



A Chaotic System With Infinite Attractors Based on Memristor

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In this article, a memristor chaotic system is constructed by introducing a cosine function flux control memristor. By analyzing the balance of the system, it is found that there are coexisting attractors, and because of the periodicity of cosine function, the chaotic system has infinite coexisting attractors. The complexity analysis of Spectral Entropy (SE) and CO is used in this paper to intuitively show the complex dynamic characteristics of the system. In addition, the introduction of paranoid propulsion also provides more possibilities for the system in engineering applications. Finally, the digital signal processing (DSP) experiment verifies the correctness of theoretical analysis and numerical analysis.

Keywords: memristor, infinite coexisting attractors, complexity analysis, offset boosting, DSP

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1 INTRODUCTION

In the 1960s, due to the birth of the Lorentz system [1], people started a period of a widespread upsurge in the study of chaos theory. Chaos is a physical phenomenon highly sensitive to initial values. Therefore, in the past decades, people have shown great interest in the creation of chaos. Chaos [2–10] has gradually developed from the study of climate to other fields, such as information science, biology, engineering, finance, and so on. So far, under the research of people, the chaotic system has been widely developed, from the original Lorentz system, to produce many different chaotic systems, including hyperchaotic system [11–17], discrete chaotic system [18–22], memristor chaotic system [23–25] and so on.

As the fourth basic circuit element, the memristor has aroused great interest in the nonlinear field since its birth. In 1971, Professor Chua made a prediction about the existence of the memristor based on the symmetrical structure of circuit elements. In 2008, HP company successfully developed a solid-state memristor that proved Chua's prediction. Therefore, the study of memristor has become a new theoretical branch. In recent years, the combination of the memristor and chaotic system has formed a new memristor chaotic system [26–32], which is also called a new research hotspot.

In addition, the multi-steady state [33–37] of dynamic systems has aroused considerable interest in the nonlinear field. When the chaotic system has the property of a periodic state, the system may have periodic attractor pairs, which are constantly shifted in the same direction, so as to achieve an infinite multi-steady state. This discovery provides a new method for studying the multi-steady state of dynamic systems in the future.

In this paper, we choose an offset boosting [38–43] chaotic system as the basis, and introduce a memristor based on the cosine function to construct a new 5-D chaotic system, which has the following properties: 1) chaos generation; 2) infinite coexistence attractor; 3) and offset boosting. In **Section 2**, the memristor in this system is analyzed, and obtain the “8” curve of proving the memristor. In **Section 3**, the corresponding chaotic phase diagram of the system is obtained through simulation analysis, and the related dynamics analysis of different parameters of the chaotic system is carried out, including the bifurcation diagram, Lyapunov exponential spectrum, infinite coexistence attractors, offset boosting, complexity analysis. In **Section 4**, the chaotic system is simulated by digital simulation and DSP experiment platform, and the reality of the physical existence of the system is verified. Finally, the conclusion of the study is given in **Section 5**.

