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No-wait inversion—a novel model for T1 mapping from inversion recovery measurements without the waiting times

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Introduction: Quantification of longitudinal relaxation time T1 gained interest as an important MR-inducible tissue property for tissue characterization. Standard inversion recovery (IR) measurements for T1 determination take a prohibitively long time, and signal models assume a perfect inversion. Acceleration is possible by using the Look–Locker (LL) technique or other accelerated, model-based algorithms. However, the calculation of real T1 values from LL acquisitions necessitates the knowledge of equilibrium magnetization M0. Thus, usually, a waiting time to allow for free relaxation between global inversion pulses must be implemented. This study aims to introduce a novel model-based fitting approach for T1 mapping without the need for such waiting times.

Methods: Single-inversion spiral LL spoiled gradient echo acquisitions were performed in a phantom and eight healthy volunteers using a 1.5T magnetic resonance imaging (MRI) scanner. The measurements comprised two parts, one without magnetization preparation and a second featuring a global inversion pulse preparation before each of the 35 slices. Acquisition was performed without any waiting time in between slices, i.e., before the inversion pulses. T1 maps were calculated based on an iterative model-based reconstruction algorithm which combines the information from these two measurements, with and without inversion.

Results: Accurate T1 maps were obtained in phantom and volunteer measurements. ROI-based mean T1 values differ by an average of 1.5% in the phantom and 5% *in vivo* between reference measurements and the proposed method. The combined fit benefits from both the information obtained in the inversion prepared and the unprepared measurements. The former provides a large dynamic range for accurate model-based fitting of the relaxation process, while the latter provides equilibrium magnetization M0, necessary to obtain accurate T1 values from a LL-like acquisition.

Conclusion: The proposed model of a combined fit of an inversion-prepared and an unprepared measurement allows for robust fast T1 mapping, even in cases of imperfect inversion due to skipped waiting times for magnetization recovery. Thus, it can render long waiting times in between inversion pulses redundant.

KEYWORDS

T1 mapping, quantitative MRI, model-based reconstruction, inversion recovery, non-Cartesian trajectory, brain

1 Introduction

Quantitative magnetic resonance imaging (MRI) allows obtaining maps of specific physical parameters in MRI. When compared to qualitative MRI (or weighted imaging), it additionally allows

- comparison of different tissues within the same individual from different locations or points in time;
- comparison between images from different individuals; and
- achievement of disease-specificity, with studies correlating abnormalities in the absolute values of certain MR-inducible tissue properties and health conditions [1–4].

The quantification of longitudinal relaxation time T1 gained interest for tissue characterization, and T1 mapping is already used in cardiac [1, 2, 5] and brain MRI [3, 4, 6]. Conventional techniques for T1 determination are based on tracking magnetization after a suitable magnetization preparation, e.g., inversion recovery (IR) measurements. Following an inversion pulse, an image is acquired for a fixed delay after the inversion pulse. Typically, the k-space is filled in a segmented fashion. Consequently, for each repetition, magnetization must be restored to the original state, requiring a waiting interval to ensure full recovery of the signal before every new magnetization preparation. Repetitions of these measurements, to acquire images at different delays after the inversion pulse, are necessary to trace the relaxation of magnetization. Thus, standard IR measurements have a prohibitively long duration and are hardly feasible for clinical application [6–8]. The Look-Locker (LL) technique [9] came as a milestone for accelerated T1 quantification. Typically, an IR magnetization preparation pulse is applied, followed by a series of low-angle radiofrequency (RF) pulses used for spoiled gradient echo imaging. Data acquired after one inversion is sorted into different k-spaces, corresponding to the different delays after the inversion. Thus, the number of necessary repetitions can be reduced compared to conventional IR T1 mapping. However, the problem of the long waiting periods, on the order of $5 \times T1_{max}$, still exists for the LL technique [2, 10]. Although LL techniques are much faster than conventional IR T1 mapping, the RF pulses used for data acquisition affect the T1 recovery. Thus, the longitudinal relaxation process is not the same as for an undisturbed IR experiment, resulting in an apparent T1 (termed T1*). The longitudinal relaxation time T1 can be calculated from the apparent relaxation time T1* using the acquisition parameters repetition time and the local flip angle [11, 12]. Alternatively, T1 can be calculated from the measured quantities, T1*, equilibrium magnetization M0, and steady-state magnetization Mss [10, 12]: $T1/T1^* = M0/Mss$. To overcome the issue of long acquisition times, single shot techniques with a very restricted spatial or temporal

resolution can be used. However, both can lead to a poor accuracy of the corresponding parameter maps. Alternatively, undersampled k-spaces can be used. Yet, these approaches require sophisticated reconstruction approaches to mitigate the effects of undersampled data sets [13, 14]. For example, the model-based acceleration of LL T1 mapping (MAP [10]) uses a single inversion to generate a series of undersampled images. Then, a pixel-wise fit of the undersampled dataset is performed using a model of monoexponential relaxation. This results in fully sampled k-spaces, leading to a series of images that cover the whole IR process. As a result, accurate parameter maps are calculated [10, 15].

Based on a novel model-based fitting approach, this study aims to introduce a technique for T1 mapping without the necessity of waiting times for free relaxation in between inversion pulses.

2 Materials and methods

2.1 Measurements

Measurements were performed on eight healthy volunteers. The study was approved by the local ethics committee, and informed consent was obtained from all volunteers before scanning. Experiments were performed on a 1.5T MRI scanner (Siemens MAGNETOM Avanto Fit) using a 20-channel head coil array. Table 1 gives an overview of all measurements performed. In general, reference measurements were performed using a turbo spin echo (TSE) sequence (Table 1, Nb 1 and 3), while acquisitions for the proposed method are based on a LL [9], spoiled gradient echo acquisition with a spiral trajectory (Table 1, Nb. 2A, 2B, 4A, 4B, 5A, and 5B).

For comparison, in each volunteer and the phantom examination, one slice was acquired using a TSE sequence using the following imaging parameters: flip angle = 150°, resolution = $0.5 \times 0.5 \text{ mm}^2$, slice thickness = 5 mm, turbo factor = 10, measurement time = 22 min, and TE = 12 ms. Inversion times were as follows: TI = 30 ms, 100 ms, 250 ms, 500 ms, 1,000 ms, 1,600 ms, and 2,500 ms, and TR was TI + 1,000 ms and ranged between 1,030 and 3,500 ms. In the phantom, measurements at TI = 6,000 ms and 10,000 ms were added.

For all other acquisitions in each slice, 400 spoiled gradient echoes (GREs) using a variable-density spiral read-out trajectory were acquired. Consecutive spiral arms were spaced apart by a double golden angle. The trajectory was designed using the *variable-density spiral design function tool* [16, 17] by Brian Hargreaves and adjusted to observe a maximum gradient amplitude of 40 mT/m and a maximum gradient slew rate of 170 mT/m/s. The readout length of one spiral arm was 6.28 ms and extended in k-space to a maximum value of 0.625/

TABLE 1 Overview of measurements and their measurement parameters.

