



High Spatio-Temporal Resolution Condenser-Free Quantitative Phase Contrast Microscopy

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Specialty section:

This article was submitted to
Optics and Photonics,
a section of the journal
Frontiers in Physics

Received: 09 March 2022

Accepted: 23 March 2022

Published: 09 May 2022

Citation:

Ma Y, Ma L, Zheng J, Liu M, Zalevsky Z
and Gao P (2022) High Spatio-
Temporal Resolution Condenser-Free
Quantitative Phase
Contrast Microscopy.
Front. Phys. 10:892529.
doi: 10.3389/fphy.2022.892529

Quantitative phase-contrast microscopy (QPCM) provides an effective approach for label-free detection of transparent samples. In this study, we propose a condenser-free quantitative phase-contrast microscopy (CF-QPCM), in which several light-emitting diodes (LEDs) distributed on a ring are used for direct ultra-oblique illumination. Such condenser-free design greatly simplifies the system's structure and releases the space for installing samples. Quantitative phase maps are reconstructed by retarding the unscattered components of the object waves for a series of phases 0 , $\pi/2$, π , and $3\pi/2$ through a high-speed spatial light modulator (SLM). With this system, quantitative phase imaging of live cells has been achieved at a spatial resolution of 231 nm (lateral) and a frame rate of 250 Hz. We believe that the proposed CF-QPCM can contribute to biomedical, industrial, chemistry fields, etc.

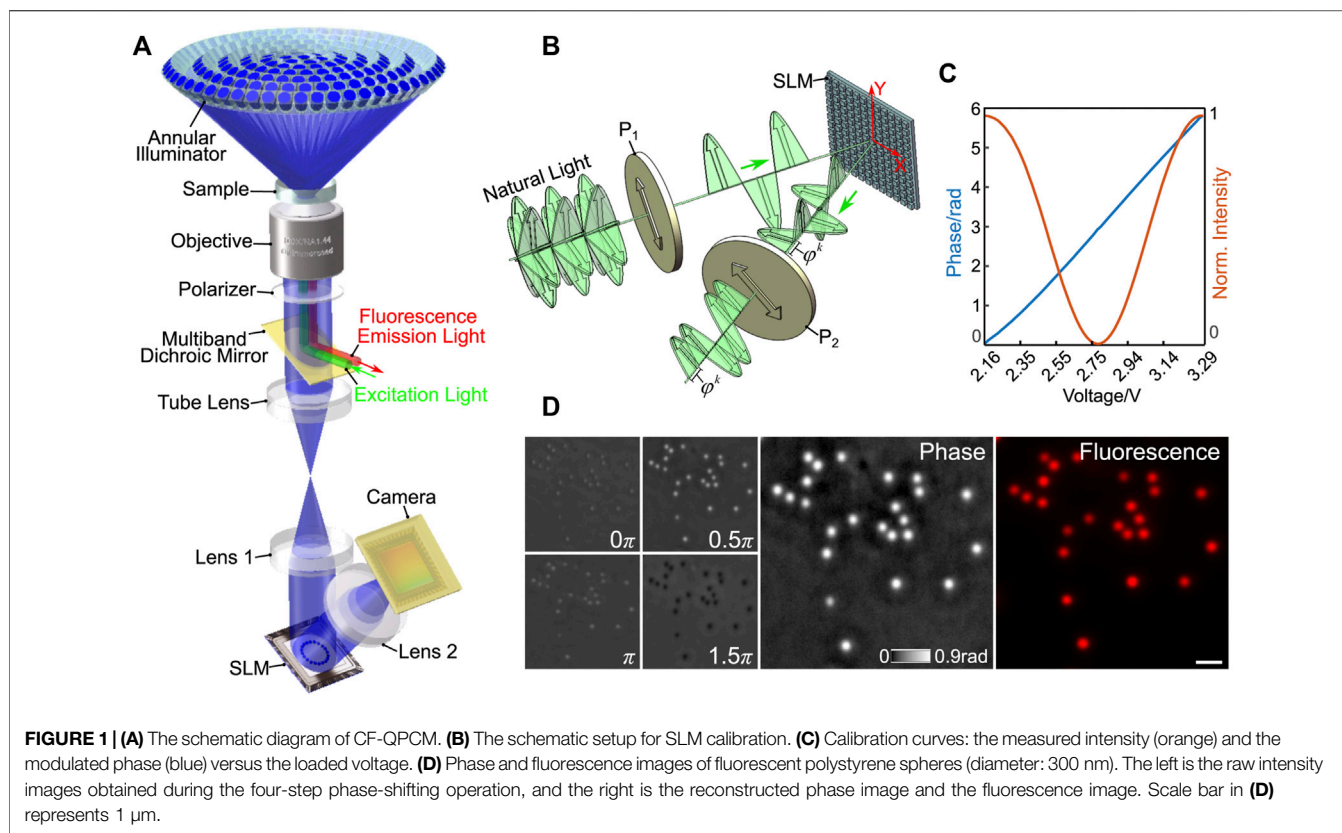
Keywords: quantitative phase imaging, phase contrast, ultra-oblique illumination, label-free, non-invasive

INTRODUCTION

The invention of microscope triggered a revolution in science and cognition, and it also greatly promoted the development of life sciences. Electron microscopes have sub-nanometer spatial resolution and play an incomparable value in structural biochemistry research [1, 2], however, the complex sample preparation hinders the application of electron microscopy in live cells. On the contrary, fluorescence microscopy has been acting as one of the mainstream techniques in life science due to the feature of being non-invasive and the capability to selectively visualize the bio-structures of interest by fluorescence labeling [3–6]. Nevertheless, the application of fluorescence microscopy is hindered by the phototoxicity and photobleaching of fluorophores used.

