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Existence of a ground-state solution for a quasilinear Schrödinger system

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In this paper, we consider the following quasilinear Schrödinger system.

$$\begin{cases} -\Delta u + u + \frac{k}{2} [\Delta |u|^2] u = \frac{2\alpha}{\alpha + \beta} |u|^{\alpha-2} u |v|^\beta, & x \in \mathbb{R}^N, \\ -\Delta v + v + \frac{k}{2} [\Delta |v|^2] v = \frac{2\beta}{\alpha + \beta} |u|^\alpha |v|^{\beta-2} v, & x \in \mathbb{R}^N, \end{cases}$$

where $k < 0$ is a real constant, $\alpha > 1, \beta > 1$, and $\alpha + \beta < 2^*$. We take advantage of the critical point theorem developed by Jeanjean (Proc. R. Soc. Edinburgh Sect A., 1999, 129: 787–809) and combine it with Pohožaev identity to obtain the existence of a ground-state solution, which is the non-trivial solution with the least possible energy.

KEYWORDS

quasilinear Schrödinger system, Pohožaev identity, ground-state solution, critical point theorem, Lebesgue dominated convergence theorem

1 Introduction

This article is concerned with the following quasilinear Schrödinger system:

$$\begin{cases} -\Delta u + u + \frac{k}{2} [\Delta |u|^2] u = \frac{2\alpha}{\alpha + \beta} |u|^{\alpha-2} u |v|^\beta, & x \in \mathbb{R}^N, \\ -\Delta v + v + \frac{k}{2} [\Delta |v|^2] v = \frac{2\beta}{\alpha + \beta} |u|^\alpha |v|^{\beta-2} v, & x \in \mathbb{R}^N, \end{cases} \tag{1.1}$$

where $k < 0$ is a real constant.

Many scholars have made significant contributions to the study of the quasilinear Schrödinger system. Wang and Huang proved the existence of ground-state solutions for a class of systems by establishing a suitable Nehari–Pohožaev-type constraint set and considering related minimization problems in [2]. The existence of infinitely many solutions was established for the quasilinear Schrödinger system by the symmetric Mountain Pass Theorem; see [3]. The existence of positive solutions was obtained by using the monotonicity trick and Morse iteration in [4]. Chen and Zhang proved the existence of ground-state solutions by minimization under a convenient constraint and concentration compactness lemma in [5].

The quasilinear Schrödinger system (1.1) is in part motivated by the following quasilinear Schrödinger equation:

$$i\epsilon\partial z = -\epsilon\Delta z + W(x)z - l(|z|^2)z - k\epsilon\Delta h(|z|^2)h'(|z|^2)z, \quad \text{for } x \in \mathbb{R}^N, N > 2, \quad (1.2)$$

where $W(x)$ is a given potential, k is a real constant, and l and h are real functions that are essentially pure power forms. The quasilinear Schrödinger Equation 1.2 describes several physical phenomena with different h ; see [6–8].

Let the case $h(s) = s, l(s) = \mu s^{\frac{p-1}{2}}$ and $k > 0$. Setting $z(t, x) = \exp(-iFt)u(x)$, one can obtain a corresponding equation of elliptic type which has the formal variational structure:

$$\epsilon\Delta u + V(x)u - \epsilon k(\Delta(|u|^2))u = \mu|u|^{p-1}u, \quad u > 0 \quad x \in \mathbb{R}^N, N > 2, \quad (1.3)$$

where $V(x) = W(x) - F$ is the new potential function. The problem (1.3) has been studied by many academics. In [9], the existence results of multiple solutions were studied via dual approach techniques and variational methods when $k > 0$ was small enough. The existence of soliton solutions was established by a minimization argument; see [10]. The Mountain Pass Theorem is combined with the principle of symmetric criticality to establish the multiplicity of solutions in [11]. In [12], the author proved the existence of soliton solutions via making a change in variables and creating a suitable Orlicz space. The minimax principles for lower semicontinuous functionals were used to find solutions in [13].

In [14], the authors used the method developed by [1, 15] to divide the energy functional into two parts and established the existence of ground-state solutions for a type of quasilinear Schrödinger equation like 1.3. Inspired by [14], we try to find the existence of ground-state solutions for system 1.1. This achievement can enrich the relatively few existing results about this system.

The main result of this paper is the following:

Theorem 1.1. *When $k < 0, \alpha > 1, \beta > 1$, and $\alpha + \beta < 2^*$, then (1.1) has a ground-state solution.*

This paper is organized as follows. In Section 2, preparation work is completed. In Section 3, we reformulate this problem and prove Theorem 1.1. In this paper, C is defined as different constants.

2 Reformulation of the problem and preliminaries

First, we explain that $L^q(\mathbb{R}^N)$ denotes the Lebesgue space with the norm

$$\|u\|_p = \left(\int_{\mathbb{R}^N} |u|^p dx \right)^{\frac{1}{p}},$$

where $1 \leq p < \infty$. $L^q = L^q(\mathbb{R}^N) \times L^q(\mathbb{R}^N)$ with the norm

$$\|(u, v)\|_p = \left(\int_{\mathbb{R}^N} |u|^p dx \right)^{\frac{1}{p}} + \left(\int_{\mathbb{R}^N} |v|^p dx \right)^{\frac{1}{p}},$$

where $1 \leq p < \infty$.

$$H^1 = \{(u, v): u, v \in L^2(\mathbb{R}^N), \nabla u, \nabla v \in L^2(\mathbb{R}^N)\}$$

with norms

$$\begin{aligned} \|(u, v)\| &= \|u\| + \|v\| \\ &= \left(\int_{\mathbb{R}^N} (|\nabla u|^2 + u^2) dx \right)^{\frac{1}{2}} + \left(\int_{\mathbb{R}^N} (|\nabla v|^2 + v^2) dx \right)^{\frac{1}{2}} \end{aligned}$$

and

$$\|(u, v)\|^2 = \|u\|^2 + \|v\|^2.$$

The embedding $H^1 \hookrightarrow L^q$ is continuous and compact for $q \in (2, 2^*)$.