Nb	Object	Sequence	Magnetization preparation	Waiting time
1	Volunteers, 1 slice each	Turbo spin echo	Inversion-prepared	-
2A	Volunteers, 35 slices each	Spoiled GRE with spiral trajectory	Non	Non
2B	Volunteers, 35 slices each	Spoiled GRE with spiral trajectory	Inversion-prepared	Non
3	Phantom, 1 slice	Turbo spin echo	Inversion-prepared	-
4A	Phantom, 35 slices	Spoiled GRE with spiral trajectory	Non	Non
4B	Phantom, 35 slices	Spoiled GRE with spiral trajectory	Inversion-prepared	Non
5A	Phantom, 35 slices	Spoiled GRE with spiral trajectory	Non	Non
5B	Phantom, 35 slices	Spoiled GRE with spiral trajectory	Inversion-prepared	20s
Common parameters				
Spoiled GRE		TSE		
TE = 0.61 ms		TE = 12 ms		
		TI = 30 ms, 100 ms, 250 ms, 500 ms, 1,000 ms, 1,600 ms, and 2,500 ms		
TR = 7.5 ms, flip angle = 10°		TR = TI + 1,000 ms, flip angle = 150°,		
Resolution: 0.8 × 0.8 mm ²		Resolution = 0.5 × 0.5 mm ²		
Slice thickness = 5 mm		Slice thickness = 5 mm		
		Turbo factor = 10		
Measurement time per slice = 3 s		Measurement time per slice = 22 min		

TABLE 2 T1 values as determined in the phantom measurements.

Phantom	T1 values (mean ± SD) [ms]							
	Measurement	Spiral acquisition with pause (Nb. 5)			Spiral acquisition with no pause (Nb. 4)			TSE (Nb. 3)
Reconstruction model	a)	b)	c)	a)	b)	c)	Reference	
ROI 1	1,363 ± 263	1,929 ± 121	1,811 ± 131	1,377 ± 271	1,302 ± 79	1,801 ± 97	1,793 ± 35	1,724
ROI 2	1,058 ± 151	1,461 ± 67	1,392 ± 69	1,129 ± 222	1,147 ± 78	1,374 ± 81	1,402 ± 24	1,451
ROI 3	904 ± 138	1,127 ± 69	1,059 ± 60	866 ± 135	1,021 ± 71	1,043 ± 73	1,017 ± 21	1,010
ROI 4	631 ± 141	765 ± 35	703 ± 33	612 ± 116	742 ± 29	702 ± 41	688 ± 8	672
ROI 5	481 ± 167	524 ± 23	492 ± 35	481 ± 112	525 ± 22	493 ± 31	488 ± 5	467
ROI 6	329 ± 167	373 ± 16	346 ± 23	358 ± 180	377 ± 25	343 ± 20	347 ± 4	333
ROI 7	252 ± 84	272 ± 14	243 ± 18	274 ± 24	271 ± 17	246 ± 18	249 ± 4	235
ROI 8	326 ± 45	204 ± 12	181 ± 14	493 ± 72	200 ± 18	184 ± 19	181 ± 3	169

Model a): unprepared measurement only.

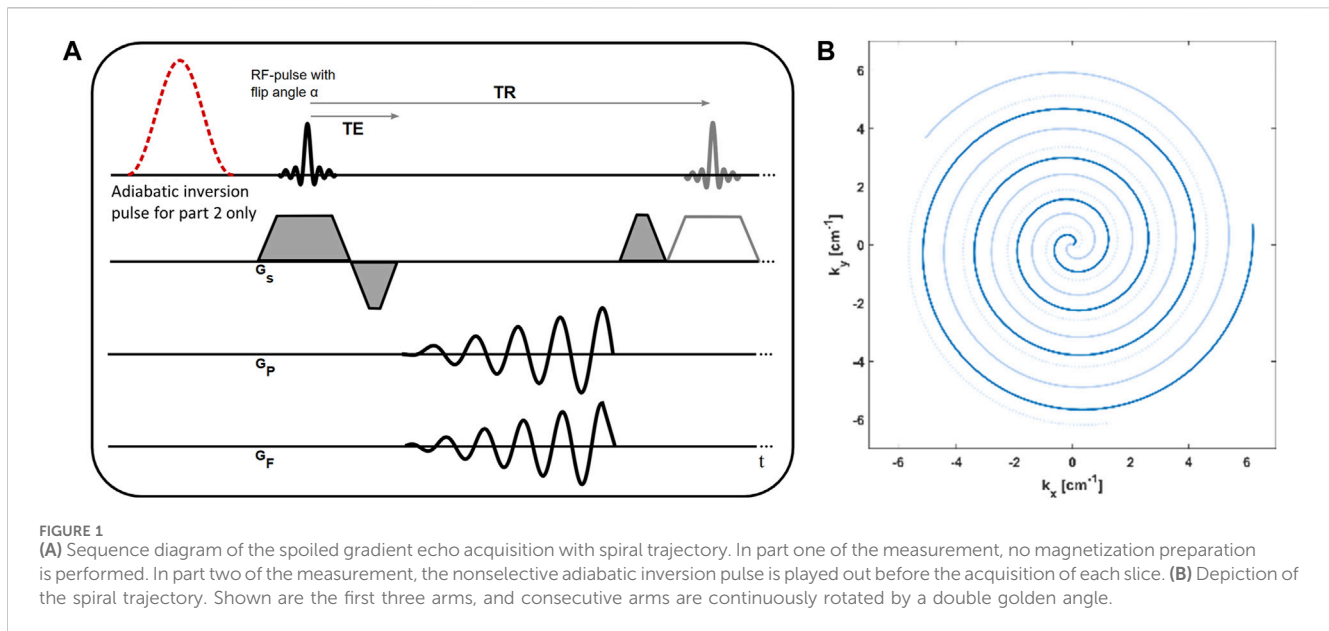
Model b): inversion-prepared measurement only.

Model c): combined fit of inversion-prepared and unprepared measurements.

Ref: TSE acquisition.

mm in order to achieve a spatial resolution of $0.8 \times 0.8 \text{ mm}^2$. By design, a total of 70 equally spaced spiral arms could fully cover the k-space for a FOV of 300 mm. However, in this study, a continuous rotation of a

double golden angle was chosen between consecutive spiral arms. Other imaging parameters include TE = 0.61 ms, TR = 7.5 ms, flip angle = 10°, slice thickness = 5 mm, and measurement time = 3s for



each slice. Each acquisition in a healthy volunteer comprised 35 slices in interleaved order and covered the whole brain.

These measurements comprised two LL-like acquisition parts: during the first part (Nb. 2A), no preparation and only slice-selective excitations were used. In the second part (Nb. 2B), a global adiabatic inversion pulse was applied before the acquisition of each slice (pulse length: 10.2 ms). No waiting time was heeded in between slices, and thus, starting magnetization was affected by the previous global inversion pulses. Both parts acquire the course of the receiver coil weighted magnetization from a starting point (M_0 or $-aM_0$) to the steady-state magnetization M_{ss} .

The same protocol was used for phantom measurements (Nb. 4A and 4B) on an *Essential System Phantom Model 106* by CaliberMRI (Handbook T1 values are given in Table 2). Here, an additional measurement with a waiting time of 20 s in between slices, and thus in between successive inversion pulses, was performed (Nb. 5A, 5B). The value of more than $10 \times$ maximum T1 in the phantom was chosen to ensure full relaxation of all phantom compartments and no signal contamination by former inversion pulses. Additionally, reference T1 values for the different vials within the phantom are provided by the vendor in the phantom manual.