Being a label-free and non-invasive imaging technique, quantitative phase microscopy (QPM) can quantitatively access the thickness or refraction index (RI) distributions of samples [7, 8]. Digital holography microscopy (DHM) can perform high-accuracy phase measurements of transparent samples by utilizing the interference between an object wave and an off-axis reference wave [9–11]. Common-path configurations were used in DHM to enhance the immunity of the system to environmental disturbances [12]. Researchers further improved the spatial resolution and realized 3D tomographic imaging by mechanically scanning the illumination [13, 14]. Nevertheless, 240 angles scanning within an annular quasi- 2π range are required and limit the imaging speed below 0.8 frames per second (FPS).

Single-beam QPM approaches, in terms of ptychographic iterative engine (PIE) [15] and Fourier ptychographic microscope (FPM) [16–18] have been developed rapidly in recent years, providing the



complex amplitude information of samples by recording a series of diffraction patterns under different illumination positions/directions, or other physical constraints. It is worth noting that the reconstruction of a high-resolution image with these methods requires recording hundreds of low-resolution diffraction patterns, and hence their imaging speed is greatly limited. QPM based on asymmetrical illumination [19] and transport of intensity equation (TIE) [20] can also retrieve the phase mapping of a sample using a relatively-complicated recovery process. Popescu and his co-workers proposed gradient light interference microscopy (GLIM), realizing 3D imaging of optically thick embryos, however it is not suitable for discontinuous samples due to its differential interference contrast essence [21]. Alternatively, phase-contrast imaging, initially proposed by Zernike in the 1940s, can convert the phase distribution of a transparent object into intensity modulation and, therefore, has also been playing an important role in life science [22]. Further, condenser-free Zernike phase-contrast microscopy using annularly-arranged or radially-scanned LEDs was proposed, simplifying the configuration and easing sample installation [23]. As an evolution from qualitative phase imaging to quantitative phase imaging, spatial light interference microscopy (SLIM) uses a spatial light modulator (SLM) instead of the conventional phase plate to perform quantitative phase imaging at temporal and spatial resolutions of 0.6 s and 350 nm, respectively, [24, 25]. Recently, the oblique angle of the annular illumination in SLIM was enlarged, forming the ultra-oblique SLIM, which enhanced the

lateral resolution to 270 nm at an imaging speed of 250 Hz [26], and which identified mitochondria inside live cells by combining a neural network with SLIM [27]. It should be noted that the annular illumination module in SLIM system is either complicated or costly since they require a water-immersed objective lens with a high numerical aperture (NA) to provide ultra-oblique illumination ($\sim 70^\circ$). Moreover, the condenser in the illumination module narrows the space where samples can be installed and hence induce the inconvenience of loading samples.

In this study, we propose condenser-free quantitative phase microscopy, termed CF-QPCM. High spatio-temporal resolution QPM imaging is performed by combining the oblique illumination from annularly-distributed LEDs and the concept of the Zernike phase contrast. The proposed system greatly simplifies the architecture, releases the complexity of sample installation, and avoids the cell contamination caused by a water-immersed condenser objective lens. CF-QPCM can be combined with various fluorescence microscopies to perform multi-modality imaging, providing complementary information for the same sample.

METHODS

Principle of CF-QPCM

CF-QPCM is expediently constructed in a wide-field fluorescence microscope frame, where a dichroic mirror wheel is equipped for integrating fluorescence imaging modality, as shown in

Figure 1A. An illuminator is comprised of dozens of LEDs, which are distributed in several concentric rings. The LEDs on different rings can be chosen independently, so that they can illuminate a sample at different angles. All LEDs used are identical and have a narrow spectrum band (470 ± 10 nm, central wavelength \pm half-width). For CF-QPCM imaging, only the LEDs on one ring are lit up simultaneously to achieve off-axis illuminations. The emitted light from the LEDs directly illuminates the sample, which is placed at the front focal plane of the objective lens (Leica, $\times 100/1.44$ Oil Immersion, WD-0.17, Germany). After passing through the sample, an object wave is formed, which carries the structural and compositional information of the sample. The object wave can be divided into scattered and unscattered waves. The scattered and unscattered waves are Fourier transformed by the objective lens and further imaged to a high-speed SLM (Meadowlark Optics, MSP1920-400-800-HSP8, United States) by a telescope system (shown in **Figure 1A**). A polarizer is placed before the SLM to maximize the phase modulation efficiency of the SLM. To perform quantitative phase imaging of CF-QPCM, the spectra of unscattered waves, which are apparently the images of all the LEDs, are retarded in phase for $0, 0.5\pi, \pi,$ and 1.5π by the SLM, as seen in **Figure 1A**. For simplicity, the sample to be studied is described with the transmittance function $t(x, y) = A(x, y) \cdot \exp(j\phi(x, y))$, where (x, y) are the coordinates on the sample plane, $A(x, y)$ and $\phi(x, y)$ are the transmittance and phase modulation of the sample on the illumination light. Then, a series of phase-contrast images that are essentially the interference between the scattered and unscattered waves are successively recorded by the camera (Andor, Zyla 4.2, United Kingdom):

$$I(x, y, \delta\varphi) = I_0(x, y) \cdot (1 + \beta^2(x, y) + 2\beta(x, y) \cdot \cos[\Delta\varphi(x, y) + \delta\varphi]) \quad (1)$$

