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# Characterisation of highly irradiated 3D trench silicon pixel sensors for 4D tracking with 10 ps timing accuracy

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3D trench silicon pixel sensors, recently developed by the TimeSPOT collaboration, have shown excellent performance in terms of spatial resolution, timing precision and detection efficiency. The combination of these three features make them one of the best candidate for inner tracking detectors operating in high luminosity hadron colliders experiments. This article presents systematic characterisations of these devices made with minimum ionising particles on irradiated sensors with neutrons up to  $2.5 \cdot 10^{16}$  1 MeV n<sub>eq</sub> cm<sup>-2</sup>. The results show that 3D trench pixels have extremely high resistance to radiation. The measured time resolution and the detection efficiency of irradiated sensors match those of non-irradiated ones if a slightly higher bias voltage, few tens of Volts, is applied to the pixels. As of today, 3D trench pixels are the only sensors capable of achieving 10 ps time resolution after being irradiated at extremely high fluences, extending by far the capabilities of future tracking systems of HEP experiments operating under extreme conditions.

#### KEYWORDS

particle tracking detectors, solid-state detectors, timing detectors, radiation-hard detectors, radiation hardness, high time resolution, high luminosity

# **1** Introduction

Future high energy physics experiments at hadronic colliders foresee operating at an increased instantaneous luminosity, significantly larger ( $\mathcal{O}(10)$ ) than the current values. As a consequence the occupancy of tracking detectors will increase, affecting the tracking and the vertex reconstruction efficiencies and the overall event reconstruction capability.

Furthermore, the detectors will face higher levels of radiation fluences, requiring superior radiation hardness to guarantee reliable performance throughout the duration of data-taking period.



The ATLAS, CMS and LHCb collaborations have shown that it is possible to recover the vertex reconstruction performance of the current detectors by measuring the time of the tracks with an accuracy of the order  $\mathcal{O}(10 \text{ ps})$  [1–3]. In this context, the TimeSPOT<sup>1</sup> R&D project [4] has developed a new family of 3D silicon pixel sensors with 55  $\mu$ m pitch that have shown excellent time resolution for Minimum Ionising Particle (MIP) detection, which make them an ideal candidate to be used in high occupancy tracking detectors operating close to the particles interaction points.

To test the radiation hardness of these sensors an irradiation campaign has been carried out exposing several TimeSPOT test structures to neutrons with a total irradiation fluence up to  $2.5 \cdot 10^{16}$  1 MeV n<sub>eq</sub> cm<sup>-2</sup>. The irradiated 3D trench sensors have been then systematically characterised with 180 GeV/c charged hadrons from the CERN SPS/H8 beamline. The timing performance as well as the detection efficiency of the irradiated sensors have been evaluated and compared to those of the non-irradiated pixels.

# 2 Materials and methods

### 2.1 Irradiated 3D trench sensors

The TimeSPOT sensors (Figure 1), referred also as 3D trench sensors, are silicon 3D pixels with the electrodes shaped in order to maximise the timing performance for MIP detection [5]. The TimeSPOT pixel has a dimension of  $55 \times 55 \,\mu\text{m}^2$  with an active thickness of  $150 \,\mu\text{m}$ . The sensor is characterised by trench shaped electrodes: two external ohmic-wall electrodes, which extend over the

entire row of a pixel matrix to provide the bias voltage to each pixel and a collecting electrode  $40 \times 5 \,\mu\text{m}^2$  with a thickness of  $135 \,\mu\text{m}$  [6].

The TimeSPOT sensors have been fully characterised in two beam test campaigns featuring both excellent time resolution, of about 10 ps, and detection efficiency of 99% [5, 7-10].