A depiction of the spoiled gradient echo pulse sequence and the spiral k-space trajectory is given in Figure 1.

2.2 Reconstruction

For all measurements, raw data was extracted from the scanner, and image reconstruction was performed offline in MATLAB. Reference TSE measurements (Table 1, Nb. 1 and 3) were reconstructed by 2D Fourier transform. Data reconstruction of spiral GRE acquisitions was performed by a modified model-based acceleration of LL T1 mapping (MAP) algorithm [10, 15]. First, the corrected spiral k-space trajectory was calculated based on the gradient system transfer function [18, 19]. Then, initial images were reconstructed by gridding the data of each spiral arm into

separate k-spaces. All reconstructions were performed on a matrix of size 512×512 and with a field of view of 40.96 cm. After 2D Fourier transform, this results in a stack of highly undersampled images, which track the time course of $T1^*$ relaxation. As one spiral arm is gridded to each k-space, the temporal resolution equals TR. This was conducted separately for both parts of the spiral measurements, with and without the inversion pulse.

This initial step was followed by the iterative reconstruction algorithm consisting of alternating steps of pixel-wise model-based fitting and restoring data consistency. Data from the two parts of the measurement (Nb. 2A and B, 4A and B, and 5A and B) was either reconstructed together using a combined model (model (c)) or, for comparison, separately only for the unprepared part (model (a) for part A) or the inversion-prepared part (model (b) for part B).

This pixel-wise fitting of the reconstruction used the following fitting models:

(a) Unprepared:

The transient phase of unprepared magnetization in spoiled gradient echo acquisition only (Nb. 2A, 4A, and 5A) with three free parameters ($M_{0_{non}}$, M_{ss} , and $T1^*$):

$$M_{non} = M_{ss} + (M_{0_{non}} - M_{ss}) * e^{(-\frac{t}{T1^*})}.$$

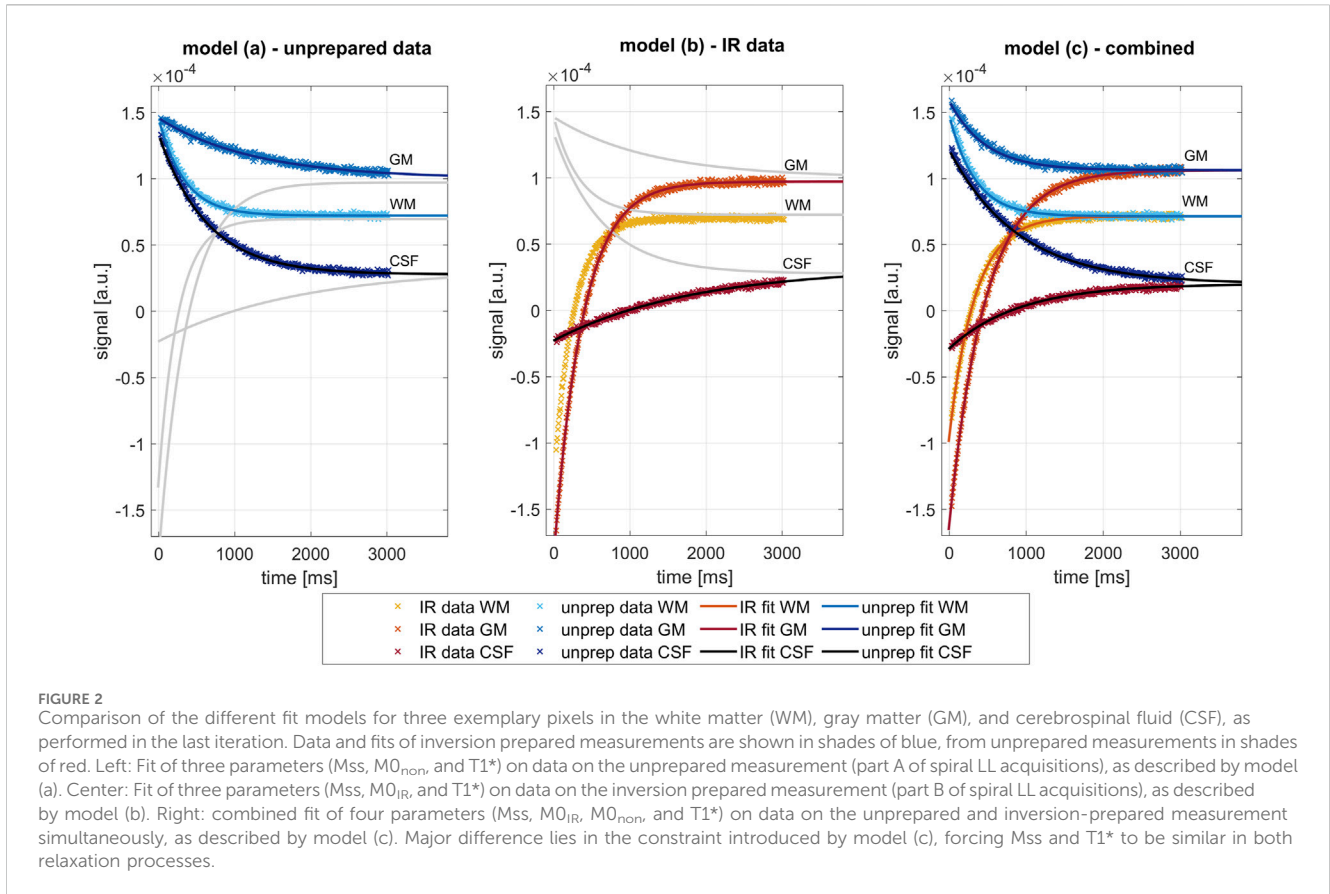
(b) Inversion-prepared:

Inversion recovery Look-Locker (IR-LL) acquisition only (Nb. 2B, 4B, and 5B) with three free parameters ($M_{0_{IR}}$, M_{ss} , and $T1^*$):

$$M_{inv} = M_{ss} - (M_{0_{IR}} + M_{ss}) * e^{(-\frac{t}{T1^*})}.$$

(c) Combined:

A combination of unprepared GRE magnetization and magnetization of IR-LL acquisitions (Nb. 2A and B, 4A and B,



and 5A and B) with four free parameters ($M_{0_{IR}}$, $M_{0_{non}}$, M_{ss} , and $T1^*$):

$$\begin{pmatrix} M_{inv} = M_{ss} - (M_{0_{IR}} + M_{ss}) * e^{-\frac{t}{T1^*}} \\ M_{non} = M_{ss} + (M_{0_{non}} - M_{ss}) * e^{-\frac{t}{T1^*}} \end{pmatrix},$$

where M_{ss} is steady-state magnetization, $T1^*$ is the LL relaxation time [9], $M_{0_{non}}$ is equilibrium magnetization, and $M_{0_{IR}}$ is magnetization directly after the inversion pulse. An illustration of the different models for data on representative voxels is given in Figure 2. Afterward, fitted images were transformed back to the k-space, and data consistency was enforced by reintroducing measured data points into these modeled k-spaces. As in the original MAP algorithm, we iteratively performed steps consisting of model-based fitting and data consistency until convergence to a solution was reached [15].