In (1.1), the Euler–Lagrange functional associated with Equation 1.1 is given by

$$I(u, v) = \frac{1}{2} \int_{\mathbb{R}^N} (1 - ku^2)|\nabla u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} |u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} (1 - kv^2)|\nabla v|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} |v|^2 dx - \frac{2}{\alpha + \beta} \int_{\mathbb{R}^N} |u|^\alpha |v|^\beta dx.$$

For (u, v) , constructing the variable like [16, 17], we have

$$\begin{aligned} dz &= \sqrt{-k}\sqrt{1 - ku^2}du, \quad z = h(u) \\ &= \frac{1}{2}\sqrt{-k}u\sqrt{1 - ku^2} + \frac{1}{2}\ln(\sqrt{-k}u + \sqrt{1 - ku^2}), \\ dw &= \sqrt{-k}\sqrt{1 - kv^2}dv, \quad w = h(v) \\ &= \frac{1}{2}\sqrt{-k}v\sqrt{1 - kv^2} + \frac{1}{2}\ln(\sqrt{-k}v + \sqrt{1 - kv^2}). \end{aligned}$$

Since h is strictly monotone, it has a well-defined inverse function f and $u = f(z), v = f(w)$. Note that

$$h(u) \sim \begin{cases} \sqrt{-k}u, & |u| \ll \sqrt{\frac{1}{-k}} \\ -\frac{k}{2}u|u|, & |u| \gg \sqrt{\frac{1}{-k}} \end{cases}, \quad h'(u) = \sqrt{-k}\sqrt{1 - ku^2}$$

and

$$\begin{aligned} f(z) &\sim \begin{cases} \frac{1}{\sqrt{-k}}z, & |z| \ll \sqrt{\frac{1}{-k}} \\ \sqrt{\frac{2}{-k|z|}}z, & |z| \gg \sqrt{\frac{1}{-k}} \end{cases}, \\ f'(z) &= \frac{1}{h'(u)} = \frac{1}{\sqrt{-k}\sqrt{1 - kv^2}} = \frac{1}{\sqrt{-k}\sqrt{1 - kf(z)^2}}. \end{aligned}$$

Similarly, the same operation holds true for $v = f(w)$.

Using the variable, (1.1) will become

$$\begin{cases} -\frac{1}{k}\Delta z + f(z)f'(z) = \frac{2\alpha}{\alpha + \beta}|f(z)|^{\alpha-2}f(z)|f(w)|^\beta, & x \in \mathbb{R}^N, \\ -\frac{1}{k}\Delta w + f(w)f'(w) = \frac{2\beta}{\alpha + \beta}|f(z)|^\alpha|f(w)|^{\beta-2}f(w), & x \in \mathbb{R}^N, \end{cases} \quad (2.1)$$

where $f: [0, \infty) \rightarrow \mathbb{R}$ and

$$f' = \frac{1}{\sqrt{-k}\sqrt{1 - kf^2}}$$

on $[0, \infty), f(0) = 0$, and $f(-t) = f(t)$ on $[0, \infty)$. From the above facts, if (z, w) is a weak solution for (2.1), then $(u, v) = (f(z), f(w))$ is a

weak solution for (1.1). The energy functional $I(u, v)$ reduces to the following functional:

$$\begin{aligned} \phi(z, w) = & \frac{1}{2} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-k}} |\nabla z|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} f^2(z) dx \\ & + \frac{1}{2} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-k}} |\nabla w|^2 dx + \frac{1}{2} \int_{\mathbb{R}^N} f^2(w) dx - \frac{2}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z)|^\alpha |f(w)|^\beta dx. \end{aligned} \tag{2.2}$$

There are some properties of $f: \mathbb{R} \rightarrow \mathbb{R}$ as follows, which are proved in [16, 17].

Lemma 2.1. *The function $f(t)$ and its derivative satisfy the following properties:*

- (i) $\frac{f(t)}{t} \rightarrow 1$ as $t \rightarrow 0$;
- (ii) $f(t) \leq |t|$ for any $t \in \mathbb{R}$;
- (iii) $f(t) \leq 2^{\frac{1}{2}} \sqrt{|t|}$ for all $t \in \mathbb{R}$;
- (iv) $\frac{f^2(t)}{2} \leq t f(t) f'(t) \leq f^2(t)$ for all $t \in \mathbb{R}$;
- (v) there exists a positive constant C such that

$$|f(t)| \geq \begin{cases} C|t|, & \text{if } t \leq 1, \\ C|t|^{\frac{1}{2}}, & \text{if } t > 1; \end{cases}$$

- (vi) $|f(t)f'(t)| \leq \frac{1}{\sqrt{2}}$ for all $t \in \mathbb{R}$.

3 Proof of theorem 1.1

In this section, we will complete the proof of Theorem 1.1. First, we will recall the critical point theorem in [1], which is crucial for proving Theorem 1.1.

