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Realization and dosimetric characterization of a mini-beam/flash electron beam

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The use of the flash effect and mini-beams have demonstrated the ability to spare healthy tissue while maintaining the same effectiveness in controlling tumors. In this study, we present the implementation and comprehensive dosimetric characterization of low-energy mini-beam radiation therapy at both conventional and ultra-high dose rates. These beams possess important features that allow for a wide range of spatial and temporal parameter variations, independently or simultaneously, for both effects. This novel capability enables the performance of *in vivo/vitro* radiobiological experiments, which are crucial for understanding the underlying mechanisms and quantitative dependencies of these effects on their respective parameters. This understanding is essential for evaluating the potential clinical applications of the two effects both individually and in terms of their potential synergistic actions.

KEYWORDS

mini-beam, flash, UHDR, UHDP, SFRT, passive collimation

1 Introduction

The initial intuition that led to the development of spatially fractionated radiotherapy (SFRT) was credited to Dr. Alban Köhler in 1909. Köhler discovered that by introducing a grid-shaped shielding system for X-ray beams (at that time, X-ray tubes were the only radiation source), he was able to treat skin tumors more effectively while significantly reducing damage to healthy tissue. Erythema and necrosis healed within a few weeks, leaving the new epidermal tissue completely healthy [1]. Although this approach was initially used for the treatment of some superficial tumors, it remained on the sidelines of radiotherapy development for a long time. However, since the 1990s, several research groups in Europe and the United States have begun to consider the possibility of implementing spatial fractionation in radiotherapy. Multiple radiobiology studies were conducted, initially using X-rays, including synchrotron light [2–5] and later with proton beams [2, 3, 6–10]. These studies have demonstrated that this highly unconventional approach has the potential to be revolutionary, allowing significant sparing of healthy tissue while

maintaining local tumor control. This outcome presents evident clinical prospects once radiation therapy accelerators (medical devices) capable of delivering such treatment become available. Most SFRT techniques spare normal tissues with conventional dose rates [11, 12]. Experiments combining Ultra High Dose Rates (UHDR) with SFRT using Microbeam Radiotherapy (MRT) and mini-beam Radiotherapy (MBRT) have been conducted [13, 14], however, establishing the additive or synergistic effects of combining these techniques is still pending, and separating the flash effect from spatial fractionation in MRT studies has yielded inconclusive results [15].

Conventional radiotherapy has always followed the paradigm of using a uniform beam to deliver the same dose to the target. In contrast, SFRT proposes the possibility of achieving a spatially periodic structure where the transverse profile is articulated in a recurring series of “peaks and valleys” [3]. The optimal structure of this pattern, which refers to the peak-to-valley dose ratio (PVDR) and the spatial extent of the “light and shadow” regions, is still the subject of ongoing research. SFRT can be implemented in two different modes, irrespective of the radiation source employed. The first mode involves using a spatially homogeneous beam and positioning a periodic screen with alternating slits and solid elements, known as the GRID technique [16–19], between the source and the target. This approach, initially proposed by Köhler in 1909 and replicated in experiments using synchrotron light with microbeams, has demonstrated promising results. The second mode utilizes the pencil beam technique, wherein multiple narrow beams, referred to as “pencil beams,” are employed to paint the desired dose distribution pattern. This requires charged particles. Proton accelerators have been particularly suitable for implementing this approach, even though most radiobiological studies have used passive collimation [20–25].

While GRID therapy has found clinical success in palliative applications and mini beam RT is approaching clinical trials with enormous potential, SFRT has not gained widespread adoption not only due to technological limitations of current available devices, as detailed in [21], but also due to the high heterogeneity of tumor coverage, which is in stark contrast to conventional radiotherapy.

Furthermore, it is important to emphasize that the optimal use of SFRT with low energy electron beams is limited to the context of potential flash delivery, as extensively discussed in [2]. Therefore, the study of the combined flash and mini-beam effect is a key aspect, and this present work serves as a foundation for future research.

A dedicated beam with the ability to independently adjust its parameters is essential not only for advancing radiobiology research but also for facilitating accurate dosimetric studies, as was the case of flash [26]. This capability has been demonstrated to be valuable in previous research involving electron Ultra High Dose per Pulse (UHDP) beams, where fundamental dosimetry solutions have been developed [27–30].

In fact, for low energy electrons an immediate application in the clinic would be Intra Operative Radiotherapy (IORT). However, the study of the mini-beam effect is fundamental in the perspective of Very High Energy Electron (VHEE) irradiation, since the mini-beam effect may play a major role since VHEE will most likely be delivered via pencil beams, that allow the possibility of “painting” the spatial dose distribution [18].

In this paper, we explore the potential of using electron mini-beams in conjunction with flash irradiation. The flash effect, as

described in previous studies [31–33], is observed when irradiation times are shorter than 0.1–0.2 s, and the average dose rate exceeds 40–100 Gy/s. The mini-Beam distribution, correlated to a tissue sparing effect [3], is characterized by a spatial dose distribution with alternating peaks and valleys, where the full width at half maximum (FWHM) of the peaks is less than 1 mm, and the distance between successive peaks is 2–4 times the FWHM. Although the combination of these two techniques has not been experimentally evaluated yet, it has the potential for a synergistic action, resulting in a significant reduction in the side effects of radiation therapy while maintaining therapeutic efficacy. Therefore, we can deliver irradiation in four different modalities (Conventional, flash, mini-Beam, and mini-beam-flash), independently varying their main parameters, to study the individual and combined effects of these techniques.

2 Materials and methods

2.1 Mini-beam design and generation

To explore the possibility of creating UHDR mini beams, Monte Carlo simulations were carried out reproducing different devised suitable collimator templates, which are passive spatial distribution modulators. These templates are integrated into the beam optic system to achieve mini-beam spatial distributions.

The mini-beam templates have been designed with the GEANT4 [34] version 10.7.2 Monte Carlo code. In particular, the Sordina IORT Technologies S.p.A. (SIT, Italy) ElectronFlash linear accelerator (linac) available at the Centro Pisano for Flash Radiotherapy (CPFR) was fully modelled, in terms of beam optics, geometry and input energy spectra with the Geant4 code, and a dedicated application was recently inserted within the official advanced examples of the Geant4 distribution with the name “eFLASH_radiotherapy.” For each simulation, the GEANT4 standard_opt4 physics list with 10^8 primary particles and 0.1 mm production cuts were set. A water phantom was implemented just after the template to score the dose, using a $0.125 \times 0.125 \times 0.5$ mm³ voxel size.

A fundamental requisite for a useful mini-beam structure, defined as having distinct peaks and valleys with high peak-to-valley dose ratio, is avoiding bleed-through of electrons through the septa, thus high Z materials are needed. In fact, by keeping the thickness equal to the practical range of the electron energy (9 MeV nominal energy), we noticed a significant reduction of PVDR for plastic materials, such as Teflon, as opposed to Tungsten due to septa bleed-through, as shown in Figure 1. Similar results have been reported also in [35] in the context of experimental IORT beam limiting devices.

Thus, we designed the collimators to have a thickness of 5 mm and be made of Tungsten. Each of the designed templates differs mainly in hole structure (grid or planar slits) and center-to-center distance (ctc), so that it is possible to study the effect of varying these parameters.