However, as this procedure converges rather slowly, we used a modified algorithm. The original MAP reconstruction algorithm can be seen as alternating projections onto two different subspaces: the data-consistent solutions and the subspace of exponential signal evolutions in all pixels. A combination of two MAP projections, either from the data-consistent solution subspace back onto the data-consistent solution subspace or from the exponential model subspace back to the model subspace, can be seen as steps into the right direction with a non-optimized step size. In other words, from two solutions in each of these subspaces, a one-dimensional (1D) consistency subspace and a 1D model subspace are defined. In the modified MAP algorithm, we calculate the point in the 1D data-

consistent subspace, which is closest to the 1D model subspace, and use it as the starting points for two further MAP iterations. These MAP steps were used again to determine new 1D subspaces and a new step size estimate. If the reconstruction problem would be linear and the measured data without noise, the exact solution would be found after a single step size adjustment. Numerical experiments suggest that, for our data, this approach reduces the number of steps before convergence by a factor of approximately 50 compared to the original MAP approach while arriving at a similar solution. For the reconstructions presented below, we performed six MAP steps and two step size adjustments (after MAP steps 2 and 4) in total.

2.3 Calculation of parameter maps

From the results of the last iteration in the model-based subspace, absolute T1 values were calculated from the apparent relaxation time $T1^*$ in LL acquisitions as [10]: $T1 = \frac{T1^* * M_{0_{IR}/non}}{M_{ss}}$ using the latest fitting values. Additionally, the inversion efficiency was calculated as $InvEff = M_{0_{IR}}/M_{0_{non}}$ from the combined model (c).

Reference TSE data were reconstructed by 2D FFT, and an adaptive coil combination was performed [20]. The average phase of the last three measurements in each pixel was subtracted from all measurements, and subsequently, the real part of the data was fitted pixel-wise using the following equation (adapted from [21]):

$$S = M_0 * \left(1 - (1 - b) * e^{-\frac{t}{T1}} + c * e^{-\frac{t}{T1}} \right),$$

with $c = (e^{TE_{last}/T1})$,

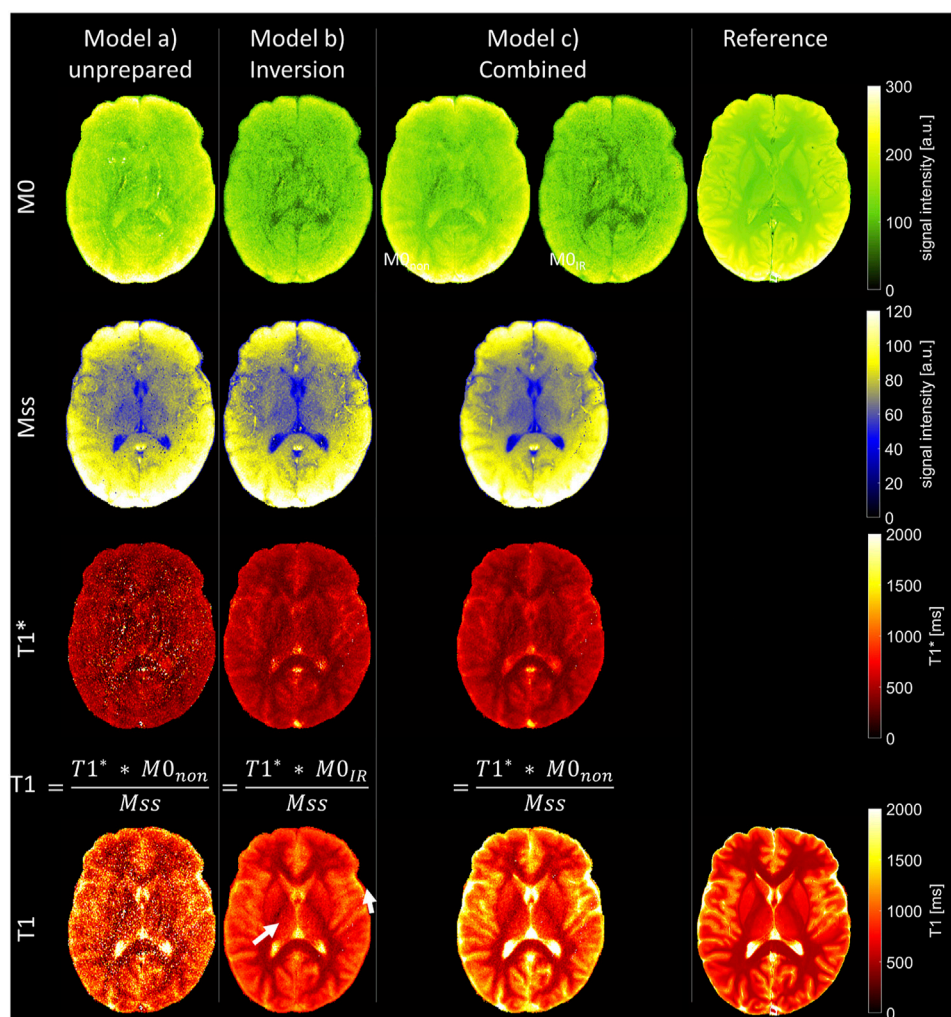


FIGURE 3

Parameter maps as determined by the three different fit models and from the reference measurement. Top: equilibrium magnetization (M_0). In the combined fit, two independent parameters (M_{0IR} and M_{0non}) are considered to allow for imperfect inversion. Center: steady-state magnetization (M_{ss}) and apparent longitudinal relaxation times ($T1^*$). Bottom: longitudinal relaxation times (T_1), as calculated from fit parameters of unprepared acquisition only [model (a)], inversion-prepared data only [model (b)], and the combined fit of both parts of the measurements [model (c)]. Arrows point to areas of T_1 underestimation.

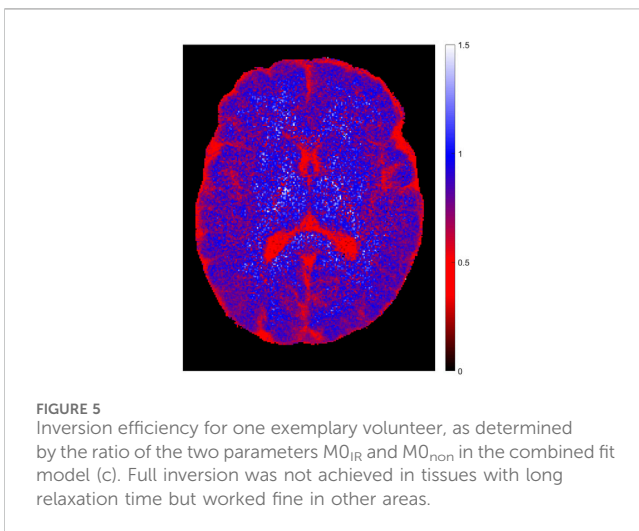
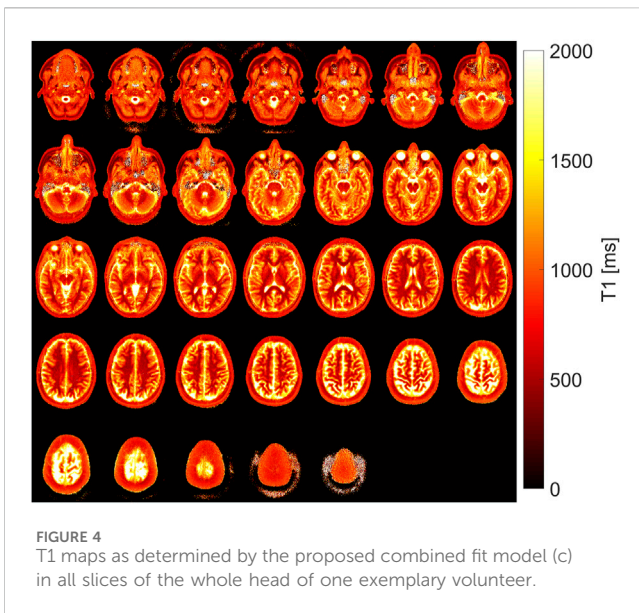
where TR is the repetition time and TI is time after inversion. Parameters b and c account for a potential non-ideal inversion and a constant effect of the acquired multi-echo train, respectively.