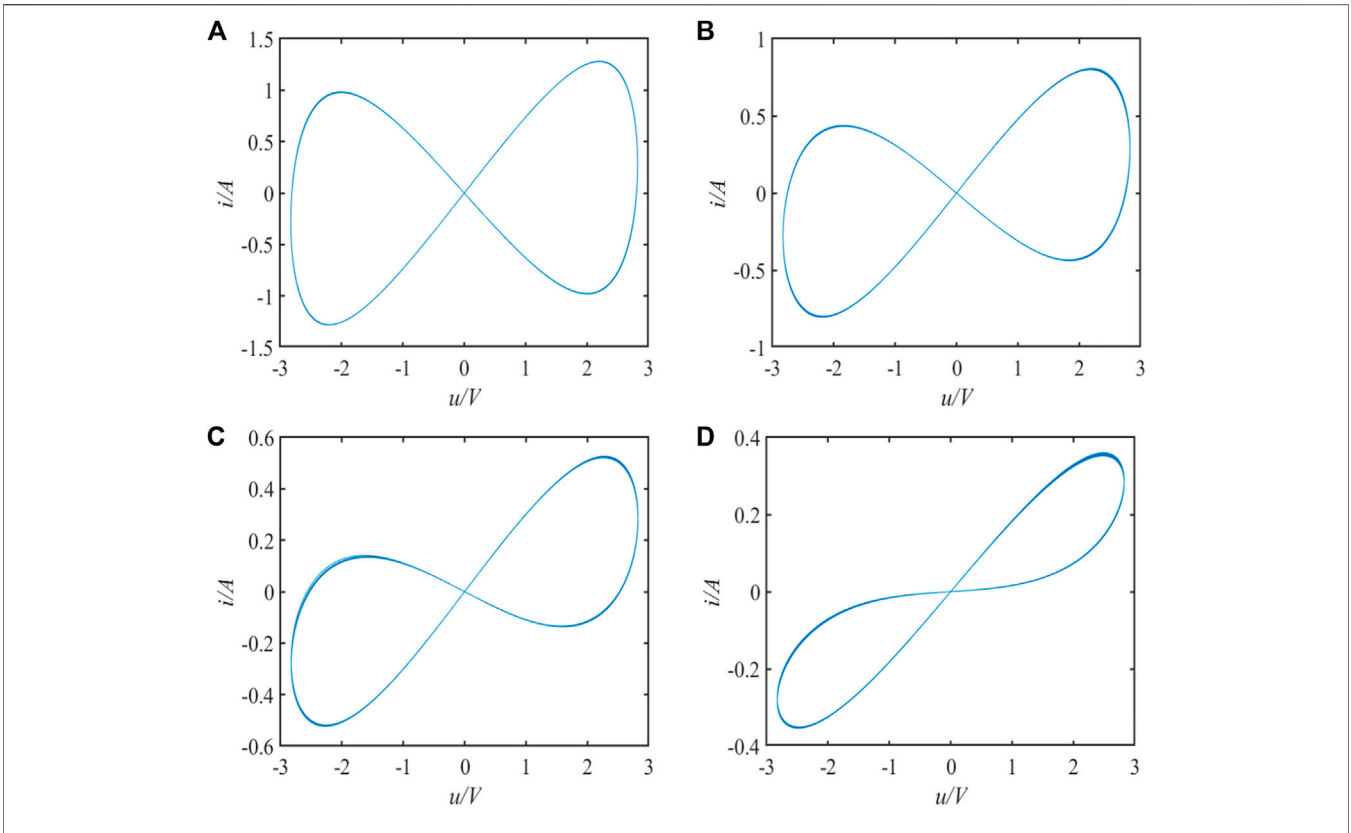


FIGURE 1 | The u - i curve of memristor: **(A)** $f = 0.5$ kHz with $u = 3$ V, **(B)** $f = 1$ kHz with $u = 3$ V, **(C)** $f = 2$ kHz with $u = 3$ V, **(D)** $f = 5$ kHz with $u = 3$ V.

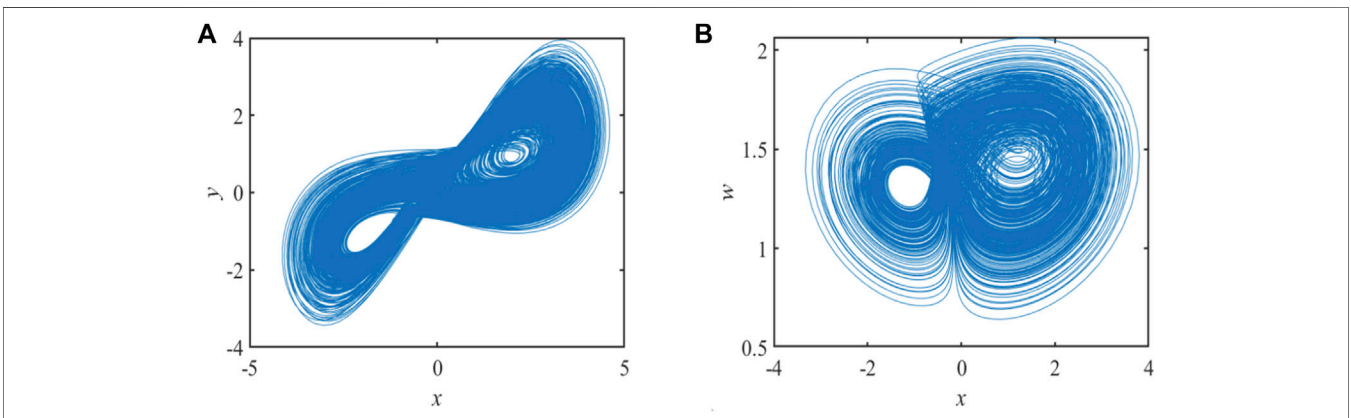


FIGURE 2 | Chaotic attractor of system (3): **(A)** X-Y plane, **(B)** X-W plane.

2 MATHEMATICAL MODEL

2.1 Memristor Model

The memristor model is expressed as follows:

$$\begin{cases} \dot{x} = G(x, y, t)y \\ \dot{y} = H(x, y, t) \end{cases}, \quad (1)$$

where x is this output of the memristor, y represents the state of the memristor, and the functions $G(\cdot)$ and $H(\cdot)$ are particularly relevant to the memristor.

The flux-controlled memristor is from cosine function, and the specific equation is shown as follows:

$$\begin{cases} W(w) = \frac{dq(w)}{d(w)} = \cos(w) \\ i = W(w)y = \cos(w)y \\ \frac{d(w)}{dt} = y^2 - w \end{cases}, \quad (2)$$

Here, the following the “8” curve of the memristor is obtained by inputting different frequency parameters as shown in **Figure 1**.