In general, 3D sensors are expected to have an higher radiation hardness with respect to planar silicon sensors due to the shorter inter-electrodes distance which guarantees a lower trapping probability of charge carriers during their motion towards the electrodes [11]. Radiation tolerance of 3D columnar sensors has been already proved up to  $3 \cdot 10^{16}$  1 MeV  $n_{eg}$  cm<sup>-2</sup> [12] for what concerns the charge carrier collection efficiency. A similar behaviour is expected also for 3D trench sensors for which however the radiation hardness was never tested. To verify the radiation tolerance of the TimeSPOT sensors an irradiation campaign has been carried out at the TRIGA Mark II Reactor<sup>2</sup> at the Jožef Stefan Institute in Ljubljana, Slovenia. Several test structures of the second TimeSPOT batch, produced in 2020 by Fondazione Bruno Kessler (FBK) in Trento, Italy, were irradiated at two different neutron irradiation fluences of  $1.0\cdot 10^{16}\,1$  MeV  $n_{eq}\,cm^{-2}$  and  $2.5\cdot 10^{16}\,1$  MeV  $n_{eq}$  cm<sup>-2</sup>. After irradiation the test structures were stored in a temperature controlled box at about -20 °C, to prevent radiation damage mitigation due to the reverse annealing. Two types of test structures, the single pixel and the triple-strip (Figure 2), were then wire bonded to the readout boards and stored back at -20 °C. In the single pixel structure the innermost of seven adjacent pixels in a row is readout, while the two adjacent pixels are connected to ground, to guarantee the proper electric field configuration. This structure has been used for the charge collection (amplitude) and timing

<sup>1</sup> https://web.infn.it/timespot/index.php

<sup>2</sup> https://www.tuwien.at/en/trigacenter/trigareactor



FIGURE 2

Pictures of some of the 3D pixel test structures used in this work. For each structure the active area is outlined in red. (A) Single pixel sensor; (B) triple strip sensor (30 pixels located in three adjacent rows).



**FIGURE 3** The TimeSPOT single channel board [14]. A magnification of the test structure wire bonded to the board input is also reported.

characterizations. The triple-strip structure consists of 30 pixels placed in three adjacent rows which have the readout electrodes shorted together. This device, having a larger active area, has been used for the efficiency measurements.

The test structures are readout through the TimeSPOT single channel front-end electronic boards detailed in [13]. This circuit features a signal-to-noise ratio (S/N) of about 20 and an electronic jitter below 7 ps at 2 fC input charge. The power consumption is about 70 mW/channel. The board directly hosts the test structures and provides the bias voltage to the sensors by means of a conductive tape placed on the back side of the test structures (Figure 3).

## 2.2 Test beam setup

Irradiated TimeSPOT sensors have been tested in May 2022 at the CERN SPS/H8 beamline with a 180 GeV/c positive hadrons beam. The beam, on average, consists of  $10^6$  particles which are extracted every 30 s in a 4 seconds-long spill and focused on an approximately circular spot of 8 mm radius measured at the setup location. The setup (Figure 4) consists of a total of four detectors precisely aligned with respect to the beam line: two 3D trench sensors and two microchannel-plate photo-multipliers (MCP-PMTs). The two 3D trench silicon sensors, along with their respective front-end electronics boards, have been securely installed within an electromagnetically shielded light-tight enclosure.

One of the sensors, the irradiated structure, is operated at low temperature. To do that the sensor is placed in a cold, temperature-monitored environment, which is isolated from the other part of the light-tight enclosure by means of a polystyrene box. Dry ice has been regularly filled inside two 3D printed containers positioned inside the polystyrene box. This relatively simple system effectively enables the operation of the irradiated sensors within a temperature range  $[-40, -20]^{\circ}$ C for about 10 h. The sensor positioned within the polystyrene box is fixed to a stationary mount which only allows to manually rotate the sensor around the vertical direction to measure the sensor response for non-normal beam incidence. The other (non-irradiated) sensor is mounted on a movable holder outside of the thermally isolated box. Such holder is driven by two closed-loop piezoelectric linear stages, allowing movements of up to 16 mm with a position accuracy of 10 nm in the transverse direction relative to the beam line. The second sensor, mounted on the piezoelectric holder, provides a signal for the acquisition trigger allowing unbiased measurements of the device under test (DUT). Two microchannel-plate photomultipliers are employed for the precise measurement (3-4 ps) of the time of arrival (TOA) of each particle. Both photomultipliers feature a 18 mm diameter and a 5.5 mm thick quartz input window.