2.4 ROI-based analysis

In each multi-slice measurement, the central slice was identified, which also corresponds to the slice chosen in single-slice reference measurements. ROIs were then manually placed in corresponding areas in all measurements. In the phantom, a total of eight circular ROIs were added to comprise each of the separate vials inside the phantom. These ROIs were labeled ROI 1 to ROI 8 in the order of decreasing T_1 values. In each volunteer, three different polygonal ROIs were drawn within that selected slice. The first ROI is placed in the frontal white matter area, the second ROI in the gray matter in the sulci in the right posterior brain, and the third ROI inside the

frontal part of one ventricle. For each ROI in the phantom, the mean T_1 values and the standard deviation (SD) within the ROI are evaluated. In each volunteer mean T_1 values and SD within the ROI are evaluated, as well as mean T_1 values and SD in each tissue over the whole group of volunteers.

Additionally, equivalence testing was performed to assess whether the phantom measurements with and without the waiting time were statistically equivalent. Specifically, the two one-sided t -test (TOST) procedure was conducted to test whether the difference in means between two groups falls within a predefined equivalence margin [22]. The equivalence bounds were set to be 3% of the mean T_1 value in the reference measurement (i.e., 59 ms in ROI 1, 44 ms in ROI 2, 32 ms in ROI 3, 21 ms in ROI 4, 15 ms in ROI 5, 10 ms in ROI 6, 7 ms in ROI 7, and 5 ms in ROI 8). Practically, the equivalence of the means was assessed by examining whether the 90% confidence interval fell entirely within the predetermined equivalence bounds. Values



within each ROI were considered as paired as the same ROIs were applied and pixel values were listed in the same order. The significance level was set at $\alpha = 0.05$.

3 Results

3.1 Signal evolution in the different models

Exemplarily, Figure 2 shows the signal evolution in three pixels located in the white matter area in the center of the brain, gray matter area, and cerebrospinal fluid (CSF) in the ventricle. Shown are the data points after data consistency was restored and the corresponding exponential curves fitted in the last iteration of each of the separate reconstruction processes (model (a), model (b), and model (c)). Note that the combined fitting, as proposed in model (c), forces M_{ss} and $T1^*$ to be similar for both relaxation processes as data from identical tissue should feature identical tissue properties.

Separate fitting of either inversion-prepared data (model (a)) or unprepared data (model (b)) imposes no such constraint, resulting in differing values.

3.2 Exemplary volunteer

Figure 3 shows parameter maps, as determined by the three different models after eight iteration steps. The combined model (c) yields two values for equilibrium magnetization M_{0IR} and M_{0non} , from the inverted and unprepared relaxation processes respectively, while steady-state magnetization M_{ss} and $T1^*$ times are preordained to be identical. Steady-state magnetization M_{ss} (Figure 3 M_{ss}) is comparable in all models. Maps of the equilibrium magnetization $M_{0IR/non}$ are shown in Figure 3 M_0 . Using model (b) for fitting of inversion-prepared data only (part B of LL measurements) enables robust fitting. Challenges to this model (b) arise in cases of imperfect inversion or altered magnetization due to the lack of waiting time between consecutive global inversions. The latter is especially apparent in areas with long relaxation times like in CSF, where the calculated values of M_{0IR} are low. Low M_{0IR} values are also determined by the combined model, but in contrast, it can robustly determine M_{0non} , and the resulting maps show a similar distribution as in the reference measurement.

$T1^*$ maps (Figure 3 $T1^*$) are noisy when determined by model (a) from unprepared data only. The maps determined by model (b) or (c) show reduced tissue contrast as usual for $T1^*$ in LL acquisitions [23–26].

$T1$ maps, as calculated from these fitting parameters, are shown in Figure 3 $T1$. The calculation of $T1$ depends on the accurate determination of $T1^*$, M_{ss} , and M_0 . Although the first is challenging in model (a), the last poses a problem in model (b). In model (c), equilibrium magnetization is determined by M_{0non} , allowing robust $T1$ calculation, as reflected in the detailed $T1$ map (Figure 3 $T1$).

The results of all 35 slices in the exemplary volunteer are shown in Figure 4. Reasonable $T1$ maps were acquired at all positions in the head and were independent of the chronological position of the slice in the measurement setup and, therefore, independent of its history of inversion pulses.

The efficiency of the inversion preparation is shown in Figure 5. Clearly, the relaxation process does not start at 100% inversion of equilibrium magnetization M_0 in tissues with long relaxation time but in many other areas.

3.3 Phantom measurements

Figure 6 shows the $T1$ maps, as determined in the phantom measurements with and without a deliberate waiting time in between slices and, therefore, before all global inversion pulses. Corresponding $T1$ values from the ROI-based analysis are collected in Table 2. The parameter maps (M_{ss} , M_0 , $T1^*$, and $T1$), as determined by the three different fit models after the eighth iteration step are presented in Supplementary Figure S1.

As in the volunteer, maps calculated from the unprepared measurement (Nb. 2A) with model (a) are noisy, while robust

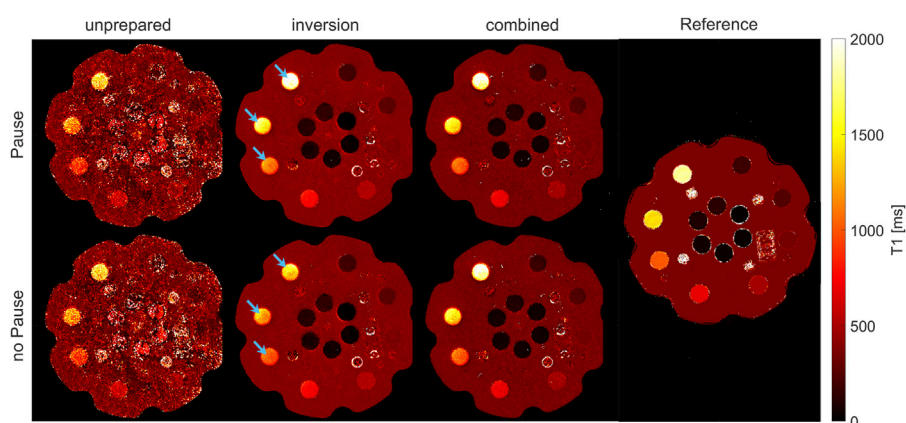


FIGURE 6

T1 maps of phantom measurement, with and without a waiting time to allow for free relaxation of magnetization. Maps calculated from the unprepared measurement by model (a) are noisy, while robust fitting was possible in the inversion-prepared case (model (b)) and with the proposed combined model (c). The major difference between measurements with and without waiting time emerges when regarding vials with long T1 values (indicated by arrows) and the inversion measurement only (model (b) for measurement part B). Measurements without waiting time result in a variable underestimation of absolute T1 values, which can be mitigated by either introducing the waiting time of 20 s or by using the combined model (c).

fitting was possible in the inversion-prepared case (Nb. 2B) with model (b) and with the proposed combined reconstruction (Nb. 2A and B) by model (c).