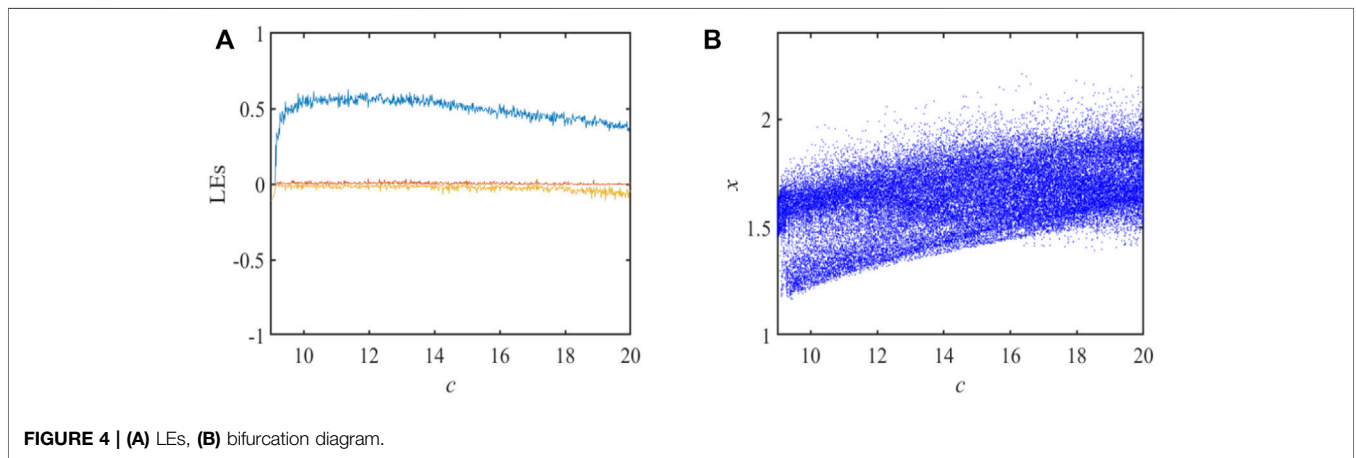
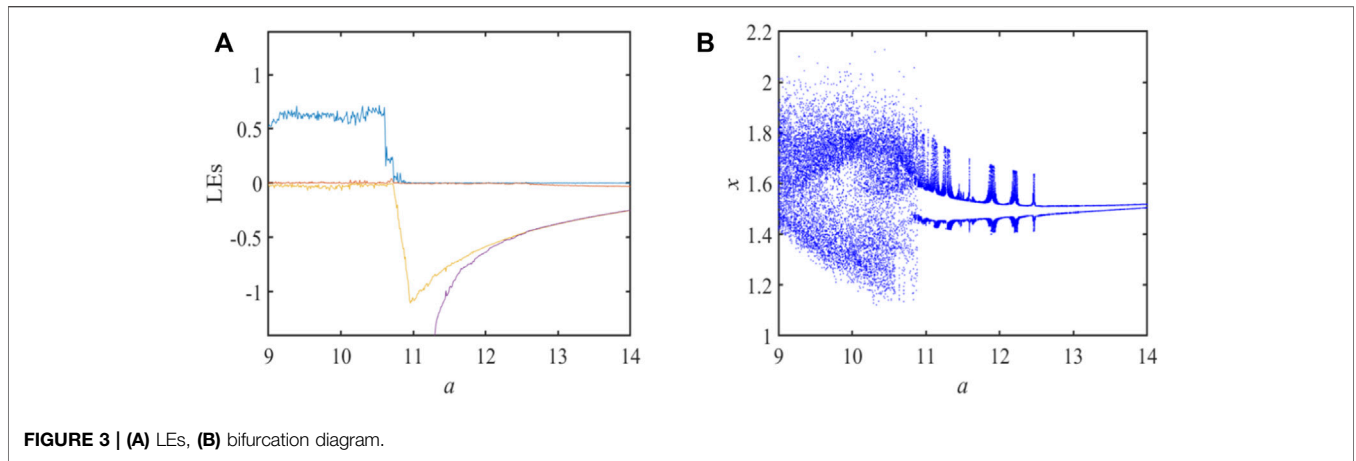


TABLE 1 | System state corresponding to different parameter a when the original conditions is (1, 1, 1, 1, 1).

Range	LEs	State	Range	LEs	State
—	0 - - - -	Divergence	10.85	+0 - - -	Weak chaos
9.00–10.72	+0 - - -	Chaos	10.86	0 - - - -	Period
10.73	0 - - - -	Period	10.87	+0 - - -	Weak chaos
10.74–10.78	+0 - - -	Weak chaos	10.88	0 - - - -	Period
10.79	0 - - - -	Period	10.89	+0 - - -	Weak chaos
10.80–10.83	+0 - - -	Weak chaos	10.90–14.00	0 - - - -	Chaos
10.84	0 - - - -	Period	—	—	—

TABLE 2 | System state corresponding to different parameter c when the original conditions is (1, 1, 1, 1, 1).