Theorem 3.1. *Let $(X, \|\cdot, \cdot\|)$ be a Banach space and $L \subset \mathbb{R}_+$ an interval. Consider the following family of C^1 -functionals on X :*

$$\Phi_\lambda(z, w) = A(z, w) + \lambda B(z, w), \quad \lambda \in L,$$

with B being non-negative and either $A(z, w) \rightarrow +\infty$ or $B(z, w) \rightarrow +\infty$ as $\|(z, w)\| \rightarrow \infty$. Assume that there are two points $(z_1, w_1), (z_2, w_2) \subset X$ such that

$$\begin{aligned} c_\lambda = & \inf_{\gamma \in \Gamma_\lambda} \max_{(t_1, t_2) \in [0, 1] \times [0, 1]} \Phi_\lambda(\gamma(t_1, t_2)) \\ & > \max\{\Phi_\lambda(z_1, w_1), \Phi_\lambda(z_2, w_2)\} \quad \text{for all } \lambda \in L, \end{aligned}$$

where $\Gamma_\lambda = \{\gamma \in C([0, 1] \times [0, 1], X): \gamma(0, 0) = (z_1, w_1), \gamma(1, 1) = (z_2, w_2)\}$. Then, for almost every $\lambda \in L$, there is a sequence $\{(z_n, w_n)\} \subset X$ such that

- (i) (z_n, w_n) is bounded;
- (ii) $\Phi_\lambda(z, w) \rightarrow c_\lambda$;
- (iii) $\Phi'_\lambda(z_n, w_n) \rightarrow 0$ in the dual X^{-1} of X .

Moreover, the map $\lambda \rightarrow c_\lambda$ is non-increasing and continuous from the left.

Let $\lambda \in L$ be an arbitrary but fixed value where c'_λ exists, where c'_λ is the derivative of c_λ with respect to λ . Let $\{\lambda_n\} \subset L$ be a strictly

increasing sequence such that $\lambda_n \rightarrow \lambda$. To prove Theorem 3.1, we will show the following lemmas:

Lemma 3.1. *There exists a sequence of path $\{\gamma_n\} \subset \Gamma$ and $K = K(c'_\lambda) > 0$ such that*

- (i) $\|\gamma_n(t_1, t_2)\| \leq K$ if $\gamma_n(t_1, t_2)$ satisfies

$$\Phi_\lambda(\gamma_n(t_1, t_2)) \geq c_\lambda - (\lambda - \lambda_n); \tag{3.1}$$
- (ii) $\max_{(t_1, t_2) \in [0, 1] \times [0, 1]} \Phi_\lambda(\gamma_n(t_1, t_2)) \leq c_\lambda + (-c'_\lambda + 2)(\lambda - \lambda_n)$.

Proof. The proof is standard; see [1].

Lemma 3.1. means that there exists a sequence of paths $\{\gamma_n\} \subset \Gamma$ such that

$$\max_{(t_1, t_2) \in [0, 1] \times [0, 1]} \Phi_\lambda(\gamma_n(t_1, t_2)) \rightarrow c_\lambda,$$

for all $n \in \mathbb{N}$ sufficiently large; starting from a level strictly below c_λ , all the “top” of the path is contained in the ball centered at the origin of fixed radius $K = K(c'_\lambda) > 0$. Now, for $\alpha > 0$, we define

$$F_\alpha = \{(z, w) \in X: \|(z, w)\| \leq K + 1 \text{ and } |\Phi_\lambda(z, w) - c_\lambda| \leq \alpha\},$$

where K is given in lemma 3.1.

Lemma 3.2. *For all $\alpha > 0$,*

$$\inf\{\|\Phi'_\lambda(z, w)\|: (z, w) \in F_\alpha\} = 0. \tag{3.2}$$

Proof. We assume that (3.2) does not hold. Then, there exists $\alpha > 0$ such that for any $(z, w) \in F_\alpha$, we obtain

$$\|\Phi'_\lambda(z, w)\| \geq \alpha. \tag{3.3}$$

Without loss of generality, we can assume that

$$0 < \alpha < \frac{1}{2} [c_\lambda - \max\{\Phi_\lambda(z_1, w_1), \Phi_\lambda(z_2, w_2)\}].$$

A classical deformation argument then says that there exists $\epsilon \in [0, \alpha]$ and a homeomorphism $\eta: X \rightarrow X$ such that

$$\eta(u) = u, \quad \text{if } |\Phi_\lambda(z, w) - c_\lambda| \geq \alpha, \tag{3.4}$$

$$\Phi_\lambda(\eta(z, w)) \leq \Phi_\lambda(z, w), \quad \forall (z, w) \in X, \tag{3.5}$$

$$\Phi_\lambda(\eta(z, w)) \leq c_\lambda - \epsilon, \quad \forall (z, w) \in X, \text{ satisfying } \|(z, w)\| \leq K \text{ and } \Phi_\lambda(z, w) \leq c_\lambda + \epsilon. \tag{3.6}$$

Let $\{\gamma_n\} \subset \Gamma$ be the sequence obtained in lemma 3.1. We choose and fix $m \in \mathbb{N}$ sufficiently large in order that

$$(-c'_\lambda + 2)(\lambda - \lambda_m) \leq \epsilon. \tag{3.7}$$

By lemma 3.1 and (3.4), $\eta(\gamma_m) \in \Gamma$. Now if $(z, w) = \gamma_m(t_1, t_2)$ satisfies

$$\Phi_\lambda(z, w) \leq c_\lambda - (\lambda - \lambda_m),$$

then (3.5) implies that

$$\Phi_\lambda(\eta(z, w)) \leq c_\lambda - (\lambda - \lambda_m). \tag{3.8}$$

If $(z, w) = \gamma_m(t_1, t_2)$ satisfies

$$\Phi_\lambda(z, w) > c_\lambda - (\lambda - \lambda_m),$$

by lemma 3.1 and (3.7), we obtain (z, w) such that $\|(z, w)\| \leq K$ with $\Phi_\lambda(z, w) \leq c_\lambda + \epsilon$. From (3.6), we obtain

$$\Phi_\lambda(\eta(z, w)) \leq c_\lambda - \epsilon \leq c_\lambda - (\lambda - \lambda_m). \tag{3.9}$$

Combining (3.8) with (3.9), we obtain

$$\max_{(t_1, t_2) \in [0, 1] \times [0, 1]} \Phi_\lambda(\eta(\gamma_m(t_1, t_2))) \leq c_\lambda - (\lambda - \lambda_m),$$

which contradicts the variational characterization of c_λ .