Here, $I_0(x, y)$ is the intensity distribution of the unscattered waves, $\Delta\varphi(x, y)$ and $\beta(x, y)$ are the phase difference and amplitude ratio between the scattered and unscattered waves, respectively. And $\delta\varphi = m \cdot \pi/2$ ($m = 0, 1, 2, 3$) is the phase retarding value to the unscattered ones. With $\delta\varphi$ being switched four times, the phase distribution of the sample can be easily recovered by:

$$\varphi(x, y) = \tan^{-1}\left(\frac{\beta(x, y) \cdot \sin(\Delta\varphi(x, y))}{1 + \beta(x, y) \cdot \cos(\Delta\varphi(x, y))}\right) \quad (2)$$

Here $\tan^{-1}\{\cdot\}$ denotes the arctangent function. Obviously, the temporal resolution of the system is determined by the maximal frame rate of the sCMOS camera and the SLM (the lower one of the two). Considering four frames are needed to reconstruct a quantitative phase image, the temporal resolution of the system is 4 ms. Ulteriorly, the temporal resolution of the system can reach up to 250 FPS by reusing the raw phase-shifting images [28].

SLM Calibration

SLM is the key component in CF-QPCM, whose nonlinear optical response of liquid crystal to loaded voltage needs to be calibrated for a linear response between the applied voltage and the generated phase. We propose a simple method to calibrate the

SLM accurately, which allows for measuring the response property of each pixel and avoiding the wavefront curvature induced by externally induced devices. The schematic diagram of the proposed SLM calibration is shown in **Figure 1B**. By nature, the liquid crystal in SLMs only respond to the light with the polarization orientation along the y direction (Y -axis), while it has no response to the light with the polarization orientation along the x direction (X -axis). Therefore, we set two polarizers, P_1 and P_2 , with their polarization direction 45° with respect to the x -direction. After traversing through the first polarizer P_1 and modulated by the SLM, the resulted light beams can be written as

$$\begin{cases} E_{m,n}^x = A_{m,n}^x \cdot \exp(j \cdot \theta_{m,n}) \\ E_{m,n}^y = A_{m,n}^y \cdot \exp[j \cdot (\theta_{m,n} + \varphi_{m,n}^k)] \end{cases} \quad (3)$$

Here, m and n represent the pixel coordinates on the SLM along the x and y directions. $\varphi_{m,n}^k$ ($k = 1, 2, \dots, 255$) is the generated phase on a certain pixel when the k th voltage is applied. Then, the final light intensity generated at each pixel after passing through the second polarizer P_2 is calculated as:

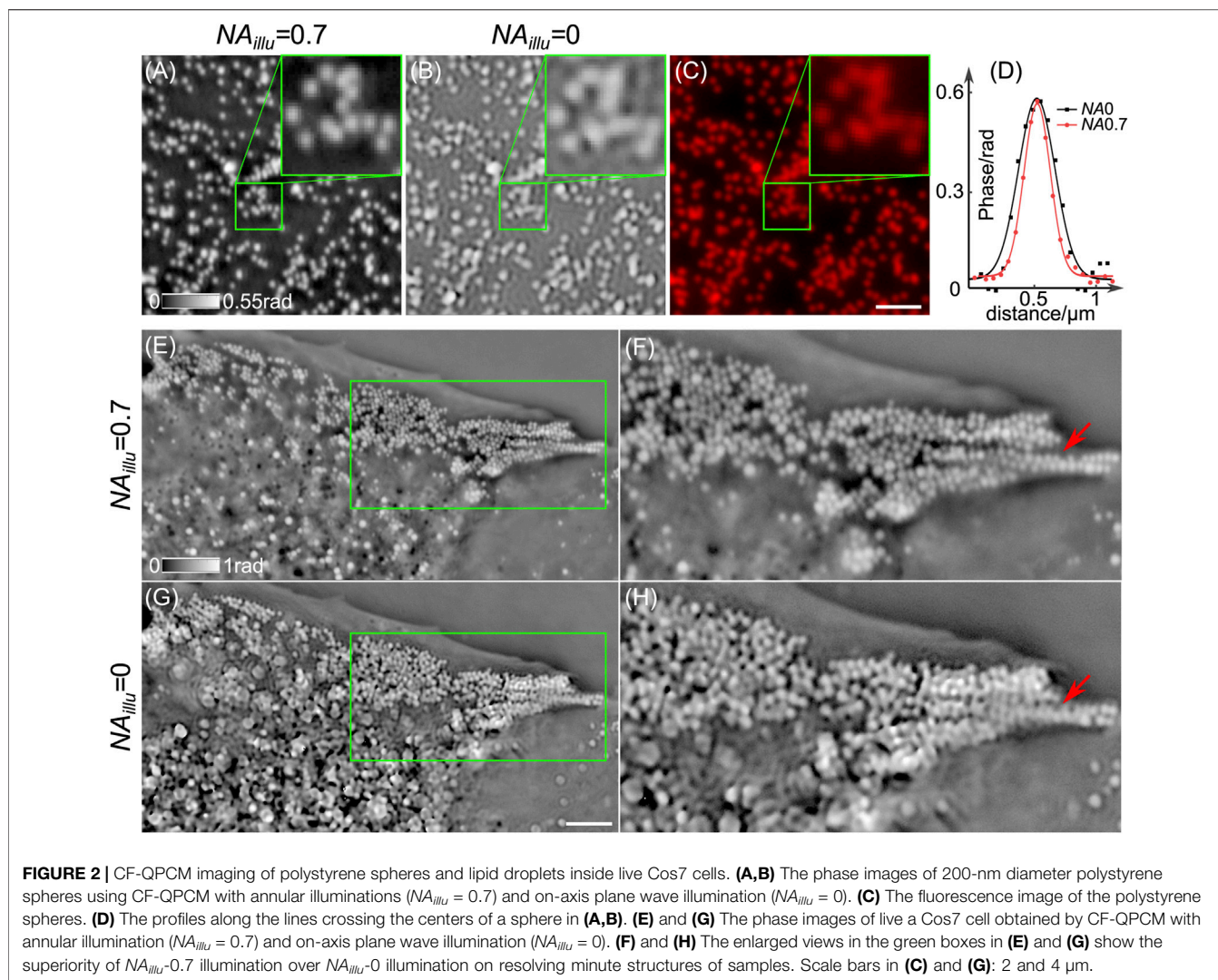
$$I_{m,n}^k = 0.5 \cdot (A_{m,n}^x)^2 + 0.5 \cdot (A_{m,n}^y)^2 + A_{m,n}^x \cdot A_{m,n}^y \cdot \cos(\varphi_{m,n}^k) \quad (4)$$