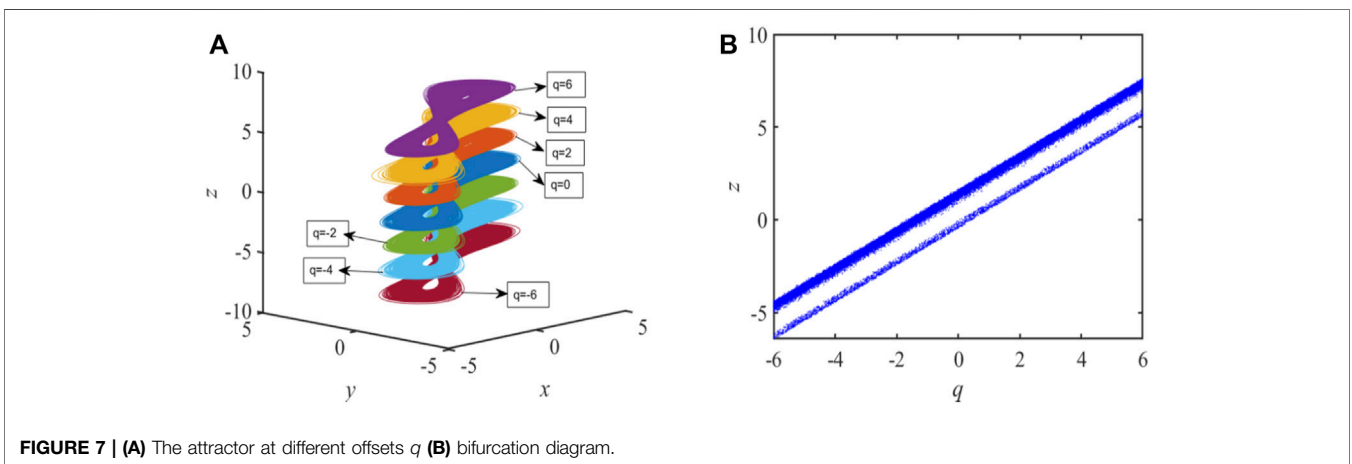
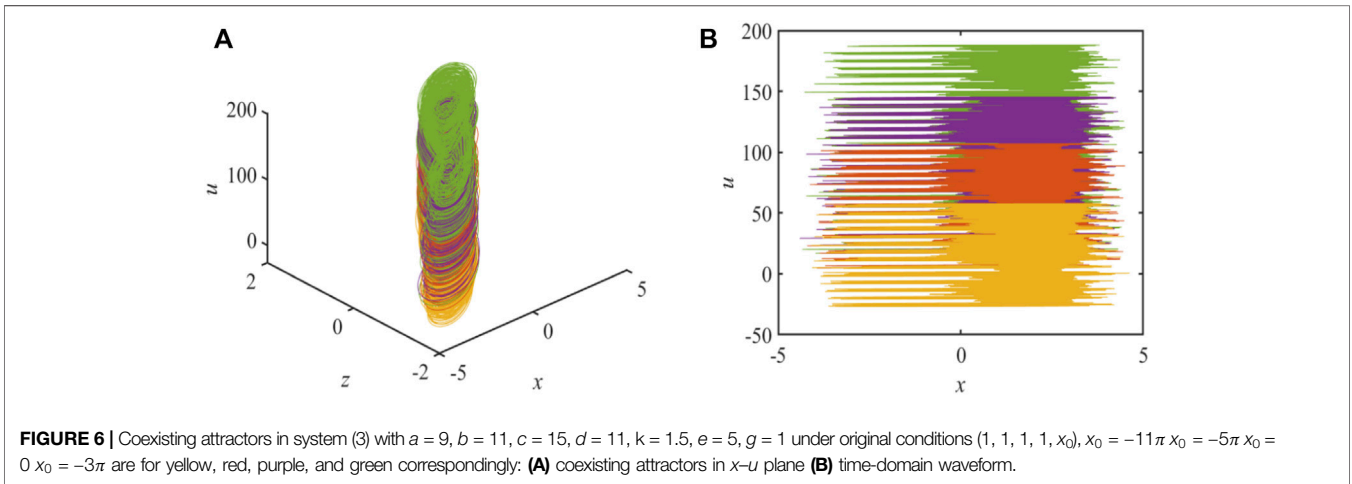
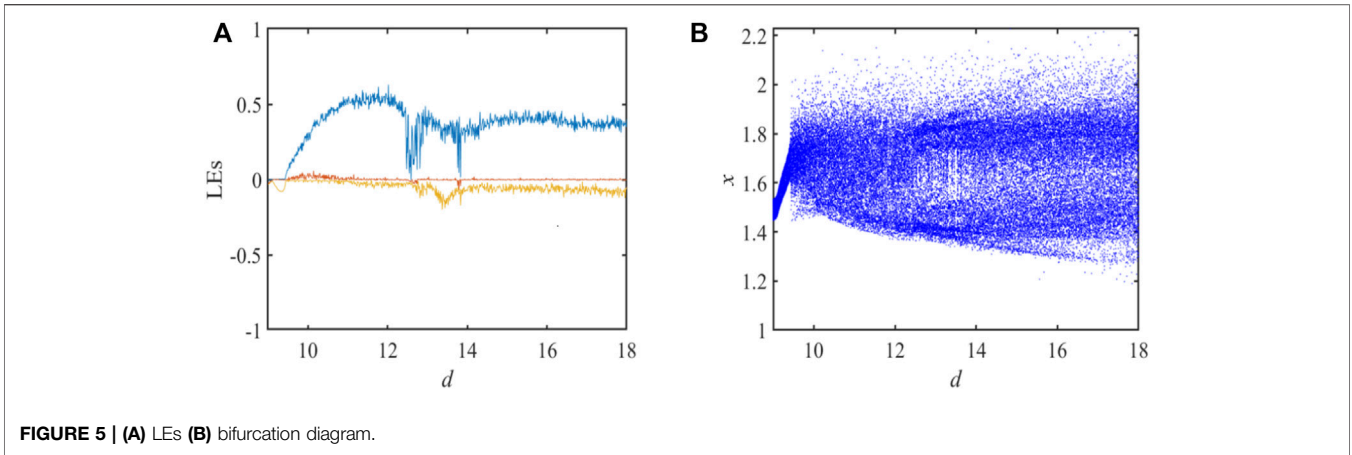
Range	LEs	State	Range	LEs	State
—	0 - - - -	Divergence	9.13–20.00	+0 - - -	Chaos
9.00–9.12	0 - - - -	Period	—	—	—

TABLE 3 | System state corresponding to different parameter d when the original conditions is (1, 1, 1, 1, 1).