Next, we prove theorem 3.1.

Proof. Since lemma 3.2 is true, there exists a PS sequence for Φ_λ at the level $c_\lambda \in \mathbb{R}$, which is contained in the ball of radius $K + 1$ centered at the origin. Hence, this is proved.

Let $L = [\frac{1}{2}, 1]$, we define the following energy functional:

$$\begin{aligned} \Phi_\lambda(z, w) = & \frac{1}{2} \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla z|^2 + z^2 + \frac{1}{\sqrt{-k}} |\nabla w|^2 + w^2 \right) dx \\ & - \lambda \int_{\mathbb{R}^N} \left(\frac{1}{2} (z^2 - f^2(z) + w^2 - f^2(w)) + \frac{2}{\alpha + \beta} |f(z)|^\alpha |f(w)|^\beta \right) dx, \end{aligned} \tag{3.10}$$

where $\lambda \in L$. Moreover, let

$$A(z, w) = \frac{1}{2} \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla z|^2 + z^2 + \frac{1}{\sqrt{-k}} |\nabla w|^2 + w^2 \right) dx$$

and

$$B(z, w) = \lambda \int_{\mathbb{R}^N} \left(\frac{1}{2} (z^2 - f^2(z) + w^2 - f^2(w)) + \frac{2}{\alpha + \beta} |f(z)|^\alpha |f(w)|^\beta \right) dx.$$

Letting $\|(z, w)\| \rightarrow +\infty$, then $A(z, w) \rightarrow +\infty$ and $B(z, w) \geq 0$.

By a standard argument in [18, 19], we have the following Pohožaev-type identity:

Lemma 3.3. *If $(z, w) \in H^1$ is a critical point of (3.10), then (z, w) satisfies $P_\lambda(z, w) = 0$, where*

$$\begin{aligned} P_\lambda(z, w) = & \frac{N-2}{2} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-k}} (|\nabla z|^2 + |\nabla w|^2) dx \\ & + \frac{N}{2} \int_{\mathbb{R}^N} (f^2(z) + f^2(w)) dx - \frac{2N\lambda}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z)|^\alpha |f(w)|^\beta dx. \end{aligned} \tag{3.11}$$

Similar to [9], we obtain the following lemma:

Lemma 3.4. $\Phi_\lambda(z, w)$ meet the conditions as follows:

(i) *there exists $(z, w) \in H^1 \setminus \{(0, 0)\}$ such that $\Phi_\lambda(z, w) < 0$ for all $\lambda \in L$;*

(ii) *for c_λ , we obtain*

$$c_\lambda = \inf_{\gamma \in \Gamma} \max_{(t_1, t_2) \in [0, 1] \times [0, 1]} \Phi_\lambda(\gamma(t_1, t_2)) > \max\{\Phi_\lambda(0, 0), \Phi_\lambda(z, w)\},$$

for all $\lambda \in L$, where

$$\Gamma = \{\gamma \in C([0, 1] \times [0, 1], H^1) : \gamma(0, 0) = (0, 0), \gamma(1, 1) = (z, w)\}.$$

Proof. (i) Let $(z, w) \in H^1 \setminus \{(0, 0)\}$ be fixed. For any $\lambda \in L = [\frac{1}{2}, 1]$, we obtain

$$\begin{aligned} \Phi_\lambda(z, w) & \leq \Phi_{\frac{1}{2}}(z, w) \\ & = \frac{1}{2} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-k}} (|\nabla z|^2 + |\nabla w|^2) dx \\ & \quad + \frac{1}{4} \int_{\mathbb{R}^N} (z^2 + f^2(z) + w^2 + f^2(w)) dx - \frac{1}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z)|^\alpha |f(w)|^\beta dx. \end{aligned}$$

As [20, 21], we consider $\phi, \varphi \in C_0^\infty(\mathbb{R})$ such that $0 \leq \phi(x) \leq 1, 0 \leq \varphi(x) \leq 1$ and

$$\phi(x) = \begin{cases} 1, & \text{if } |x| \leq 1, \\ 0, & \text{if } |x| \geq 1, \end{cases} \quad \varphi(x) = \begin{cases} 1, & \text{if } |x| \leq 1, \\ 0, & \text{if } |x| \geq 1. \end{cases}$$

By Lemma 2.1 (ii) and (v), we obtain

$$|f(t\phi)| \geq C|t\phi| \geq C\varphi(t)\phi.$$

By Lemma 2.1 (ii),

$$\begin{aligned} \Phi_\lambda(t_1\phi, t_2\varphi) & \leq \frac{1}{2} \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla t_1\phi|^2 + t_1^2\phi^2 \right) dx \\ & \quad + \frac{1}{2} \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla t_2\varphi|^2 + t_2^2\varphi^2 \right) dx \\ & \quad - \frac{1}{\alpha + \beta} \int_{\mathbb{R}^N} |f(t_1\phi)|^\alpha |f(t_2\varphi)|^\beta dx \\ & \leq \frac{t_1^2}{2} \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla\phi|^2 + \phi^2 \right) dx \\ & \quad + \frac{t_2^2}{2} \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla\varphi|^2 + \varphi^2 \right) dx \\ & \quad - \frac{C(|f(t_1)|^\alpha + |f(t_2)|^\beta)}{\alpha + \beta} \int_{\mathbb{R}^N} |\phi|^\alpha |\varphi|^\beta dx. \end{aligned}$$