2.2 Monte Carlo simulations

To further strengthen the reliability of dose evaluations in terms of spatial distribution, we compared the Monte Carlo

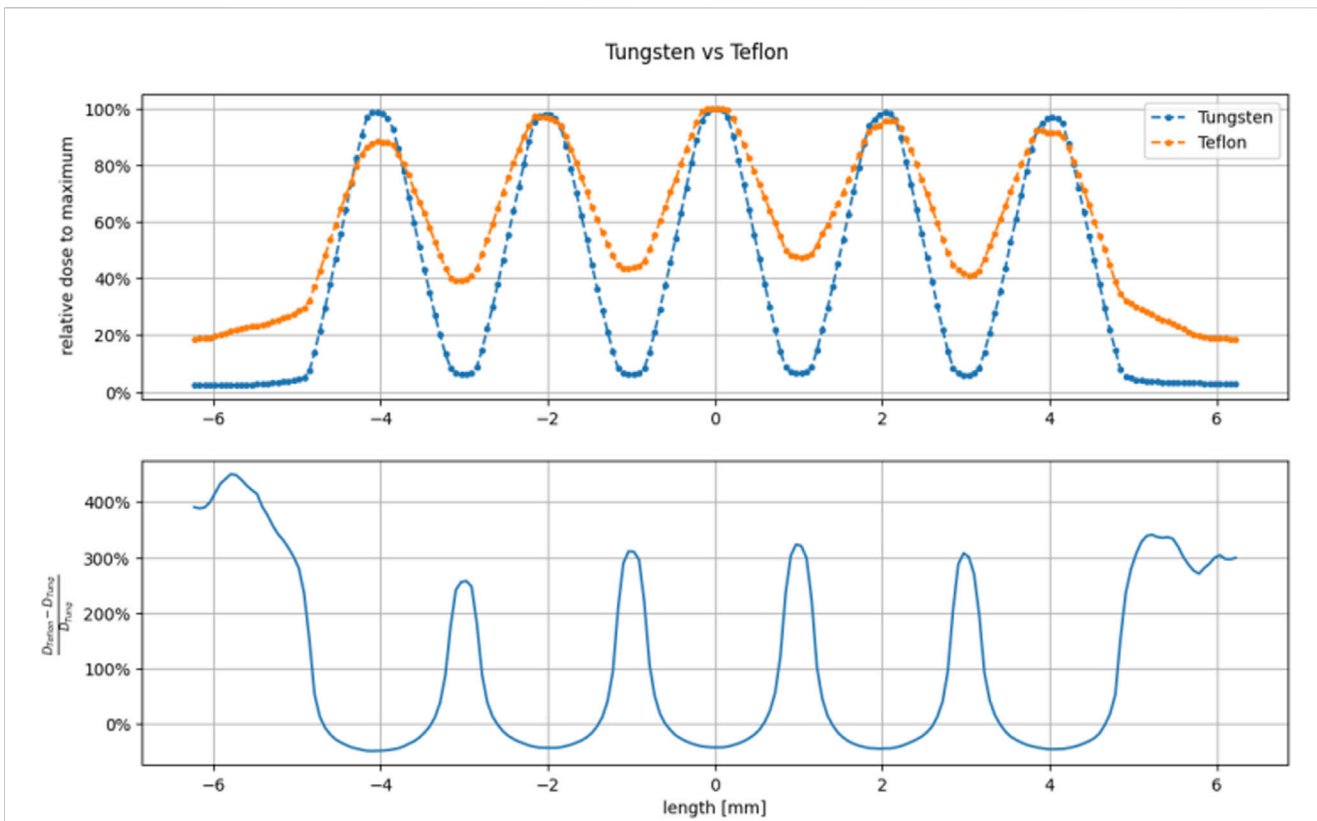


FIGURE 1 Comparison between Tungsten (thickness 5 mm) and Teflon (thickness 25 mm) grid-hole template superficial dose profiles. The thickness has been calculated as the practical range extrapolated from the PDD in the selected material. We can see that Teflon cannot achieve useful peak to valley ratio since the thickness of the walls of the holes (1 mm) are much smaller than the practical range.

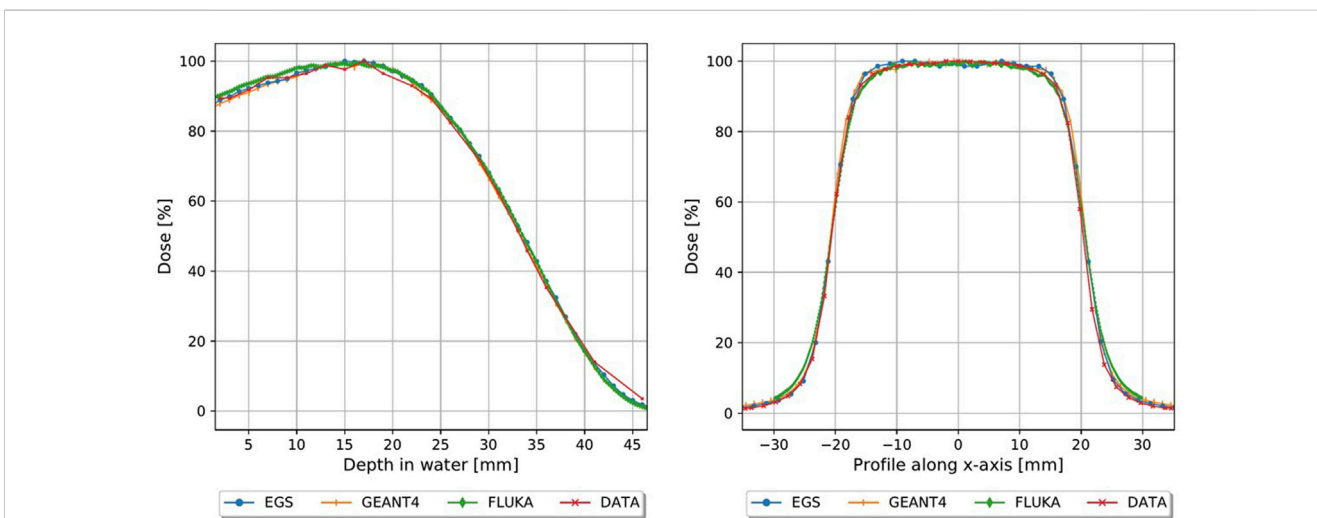


FIGURE 2 Monte Carlo simulations plotted with FlashDiamond measurements, PDD (left) and profiles at R100 (right).

simulations with experimental data acquired with radiochromic gafchromic films, which are currently considered the most suitable dosimeters with high spatial resolution (below 25 μm [36]). However, it is important to acknowledge that these

dosimeters have certain limitations, mostly related to the accuracy of dose reading [37]. To address this, we conducted a comparative analysis by comparing the gafchromic data with results obtained from three independent Monte Carlo simulation

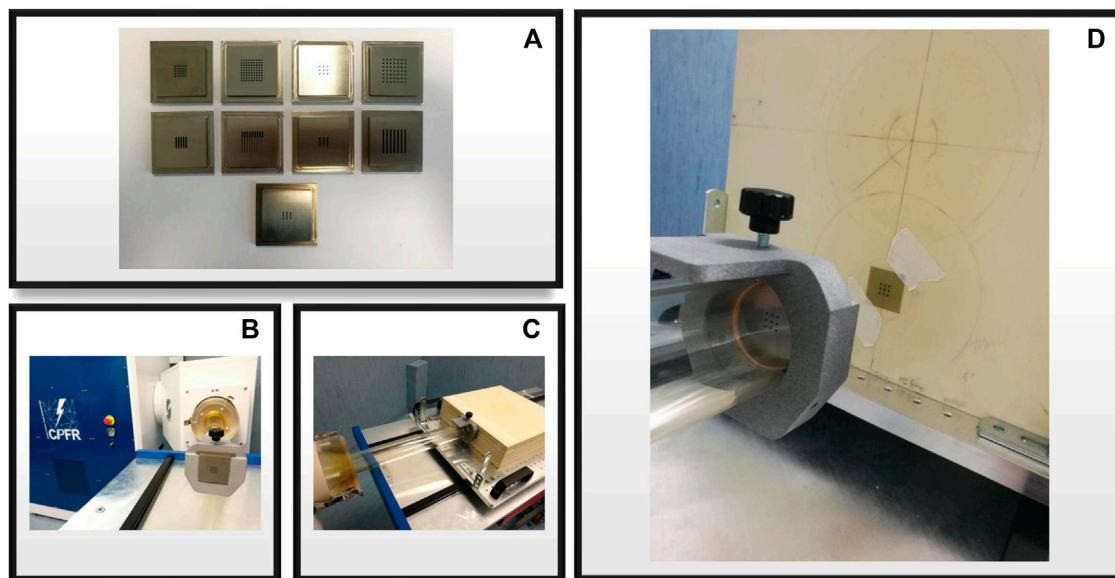


FIGURE 3

(A) All nine template configurations for grid and planar collimators; (B) template mounted on the accelerator's applicator with a nylon 3d printed holder; (C) experimental setup close-up, applicator, template and phantom are shown; (D) close-up of the beam optics, template and irradiated gafchromic film after dose delivery with a visible mini-beam pattern.