The statistical equivalence test for the combined reconstruction model (model (c)) revealed that the mean difference between T1 values from measurements with and without pause was below the equivalence bounds in all ROIs. Tests on the T1 values gained by model (b) from IR measurements showed that the mean difference was below the equivalent bounds in ROIs 5–8. T1 values in the ROIs 1–4, featuring the longest T1 values, cannot be considered equivalent within the given bounds. Tests for reconstruction model (a), unprepared measurement, revealed that the mean difference between T1 values from measurements with and without pause was outside the equivalence bounds in all ROIs, except for ROI 5.

Deviations of ROI-based mean T1 values between the combined fit model (c) and the TSE reference measurement were 0.4%, 2.0%, 2.5%, 2.0%, 1.0%, 1.2%, 1.2% and 1.6%, respectively, for the eight ROIs in the decreasing T1 order. Deviations to phantom manual values were 4.3%, 5.6%, 3.2%, 4.3%, 5.3%, 2.9%, 4.5%, 8.2%, respectively, for ROIs 1–8 (see also [Supplementary Table S1](#)).

3.4 Volunteer ensemble

T1 maps were calculated for all slices in all eight healthy volunteers and by all three reconstruction models. The T1 maps acquired from the TSE reference measurement and T1 maps calculated from the combined model (c) in a corresponding slice are shown in [Supplementary Figure S2](#). Mean T1 and SD, as calculated by model (a), were $1,567 \pm 139$ ms in CSF, 967 ± 37 ms in the gray matter, and 529 ± 43 ms in the white matter. Mean T1 and SD, as calculated by model (b), were $2,278 \pm 240$ ms in CSF, $1,022 \pm 44$ ms in the gray matter, and 559 ± 84 ms in the white matter. For the combined model (c), mean T1 values and SD of $2,911 \pm 240$ ms in CSF, $1,198 \pm 58$ ms in the gray matter, and 557 ± 40 ms in the white matter were found. Separate results for each volunteer and area are collected in [Table 3](#),

relative deviations are collected in [Supplementary Table S2](#). In general, model (a) underestimates T1 values and has high SD within the ROIs in each volunteer. Model (b) has lower SD within ROIs and shows underestimation of T1 values, especially in tissue with slow relaxation. The results from model (c) agree within the error bounds with the T1 times as determined by the reference measurement and are in good accordance with the literature values [27–29]. In addition, they feature low SD within each of the evaluated areas, indicating robust fitting.

4 Discussion

The proposed method aims to accomplish fast T1 mapping without the need for waiting times in between inversion pulses. Therefore, the combination of information from two parts of a measurement, with and without inversion, is proposed. To obtain T1 values from a LL acquisition, information on the apparent relaxation time $T1^*$ and on equilibrium magnetization M_0 and steady-state magnetization M_{ss} is necessary [9, 12]. Inversion-prepared measurements provide a large dynamic range for accurate model-based fitting of the relaxation process (and thus $T1^*$), but without considering the necessary time for magnetization recovery, they cannot provide M_0 . In contrast, unprepared measurements can provide the equilibrium magnetization M_0 , but determination of $T1^*$ is challenging, due to the low dynamic range of the relaxation process. Information on M_0 could, otherwise, be procured by considering long waiting times in between inversion pulses to allow full relaxation before each such pulse. The main advantage of the proposed approach is the omission of such waiting times, which significantly decreases acquisition time.

4.1 Sequence design

The design of the spiral trajectory was driven by two limitations. First, the desired resolution and second a limited readout length, in

TABLE 3 T1 values determined from ROIs placed in CSF, gray matter, and white matter for the three different fitting models.

Volunteer	T1 values (mean \pm SD) [ms]											
ROI location	CSF				Gray matter				White matter			
Reconstruction model	a)	b)	c)	Reference	a)	b)	c)	Reference	a)	b)	c)	Reference
1	1,619 \pm 60	2,549 \pm 174	2,917 \pm 55	2,387 \pm 158	947 \pm 34	969 \pm 71	1,162 \pm 57	989 \pm 51	546 \pm 33	664 \pm 161	595 \pm 62	571 \pm 26
2	1,467 \pm 55	1,751 \pm 148	2,736 \pm 88	2,434 \pm 130	965 \pm 33	1,019 \pm 78	1,185 \pm 55	1,065 \pm 69	508 \pm 18	547 \pm 108	536 \pm 25	527 \pm 6
3	1,516 \pm 118	2,199 \pm 219	2,815 \pm 104	2,723 \pm 35	1,051 \pm 39	1,094 \pm 136	1,331 \pm 63	1,129 \pm 50	535 \pm 20	596 \pm 133	575 \pm 27	547 \pm 10
4	1,520 \pm 69	2,228 \pm 226	2,903 \pm 59	2,908 \pm 184	973 \pm 21	979 \pm 54	1,169 \pm 22	1,122 \pm 32	528 \pm 25	597 \pm 147	547 \pm 30	539 \pm 6
5	1,702 \pm 103	2,306 \pm 209	2,998 \pm 5	3,130 \pm 55	952 \pm 25	1,008 \pm 83	1,210 \pm 34	1,176 \pm 59	531 \pm 16	478 \pm 71	569 \pm 24	546 \pm 15
6	1,449 \pm 98	2,406 \pm 169	2,962 \pm 35	2,699 \pm 154	936 \pm 24	1,022 \pm 93	1,173 \pm 35	1,241 \pm 68	580 \pm 21	484 \pm 82	590 \pm 30	565 \pm 7
7	1,828 \pm 93	2,413 \pm 157	2,998 \pm 6	2,781 \pm 135	937 \pm 25	1,007 \pm 68	1,146 \pm 28	1,063 \pm 51	438 \pm 19	659 \pm 105	470 \pm 19	519 \pm 12
8	1,436 \pm 102	2,369 \pm 164	2,958 \pm 49	2,677 \pm 172	975 \pm 29	1,079 \pm 75	1,208 \pm 51	1,122 \pm 57	566 \pm 23	445 \pm 158	573 \pm 27	560 \pm 10
Mean	1,567 \pm 139	2,278 \pm 240	2,911 \pm 92	2,717 \pm 240	967 \pm 37	1,022 \pm 44	1,198 \pm 58	1,113 \pm 76	529 \pm 43	559 \pm 84	557 \pm 40	547 \pm 18

Model a): unprepared measurement only.

Model b): inversion-prepared measurement only.

Model c): combined fit of inversion-prepared and unprepared measurements.