Range	LEs	State	Range	LEs	State
—	0 - - - -	Divergence	12.75–12.80	+0 - - -	Chaos
9.00–9.43	0 - - - -	Period	12.81	+0 - - -	Weak chaos
9.43–12.47	+0 - - -	Chaos	12.82–13.75	+0 - - -	Chaos
12.48	+0 - - -	Weak chaos	13.76	+0 - - -	Weak chaos
12.49–12.54	+0 - - -	Chaos	13.77–13.78	+0 - - -	Chaos
12.55	+0 - - -	Weak chaos	13.79	+0 - - -	Weak chaos
12.56–12.58	+0 - - -	Chaos	13.80–13.82	+0 - - -	Chaos
12.59	0 - - - -	Period	13.83	+0 - - -	Weak chaos
12.60–12.73	+0 - - -	Chaos	13.84–20	+0 - - -	Chaos
12.74	+0 - - -	Weak chaos	—	—	—

2.2 Equilibrium Points Set and Stability

By combining the memristor with the 4-D chaotic system, the specific expression of the chaotic system is shown below:



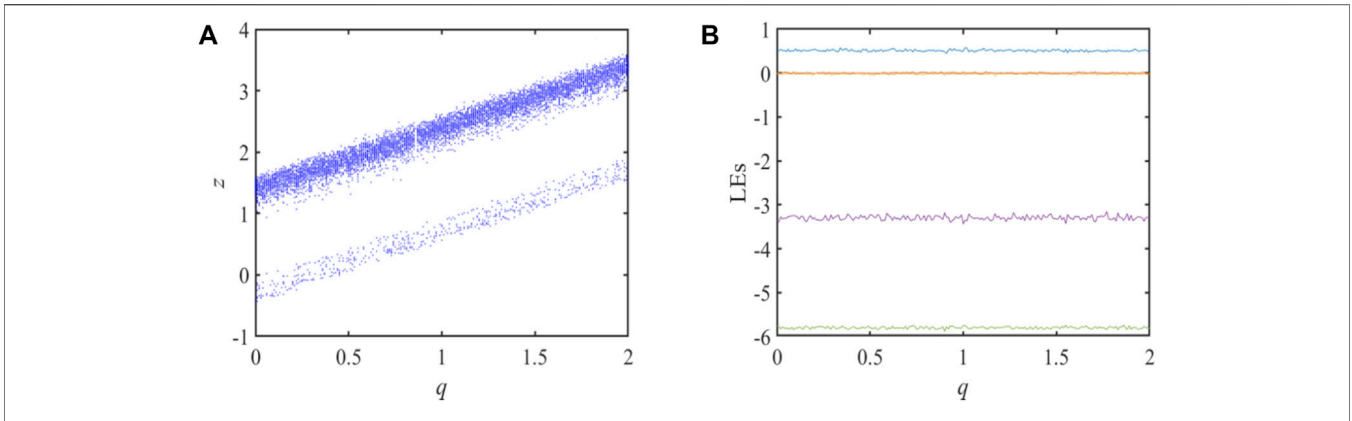


FIGURE 8 | (A) bifurcation diagram **(B)** Lyapunov exponent spectrum.