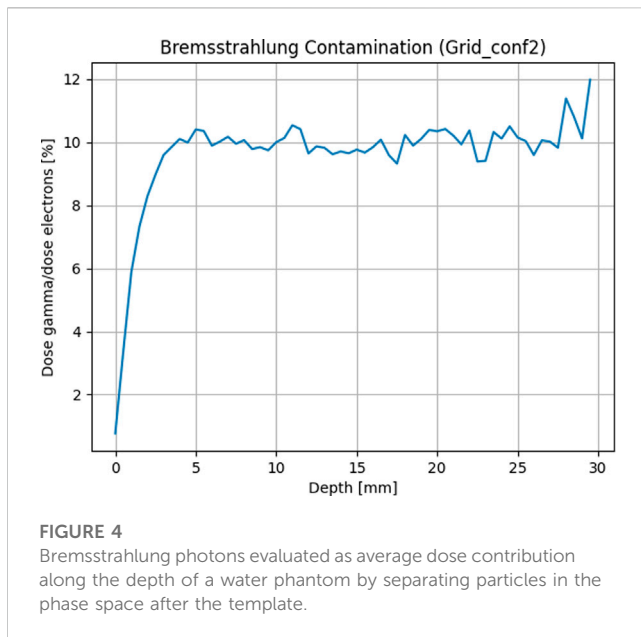
TABLE 1 Mini-beam template geometric parameters.

Configuration name	Hole/slit dimensions (mm ²)	# of holes/slits	Center to center distance (ctc) (mm)
Grid_conf1	1 × 1	25	2
Grid_conf2	1 × 1	9	3
Grid_conf3	1 × 1	81	2
Grid_conf4	1 × 1	49	3
Planar_conf1	1 × 9	5	2
Planar_conf2	1 × 7	3	3
Planar_conf3	1 × 17	9	2
Planar_conf4	1 × 17	7	3
Planar_conf5	0.5 × 6	3	3

codes: Electron Gamma Shower by the National Research Council Canada (EGSnrc) [38], Geant4, and FLUKA [39, 40]. The realistic input energy spectrum and spatial distribution before the Titanium window were provided by SIT Sordina company. To validate the accuracy of the simulations, we compared them with measurements obtained using the FlashDiamond [41] detector, including PDD and dose profile measurements. In order to obtain absolute dose values from the simulations, we normalized the mini-beam dose percentage with respect to an open field measurement at the buildup region. This normalization was then scaled to the desired depth using the PDD curve.

2.2.1 EGSnrc

The simulations performed using EGSnrc consisted of two parts, using the beamline simulation code (BEAMnrc) and the dose deposition simulation code (DOSXYZnrc). At first, a phase space was obtained using BEAMnrc, then the dose was calculated using DOSXYZnrc in the water phantom. The voxel size was adjusted to closely match the dots per inch (DPI) scan setting of the gafchromic film and the film thickness. With BEAMnrc, we modeled the accelerator optics starting from the titanium window down to the tungsten template. To include all the beam optic materials, we expanded the default International Commission on Radiation Units and Measurements (ICRU) material data file (ICRU521) using the EGSnrc graphic interface.



2.2.2 GEANT4

We also recreated the EGSnrc setup using GEANT4 to compare it with the gafchromic measurements. We used the same input spectrum and designed the simulated setup to mimic the modeling approach in EGSnrc. However, in GEANT4, we directly calculated the dose in a water phantom without generating a phase space.

2.2.3 FLUKA

In addition, we replicated the EGSnrc configuration using FLUKA as an alternative comparison to the measurements. Keeping the input spectrum unchanged, the simulated setup closely emulates the modeling methodology employed in EGSnrc. Like in GEANT4, in FLUKA, the dose calculation was performed directly in a water phantom, eliminating the need for generating a phase space.

2.3 UHDP-mini-beam by using the triode-gun electronflash linac

The flash linac adopted is the Triode-Gun equipped ElectronFlash manufactured by SIT, available at the CPFR in Pisa, Italy. It operates at energies of 7 and 9 MeV, delivering a dose-rate of up to 5,000 Gy/s. Field size and Normal Treatment Distance (NTD) is achieved by means of passive collimation with polymethyl methacrylate (PMMA) cylindrical tubes, called applicators. Each applicator is uniquely identified by its diameter (ex. Applicator 100 mm, 50 mm, 40 mm, etc.). With this linac it's possible to achieve flash conditions and vary the main parameters (Dose per pulse, pulse width, pulse repetition frequency ecc.) independently from one and another, without changing experimental setup or beam collimation [42].

2.4 Experimental setup

Measurements were performed using the ElectronFlash linac in 9 MeV mode with a 40 mm diameter applicator. For electron

flash, beam monitoring is critical as conventional systems based on ionization chambers are not compatible with the high beam current [43]. As in the case of flash with protons [44], beam monitoring is a topic of great interest which requires non-conventional solution. In fact, ElectronFlash comes equipped with a IEC 60601-2-1 [45] compliant beam current transformer (ACCT) based monitoring system, which correlates the beam current to delivered dose in the form of monitor units (MU) displayed on the machine human interface (HMI) system. After an initial MU calibration check without the mini-beam template, we fixed the pulse width (t_p), pulse repetition frequency (PRF), and dose per pulse to ensure each gafchromic film received the same dose at the open field electron beam build up depth (R100). The EBTXD gafchromic films [36] were analyzed with an Epson Expression scanner 10000XL after 48 h, 254 DPI with the background subtracted from a pre-irradiation scan. From the scan, we converted the optical density (OD) to dose with previously measured dose calibration curves. Each film was irradiated accumulating a total dose of 23.5 Gy at R100, using a dose per pulse of about 0.2 Gy/p, t_p of 4 μ s and a PRF of 50 Hz.

The mini-beam templates were attached to the applicator using a 3D-printed nylon holder, and the films were placed between slabs of a plastic water phantom to measure the dose profile at each depth. The total phantom dimensions are $30 \times 30 \times 15$ cm³, with singular slab thickness ranging from 1 mm to 1 cm. For each template we evaluated the main characteristics, such as ctc, PVDR, irradiated surface, and the effect of hole type and dimensions.

As by design, the nylon holder does not extend beyond the template, ensuring that only the tungsten is in contact with the phantom. For each template, we measured the dose profile at various depths ranging from 0 up to 4 mm water depth, using 25.6×25.6 mm² films positioned perpendicular to the beam axis in the water equivalent phantom. In the case of GRID templates, we sampled additional points along the depth to evaluate the PDD (percentage depth dose). For the PLANAR templates, a 51.2×51.2 mm² film was placed parallel to the beam axis at a depth of 4 mm, with normalization to the entrance dose measured using a 25.6×25.6 mm² film positioned perpendicular to the beam axis during the same irradiation.

3 Results

3.1 Monte Carlo validation

Each of the Monte Carlo codes show good agreement with the open field experimental data in the water phantom, as shown in Figure 2. Gamma index analysis [46] was performed for the Monte Carlo curves, comparing each one with the FlashDiamond data. With a Dose Difference (DD) of 3% and a Distance to Agreement (DTA) of 3 mm, we have obtained over a 90% agreement for the curves across all Monte Carlo codes. We decided on these values of DD and DTA due to the dose response and spatial resolution of the FlashDiamond detector [41]. Thus, we can reasonably use the Monte Carlo codes for a comparison with the experimental data.