Ref: TSE acquisition.

order to prevent strong effects of T_2^* relaxation along the readout. Furthermore, a variable density factor of two was chosen to optimize the acquisition of the k -space where most of the energy in an image is concentrated near the origin of the k -space [17]. Consecutive spiral arms were rotated in the spiral plane by a double golden angle. One advantage of golden angle sampling is the uniform coverage of the k -space along time. For optimized fitting of the exponential T_1 relaxation, ideally, three full k -spaces are available at set points in time. For a fixed T_1 value, the ideal position of these in time is known. For unknown multiple T_1 values, golden angle sampling provides a good compromise. For center-out trajectory designs, the double golden angle ($137.51^\circ \approx 360^\circ - 2 \times 111.25^\circ$) features the same properties as the standard golden angle (111.25°) for full spoke acquisitions [30].

4.2 Reference measurement

Reference multi-inversion TSE measurements were performed in one slice in each volunteer and in the phantom. The settings were adjusted to minimize SAR in all acquisitions and to keep overall scan time within reasonable limits. A variable TR of $TI + 1,000$ ms was chosen for each acquisition in order to realize equal relaxation time in between the last excitation for acquisition and the subsequent inversion pulse. A comprehensive overview of IR signal modeling can be found in MRI handbooks and online (e.g., [21, 31, 32]), including the signal dependencies on TR. In contrast, other measurement setups use constant long TR values (e.g., $5 \times T_1$) and include a long TR assumption in the fitting process.

4.3 Comparison of T_1 measurements

In this study, T_1 was determined by multi-inversion TSE measurements, as well as by inversion-prepared and unprepared spiral LL acquisitions in combination with three different fit models. Additionally, handbook values for the phantom are given.

Values given in the phantom handbook are generally lower than measured values (with the exception of ROI 2) independent of the acquisition type and fitting model. For our proposed fitting model (c), deviations are between 2.9% and 8.2%. Differences between the reference measurement and handbook values are in a similar range, between 0.7% and 7.1%. The main reason might be a difference in temperature as handbook values are stated at 20°C , whereas both reference and LL measurements were performed at 22°C – 23°C . The temperature dependence of T_1 is known to be between 1% and 3% per degree temperature change [33, 34]. Additionally, the reference measurement was performed as a TSE sequence, where the prolonged echo trains average signal around the set TI .

Reconstruction of LL acquisitions by model (a) has low mean T_1 values and maps are very noisy, which is also reflected in the high SD within each ROI in the phantom and in all volunteers. Model (b) has lower SD within ROIs but also shows an underestimation of T_1 values in comparison to reference measurements or other fitting models, especially in areas with slow relaxation (e.g., ROIs 1 and 2 in phantom or CSF and gray matter *in vivo*). Due to that, both are not considered a recommended technique.

T_1 values, as determined by the proposed combined fit model (c), are in good agreement with reference values obtained by multiple

inversion TSE measurements for phantom and *in vivo* cases. In *in vivo* measurements, T_1 values deviate by 1.8% in the white matter, 7.1% in the gray matter, and 6.6% in CSF, and are in a range of literature values given for the evaluated tissues [27–29]. In the phantom, a T_1 range of 150–2,000 ms is covered. Within these ROIs, T_1 values deviate by no more than 2.5% from the reference measurement.

For very slow relaxing components (like CSF), inaccuracies can occur due to the short overall LL acquisition time per slice. Although main features of the relaxation process are covered within the 3 s acquisition, the magnetization level of very slow components has not necessarily reached the steady-state magnetization yet.

4.4 Model considerations

The main advantage of model-based reconstructions lies in the usage of prior knowledge. The MAP approach applied in this study combines prior knowledge about the relaxation behavior with the contrast information acquired using a spiral sampling scheme [10, 15]. The iterative reconstruction algorithm introduces more and more information about the relaxation process into the initial k -spaces, which originally contain only one spiral arm each. Although, for example, fingerprinting assumes that incoherent artifacts do not influence the fitting, MAP uses the information from other time points to consistently fill each k -space. Thus, over the iterations, the k -space of each point in time is filled and undersampling artifacts are eliminated. As a result, one complete image is generated for every time point TI of the relaxation process. Specifically, for the study presented here, a total of 400 full images spaced apart by TR were obtained.

The relaxation process from equilibrium magnetization (or inversion) into steady-state magnetization depends on T_1^* and can be fitted by corresponding exponential equations of T_1^* in inversion-prepared and unprepared measurements. The main advantage of the former is the large dynamic range of signal evolution, which benefits robust fitting, especially in accelerated methods. Unfortunately, to determine equilibrium magnetization M_0 , it also demands perfect inversion and, thus, a waiting time to allow free relaxation in between pulses. By combining the inversion acquisition with an unprepared measurement, M_0 can be adequately recovered in a combined fit, as given by model (c). Thus, no waiting time is necessary. As an additional benefit, the combined fit model allows an investigation of the performance of the inversion pulse by comparing $M_{0,\text{non}}$ and $M_{0,\text{IR}}$.

The main advantage of such a model-based approach is the efficient use of information gathered in the shortest measurement times. Reconstructions can be performed on a fine temporal grid as no binning of phase-encoding steps is used, but each acquired spiral arm provides information on one time point along the relaxation curves. In combination with pixel-wise fitting, no spatial or temporal filtering is applied.

4.5 Inversion pulses

The presented study used a global inversion pulse to provide maximum insensitivity to B_1 inhomogeneity [35]. The inversion was then followed by the acquisition of a single slice. Consequently, all slices are affected each time an inversion is played out. In other

words, all slices, but the first, have a history of multiple inversions. If no waiting time for free relaxation in between the inversion pulses is heeded, the inversion acts on an arbitrary magnetization vector and not on the equilibrium magnetization M_0 . Thus, the start point of the relaxation process shifts from $-M_0$, as in a perfect inversion, to a different value, and M_0 can no longer be determined. Loss of the M_0 information subsequently prevents the necessary $T1^*$ correction for LL acquisitions [9, 10, 12].

An additional advantage of a global inversion is that it allows generalization of the method, for example, for segmented 3D [36–38] acquisitions or simultaneous multi-slice measurements [39].

4.6 Other T1 mapping techniques

T1 mapping is often included in multiparametric acquisition frameworks [40], like fingerprinting [41–44], the MR multitasking framework [45, 46], or SyMRI [47, 48], which can achieve clinically reasonable scan times below 10 min. Although the additional information gained might be a benefit, the faster acquisition in our proposed method is favorable for dedicated T1 mapping applications.

Another common technique to determine T1 is the variable flip angle method [49]. Several approaches were presented, which can acquire whole-head T1 maps under 10 min [3, 6, 50]. Although they are very fast, they require exact knowledge of the flip angle. Additionally, problems with the B1 inhomogeneity and slice profile may arise, especially for thin slices and fast pulses [51]. T1 values determined by the variable flip angle method are also dependent on the RF spoiling procedure [52].

In techniques based on IR-LL methods, T1 can be calculated from $T1^*$ either by using the acquisition parameters repetition time TR and the local flip angle or from the measured quantities, $T1^*$, M_0 , and M_{ss} [9, 12]. By choosing optimal settings for a given experiment and incorporating flip angle information, a bias from imperfect inversion can be minimized. Shin et al. [11] performed whole brain T1 mapping in approximately 4 min with a resolution of $1 \times 1 \times 4 \text{ mm}^3$. Within an optimized setup, an average error of T1 of 1.2% was achieved, without requiring a specific waiting time in between inversion pulses. In comparison to our approach, this study achieved comparable measurement times, albeit slightly lower in-plane resolution and volume coverage.