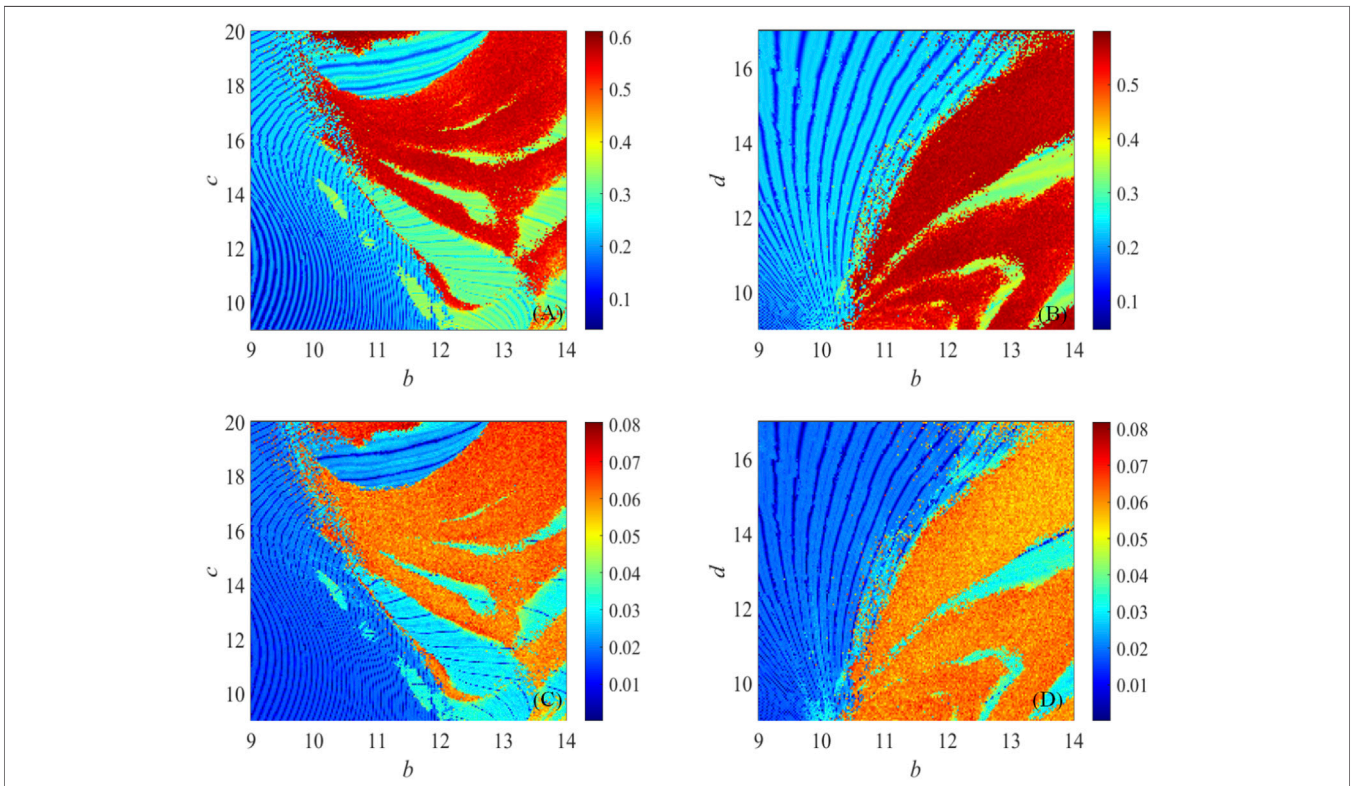


FIGURE 9 | Complexity diagram with different parameter *a* and *e*. **(A)** Parameters *a* *d* in SE complexity; **(B)** parameters *a* *e* in SE complexity; **(C)** parameters *a* *d* in C_0 complexity; **(D)** parameters *a* *e* in C_0 complexity.

$$\begin{cases} \dot{x} = ay - bz + \cos u \\ \dot{y} = cy \cos w - x^2z + k \\ \dot{z} = -dz + ex \\ \dot{u} = gy \\ \dot{w} = y^2 - w \end{cases} \quad (3)$$

The divergence of the system is shown in Eq. 4,

$$\nabla V = \frac{\partial \dot{x}}{\partial x} + \frac{\partial \dot{y}}{\partial y} + \frac{\partial \dot{z}}{\partial z} + \frac{\partial \dot{u}}{\partial u} + \frac{\partial \dot{w}}{\partial w}, \quad (4)$$

when $a = 9, b = 11, c = 15, d = 11, k = 1.5, e = 5,$ and $g = 1,$ and this original conditions $(1, 1, 1, 1, 1); \nabla V$ not greater than zero; it is proved that the system is dissipative, there could be chaotic attractors in this system.

Setting $\dot{x} = \dot{y} = \dot{z} = \dot{u} = \dot{w} = 0,$ then lead to

$$\begin{cases} ay - bz + \cos u = 0 \\ cy \cos w - x^2z + k = 0 \\ -dz + ex = 0 \\ gy = 0 \\ y^2 - w = 0 \end{cases}, \quad (5)$$

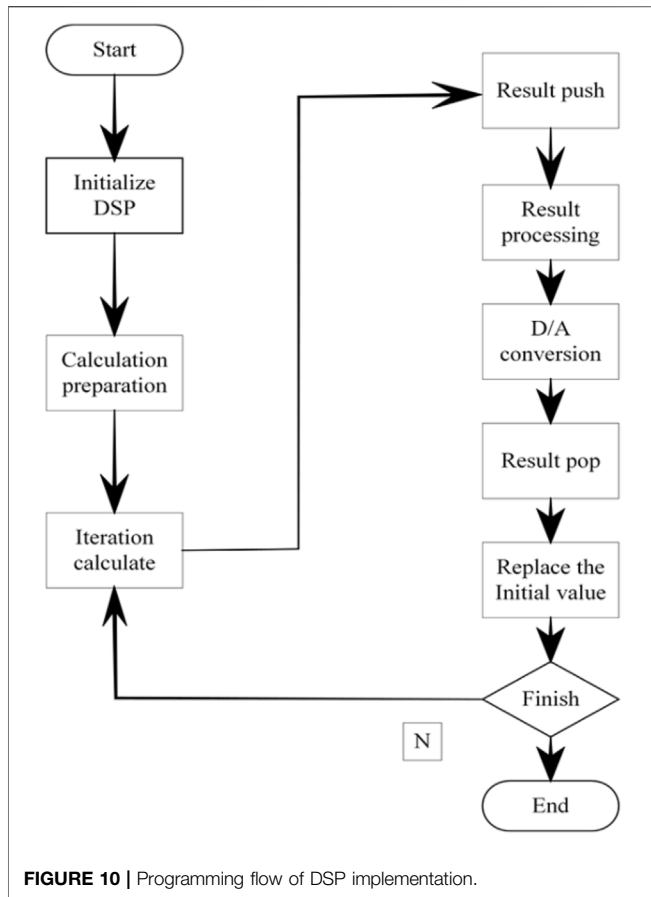


FIGURE 10 | Programming flow of DSP implementation.

from this the following equation can be obtained

$$E^* = \begin{cases} \text{None}, & k \neq 0 \\ \left(0, 0, 0, 0, n\pi + \frac{\pi}{2}\right), & k = 0 \end{cases} \quad (6)$$

If $k = 0$, the equilibrium set is set to O , Jacobi matrix J_E at the equilibrium points O

$$J_E = \begin{bmatrix} -a & z & y & b & 1 \\ -z & c & -x & 0 & 0 \\ y & x & -d & 0 & 0 \\ z & 0 & x & -e & 0 \\ 0 & g & 0 & 0 & 0 \end{bmatrix}, \quad (7)$$

then, the secular equation set

$$\lambda^5 + a_1\lambda^4 + a_2\lambda^3 + a_3\lambda^2 + a_4\lambda + a_5 = 0, \quad (8)$$

where $a_1 = 13.9$, $a_2 = -59.1$, $a_3 = 408.5$, $a_4 = 1,324.4$, and $a_5 = -39.5$.