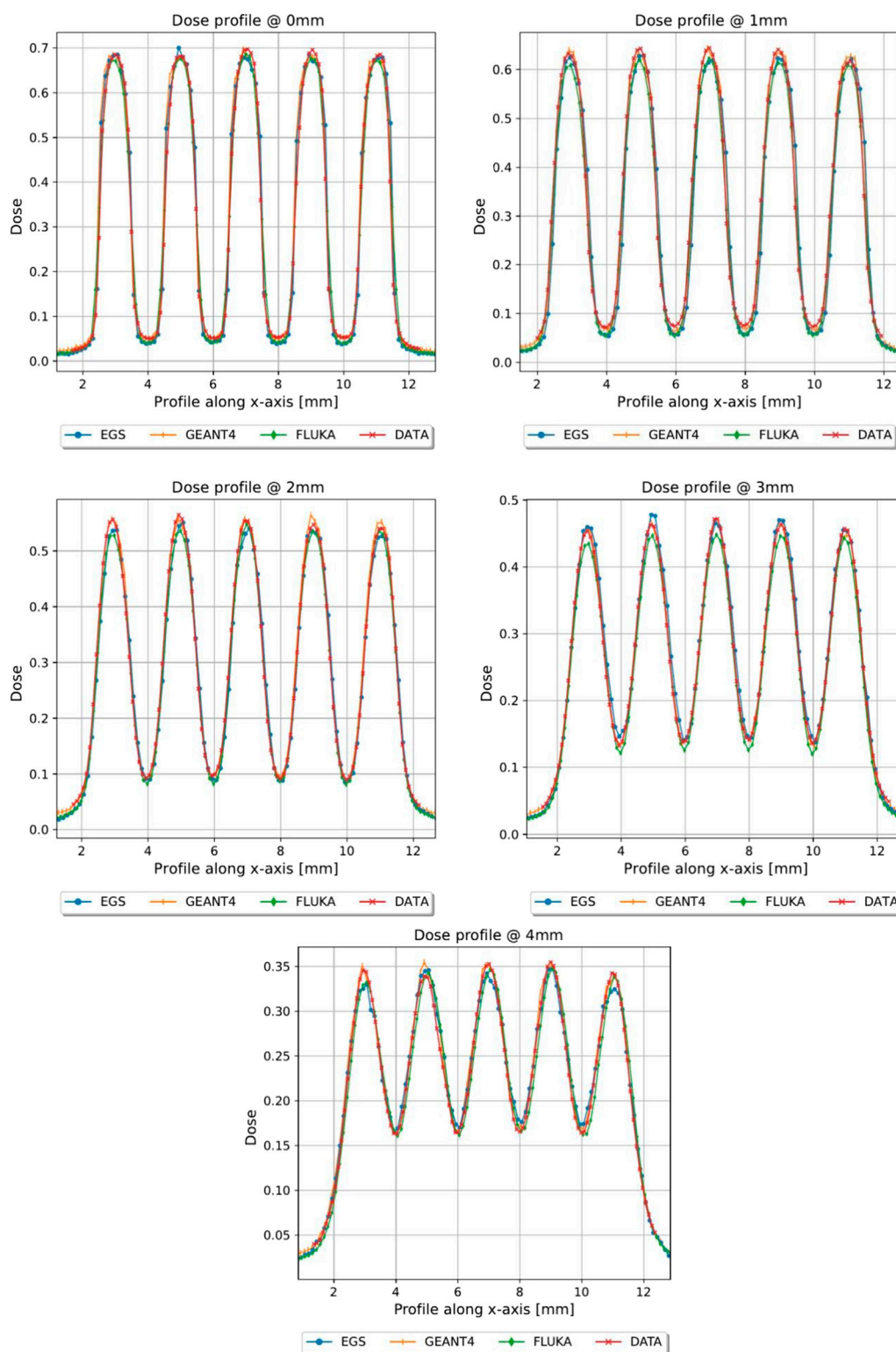


FIGURE 5

Dose profiles, normalized to the dose value of the open field at R100, of the first 4 mm in depth of grid conf1, Monte Carlo vs. Gafchromic film.