The signals from the silicon sensors and the two MCP-PMTs are processed and recorded by a 8 GHz analog bandwidth 20 GSa/s 4 channels digital oscilloscope. The connection of the sensors and the MCP-PMTs to the oscilloscope is made through low-loss RF cables. The typical trigger requires the coincidence in a time window of 20 ns of two signals, one of the MCP-PMTs and one of the non-irradiated 3D trench sensor, with a minimum amplitude. This trigger condition avoids bias on the DUT measurements, which is not involved in the acquisition



#### FIGURE 4

Pictures of the beam test setup. (A) Front view of the light-tight box containing the two silicon sensors aligned; the irradiated sensor is placed between two dry ice containers used for the cooling. (B) Back view of the two silicon sensors in which the temperature sensor (PT100) is visible. (C) Global view of the setup showing the black light-tight box containing the polystirene box, where the irradiated sensor is placed, and the two MCP-PMTs providing the time of arrival of the particles. The oscilloscope used for the signal acquisition is also visible in the picture.



#### FIGURE 5

A typical event in which a charged particle crosses the four detectors. The signals from the two silicon sensors are shown in yellow (non-irradiated) and green (irradiated), while the signals of the two MCP-PMTs are reported in orange and blue.



requirements. A representative event recorded with the described setup is reported in Figure 5.

## 2.3 Analysis method

The analysis method strictly follows what discussed in previous publications [5, 8]. Each waveform is processed to determine the signal amplitude and time of arrival (TOA).

A signal is identified by looking for a minimum in the waveform in a time window where signals are expected to contribute. The signal amplitude is determined by subtracting the minimum value found from the corresponding extrapolated value of the baseline, which is determined event-by-event by fitting the first samples of the waveform.

Among the different methods already detailed in [5, 8], the reference and the *spline* are used in this article. The reference method embodies the amplitude and rise time compensation (ARC)



FIGURE 7

(A) Detection efficiency versus the bias voltage for the irradiated and non-irradiated 3D trench sensors operated at normal beam incidence condition. (B) Detection efficiency versus tilt angle for the irradiated ( $0^\circ$ ,  $5^\circ$ ,  $20^\circ$ ) and non-irradiated ( $0^\circ$ ,  $5^\circ$ ,  $10^\circ$ ,  $20^\circ$ ) 3D trench sensor operated at a bias voltages of -130 V and -100 V, respectively.





method [15]. In this approach, the signal waveform undergoes a delay of approximately half its rise time, followed by subtracting the delayed waveform from the original. Subsequently, the TOA is determined as the time where the resulting waveform exceeds half of its maximum amplitude. On the other hand, the *spline* method interpolates the waveforms using cubic splines. The TOA is then set to the time at which the signal exceeds a specified fraction (20%) of its maximum amplitude. While the reference method has shown the best performance in term of time resolution, the *spline* method has a wider application being a refined, off-line implementation of the constant fraction discriminator.

The analysis method utilised is proved not to bias the results, which is a fundamental requirement for the measurements that are discussed in the following. In particular, it is capable to reconstruct the observables also in case of very low charge deposits, which is



Time resolution of the single pixels as a function of the bias voltage for the different irradiation fluences obtained with the (A) reference and (B) spline methods. The contribution due to the resolution of the time reference is subtracted. Irradiated pixels are operated in a temperature range [-40, -20]°C, while non-irradiated pixel is operated at room temperature.