The dc term $0.5 \cdot (A_{m,n}^x)^2 + 0.5 \cdot (A_{m,n}^y)^2$ and the modulation depth $A_{m,n}^x \cdot A_{m,n}^y$ can be determined when applying the voltage from 0 to the maximal voltage allowed and determining the maximal and minimal intensities for the intensity series obtained. Then the relation between the $\varphi_{m,n}^k$ and the voltage can be obtained with the $\cos^{-1}\{\cdot\}$ function, as is shown with the orange curve in **Figure 1C**. To verify the accuracy of the SLM calibration, a four-step phase measurement by CF-QPCM was performed on 300 nm fluorescent polystyrene spheres (RF300C, Huge Biotechnology), as shown in **Figure 1D**. Notably, the phase and fluorescence imaging were performed simultaneously, and the results are shown in **Figure 1D**-right. Considering the refractive indices of polystyrene (n) and the nutrient solution (n_m) are, respectively, 1.55 and 1.33, the induced maximum phase value by an individual polystyrene sphere can be estimated as $2\pi \cdot (n - n_m) \cdot d/\lambda = 2\pi \cdot (1.55 - 1.33) \cdot 300/470 = 0.88 \text{ rad}$, which is consistent with the measured value of 0.9 rad in the center of a sphere.

EXPERIMENTS AND RESULTS

Spatial Resolution Verification of CF-QPCM

In the CF-QPCM system, we use an annular illuminator circled by 24 LEDs, generating oblique illuminations ($NA_{illu} = 0.7$) to achieve higher spatial resolution. Under such illumination, a sample is imaged by an oil-immersed objective lens ($\times 100/1.44$, Leica). According to the principle of partially coherent imaging, the theoretical lateral resolution of the system is $\delta = \lambda/(NA_{imag} + NA_{illu}) = 220$ nm under the oblique illuminations and 326 nm under on-axis plane wave illumination, respectively. Then, we use 200 nm diameter fluorescent polystyrene spheres (RF200C, Huge Biotechnology) to measure the lateral resolution of the CF-QPCM system. **Figure 2** shows the images of the