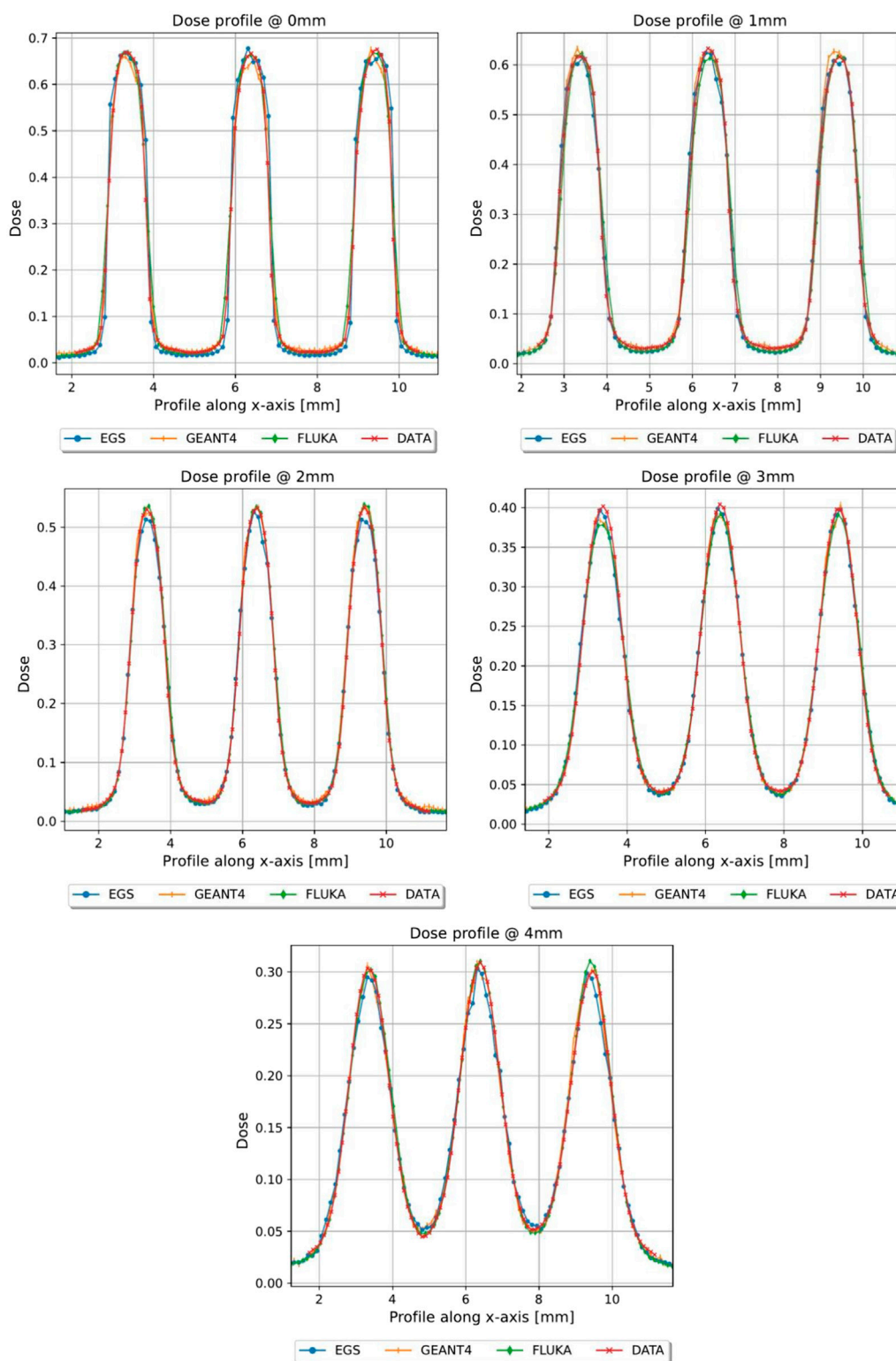
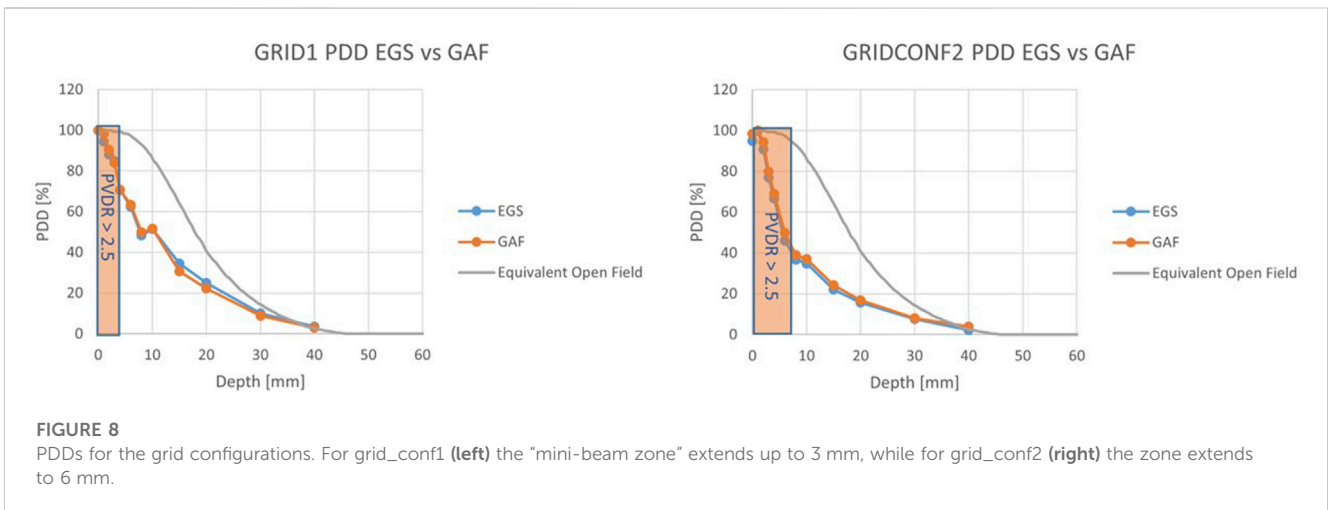
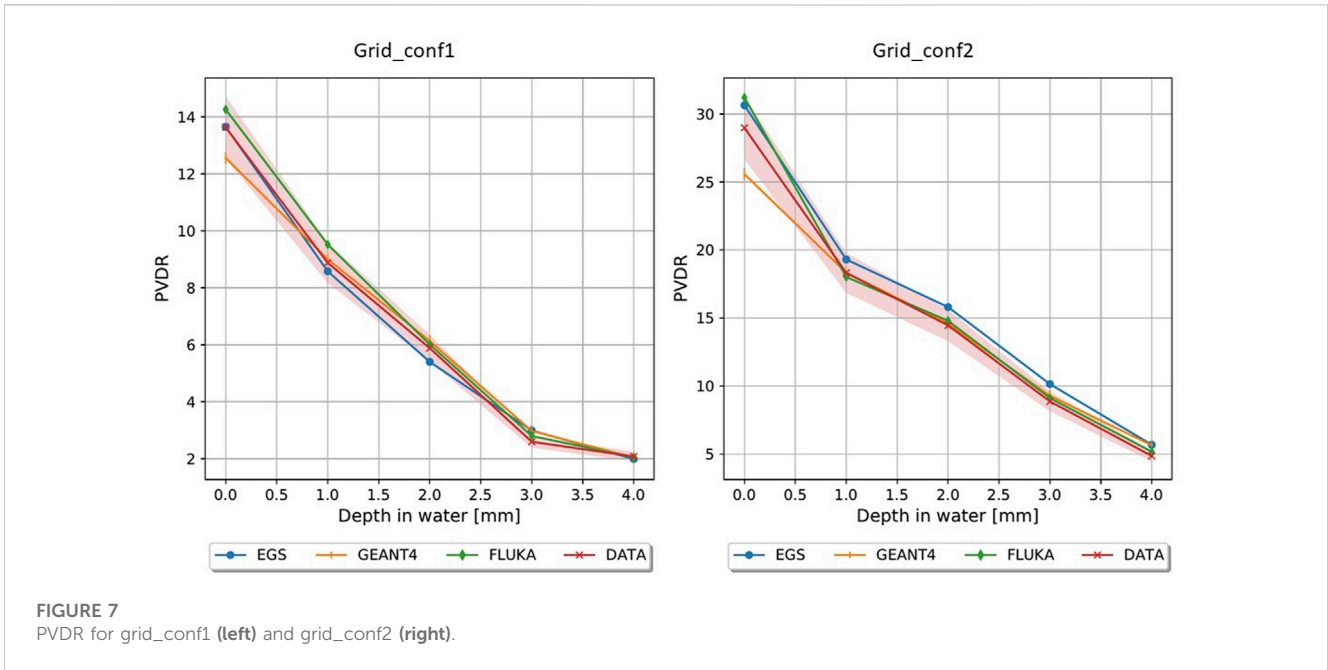


FIGURE 6

Dose profiles, normalized to the dose value of the open field at R100, of the first 4 mm in depth of grid conf2, Monte Carlo vs. Gafchromic film.



3.2 Template construction and characteristics

The final collimators which have been constructed are shown in Figure 3, which also shows a picture of the template attachment on the PMMA applicator and a close-up of the setup, and in Table 1 the main geometric characteristics of each one is summarized. Each one is $5 \times 5 \times 0.5 \text{ cm}^3$ and is made of Tungsten, to ensure minimal electron septa bleed-through and maintain a useful mini-beam structure. Since the templates are made of a high Z material, it is useful to evaluate secondary radiation and its effect on the mini-beam distribution. As shown in Figure 4, Bremsstrahlung photons, evaluated in the worst-case scenario of the grid_conf2 template (the one with least open volume), constitute less than 10% for depth up to 5 mm of the dose, which is deposited mostly at depths beyond the desired mini-beam structure. A rough approximation can be made by utilizing the radiative stopping powers listed in the National Institute of Standards and Technology (NIST) tables for Tungsten

[47]. This calculation overlooks any directional or geometric aspects of the beam and yields an estimate of approximately 10% of the maximum dose attributable to radiative loss.

3.3 Experimental measurements and comparison with Monte Carlo simulations

The templates used for measurement, referenced in Table 1, are as follows:

1. Grid_conf1
2. Grid_conf2
3. Planar_conf1
4. Planar_conf3

In Figures 5, 6, the transversal dose profiles obtained experimentally, at the center of the beam from 0 mm up to

