

Calculation of the Coupling Coefficient in Step-Index Multimode Polymer Optical Fibers Based on the Far-Field Measurements

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Using the power flow equation (PFE), this article investigates mode coupling in step-index (SI) multimode (MM) polymer optical fiber (POF). This equation's coupling coefficient was initially fine-tuned so that it could appropriately reconstruct previously recorded far-field (FF) power distributions. The equilibrium mode distribution (EMD) and steady-state distribution (SSD) in the SI MM POF were found to be obtained at lengths $L_c = 15$ m and $z_s = 41$ m, respectively. These lengths are substantially shorter than their glass optical fiber counterparts. Such characterization of the investigated POF can be used in its employment as a part of the communication or sensory system. Namely, the POF's bandwidth is inverse linear function of fiber length (z^{-1}) below the coupling length L_c . However, it has a $z^{-1/2}$ dependence beyond this equilibrium length. Thus, the shorter the coupling length L_c , the sooner transition to the regime of slower bandwidth decrease occurs. It is also important to be able to determine a modal distribution at a certain length of the POF employed as a part of optical fiber sensory system.

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INTRODUCTION

Silica optical fiber has a lot of advantages such as a low loss, lightweight, and high bandwidth [1–4]. Although silica optical fiber has a lot of advantages, is not appropriate for short-distance applications such as automotive [5], local area networks [6], and visible light communication (VLC) applications [7]. On the other hand, POFs are made of poly (methylmethacrylate) (PMMA) [8, 9], polycarbonate (PC) [10, 11], ZEONEX [®] [12, 13], TOPAS [®] [14, 15], poly dimethyl siloxane (PDMS) [16, 17] or hydrogel [18, 19].

Apart from optical fiber material, optical fibers usually have step-index or graded-index refractive index distribution, and can operate in a single-mode or multimode regime. There are several typical commercial POFs, such as Mitsubishi Rayon's Eska Extra EH 4001 SI MM POF [20] and CYTOP[®] graded index (GI) POF [21], both of them could transmit a large number of modes, and have been used for short-range communications and sensing area in fiber optic sensory systems [22]. Additionally, micro-structured polymer optical fiber which was first produced by Argyros' group, is one very specific type of POF [23]. However, due to the transmission loss and butcoupling with commercial silica optical fiber, only one Start-up Company supplied the commercial

1



product [24]. PMMA SI MM POF is still the most common one with lots of commercial products on the market.

Mode coupling has a significant impact on the transmission performance of PMMA SI MM POF. The transfer of power between nearby modes is represented by such coupling. It causes the launched light's angular power distribution to be disrupted [25, 26]. It is caused by the optical fiber's intrinsic perturbation effects (for example, variations of the refractive index distribution and microscopic bends). The angular power distribution is predicted to be influenced by the launch conditions and mode coupling characteristics. Thus, a Gaussian beam launched at an angle θ_0 in respect to the optical fiber axis, at the output end of a short piece of optical fiber, is seen as a sharp ring pattern (the ring diameter depends on θ_0). Mode coupling, on the other hand, causes such a ring pattern to be distorted in longer fibers and finally transformed into disk. The coupling length L_c marks the fiber length at which the ring pattern of the highest order guiding mode evolved into disk, indicating that an EMD is established. Coupling can fully complete at optical fiber length z_s ($z_s > L_c$), which is referred to as an SSD.

The light transmission performance of PMMA SI MM POF is investigated in this work. The modal power diffusion model is used to obtain the coupling coefficient D. In this way, we computed the L_c required to accomplish the EMD and the length z_s to achieve the SSD. By such a characterization of mode coupling process in an optical fiber one can predict at which length $z = L_c$ one can expect that bandwidth decrease with length would start to decelerate, as explained in more details in Results and Discussion. It is also important to know the state of mode coupling in an optical fiber employed as a part of optical fiber sensory system, especially in terms of the modal distribution at certain fiber length.

POWER FLOW EQUATION

The Gloge's PFE has the following form [27]:

$$\frac{\partial P(\theta, z)}{\partial z} = -\alpha(\theta)P(\theta, z) + \frac{D}{\theta} \frac{\partial}{\partial \theta} \left(\theta \frac{\partial P(\theta, z)}{\partial \theta}\right) \quad (1)$$

where $P(\theta, z)$ is the angular power distribution, θ is the angle of propagation, z is the distance of the propagation, D is the coupling coefficient (assumed constant [27, 28]) and $\alpha(\theta)$ is the modal attenuation. **Eq. 1** can be reduced to [29]:

$$\frac{\partial P(\theta, z)}{\partial z} = \frac{D}{\theta} \frac{\partial P(\theta, z)}{\partial \theta} + D \frac{\partial^2 P(\theta, z)}{\partial \theta^2}$$
(2)

Steady-state solution of Eq. 2 is given as [28]:

$$P(\theta, z) = J_0\left(2.405\frac{\theta}{\theta_c}\right) \exp\left(-\gamma_0 z\right)$$
(3)

where J_0 is the Bessel function of the first and zero-order, $\gamma_0 [m^{-1}] = 2.405^2 D/\theta c^2$ is the attenuation coefficient. In this work, solved **Eq. 2** using the explicit finite difference method [29].