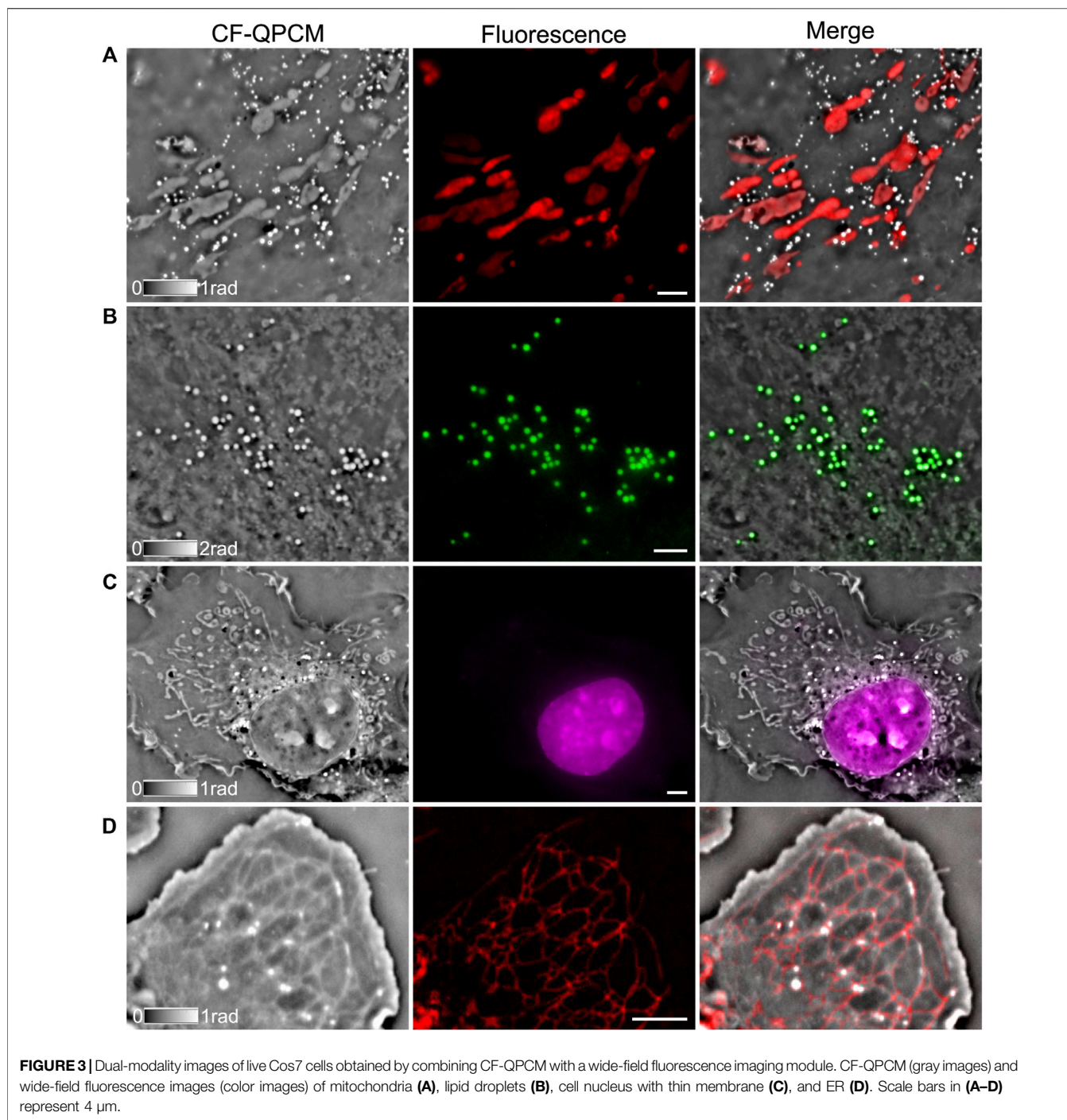


polystyrene spheres obtained with CF-QPCM in comparison with fluorescence imaging (**Figure 2C**). Specifically, **Figures 2A,B** are the phase images of the polystyrene spheres obtained with annular illumination ($NA_{illum} = 0.7$) and on-axis plane wave illumination ($NA_{illum} = 0$). The phase distributions along the line that crosses the center of the same individual sphere in **Figures 2A,B** were extracted and fitted with Gaussian functions in **Figure 2D**. It is apparent that the prior one has a higher spatial resolution than the later one. When dozens of spheres were analyzed like the way shown in **Figure 2D**, the statistics reveal that the measured lateral resolutions (in terms of full width at half maximum, FWHM) are $231 \pm 5 \text{ nm}$ and $330 \pm 10 \text{ nm}$ (mean \pm standard deviation) for NA_{illum} -0.7 and NA_{illum} -0 illumination cases. The enlarged views in **Figures 2A,D** show the superiority of NA_{illum} -0.7 illumination over NA_{illum} -0 illumination on resolving minute structures of samples. Moreover, we also used the decorrelation analysis [29] to quantitatively determine the lateral resolution of **Figures 2A,B**, and the results are $239 \pm 10 \text{ nm}$ and $319 \pm 10 \text{ nm}$ for the two illumination cases. The measured resolution is a little lower than

the theoretical resolution, as the imaging system is unavoidably accompanied by optical aberration. Further, we imaged lipid droplets inside live Cos7 cells with the above two illumination angles, as shown in **Figures 2E–H**. It is clear from **Figures 2F,H** that lipid droplets can be separated under NA_{illum} -0.7 illumination while failed under NA_{illum} -0 illumination. The comparison further confirms that the proposed CF-QPCM under oblique illumination renders the higher resolution images for fine structures inside live cells, compared with the QPCM with plane-wave illumination.

Dual-Modality Imaging of Live Cells

CF-QPCM can visualize fine structures of sub-cellular organelles inside live cells once for all with high contrast and high resolution, but it lacks the specificity of organelles. As a remedy, we integrated both the CF-QPCM module and fluorescence imaging module into a microscope frame to obtain complementary information simultaneously for the same organelles, as shown in **Figure 3**. First, we demonstrated phase and fluorescence dual-modality imaging on mitochondria in live