It follows that $\Phi_\lambda(t_1\phi, t_2\varphi) \rightarrow -\infty$ as $(t_1, t_2) \rightarrow (+\infty, +\infty)$. Thus, there exists $(t_3, t_4) > 0$ such that $\Phi_\lambda(t_3\phi, t_4\varphi) < 0$. Thus, taking $(z, w) = (t_3\phi, t_4\varphi)$, we obtain $\Phi_\lambda(z, w) < 0$ for all $\lambda \in L$.

(ii) As [20, 22], there exists $C > 0$ and $\rho_1 > 0$ small enough such that

$$\int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla z|^2 + f^2(z) + \frac{1}{\sqrt{-k}} |\nabla w|^2 + f^2(w) \right) dx \geq C\|(z, w)\|,$$

for $\|(z, w)\| \leq \rho_1$. From Lemma 2.1 (iii) and Hölder inequality, we obtain

$$\begin{aligned} \Phi_\lambda(z, w) & \geq \frac{1}{2} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-k}} |\nabla z|^2 + f^2(z) dx + \frac{1}{2} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-k}} |\nabla w|^2 + f^2(w) dx \\ & \quad - \frac{1}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z)|^\alpha |f(w)|^\beta dx \\ & \geq C\|(z, w)\| - \frac{1}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z)|^\alpha |f(w)|^\beta dx \\ & \geq C\|(z, w)\| - C\|z^{\alpha_1}\|_p \|w^{\beta_1}\|_p, \quad \text{for all } \|(z, w)\| \leq \rho_1, \end{aligned}$$

where $\alpha_1 = \alpha$ or $\frac{\alpha}{2}, \beta_1 = \beta$ or $\frac{\beta}{2}$, and $(\frac{1}{p} + \frac{1}{p'}) = 1$. It can conclude that Φ_λ has a strict local minimum at 0, and hence, $c_\lambda > 0$.

By Theorem 3.1, it is easy to know that for every $\lambda \in [\frac{1}{2}, 1]$, there exists a bounded sequence $(z_n, w_n) \subset H^1$ such that $\Phi_\lambda(z_n, w_n) \rightarrow c_\lambda$ and $\Phi'_\lambda(z_n, w_n) \rightarrow 0$.

Lemma 3.5. *If $(z_n, w_n) \subset H^1$ is the sequence obtained above, then for almost every $\lambda \in L = [\frac{1}{2}, 1]$, there exists $(z_\lambda, w_\lambda) \in H^1 \setminus \{(0, 0)\}$ such that $\Phi_\lambda(z_\lambda, w_\lambda) \rightarrow c_\lambda$ and $\Phi'_\lambda(z_\lambda, w_\lambda) \rightarrow 0$.*

Proof. Since (z_n, w_n) is bounded in H^1 , up to a subsequence, there exists $(z_\lambda, w_\lambda) \in H^1$ such that

$$\begin{aligned} &(z_n, w_n) \rightharpoonup (z_\lambda, w_\lambda) \text{ in } H^1, \\ &(z_n, w_n) \rightarrow (z_\lambda, w_\lambda) \text{ in } L^s \text{ for all } 2 < s < 2^*, \\ &(z_n(x), w_n(x)) \rightarrow (z_\lambda(x), w_\lambda(x)) \text{ a.e. in } \mathbb{R}^N. \end{aligned}$$

Since $\Phi'_\lambda(z_n, w_n) \rightarrow 0$, by the Lebesgue dominated convergence theorem, it is easy to get $\Phi'_\lambda(z_n, w_n) \rightarrow \Phi'_\lambda(z_\lambda, w_\lambda)$, that is, $\Phi'_\lambda(z_\lambda, w_\lambda) = 0$, as shown in [23]. Similar to [22, 24, 25], there exists $C > 0$ such that

$$\int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla(z_n - z_\lambda)|^2 + (f(z_n)f'(z_n) - f(z_\lambda)f'(z_\lambda))(z_n - z_\lambda) \right) dx \geq C \|z_n - z_\lambda\|^2, \tag{3.12}$$

$$\int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla(w_n - w_\lambda)|^2 + (f(w_n)f'(w_n) - f(w_\lambda)f'(w_\lambda))(w_n - w_\lambda) \right) dx \geq C \|w_n - w_\lambda\|^2. \tag{3.13}$$

By Hölder inequality and Lemma 2.1(ii) and (iv), we deduce that

$$\begin{aligned} &\frac{2\alpha}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z_n)|^{\alpha-2} f(z_n) f'(z_n) |f(w_n)|^\beta (z_n - z_\lambda) dx \\ &+ \frac{2\beta}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z_n)^\alpha |f(w_n)|^{\beta-2} f(w_n) f'(w_n) \\ &\times (w_n - w_\lambda) dx \leq \frac{2\alpha}{\alpha + \beta} \int_{\mathbb{R}^N} |z_n|^{\alpha-1} |w_n|^\beta (z_n - z_\lambda) dx \\ &+ \frac{2\beta}{\alpha + \beta} \int_{\mathbb{R}^N} |z_n|^\alpha |w_n|^{\beta-1} (w_n - w_\lambda) dx \leq \frac{2\alpha}{\alpha + \beta} \left(\int_{\mathbb{R}^N} |z_n|^{\frac{1}{\beta}} |w_n|^{\frac{1}{\alpha}} dx \right)^{(\alpha-1)\beta} \|z_n - z_\lambda\|_{p_1} \\ &+ \frac{2\alpha}{\alpha + \beta} \left(\int_{\mathbb{R}^N} |z_n|^{\frac{1}{\beta}} |w_n|^{\frac{1}{\alpha}} dx \right)^{(\beta-1)\alpha} \|w_n - w_\lambda\|_{p_2} \rightarrow 0, \end{aligned} \tag{3.14}$$