In our earlier published work, we proposed a method which enables that the coupling coefficient D can be obtained from just two output angular power distributions $P(\theta, z)$ in the case of centrally launched beam [30]. As an alternative, in this work we propose a method for calculating the coupling coefficient Don the basis of the measured FFPs p(x, z), illustrated in **Figure 1**.

RESULTS AND DISCUSSION

We used the PFE (2) to calculate the lengths L_c and z_s for the POF that Ribeiro et al. [31] studied experimentally. This POF (Mitsubishi Rayon's Eska Extra EH 4001) had NA = 0.47, the inner critical angle $\theta_c = 18^{\circ}$ ($\theta_c = 27.4^{\circ}$ measured in air), core









FIGURE 4 The normalized angular power distribution at the end of PMMA SI MM POF, obtained by solving the PFE for four Gaussian launch distributions with input angles $\theta_0 = 0^\circ$ (-), $\theta_0 = 5^\circ$ (- -), $\theta_0 = 10^\circ$ ($\bullet \bullet \bullet$) and $\theta_0 = 15^\circ$ (- $\bullet -$), with (FWHM)_{z=0} = 7° (g represent the analytical SSD).

refractive index $n_{core} = 1.4897$, numerical aperture NA = 0.46 and fiber diameter d = 1 mm. Ribeiro et al. focused a He-Ne laser beam at 633 nm onto the input end of the fiber in their experiment (**Figure 1**). They used a CCD to detect the FF patterns of the fiber output.

The coupling coefficient D is obtained by its fine-tuning in the PFE, thus recreating the Ribeiro et al.'s measured FF patterns in this fiber for lengths of z = 1 m and z = 19 m (**Figure 2**). This necessitated numerically calculating the PFE for various values of D. The value of $D = 9.2 \times 10^{-4}$ rad²/m enabled the best fit between the calculated and measured output power distributions. A Gaussian launch beam with $(FWHM)_{z=0} = 7^{\circ}$ is used in the computations. **Figure 3** shows the output power distribution for a z = 1 and 19 m length POF with an input angle of $\theta_0 = 0^{\circ}$ obtained as numerical solution of **Eq. 2**. The distributions P(x, z = 1 and 19 m, L = 110 mm) are generated from the distributions $P(\theta, z = 1 \text{ and } 19 \text{ m}, L = 110 \text{ mm})$, where $x = L \cdot tg\theta$ and L is the receiving distance.

In **Figure 4**, our numerical results for the normalized output angular power distribution for input angles $\theta_0 = 0, 5, 10, \text{ and } 15^{\circ}$ are shown. At short fiber lengths, a significant mode coupling is observed for low order modes, as can be seen in **Figure 4A**. The EMD is achieved at length $z = L_c = 15$ m (**Figure 4C**), while the SSD is established at $z_s = 41$ m in **Figure 4D**.

The coupling coefficient *D* for the PMMA SI MM POF evaluated in this work is similar to those which we obtained for other investigated POFs ($\sim 10^{-4} \text{ rad}^2/\text{m}$), so the characteristic lengths for coupling, L_c and z_s , are similar (L_c is between $\approx 15-35 \text{ m}$ and z_s is between ≈ 40 and 100 m [25, 29, 30]). We have previously reported that glass optical fibers show the weakest strength of mode coupling ($D \sim 10^{-7}$ to $10^{-6} \text{ rad}^2/\text{m}$), so their SSD lengths z_s are between $\approx 1-10 \text{ km}$ [32].

It is worth noting that the length dependence of the bandwidth of POF is determined by mode coupling behavior. The bandwidth is inverse linear function of fiber length (z^{-1}) below the coupling length L_c . However, it has a $z^{-1/2}$ dependence beyond this equilibrium length. Thus, the shorter the coupling length L_c , the sooner transition to the regime of slower bandwidth decrease occurs [26]. It is obvious that mode coupling has a beneficial influence on bandwidth. Therefore, it is of great interest to characterize mode coupling in an optical fiber in order to predict its transmission characteristics, especially to determine at which coupling length L_c one can expect that bandwidth would start to improve. It is also important to know the state of mode coupling in an optical fiber employed as a part of optical fiber sensory system, especially in terms of the modal distribution at certain fiber length.

CONCLUSION

We investigate a mode coupling along a PMMA SI MM POF previously investigated experimentally by Ribeiro et al. [31]. To appropriately recreate the measured FF patterns reported before, the coupling coefficient *D* in the PFE is tweaked. As a result, the lengths L_c and z_s that characterize the coupling process are obtained. Such characterization of the investigated PMMA SI MM POF can be used in its employment as a part of communication or sensory system. In practice, by characterization of mode coupling in an optical fiber one can predict its transmission characteristics, especially to determine length-dependent bandwidth behavior. Since POF's bandwidth is inverse linear function of fiber length (z^{-1}) below the coupling

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length L_c , while it has a $z^{-1/2}$ dependence beyond this equilibrium length, it is obvious that the shorter the coupling length L_c leads to the sooner transition to the regime of slower bandwidth decrease. It is also important to be able to determine a modal distribution at a certain length of the fiber employed as a part of optical fiber sensory system.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/supplementary material, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

SS: methodology, conceptualization, formal analysis, validation, writing—original draft. RM: conceptualization, formal analysis, funding acquisition, writing—review and editing. AD: formal analysis, validation, writing—review and editing. BD: validation, writing—review and editing. AS: validation, writing—review and editing.

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