TABLE 1 Time resolution for non-irradiated and irradiated 3D trench pixels operated at different tilt angles. The contribution due to the resolution of the time reference is subtracted. The sensors were operated at a bias voltage of -100 V.

Fluence	σ <sup>Si</sup> (reference) [ps]		$\sigma_{ m eff}^{ m Si}$ (spline) [ps]	
	0°	20°	0°	20°
Not irradiated	$12.0 \pm 0.7$	$17.5 \pm 0.9$	17.8 ± 1.0	21.7 ± 0.9
$2.5 \cdot 10^{16} \ 1 \ \text{MeV} \ n_{eq} \ \text{cm}^{-2}$	9.8 ± 0.4	13.5 ± 0.6	13.6 ± 1.1	16.9 ± 0.9

important for the characterisation of the performance of the whole detector under test that typically is not part of the triggering logic.

# **3** Results

## 3.1 Single pixel signal amplitude

Radiation generates traps in the silicon bulk that block the charge carriers for a time by far larger than the typical sensor charge collection time. As a consequence, irradiated sensors feature a reduction of both the collected charge and the signal amplitude, which can be mitigated by increasing the bias voltage applied to the sensor. To study this effect, the two 3D trench silicon pixel structures, irradiated at  $1.0 \cdot 10^{16}$  and  $2.5 \cdot 10^{16}$  1 MeV n<sub>eq</sub> cm<sup>-2</sup>, are operated at different bias voltages (-20 V, -40 V, -100 V, -120 V and -150 V).

Figure 6 shows the resulting amplitude distributions which are compared to the measurements of the non-irradiated pixel from [8].

Both irradiated sensors feature a clear dependency of the most probable amplitude value on the bias voltage. Notice that since the amplifier boards used for the two irradiated sensors have slightly different gains, an exact quantitative comparison between the most probable amplitude values cannot be done. Nevertheless, it is evident that both irradiated sensors recover a pulse height distribution similar to those of the non-irradiated one when a bias voltage slightly lower than -100 V is applied. Moreover, the fact that the sensor irradiated at  $2.5 \cdot 10^{16}$  1 MeV n<sub>eq</sub> cm<sup>-2</sup> still shows a regular, Landau-shaped distribution suggests that the radiation hardness limit for such sensors could be even higher.

These measurements clearly prove the radiation hardness of such devices and the unchanged performance with respect to non-irradiated sensors indicate that they can be operated even at fluences higher than  $2.5 \cdot 10^{16}$  1 MeV  $n_{eq}$  cm<sup>-2</sup>.

# 3.2 Detection efficiency

The detection efficiency of irradiated sensors is studied in different conditions and compared to the published results of the non-irradiated sensors [8]. The measurement is performed using a triple-strip structure as DUT, and by triggering on tracks crossing a 3D trench single pixel precisely centred, projectively along the beam line, in front of the DUT. A particle detected by the triggering pixel will also cross the triple strip producing a signal unless it ends up inside the trenches or gets trapped in the bulk. The two MCP-PMTs are also readout to obtain a precise charged particle time reference.

The efficiency is given by the ratio  $\eta = N_{\rm ts}/N_{\rm tr}$ , where  $N_{\rm ts}$  is the number of events detected by the triple-strip sensor and  $N_{\rm tr}$  is the

number of triggers. Additional requirements on the TOA of the signal in the DUT with respect to the time reference given by the two MCP-PMTs as well as corrections due to small beam divergence are applied following the same methods reported in [8].

Different measurements are performed using the two irradiated 3D trench triple-strip structures: by varying either the bias voltages or the tilt-angle with respect to the beam line. Figure 7A shows the results for the triple-strip structure irradiated at  $1.0 \cdot 10^{16}$  1 MeV  $n_{eq}$  cm<sup>-2</sup> as a function of the bias voltage (green points). For  $V_{bias} \leq -60$  V the detection efficiency is compatible with the value measured for the non-irradiated sensor (black points). This result is consistent with the recovery of signal amplitude studied in the previous section.