where $p_1 = \frac{1}{(\alpha-1)\beta}$ and $p_2 = \frac{1}{(\beta-1)\alpha}$. Similarly, we obtain

$$\begin{aligned} &\frac{2\alpha}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z_\lambda)|^{\alpha-2} f(z_\lambda) f'(z_\lambda) |f(w_\lambda)|^\beta (z_n - z_\lambda) dx \\ &+ \frac{2\beta}{\alpha + \beta} \int_{\mathbb{R}^N} |f(z_\lambda)^\alpha |f(w_\lambda)|^{\beta-2} f(w_\lambda) f'(w_\lambda) (w_n - w_\lambda) dx \rightarrow 0. \end{aligned} \tag{3.15}$$

Following (3.12), (3.13), (3.14), and (3.15), we obtain

$$\begin{aligned} 0 &\leftarrow \langle \Phi'_\lambda(z_n, w_n) - \Phi'_\lambda(z_\lambda, w_\lambda), (z_n - z_\lambda, w_n - w_\lambda) \rangle \\ &= \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla(z_n - z_\lambda)|^2 + f(z_n - z_\lambda) f'(z_n - z_\lambda) (z_n - z_\lambda) \right) dx \\ &+ \int_{\mathbb{R}^N} \left(\frac{1}{\sqrt{-k}} |\nabla(w_n - w_\lambda)|^2 + f(w_n - w_\lambda) f'(w_n - w_\lambda) (w_n - w_\lambda) \right) dx \\ &- \frac{2\alpha}{\alpha + \beta} \int_{\mathbb{R}^N} [|f(z_n)|^{\alpha-2} f(z_n) f'(z_n) |f(w_n)|^\beta \\ &- |f(z_\lambda)|^{\alpha-2} f(z_\lambda) f'(z_\lambda) |f(w_\lambda)|^\beta] (z_n - z_\lambda) dx \\ &- \frac{2\beta}{\alpha + \beta} \int_{\mathbb{R}^N} [|f(z_n)^\alpha |f(w_n)|^{\beta-2} f(w_n) f'(w_n) \\ &- |f(z_\lambda)^\alpha |f(w_\lambda)|^{\beta-2} f(w_\lambda) f'(w_\lambda)] (w_n - w_\lambda) dx \\ &\geq C \|z_n - z_\lambda\|^2 + C \|w_n - w_\lambda\|^2 + o_n(1), \end{aligned} \tag{3.16}$$

which implies that $(z_n, w_n) \rightarrow (z_\lambda, w_\lambda)$ in H^1 . Thus, (z_λ, w_λ) is a non-trivial critical point of $\Phi_\lambda(z, w)$ with $\Phi_\lambda(z_\lambda, w_\lambda) = c_\lambda$.

Next, we prove Theorem 1.1.

Proof. At first, using Theorem 3.1, for arbitrary $\lambda \in L = [\frac{1}{2}, 1]$, there is a $(z_\lambda, w_\lambda) \in H^1$ such that

$$\begin{aligned} &(z_n, w_n) \rightharpoonup (z_\lambda, w_\lambda) \neq (0, 0) \text{ in } H^1, \\ &\Phi_\lambda(z_n, w_n) \rightarrow c_\lambda \text{ and } \Phi'_\lambda(z_n, w_n) \rightarrow 0. \end{aligned}$$

By Lemma 3.5, we obtain

$$\Phi_\lambda(z_\lambda, w_\lambda) \rightarrow c_\lambda \text{ and } \Phi'_\lambda(z_\lambda, w_\lambda) = 0.$$

Thus, there exists $\lambda_n \in [\frac{1}{2}, 1]$ such that

$$\begin{aligned} &\lambda_n \rightarrow 1, (z_{\lambda_n}, w_{\lambda_n}) \in H^1, \\ &\Phi'_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) = 0 \text{ and } \Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) = c_{\lambda_n}. \end{aligned}$$

Next, we prove that $\{(z_{\lambda_n}, w_{\lambda_n})\}$ is bounded in H^1 . From Lemma 3.4

$$\Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) = c_{\frac{1}{2}}, \quad \Phi'_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) = 0,$$

it follows that

$$\begin{aligned} c_{\frac{1}{2}} &\geq \Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) \\ &= \Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) - \frac{1}{N} P_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) \\ &= \frac{N-2}{2N} \int_{\mathbb{R}^N} \left(\frac{2}{N-2} \frac{1}{\sqrt{-k}} (|\nabla z_{\lambda_n}|^2 + |\nabla w_{\lambda_n}|^2) + f^2(z_{\lambda_n}) + f^2(w_{\lambda_n}) \right) dx. \end{aligned} \tag{3.17}$$