As shown in Eq. 6, the system is unstable according to the Rous criterion. Hence, when n be an any constant, and any equilibrium points within $O(0, 0, 0, 0, n\pi + \frac{\pi}{2})$, other parameters are $n = 0$, $a = 9$, $b = 11$, $c = 15$, $d = 11$, $k = 1.5$, $e = 5$, and $g = 1$, it results $a_2 = -59.1 < 0$, $a_5 = -39.5 < 0$, and $\lambda_1 = -18.1636$, $\lambda_2 = 3.2391 + 4.7287i$, $\lambda_3 = 3.2391 - 4.7287i$, $\lambda_4 = -2.2397$, $\lambda_5 = 0.0296$. At this point, the system is chaos.

3 NUMERICAL DIAGRAMS OF THE DYNAMICAL BEHAVIOR

3.1 Chaotic Attractor

Let $a = 9$, $b = 11$, $c = 15$, $d = 11$, $k = 3$, $e = 5$, $g = 1$, and original conditions $(1, 1, 1, 1, 1)$. Through simulation of the system, the phase diagrams of different chaotic attractors are shown in Figures 2, 3.

3.2 LEs and Bifurcation Diagram of the System

System (3) sets three parameters, such as a , c , and d . The specific parameter range is as follows $a \in [10, 15]$, $c \in [4, 6]$, and $d \in [3, 12]$.

After fixed the remaining parameters, set a , c , and d as variables, respectively, to determine their initial values. When analyzing the state of the system when the variable changes, the bifurcation diagram is usually combined with Lyapunov

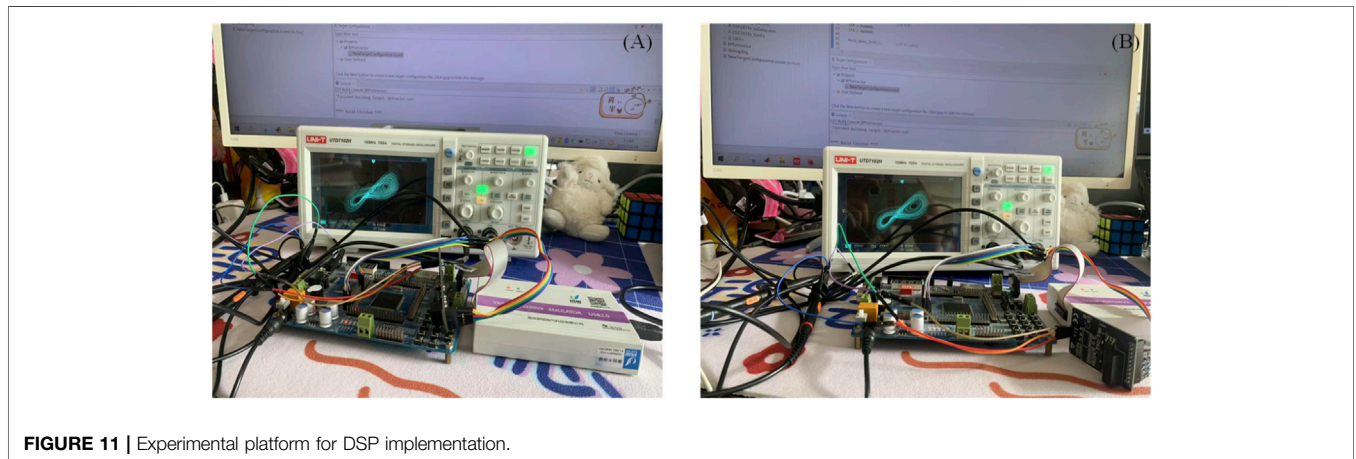


FIGURE 11 | Experimental platform for DSP implementation.

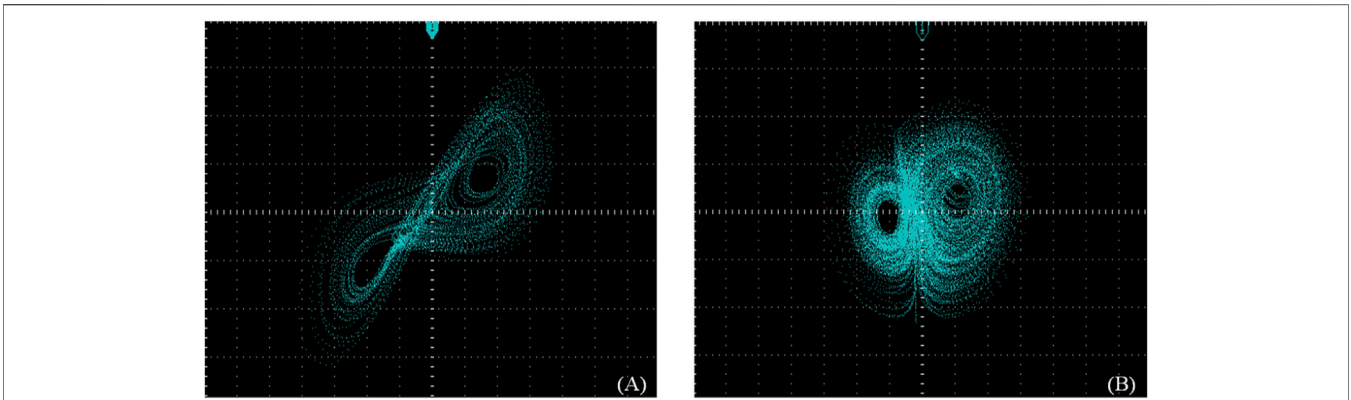


FIGURE 12 | Chaotic attractor on the DSP platform: **(A)** X-Y plane, **(B)** X-W plane.

exponential spectrum, so that the changing state of the system can be analyzed more specifically.