Other techniques based on IR-LL methods mostly rely on segmented acquisitions with multiple inversions, especially for 3D acquisitions, which require a high number of readout lines. For example, Henderson et al. [36] performed whole-head T1 mapping with a resolution of $1.4 \times 1.4 \times 6 \text{ mm}^3$ in 8 min. The volume was divided into 128 segments, each starting with an inversion pulse and a delay time of 2 s in between. Similar measurement times were achieved by Maier et al. [38] for 1-mm isotropic resolution by dividing the k-space into 60 segments. Building on that, interleaved slab-selective inversion preparations were proposed. The necessary delay times can there be used to acquire different slabs. For example, Li et al. [37] acquired whole-head T1 maps with 1-mm isotropic resolution within 4 min 21 s by separating the acquisition into six slabs.

The introduction of saturation pulses prior to the global inversion pulse of each segment can shorten necessary delay times. Such a setup was, for example, implemented by

Deichmann et al. [53] to minimize waiting times from 15 s to 3 s in a segmented acquisition. The so-called TAPIR sequence [7, 54] uses a saturation pulse and a delay of 2–2.5 s before the inversion pulse. Similar to our technique, the relaxation process after the inversion does not start from inverted equilibrium magnetization. Although in TAPIR, the additional preparation step allows control over this start point, our approach assumes a random start point. The necessary information on equilibrium magnetization is then gained either by modeling (TAPIR) or by an additional part in the measurement (as proposed here). A 3D extension of the TAPIR algorithm was proposed by Claeser et al. [55] who acquired the whole head in five segments with a resolution of $0.94 \times 0.94 \times 2.5 \text{ mm}^3$ in 3 min 25 s. This method can achieve acquisition times similar to our proposed method.

The shortest scan times are generally achieved by single-shot, inversion-prepared LL techniques, with high undersampling rates in combination with a sophisticated reconstruction algorithm. Feng et al. [13] acquired 32 slices in the human head with a resolution of $1.25 \times 1.25 \times 3 \text{ mm}^3$ in 2 min 32 s and with a resolution of $0.875 \times 0.875 \times 3 \text{ mm}^3$ in 2 min 49 s. In addition, Müller et al. [56] achieved a resolution of $0.5 \times 0.5 \times 4 \text{ mm}^3$ or $0.5 \times 0.5 \times 3 \text{ mm}^3$ in 1 min 55 s and 2 min 36 s, respectively. Both methods provide in-plane resolutions similar to our proposed method ($0.8 \times 0.8 \times 5 \text{ mm}^3$). Measurement times are even shorter, albeit the total volume covered is also smaller ($280 \times 280 \times 96 \text{ mm}^3$ [13] and $192 \times 192 \times 115/117 \text{ mm}^3$ [56] vs $410 \times 410 \times 175 \text{ mm}^3$ in our method). Wang et al. [39] combined such an approach with a simultaneous multi-slice acquisition, thus allowing even further acceleration of whole-head acquisitions. A total of five stacks, comprising five slices each, were acquired in approximately 1 min. Out of the total time, 4 s each were spent on the acquisition of the five stacks, and in between those, a 10-s waiting period was heeded. Considering that, an evaluation of the combination of our proposed method without waiting times, with a simultaneous multi-slice acquisition, is a very interesting direction for future work.

Another recent study also aims at eliminating the waiting time in between inversion pulses, by using an AI-based reconstruction method to gain quantitative T1 values [57]. In comparison to our method, the main benefit is the extremely short reconstruction time once training of the data is accomplished. Nevertheless, our method does not require a matching training dataset and does not rely on specific acquisition parameters.

4.7 Measurement time

In this study, the whole brain was covered by acquiring 35 slices. As the acquisition of each slice only took 3 s, total measurement time for both inversion and unprepared measurement together resulted in 210 s. To the best of our knowledge, only few other techniques can provide T1 mapping in the whole head at submillimeter resolution in such short acquisition times [13, 39, 55, 56].

If the traditional waiting period of $5 \cdot T1$ (e.g., approx. 1.2 s in the gray matter) after each inversion pulse was observed, total measurement time for an inversion prepared acquisition would have been 309 s, assuming a similar sequence setup. Over this reference, the proposed algorithm presents an acceleration, which is even more pronounced if longer T1 values are considered (e.g., in

CSF). In future implementations, further acceleration by shortening the unprepared acquisition might be possible and should be evaluated.

In the phantom, two sets of measurements were performed, with and without the waiting time, to evaluate the influence of the omission of waiting time for free relaxation on T1 values. Statistical equivalence testing revealed that the mean difference of T1 values, acquired from the combined reconstruction model (c), was below the equivalence bounds for the whole range of T1 values covered by the phantom (175–1900 ms). In contrast, the results from IR measurements without waiting time are not equivalent to the results from measurements with waiting time for free relaxation. In conclusion, by applying the combined reconstruction model, measurement time can be reduced without compromising T1 mapping.

4.8 Limitations

Common limitations of iterative reconstruction procedures are long reconstruction times and a high computational load. Although, here, the introduction of the step size adjustment into the reconstruction pipeline could significantly reduce the number of steps before convergence (approximately by a factor of 50), the offline implementation and overall time requirements do not yet allow an easy introduction into clinical practice. Further acceleration could be gained in future by using time-optimized fitting procedures, for example, by a reduced dimension nonlinear least-square fitting, as proposed in Barral et. al [58].

Furthermore, our approach is sensitive to model mismatches, which can occur, for example, due to motion, inflow, slice profile, partial volume effects, or multiple T1 components. Here, a rather simple straightforward exponential behavior is used as a model, which does not consider any such effects but, on the other hand, is very robust. Nevertheless, the introduction of such aspects into the model is possible [59].

The number of iterations and the setup of MAP steps and step size adjustments were empirically determined. The optimal settings might differ for different acquisition parameters. Then, other and maybe more flexible stop criteria could additionally be evaluated.

5 Conclusion

The proposed model of a combined fit of an inversion-prepared and unprepared measurement allows for robust fast T1 mapping, even in cases of insufficient magnetization relaxation before consecutive inversion pulses. It can thus render long waiting times in between inversion pulses redundant.

Data availability statement

The raw data supporting the conclusions of this article will be made available by the authors, without undue reservation. The reconstruction code developed for this study can be found here: <https://github.com/expRad/nowaitinversion>.

Ethics statement

The studies involving humans were approved by Klinische Ethikkomitee (KEK), University Hospital Würzburg. The studies were conducted in accordance with the local legislation and institutional requirements. The participants provided their written informed consent to participate in this study.

Author contributions

JB: conceptualization, data curation, formal analysis, investigation, methodology, software, validation, and writing—original draft. JK: formal analysis, investigation, methodology, software, writing—original draft, and writing—review and editing. OS: data curation, software, and writing—review and editing. AW: data curation, investigation, resources, and writing—review and editing. HK: conceptualization, data curation, methodology, project administration, resources, supervision, and writing—review and editing. AS: conceptualization, formal analysis, investigation, methodology, project administration, resources, software, supervision, validation, visualization, writing—original draft, and writing—review and editing.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Supplementary material

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fphy.2023.1299522/full#supplementary-material>

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