Since the trenches are a non-active volume, 3D trench sensors reach the full detection efficiency only if operated tilted with respect to particle flight directions. Therefore, the detection efficiency of the triple-strip structure irradiated at  $2.5 \cdot 10^{16}$  1 MeV  $n_{eq}$  cm<sup>-2</sup> is evaluated as a function of the tilt angle and compared to the results obtained for the non-irradiated sensor [8]. The results reported in Figure 7B show that the efficiency of the irradiated sensor (red points) is compatible with that of a non-irradiated sensor (black points) reaching about 99% detection efficiency when tilted at angles larger than 10°. The reported measurements are obtained for a bias voltage of -130 V and -100 V for the irradiated and non-irradiated sensors, respectively.

These results show that the irradiated sensors have the same behaviour of non-irradiated ones by providing higher absolute bias voltage, in agreement to what is observed for the amplitude measurements.

#### 3.3 Time resolution

Another key aspect to characterise the detection capabilities of the 3D trench silicon sensor in high radiation environment is given by its timing performance after irradiation.

In this section we report the timing performance of the two 3D trench single pixel structures, irradiated at  $1.0 \cdot 10^{16}$  and  $2.5 \cdot 10^{16}$  1 MeV  $n_{eq}$  cm<sup>-2</sup>, operated at different bias voltages ([-150, -20]V) at two different tilt angles with respect to the beam line (0° and 20°). For these measurements the acquisition was triggered by the coincidence of signals detected on a non-irradiated silicon strip sensor<sup>3</sup> placed upstream the pixel and two MCP-PMTs located downstream. The pixel was carefully aligned in the beam transverse plane with respect to the strip in order to maximise the rate of detected coincidences.

The time of arrival of the tested sensor ( $t_{Si}$ ) is determined by means of the reference and *spline* methods discussed in Section 2.3. The timing performance of the sensor are measured with respect to the time reference given by the mean of the two MCP-PMTs TOAs ( $\langle t_{\text{MCP-PMTs}} \rangle$ ), which has a precision in the range of 3 ps-4 ps.

Figure 8 shows the distributions of  $t_{Si} - \langle t_{MCP-PMTs} \rangle$  for the two irradiated pixels operating at bias voltage of -150 V. Minimal requirements on the amplitudes of the MCP-PMTs and on the triggering sensor are applied to reject badly reconstructed events. The two distributions are very similar to each other and have the same characteristics as the distribution of non-irradiated pixels (see Figure 9

of Ref; [8]), featuring a narrow peak and a small tail due to late signals. The distribution is fit by the sum of two Gaussian functions describing the signal, and a constant modelling the background. The time resolution of the single pixel,  $\sigma_{\rm eff}^{\rm Si}$  is obtained by combining the parameters of the two Gaussian functions and by subtracting the uncorrelated contribution due to the time reference jitter.

Figure 9 shows the resulting values of the time resolution as a function of the sensor bias voltage for the two irradiated pixels and, as a reference, for the non-irradiated one, operated respectively at low (in the range [-40, -20]°C) and room temperature. Both results obtained with the reference and the spline methods are reported. The irradiated sensors have very similar timing performance to the non-irradiated one. The time resolution obtained with the reference method at a bias voltage of -100 V are 9.8  $\pm$  0.4 ps for the 2.5  $\cdot$  10<sup>16</sup> 1 MeV n<sub>eq</sub> cm<sup>-2</sup> pixel and 11.3  $\pm$  0.7 ps for the 1.0  $\cdot$  10  $^{16}$  1 MeV  $n_{eq}$  cm  $^{-2}$  pixel to be compared to the 12.0  $\pm$  0.7 ps of the non-irradiated sensor operated at the same bias voltage. These results clearly show that the time resolution of 3D trench sensor is not affected by the radiation damage up to the considered fluences. The slightly better timing performance of the irradiated sensors with respect to the non-irradiated ones are related to the different performance of the front-end electronic boards used for the different sensors and to the different temperature condition. The measurements also show that the time resolution of both the irradiated and non-irradiated sensors does not improve significantly below a certain value of bias voltage, which is about -25, -60 and -80 V for the non-irradiated, for the  $1.0 \cdot 10^{16}$  and  $2.5 \cdot 10^{16}$  1 MeV  $n_{eq}$  cm<sup>-2</sup> sensors, respectively. This effect can be explained by the amplitude decrease of the irradiated sensors when operated at lower absolute bias voltages as discussed in Sec. 3.1.