By Lemma 2.1 (v) and Sobolev inequality, it follows that

$$\int_{|z_{\lambda_n}| \leq 1} z_{\lambda_n}^2 dx \leq C \int_{\mathbb{R}^N} f^2(z_{\lambda_n}) dx, \quad \int_{|w_{\lambda_n}| \leq 1} w_{\lambda_n}^2 dx \leq C \int_{\mathbb{R}^N} f^2(w_{\lambda_n}) dx$$

and

$$\begin{aligned} &\int_{|z_{\lambda_n}| > 1} z_{\lambda_n}^2 dx \leq \int_{|z_{\lambda_n}| > 1} z_{\lambda_n}^{2^*} dx \leq C \left(\int_{\mathbb{R}^N} |\nabla z_{\lambda_n}|^2 dx \right)^{\frac{2}{2^*}}, \\ &\int_{|w_{\lambda_n}| > 1} w_{\lambda_n}^2 dx \leq \int_{|w_{\lambda_n}| > 1} w_{\lambda_n}^{2^*} dx \leq C \left(\int_{\mathbb{R}^N} |\nabla w_{\lambda_n}|^2 dx \right)^{\frac{2}{2^*}}. \end{aligned}$$

Therefore,

$$\begin{aligned} &\int_{\mathbb{R}^N} (z_{\lambda_n}^2 + w_{\lambda_n}^2) dx \\ &= \int_{|z_{\lambda_n}| \leq 1} z_{\lambda_n}^2 dx + \int_{|z_{\lambda_n}| > 1} z_{\lambda_n}^2 dx + \int_{|w_{\lambda_n}| \leq 1} w_{\lambda_n}^2 dx + \int_{|w_{\lambda_n}| > 1} w_{\lambda_n}^2 dx \\ &\leq C \int_{\mathbb{R}^N} f^2(z_{\lambda_n}) dx + C \int_{\mathbb{R}^N} f^2(w_{\lambda_n}) dx \\ &+ C \left(\int_{\mathbb{R}^N} |\nabla z_{\lambda_n}|^2 dx \right)^{\frac{2}{2^*}} + C \left(\int_{\mathbb{R}^N} |\nabla w_{\lambda_n}|^2 dx \right)^{\frac{2}{2^*}}. \end{aligned} \tag{3.18}$$

Combining (3.17) and (3.18), we infer that there exists $C > 0$ such that

$$\int_{\mathbb{R}^N} (z_{\lambda_n}^2 + w_{\lambda_n}^2) dx \leq C.$$

Thus, there exists $C > 0$ independent of n such that

$$\|(z_{\lambda_n}, w_{\lambda_n})\|^2 = \int_{\mathbb{R}^N} (|\nabla z_{\lambda_n}|^2 + z_{\lambda_n}^2) dx + \int_{\mathbb{R}^N} (|\nabla w_{\lambda_n}|^2 + w_{\lambda_n}^2) dx \leq C.$$

Next, we can assume that the limit of $\Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n})$ exists. By Theorem 3.1, we know that $\lambda \rightarrow c_\lambda$ is continuous from the left. Thus, we obtain

$$0 \leq \lim_{n \rightarrow \infty} \Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) \leq c_{\frac{1}{2}}.$$

Then, by using the fact that

$$\begin{aligned} \Phi(z_{\lambda_n}, w_{\lambda_n}) &= \Phi_{\lambda_n}(z_{\lambda_n}, w_{\lambda_n}) \\ &+ \frac{(\lambda_n - 1)}{\alpha\beta} \int_{\mathbb{R}^N} \frac{2}{\alpha + \beta} |f(z_{\lambda_n})|^\alpha |f(w_{\lambda_n})|^\beta dx \end{aligned}$$

and

$$\begin{aligned} \langle \Phi'(z_{\lambda_n}, w_{\lambda_n}), (\phi, \psi) \rangle &= \langle \Phi_{\lambda_n}'(z_{\lambda_n}, w_{\lambda_n}), (\phi, \psi) \rangle \\ &+ \frac{(\lambda_n - 1)}{\beta} \int_{\mathbb{R}^N} \frac{2}{\alpha + \beta} |f(z_{\lambda_n})|^{\alpha-1} f'(z_{\lambda_n}) \phi |f(w_{\lambda_n})|^\beta dx \\ &+ \frac{(\lambda_n - 1)}{\alpha} \int_{\mathbb{R}^N} \frac{2}{\alpha + \beta} |f(z_{\lambda_n})|^\alpha |f(w_{\lambda_n})|^{\beta-1} f'(w_{\lambda_n}) \psi dx, \end{aligned}$$

for any $\phi, \psi \in C_0^\infty(\mathbb{R}^N)$ and $\|(z_{\lambda_n}, w_{\lambda_n})\| \leq C$, it follows that

$$\lim_{n \rightarrow \infty} \Phi(z_{\lambda_n}, w_{\lambda_n}) = c_1, \quad \lim_{n \rightarrow \infty} \Phi'(z_{\lambda_n}, w_{\lambda_n}) = 0.$$

Up to a subsequence, there exists a subsequence $(z_{\lambda_n}, w_{\lambda_n})$ denoted by (z_n, w_n) and $(z_0, w_0) \in H^1$ such that $(z_n, w_n) \rightharpoonup (z_0, w_0)$ in H^1 . Using the same method as Lemma 3.5, we will obtain the existence of a non-trivial solution (z_0, w_0) for Φ and $\Phi'(z_0, w_0) = 0$ and $\Phi(z_0, w_0) = c_1$.

To find ground-state solutions, we need to define that

$$m := \inf\{\Phi(z, w) : (z, w) \neq (0, 0), \Phi'(z, w) = 0\}.$$

By Lemma 3.3, it follows that

$$P(z, w) = P_1(z, w) = 0.$$

According to (3.17), we have $m \geq 0$. Let (z_n, w_n) be a sequence such that

$$\Phi'(z_n, w_n) = 0 \text{ and } \Phi(z_n, w_n) \rightarrow m.$$

Similar to Lemma 3.5, we can prove that there exists $(z', w') \in H^1$ such that

$$\Phi'(z', w') = 0 \text{ and } \Phi(z', w') = m,$$

which implies that $(u', v') = (f(z'), f(w'))$ is a ground-state solution of (1.1). The proof is complete.