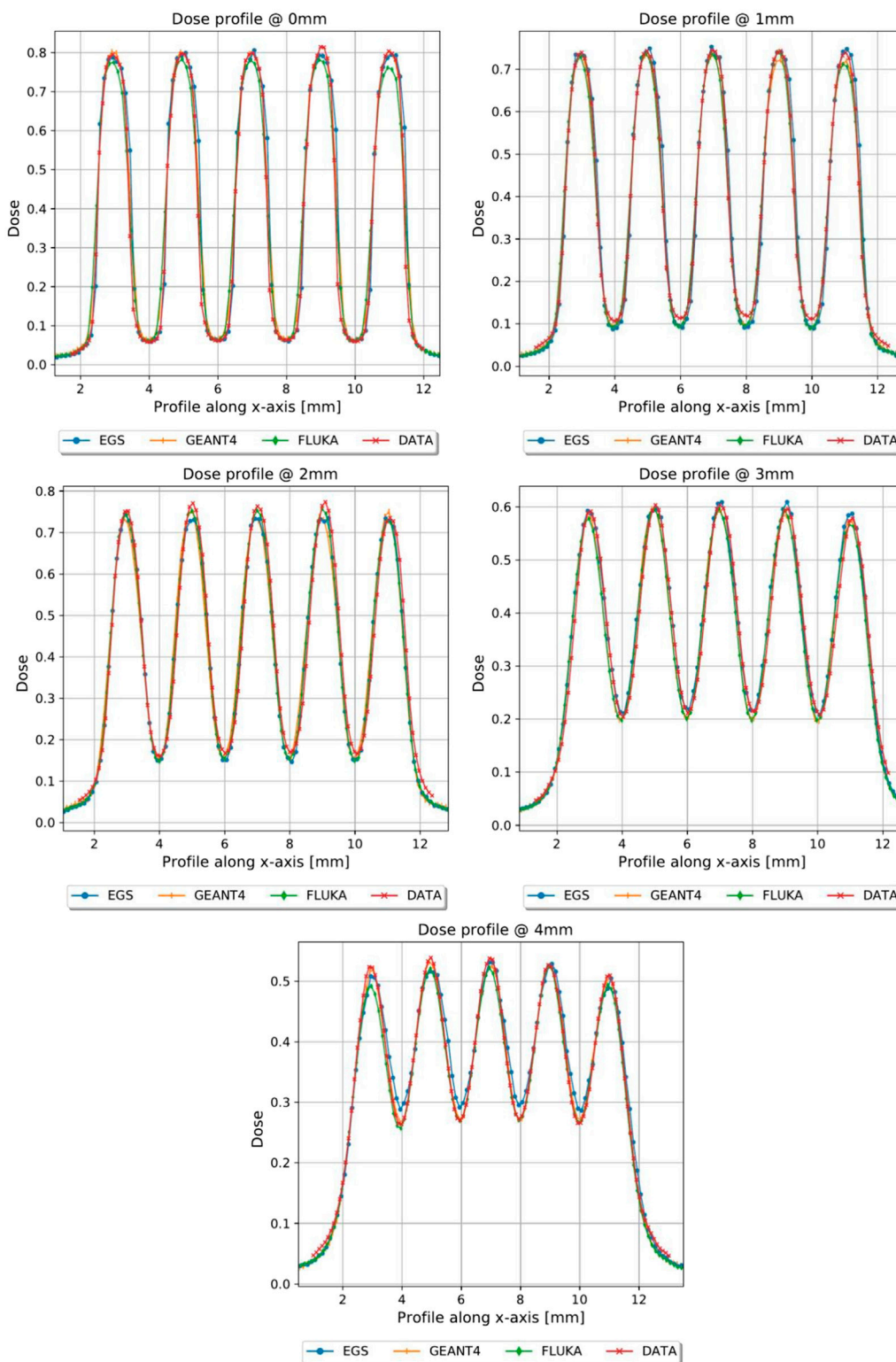


FIGURE 9
Dose profiles, normalized to the dose value of the open field at R100, of the first 4 mm in depth of planar_conf1, MC vs. GAF.

4 mm in water depth are presented for the grid configurations, grid1 and grid2, along with the expected values from the Monte Carlo (MC) simulations. For all profiles, the dose is normalized to the dose at R100 of the open 40 mm field. To evaluate the agreement

of the Monte Carlo simulations, we used gamma index analysis with DD of 6% and DTA of 2 mm to achieve an agreement greater than 93% with the experimental data at each depth. The choice of DD and DTA is due to the high dose uncertainty and high spatial resolution

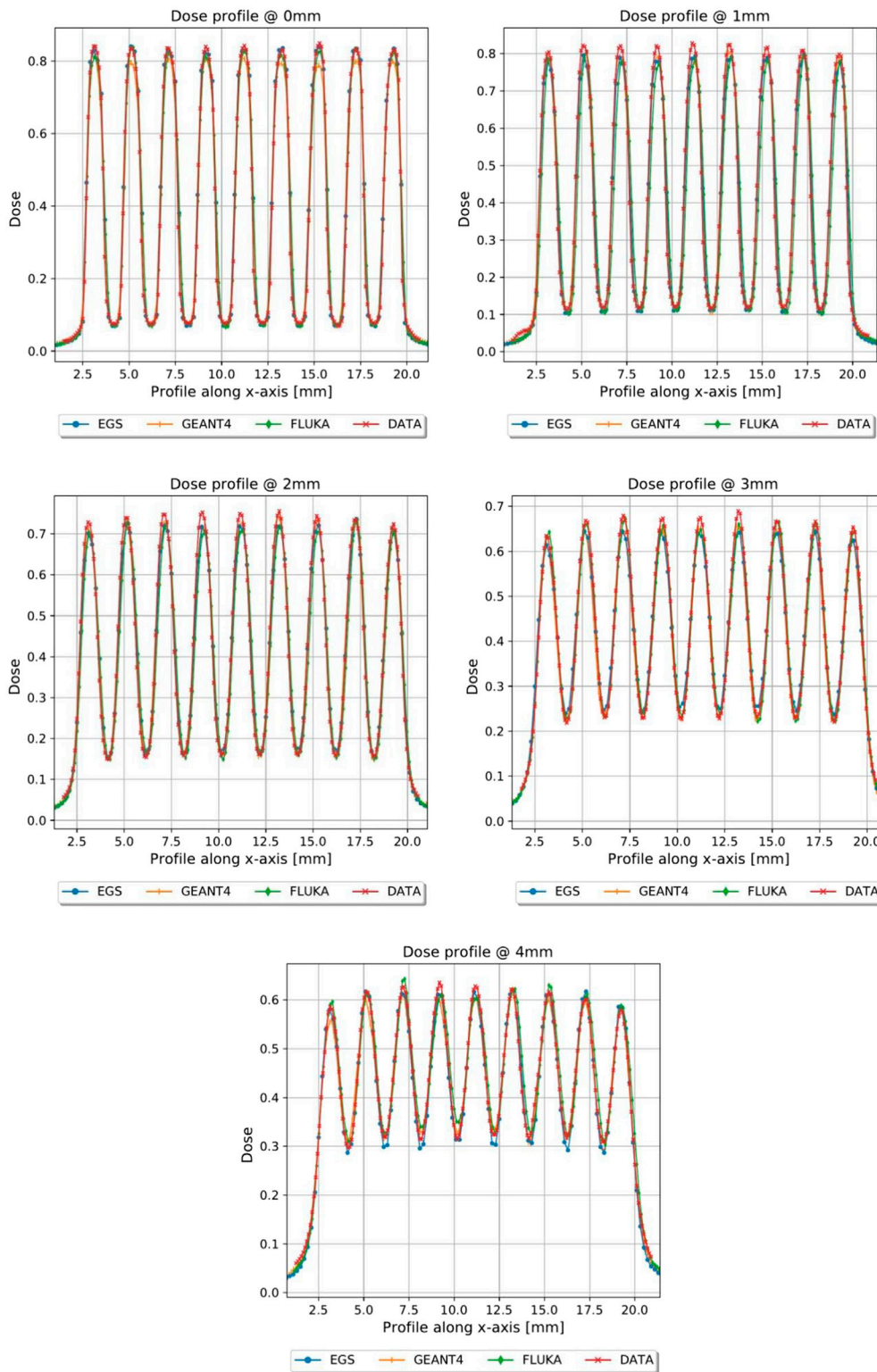


FIGURE 10
Dose profiles, normalized to the dose value of the open field at R100, of the first 4 mm in depth of planar_conf3, MC vs. GAF.

of gafchromic films, as pointed out in [36, 37, 41]. The agreement is further observed in the PVDR values reported as function of the depth depicted in Figure 7. As anticipated, a higher ctc results in a significantly greater PVDR, as the increased septa blocks adjacent

hole contamination, leading to a substantial decrease in the valley dose at the expense of a lower average dose. Figure 8 illustrates the average dose PDDs (percentage depth dose) and the depth at which PVDR > 2.5 for the grid templates. As comparison, also the