Cos7 cells stained with fluorescent dye (Thermo Fisher, MitoTracker 7512, United States), and the results are shown in **Figure 3A**. The merged image indicates that CF-QPCM can perform high-contrast imaging of mitochondria together with the surrounding organelles without fluorescent labeling. Second, we performed phase and fluorescence dual-modality imaging on lipid droplets in live Cos7 cells labeled with specific dye (Thermo Fisher, HCS LipidTOX, United States), as shown in **Figure 3B**. The phase images of the lipid droplets have high

contrast as their inner refractive index is quite different from the cytoplasm. Third, we performed phase and fluorescence dual-modality imaging on cell nucleus in live Cos7 labeled with desired dye (Biotium, NucSpot[®] Live 650 Nuclear Stain, United States), as shown in **Figure 3C**. The cell nucleus is often easily recognized since it has a clear membrane. The comparison between CF-QPCM and fluorescence images of the nucleus indicates that CF-QPCM can distinguish the nucleus from other surrounding organelles, such as mitochondria. Lastly, we performed phase

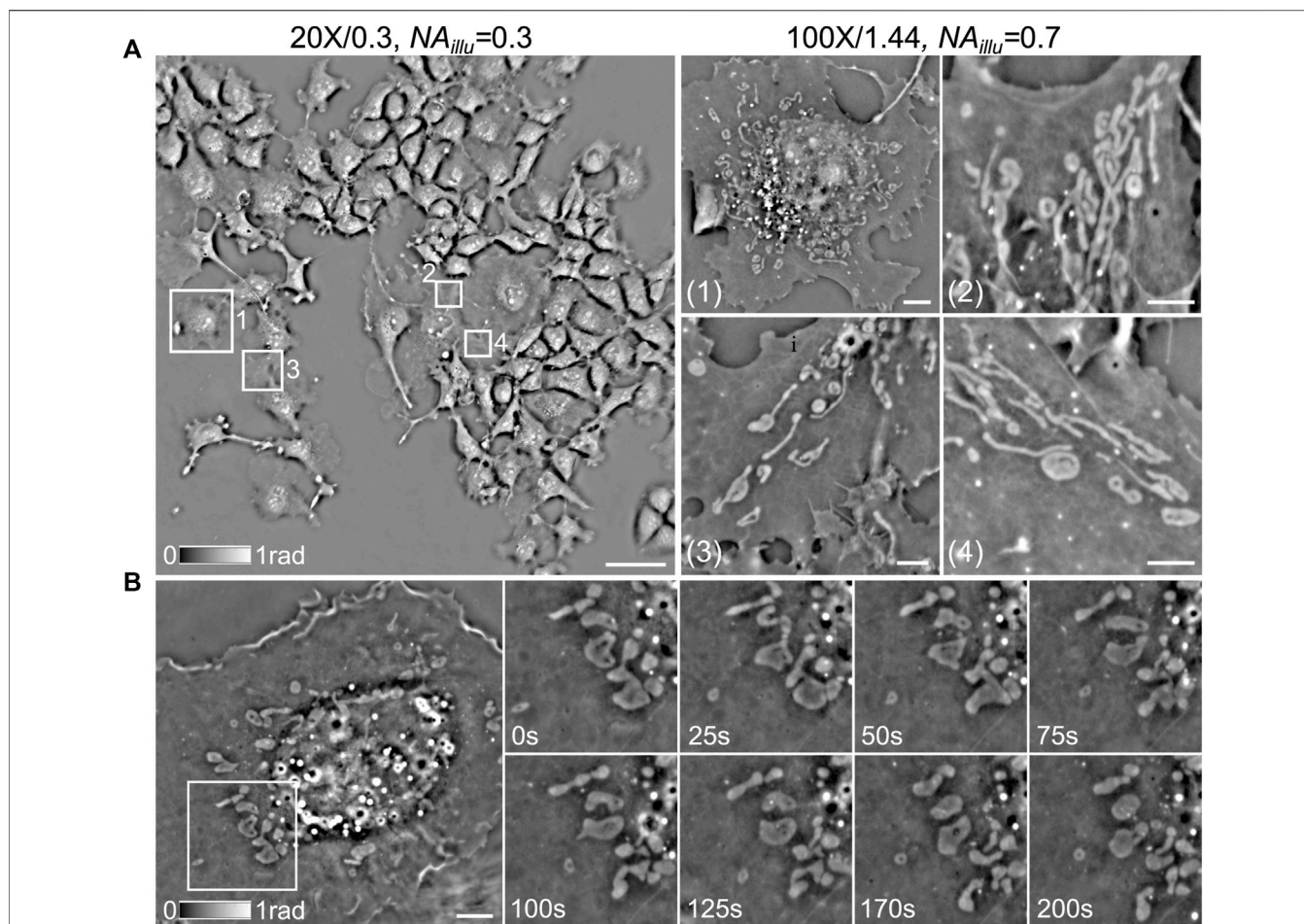


FIGURE 4 | Phase imaging of live Cos7 cells with CF-QPCM under different configurations. **(A)** left: phase image of live Cos7 cells obtained with the large-FOV imaging mode; right: phase images of several sub-regions marked with white boxes on the left, obtained with the small-FOV imaging mode. **(B)** left: phase image of a live Cos7 cell obtained with the small-FOV imaging mode; right: time series of mitochondrial fission in the white box on the left. Scale bars in **(A)**-left, **(A)**-right, and **(B)** are 65 μm , 4 μm , and 6 μm , respectively.

and fluorescence dual-modality imaging on endoplasmic reticulum (ER) in live Cos7 cells labeled with ER Tracker (Thermo Fisher, ER Tracker E34250, United States), as shown in **Figure 3D**. Similar structures of ERs can be visualized in both the phase image and the fluorescence image. It is confirmed from the above experiments that CF-QPCM has the power to image organelles without fluorescent labeling and is compatible with other imaging modalities to provide complementary information.