Data availability statement

The original contributions presented in the study are included in the article/supplementary material; further inquiries can be directed to the corresponding authors.

References

- Jeanjean L. On the existence of bounded Palais-Smale sequences and application to a Landesman-Lazer type problem set on \mathbb{R}^N . *Proc R Soc Edinb Sect A*. (1999) 129:787–809.
- Wang Y, Huang X. Ground states of Nehari-Pohozaev type for a quasilinear Schrödinger system with superlinear reaction. *Electron Res Archive* (2023) 31(4): 2071–94. doi:10.3934/era.2023106
- Chen C, Yang H. Multiple Solutions for a Class of Quasilinear Schrödinger Systems in \mathbb{R}^N . *Bull Malays Math Sci Soc* (2019) 42:611–36. doi:10.1007/s40840-017-0502-z
- Chen J, Zhang Q. Positive solutions for quasilinear Schrödinger system with positive parameter. *Z Angew Math Phys* (2022) 73–144. doi:10.1007/S00033-022-01781-1
- Chen J, Zhang Q. Ground state solution of Nehari-Pohozaev type for periodic quasilinear Schrödinger system. *J Math Phys* (2020) 61:101510. doi:10.1063/5.0014321
- Lange H, Toomire B, Zweifel PF. Time-dependent dissipation in nonlinear Schrödinger systems. *J Math Phys* (1995) 36:1274–83. doi:10.1063/1.531120
- Laedke EW, Spatschek KH, Stenflo L. Evolution theorem for a class of perturbed envelope soliton solutions. *J Math Phys* (1983) 24:2764–9. doi:10.1063/1.525675
- Ritchie B. Relativistic self-focusing and channel formation in laser-plasma interactions. *Phys Rev E* (1994) 50:687–9. doi:10.1103/physreve.50.r687
- Chen J, Huang X, Cheng B, Zhu C. Some results on standing wave solutions for a class of quasilinear Schrödinger equations. *J Math Phys* (2019) 60:091506. doi:10.1063/1.5093720
- Liu JQ, Wang ZQ. Soliton solutions for quasilinear Schrödinger equations I. *Proc Amer Math Soc* (2002) 131(2):441–8. doi:10.1090/s0002-9939-02-06783-7
- Severo UB. Symmetric and nonsymmetric solutions for a class of quasilinear Schrödinger equations. *Adv Nonlinear Stud* (2008) 8:375–89. doi:10.1515/ans-2008-0208
- Moameni A. Existence of soliton solutions for a quasilinear Schrödinger equation involving critical exponent in \mathbb{R}^N . *J Differential Equations* (2006) 229: 570–87. doi:10.1016/j.jde.2006.07.001
- Alves CO, de Moraes Filho DC. Existence and concentration of positive solutions for a Schrödinger logarithmic equation. *Z Angew Math Phys* (2018) 69–144. doi:10.1007/s00033-018-1038-2
- Chen J, Chen B, Huang X. Ground state solutions for a class of quasilinear Schrödinger equations with Choquard type nonlinearity. *Appl Math Lett* (2020) 102: 106141. doi:10.1016/j.aml.2019.106141
- Yang X, Zhang W, Zhao F. Existence and multiplicity of solutions for a quasilinear Choquard equation via perturbation method. *J Math Phys* (2018) 59:081503. doi:10.1063/1.5038762

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16. Colin M, Jeanjean L Solutions for a quasilinear Schrödinger equation: a dual approach. *Nonlinear Anal* (2004) 56:213–26. doi:10.1016/j.na.2003.09.008
17. Liu JQ, Wang Y, Wang ZQ Soliton solutions for quasilinear Schrödinger equations, II. *J Differential Equations* (2003) 187:473–93. doi:10.1016/s0022-0396(02)00064-5
18. Chen J, Zhang Q Existence of positive ground state solutions for quasilinear Schrödinger system with positive parameter. *Appl Anal* (2022) 102:2676–91. doi:10.1080/00036811.2022.2033232
19. Willem M *Minimax theorems*. Berlin: Birkhauser (1996).
20. Chen S, Wu X Existence of positive solutions for a class of quasilinear Schrödinger equations of Choquard type. *J Math Anal Appl* (2019) 475:1754–77. doi:10.1016/j.jmaa.2019.03.051
21. do Ó JM, Miyagaki OH, Soares SHM Soliton solutions for quasilinear Schrödinger equations with critical growth. *J Differential Equations* (2010) 248:722–44. doi:10.1016/j.jde.2009.11.030
22. Fang X, Szulkin A Multiple solutions for a quasilinear Schrödinger equation. *J Differential Equations* (2013) 254:2015–32. doi:10.1016/j.jde.2012.11.017
23. Li G On the existence of nontrivial solutions for quasilinear Schrödinger systems. *Boundary Value Probl* (2022) 2022:40. doi:10.1186/s13661-022-01623-z
24. Wu X Multiple solutions for quasilinear Schrödinger equations with a parameter. *J Differential Equations* (2014) 256:2619–32. doi:10.1016/j.jde.2014.01.026
25. Zhang J, Tang X, Zhang W Infinitely many solutions of quasilinear Schrödinger equation with sign-changing potential. *J Math Anal Appl* (2014) 420:1762–75. doi:10.1016/j.jmaa.2014.06.055