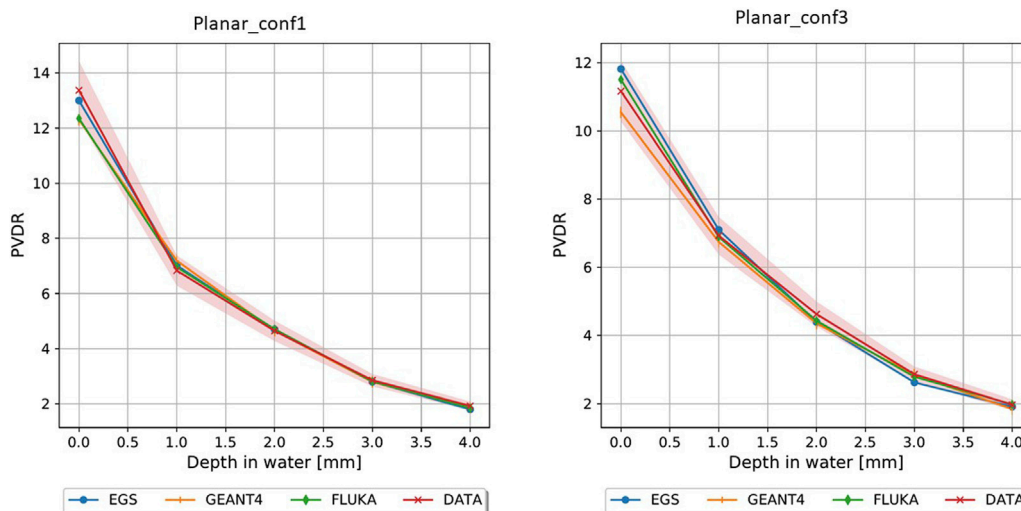


FIGURE 11 PVDR planar_conf1 (left) and planar_conf3 (right).

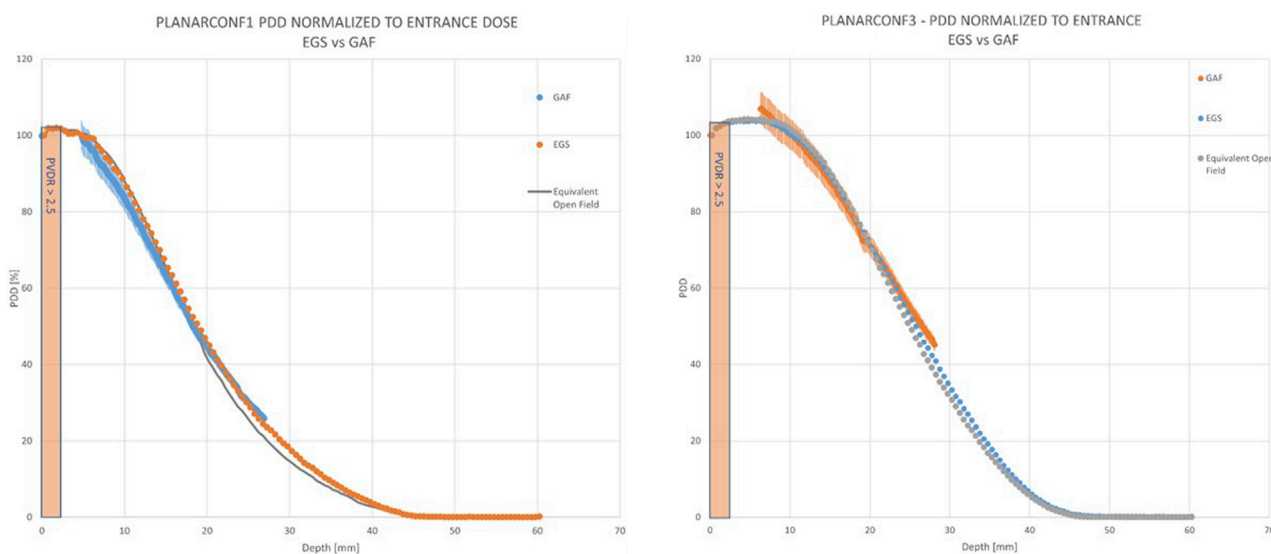


FIGURE 12 PDDs for the planar configurations. For planar_conf1 (left) the “mini-beam zone” extends up to 3 mm, which is the same as planar_conf3 (right).

equivalent open field (obtained as EGSnrc simulated tungsten slab of 5 mm with an aperture of the dimensions of the hole distribution, in this case 1 cm) is shown. We can observe that a greater ctc considerably expands the “mini-beam zone” (PVDR > 2.5) while creating a steeper PDD. This mini-beam structure covers a zone up to the depth of the 50% dose deposition (R50) in water, while reducing the total penetration of the dose distribution for greater depth. This makes the grid template particularly interesting, as it is possible to obtain a large area of “mini-beam” characteristics, with the possibility of performing *in vivo* experiments on organs and tissues of dimensions of 5–6 mm. Figures 9, 10 present the profiles for the planar templates, showing a good agreement also above 93%

of gamma index with DD 6% and DTA 2 mm. In Figure 11, the PVDR as a function of depth for the planar configuration is displayed. Apart from a slight discrepancy (15%) at the entrance, where a small increase in the valley dose can significantly impact the PVDR, the greater field size does not affect significantly the mini-beam distribution along the depth. Figure 12 shows the PDDs normalized to the entrance dose and “mini-beam zone” (PVDR > 2.5) for the templates, compared to the equivalent open beam configuration (2 × 2 cm² field). Since the positioning of the film parallel to the beam resulted in an underexposure after the R50 beyond the constructor’s tolerances, that part is omitted on the graph. As expected, with an increase in field size, we approach an

TABLE 2 Mini-beam characteristics for each template.

Template	Depth (mm)	PVDR	D _p /D _{open} (%)	D _v /D _{open} (%)	D _{avg} /D _{open} (%)	Ctc (mm)	FWHM (mm)	R50 (mm)
Grid_conf1	0	13.6	70.8	5.2	37.9	2	0.9	9
	1	8.9	63.9	7.2	35.8	2	1	
	2	5.9	56.5	9.6	33.2	2	1.1	
	3	2.6	47.4	18.2	31.6	2	1.2	
	4	2.1	30.9	14.7	25.1	2	1.5	
Grid_conf2	0	28.9	68.9	2.4	25.1	3	0.9	6
	1	18.3	60.9	3.3	24.5	3	1	
	2	15.3	49.7	3.5	23.7	3	1	
	3	9.5	54.7	4.3	20.1	3	1.2	
	4	6.5	40.2	4.6	17.1	3	1.3	
Grid_conf3	0	14	69.2	4.9	38.5	2	1	10
	1	9.6	58.5	6.1	33.9	2	1	
	2	5.7	51.1	9.0	29.4	2	1.1	
	3	3	39.3	13.1	26.0	2	1.2	
	4	1.9	30.2	15.9	22.9	2	1.5	
Grid_conf4	0	32.8	64.1	2.0	22.3	3	1	6
	1	24.2	53.5	2.2	19.8	3	1	
	2	16.6	47.2	2.8	19.0	3	1.1	
	3	9.6	35.7	3.7	16.5	3	1.2	
	4	5.4	26.1	4.8	14.2	3	1.5	
Planar_conf1	0	13.4	81.8	6.1	39.9	2	0.9	19
	1	6.8	72.9	10.7	39.6	2	1	
	2	4.7	73.6	15.7	39.4	2	1.1	
	3	2.9	59.2	20.4	38.2	2	1.3	
	4	1.9	52.2	27.5	37.8	2	1.6	
Planar_conf2	0	28.4	84.0	3.0	32.9	3	1	15
	1	18.6	74.8	4.0	31.0	3	1	
	2	13.7	66.3	4.8	29.7	3	1.1	
	3	8	64.0	8.0	29.0	3	1.3	
	4	4.7	53.2	11.3	27.8	3	1.5	
Planar_conf3	0	11.2	84.3	7.5	41.3	2	0.9	26
	1	6.9	81.9	11.9	43.7	2	1	
	2	4.6	74.0	16.1	43.1	2	1.1	
	3	2.9	67.6	23.3	42.1	2	1.3	
	4	2	63.1	31.6	42.8	2	1.5	
Panar_conf4	0	25	82.0	3.3	31.4	3	0.8	15
	1	15.1	78.6	5.2	32.4	3	0.9	
	2	10.2	70.5	6.9	31.2	3	1	

(Continued on following page)

TABLE 2 (Continued) Mini-beam characteristics for each template.