Biological Dynamics Captured by CF-QPCM

As mentioned before, the LEDs on the rings with different diameters can be chosen to generate different illumination angles, according to different field of view (FOV) of CF-QPCM imaging. For the large-FOV imaging mode, an objective ($\times 20/0.3$) and the LEDs on a ring with a smaller diameter (yielding $NA_{illu} = 0.3$) are selected to match the phase ring on the SLM. For the small-FOV imaging mode, an oil-immersed objective ($\times 100/1.44$) is used, and the LEDs on a ring with a larger diameter (yielding $NA_{illu} = 0.7$) are selected to

match the phase ring on the SLM. The large-FOV imaging mode allows us to study the overall confluency and distribution of cells, as well as the interaction between different cells in a large FOV, while the small-FOV imaging mode enables us to investigate the structure and dynamics of internal organelles with high resolution without fluorescent labeling. Live Cos7 cells were used as the sample to demonstrate the two imaging modes of CF-QPCM in the following.

To demonstrate the first imaging mode of CF-QPCM, **Figure 4A**, (left) shows a phase image of live Cos7 cells in a FOV of $500 \times 500 \mu\text{m}^2$ captured with the objective ($\times 20/0.3$) and $NA_{illu} = 0.3$. Dozens of cells can be observed in the large FOV once for all, which helps us to search for the sub-regions of interest. With the same setting, we captured the flow process of red blood cells (RBCs) as well in the blood vessels of zebrafish (**Supplementary Video S1**) with a temporal resolution of 100 frames per second (FPS). It can be found that RBCs flow in the blood vessels one by one, and the flow speed is different at different locations. To demonstrate the second imaging mode of

CF-QPCM, the sub-regions of interest (about $25 \times 25 \mu\text{m}^2$), e.g., those marked with 1, 2, 3, 4, can be further visualized with high resolution using oil-immersed objective ($\times 100/1.44$) and $NA_{\text{illu}} = 0.7$, as shown in **Figure 4A**, right). Furthermore, CF-QPCM was used to image mitochondria, which play a key role in energy metabolism and are involved in cell aging and death. **Figure 4B** shows the time series of mitochondrial fission captured with CF-QPCM, and the dynamic process can be seen in **Supplementary Video S2**. In addition, the dynamic processes (notably the interactions) of ER, mitochondria, and lipid droplets in a live Cos7 cell can be found in **Supplementary Video S3**. The results indicate that the proposed CF-QPCM can perform high-quality phase imaging for different scales of samples to meet different application requirements.

CONCLUSION AND DISCUSSIONS

In this paper, we have demonstrated a high-resolution, condenser-free quantitative phase-contrast microscopy (CF-QPCM) by combining partially coherent annular illumination with quantitative phase-contrast microscopy (QPCM). In CF-QPCM, the transparent samples are directly illuminated with circularly distributed LEDs, greatly simplifying the system configuration, easing the sample installation, and avoiding the cell contamination caused by the water-immersed condenser objective lens. CF-QPCM has a very high spatial resolution of 231 nm (lateral) and a temporal resolution of 250 Hz for 2D imaging, and it can capture the dynamics of organelles inside live cells under normal survival conditions (37°C and 5% CO_2) without fluorescent labeling. In the future, if higher spatial resolution is required, oblique illumination with an even larger angle could be used, although this may reduce the space between the sample and illuminator. CF-QPCM can be combined with various fluorescence microscopies to form the multi-modality imaging system. Due to its common path interference configuration, CF-QPCM has a strong immunity to external disturbances, which is very suitable for long-term observation of live samples. And in the proposed CF-QPCM, the annular illumination angle of LEDs can be expediently adjusted through electrical control according to different applications, providing complementary information for the same sample. Of note, the application of CF-QPCM to organelle-specific phase-contrast image can be enhanced by using a neural network, trained with phase and fluorescence image pairs of the same organelles [27]. There are also many label-free QPCM method to provide phase-contrast images for visualizing organelles inside live cells, such as optical diffraction tomography (ODT) with annular quasi- 2π illumination [13, 14]. While, CF-QPCM only needs three or four phase-shifted raw images to reconstruct a quantitative phase image and hence has a much faster imaging speed, more suitable for observing the dynamics and interactions of live organelles. Meanwhile, a circularly distributed LEDs illuminator provides a superior spatial resolution to visualize

tiny structures of organelles. The quantitative phase imaging capability of CF-QPCM also opens a gate for inspection of industrial devices, like MEMS or microlens arrays. Compared with the traditional approaches, such as digital holographic microscopy (DHM), CF-QPCM has a superior phase stability and sensibility due to its common-path configuration. CF-QPCM can also be applied to chemistry field, for example, to sense the composition of different solvents or to discriminate different substrate materials since different chemical constituent have different refractive index. We believe such a simple and versatile apparatus will be widely applied for biomedical fields and life science.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/**Supplementary Material**, further inquiries can be directed to the corresponding author.

ETHICS STATEMENT

The animal study was reviewed and approved by Animal Experiment Ethics Committee of Xidian University.