Set $a \in [9, 15]$, $b = 11$, $c = 15$, $d = 11$, $k = 1.5$, $e = 5$, $g = 1$. The LEs and bifurcation diagram as shown in **Figure 2** are obtained through software simulation. Meanwhile, smaller LEs are omitted below for easy observation. In the process of adjusting parameter a , it is found that the dynamical behavior of the system has complex dynamical characteristics, including a variety of states, such as period and chaos, and their mutual transformation. It can be observed from the figure that when parameter $a \in [9, 11.2]$, and in the following range, the system slowly transforms from chaos to period and remains stable. In addition, parameter a ranges the corresponding system status can be clearly understood through the **Table 1**.

Set the parameter $c \in [10, 30]$ and keep the other unchanged. When $c \in [9, 9.12]$, the system is in a period, as shown in **Figure 4**. At the same time, the system changes rapidly from period to chaos in the following range. Specific numerical changes are shown in **Table 2** and **Figure 5**.

Set $d \in [9, 18]$, and keep the other unchanged. In **Figure 6**, when $d \in [9, 9.43]$, at this point, the system LEs is 0, indicating that the system is in a periodic state, then the system changes to a chaos. When parameter $d = 12.74$ the system changes briefly to a periodic state. **Table 3** clearly shows the process of system change within the range of parameter d .

3.3 Infinite Coexisting Attractors

Coexistence attractor is a kind of phenomenon which mainly occurs in special nonlinear systems and has become a research hotspot. Because of the existence of trigonometric function in an equation, there exists the infinite coexisting attractor. The main phenomenon is that when parameters remain fixed and the initial values change, the trajectories could gradually tend to different states of motion. The phenomenon of the infinite coexisting attractor could be observed in **Figure 6**.

3.4 Offset Boosting Scheme

A new feedback state is introduced in the system to control the system flexibly so that the attractor and its attractor pool can move arbitrarily. This method is called offset boosting. By introducing a parameter q to boost the variable z , the system's

attractor and its attractor pool are controlled. The improved offset-boosted system is shown in **Eq. 9**:

$$\begin{cases} x = ay - b(z - q) + \cos u \\ y = cy \cos w - x^2(z - q) + k \\ z = -d(z - q) + ex \\ u = gy \\ w = y^2 - w \end{cases} \quad (9)$$

Let $a = 9$, $b = 11$, $c = 15$, $d = 11$, $k = 1.5$, $e = 5$, $g = 1$, when q is a constant. **Figure 7A** shows the 3D projection of the attractor with different offsets q , and with the increase of parameter q , the position of the attractor also shows a regular upward trend. **Figure 7B** shows the corresponding bifurcation diagram.

In **Figures 7A,B**, the local bifurcation diagram and the corresponding Lyapunov exponent spectrums are shown respectively. When parameter $q \in [0, 2]$, the state variable z of the system augment with the augment of offset variable q , while the attractor LEs of the system does not change. In **Figure 8B**, the attractor is offset accordingly by introducing new variables, which have considerable practical application value in engineering.

3.5 Complexity Analysis

The complexity analysis is one of the important methods to study chaotic system. In practical studies, complexity algorithms are generally introduced to measure chaotic sequences. In order to show the complexity of the chaotic system more clearly, as shown in **Figure 9B**, parameters b, c, b , and d are set respectively. When $b \in [10.5, 14]$ and $d \in [9, 18]$, it could be clearly observed that the region is dark, indicating that the system is in a chaotic state at this time. In this diagram, as the color gets darker, it also means that the system gets more complex. In addition, this multi-dimensional complexity analysis method also offers a deterministic foundation for the parameter pick of the systems.

4 DIGITAL CIRCUIT PLATFORM IMPLEMENTATION

Fast speed, high accuracy, and low environmental impact are the characteristics of the DSP chip F28335. The chaos is confirmed on

the experimental platform. The D/A converter needs to simulate and convert the sequence code generated by the DSP to snatch the output sequence displayed on the corresponding range (UTD7102H). By discretization of continuous chaos, correlation can be processed on DSP platform. Then, the fourth-order Runge-Kutta method is used to transform them into discrete chaotic sequences. Finally, the iterative relation is imported into DSP by C language and a corresponding simulation is carried out. The actual operation and experiment are shown in **Figure 10**.

The parameters were set as $a = 9$, $b = 11$, $c = 15$, $d = 11$, $k = 1.5$, $e = 5$, $g = 1$. The concrete experimental objects are shown in **Figure 11**, and the attractor shown in the figure corresponds to **Figures 2A,B** one by one.

The phase diagram in **Figure 12** shown on an oscilloscope is identical to the phase diagram simulated by the computer, The experiment proves that the digital simulation circuit is correct.

5 CONCLUSION

In this article, a memristor chaotic system is proposed and its dynamic behavior is analyzed. At the same time, it is proved that the appearance of infinite coexisting attractor caused by the periodicity of a trigonometric function is correct. By introducing the variable q into the variable z , the variable z changes parallel with the control variable q , while the value of LEs remains stable. It is proved that the method of flexibly changing the original sequence by introducing a new control variable is effective, and this method is called offset boosting. Finally, the corresponding digital circuit experiment is

carried out, and the correctness of the experimental results is also verified. The influence of trigonometric function on memristor chaotic system is studied in this article, which provides a new reference and idea for the future study of the interaction between the chaotic system and periodic function.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/Supplementary Materials, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

JW provided the idea of algorithm, carried out the simulations, arranged the architecture and drafted the manuscript. JW supervised the work and revised the manuscript. Both authors read and approved the final manuscript.

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