Template	Depth (mm)	PVDR	D _p /D _{open} (%)	D _v /D _{open} (%)	D _{avg} /D _{open} (%)	Ctc (mm)	FWHM (mm)	R50 (mm)
	3	6.8	64.5	9.5	28.8	3	1.1	
	4	4	55.3	13.8	28.8	3	1.3	
Planar_conf5	0	28	60.5	2.2	14.9	3	0.5	17
	1	20.1	52.0	2.6	14.8	3	0.6	
	2	15.1	38.9	2.6	12.7	3	0.7	
	3	10.3	31.6	3.1	9.5	3	0.9	
	4	6.3	25.7	4.1	9.1	3	1.2	

TABLE 3 MINI-BEAM-CONV and MINI-BEAM-FLASH characteristics for PEAK and AVERAGE dose points.

Template	Average DR @peak (Gy/s)	DPP@peak (Gy/p)	DR_pulse@peak (kGy/s)	Average DR @average (Gy/s)	DPP@average (Gy/p)	DR_pulse@average (kGy/s)
Grid_conf1	(0.09, 1404.36)	(0.09, 5.73)	(21.36, 1431.70)	(0.05, 787.44)	(0.05, 3.21)	(11.98, 802.77)
Grid_conf2	(0.08, 1346.26)	(0.08, 5.49)	(20.48, 1372.48)	(0.03, 541.49)	(0.03, 2.21)	(8.24, 552.04)
Grid_conf3	(0.08, 1301.5)	(0.08, 5.31)	(19.8, 1326.85)	(0.05, 760.0)	(0.05, 3.1)	(11.56, 774.8)
Grid_conf4	(0.07, 1206.5)	(0.07, 4.92)	(18.35, 1230.00)	(0.03, 446.5)	(0.03, 1.82)	(6.79, 455.2)
Planar_conf1	(0.10, 1613.58)	(0.10, 6.58)	(24.54, 1645.00)	(0.05, 876.43)	(0.05, 3.57)	(13.33, 893.50)
Planar_conf2	(0.10, 1649.20)	(0.10, 6.73)	(25.09, 1681.32)	(0.04, 684.00)	(0.04, 2.79)	(10.40, 697.32)
Planar_conf3	(0.11, 1746.27)	(0.11, 7.12)	(26.56, 1780.28)	(0.06, 930.08)	(0.06, 3.79)	(14.15, 948.20)
Planar_conf4	(0.10, 1683.4)	(0.10, 6.87)	(25.61, 1716.18)	(0.04, 693.50)	(0.04, 2.83)	(10.55, 707.01)
Planar_conf5	(0.07, 1178.00)	(0.07, 4.81)	(17.92, 1200.94)	(0.02, 334.40)	(0.02, 1.36)	(5.09, 340.91)

open field configuration, resulting in greater penetration, as precedingly reported by [18]. Furthermore, in comparison to the grid configuration, the PDD does not significantly differ from the equivalent open field. This makes the planar template interesting from a clinical point of view, as it has a PDD practically identical to that of the conventional beam of equal energy (same depth of treatment) with the possibility of having a mini-beam effect, for example, on the skin, potentially reducing the collateral effects to this organ.

3.4 Mini-beam characteristics for each template

In Table 2, the main mini-beam characteristics are reported for each constructed template. We reported, for each template and each depth, the PVDR, the ratio of the dose at the peak (D_p) and the dose at the buildup depth without the template (D_{open} , which is to be considered a reference dose value), the ratio of the dose at the valley (D_v) and D_{open} , the ratio of the average dose (D_{avg}) and D_{open} , the ctc, the FWHM and finally the R50 (which is the same for each template). For the sake of legibility, uncertainties on single dose measurements are assumed to be 4%, to be appropriately propagated for each derived quantity.

3.5 Mini-beam-flash irradiation parameters

Finally, we present in Table 3 the possible parameters for mini-beam-conv and mini-beam-UHDR irradiation (Average dose rate (DR), dose per pulse (DPP) and dose rate within the pulse (DR_PULSE) at the peak and average mini-beam dose), at the reference depth of 1 mm, for the peak and average dose. All values are reported as ranges in square brackets, indicating the minimum (CONV) and maximum (FLASH) values achievable. These parameters are to be considered indicative, since we can tune the ElectronFlash to obtain a wide range of characteristics in a continuous way, within the limits of proper linac function.

4 Discussion and conclusion

In this study, we have designed, simulated, and finally realized mini-beam and UHDR-mini-beams by using the triode-gun ElectronFlash low energy electrons linac and special tungsten passive template; we have then completely characterized dosimetrically our beams by means of radiochromic films. We have observed a good agreement between film measurements and Monte Carlo simulations. Due to the nature of radiochromic films, a greater discrepancy in the valley is to be expected with respect to the peaks, since the dose is much lower and the region

of interest in which the dose is calculated is quite small. Nevertheless, the comparison with the simulations with various codes gave additional robustness to the experimental measurements. Across all Monte Carlo codes, the obtained results exhibited comparable outcomes with only minor discrepancies in valley dose. The relative dose values and PVDR measurements were in almost all cases within the uncertainty range of the gafchromic films, with discrepancies most likely due to different statistics and variation between the unique code parameters. Nevertheless, these findings affirm the capability of each of the evaluated codes to accurately simulate the linac geometry and templates. We have observed that increasing the ctc results in a decrease in the valley dose, leading to higher PVDR and a larger mini-beam zone. Having both grid and planar configurations at our disposal allows for great versatility in experiments. A grid configuration reduces dose deposition along the depth and enables a higher PVDR and mini-beam zone with fixed template parameters. This can be beneficial for *in vivo* irradiations for target up to 5–6 mm of size. On the other hand, the planar geometry diminishes the PVDR but achieves a higher peak dose and deeper dose penetration. This may be promising from a clinical perspective, as the beam does not lose penetration compared to the open field and has a mini-beam component in the first few mm of depth, which could substantially reduce side effects on the skin. Regarding bremsstrahlung contamination, even in the worst-case scenario, it is not significant within the mini-beam zone, and an appropriate mini-beam structure is maintained. This allows for the use of tungsten as the material of choice, at least in the case of low-energy electrons, as the main cause of mini-beam loss is primarily electron septa crossing, which is avoided due to the high atomic number (Z) of the material. Thus, the proposed method of generating a mini-beam structure offers great versatility. This versatility is particularly important when combined with the flash capabilities of our triode-gun ElectronFlash linac, as it enables the study of both mini-beam and flash effects by independently varying the fundamental spatial and temporal parameters involved in the two effects. This is fundamental for quantitative mini-beam and flash experiments, since we can vary all the fundamental parameters independently one from the other and without altering the experimental setup, guaranteeing a wide range of investigations and great reproducibility. The mini-beam and mini-beam-flash operating beams that we have realized, can be used to carry out radiobiological experiments necessary to study the quantitative dependencies of the flash and mini-beam effects and the beam parameters that characterize them, and to understand their underlying radiobiological mechanisms. These experiments are aimed at optimizing the clinical implementation of flash, mini-beam and possibly mini-beam flash radiotherapy. This will be fundamental also for the future transition to VHEE [48, 49], as mini-beams and the flash effect will play a crucial role. We are thus able to continue working towards a robust clinical protocol and evaluate the biological modifying factors necessary for a treatment planning system, both pre-clinical and clinical.

Data availability statement

The original contributions presented in the study are included in the article/Supplementary Material, further inquiries can be directed to the corresponding author.

Author contributions

JP: Writing—original draft, Writing—review and editing. FR: Writing—review and editing. MC: Writing—review and editing. DS: Writing—review and editing. GFe: Writing—review and editing. GFr: Writing—review and editing. LM: Writing—review and editing. GM: Writing—review and editing. VP: Writing—review and editing. YP: Writing—review and editing. FD: Writing—review and editing.

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Conflict of interest

Authors JP, GFe, and LM were employed by Sordina IORT Technologies S.p.A.

The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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