Given the convenience in operating the 3D trench sensors tilted with respect to the particle flight directions in terms of detection efficiency (Sec. 3.2), the time resolution is also evaluated for a tilt angle of 20°. Table 1 show the results for both the non-irradiated and the irradiated  $(2.5 \cdot 10^{16} 1 \text{ MeV } n_{eq} \text{ cm}^{-2})$  pixels. Both pixels feature a similar worsening of the time resolution of maximum +40% when operated at 20°. Such behaviour is due to the larger fraction of tracks that crosses adjacent pixels leading to charge sharing between two pixels and thus in a reduced signal amplitude of the pixel under test (see details in [8–10]). As shown in Ref. [8], in such cases the timing performance recovers that of the single pixel at normal incidence by combining the information of the two hit pixels with a clustering algorithm.

## 4 Discussion

The innovative 3D trench silicon pixel sensors, developed by the TimeSPOT collaboration, have demonstrated exceptional performance in detecting high-energy charged particles, exhibiting high detection efficiency and precise timing resolution, as highlighted in previous works [5, 8]. This study explores their radiation tolerance, showcasing their resistance up to fluences of  $2.5 \cdot 10^{16}$  1 MeV  $n_{eq}/cm^2$  through measurements conducted during a beam test campaign on irradiated sensors.

Due to radiation damage, the detection efficiency of 3D trench sensor diminishes for  $V_{\text{bias}} \ge -40$  V. However, this reduction is offset by restoring the full efficiency when applying a slightly larger reverse bias voltage. This observation is supported by the analysis of

<sup>3</sup> A row of ten 3D trench pixels with the readout electrodes shorted together.

amplitude distributions at various bias voltages, revealing a distinctive peaking shape similar to that of the non-irradiated pixel for bias voltages slightly lower than -100 V.

Crucially, the timing performance of the irradiated sensors match those of their non-irradiated counterparts, underscoring the negligible impact of radiation damage on these sensors.

As of today, 3D trench-type silicon detectors stand out as the fastest radiation-hard charged-particles pixel sensors available. Their promising features position them as a compelling solution for future upgrades in the tracking systems of numerous High Energy Physics experiments that operate at exceptionally high instantaneous luminosity.

## Data availability statement

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

## Author contributions

FB: Formal Analysis, Investigation, Writing-review and editing. AC: Conceptualization, Investigation, Methodology, Writing-review and editing. GMC: Conceptualization, Writing-review and editing. GFDB: Conceptualization, Writing-review and editing. MG: Conceptualization, Formal Analysis, Investigation, Methodology, Writing-review and editing. LLD: Investigation, Methodology, Writing-review and editing. ALai: Conceptualization, Investigation, Writing-review and editing. ALam: Conceptualization, Formal Analysis, Investigation, Methodology, Writing-original draft, Writing-review and editing. ALam: Conceptualization, Formal Analysis, Investigation, Methodology, Writing-original draft, Writing-review and editing. ALoi: Conceptualization, Investigation, Writing-review and editing. GS: Investigation, Formal Analysis, Writing-review and editing. SV: Investigation, Formal Analysis, Writing-original draft, Writing-review and editing.

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# **Conflict of interest**

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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