AUTHOR CONTRIBUTIONS

PG conceived and supervised the project. YM and LM Performed experiments and data analysis. JZ, ML, and ZZ contributed to data analysis. YM wrote the draft of the manuscript. All the authors edited the manuscript.

FUNDING

The authors acknowledge the support from the National Natural Science Foundation of China (NSFC 62105251, 62075177, 12104354); Natural Science Foundation of Shaanxi Province (2021JQ-184); Fundamental Research Funds for the Central Universities (XJS210503 and XJS210504); the National Key Research and Development Program of China (2021YFF0700303); Exchange Program Between China-and Poland (2021–2022); Guangdong Basic and Applied Basic Research Foundation (2020A1515110590); Research Fund of State Key Laboratory of Transient Optics and Photonics.

SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fphy.2022.892529/full#supplementary-material>

REFERENCES

- Yuan W., Zhu B., Li X.-Y., Hansen T. W., Ou Y., Fang K, et al. Visualizing H₂O Molecules Reacting at TiO₂ Active Sites with Transmission Electron Microscopy. *Science* (2020) 367:428–30. doi:10.1126/science.aay2474
- Kumar A., Baker J. N., Bowes P. C., Cabral M. J., Zhang S., Dickey E. C, et al. Atomic-resolution Electron Microscopy of Nanoscale Local Structure in lead-based Relaxor Ferroelectrics. *Nat. Mater.* (2021) 20:62–67. doi:10.1038/s41563-020-0794-5
- Shao L., Kner P., Rego E. H., Gustafsson M. G. L. Super-resolution 3D Microscopy of Live Whole Cells Using Structured Illumination. *Nat. Methods* (2011) 8:1044–1046. doi:10.1038/nmeth.1734
- Hell S. W., Wichmann J. Breaking the Diffraction Resolution Limit by Stimulated Emission: Stimulated-Emission-Depletion Fluorescence Microscopy. *Opt. Lett.* (1994) 19:780–782. doi:10.1364/OL.19.000780
- Betzig E., Trautman J. K., Harris T. D., Weiner J. S., Kostelak R. L. Breaking the Diffraction Barrier: Optical Microscopy on a Nanometric Scale. *Science* (1991) 251:1468–1470. doi:10.1126/science.251.5000.1468
- Rust M. J., Bates M., Zhuang X. Sub-diffraction-limit Imaging by Stochastic Optical Reconstruction Microscopy (STORM). *Nat. Methods* (2006) 3: 793–796. doi:10.1038/nmeth929
- Micó V., Zheng J., Garcia J., Zalevsky Z., Gao P. Resolution Enhancement in Quantitative Phase Microscopy. *Adv. Opt. Photon.* (2019) 11:135–214. doi:10.1364/AOP.11.000135
- Gao P., Yuan C. Resolution Enhancement of Digital Holographic Microscopy via Synthetic Aperture: a Review. *gqxzz* (2022) 3:1–16. doi:10.37188/lam.2022.006
- Park Y., Choi W., Yaqoob Z., Dasari R., Badizadegan K., Feld M. S. Speckle-field Digital Holographic Microscopy. *Opt. Express* (2009) 17:12285–92. doi:10.1364/OE.17.012285
- Zheng J., Pedrini G., Gao P., Yao B., Osten W. Autofocusing and Resolution Enhancement in Digital Holographic Microscopy by Using Speckle-Illumination. *J. Opt.* (2015) 17:085301. doi:10.1088/2040-8978/17/8/085301
- Gao P., Pedrini G., Osten W. Structured Illumination for Resolution Enhancement and Autofocusing in Digital Holographic Microscopy. *Opt. Lett.* (2013) 38:1328–30. doi:10.1364/OL.38.001328
- Zhang M., Ma Y., Wang Y., Wen K., Zheng J., Liu L, et al. Polarization Grating Based on Diffraction Phase Microscopy for Quantitative Phase Imaging of Paramecia. *Opt. Express* (2020) 28:29775–87. doi:10.1364/OE.404289
- Dong D., Huang X., Li L., Mao H., Mo Y., Zhang G, et al. Super-resolution Fluorescence-Assisted Diffraction Computational Tomography Reveals the Three-Dimensional Landscape of the Cellular Organelle Interactome. *Light Sci. Appl.* (2020) 9:1–15. doi:10.1038/s41377-020-0249-4
- Cotte Y., Toy F., Jourdain P., Pavillon N., Boss D., Magistretti P, et al. Marker-free Phase Nanoscopy. *Nat. Photon.* (2013) 7:113–117. doi:10.1038/nphoton.2012.329
- Godden T. M., Suman R., Humphry M. J., Rodenburg J. M., Maiden A. M. Ptychographic Microscope for Three-Dimensional Imaging. *Opt. Express* (2014) 22:12513–12523. doi:10.1364/OE.22.012513
- Zhang H., Jiang S., Liao J., Deng J., Liu J., Zhang Y, et al. Near-field Fourier Ptychography: Super-resolution Phase Retrieval via Speckle Illumination. *Opt. Express* (2019) 27:7498–512. doi:10.1364/OE.27.007498
- Alexandrov S. A., Hillman T. R., Gutzler T., Sampson D. D. Synthetic Aperture Fourier Holographic Optical Microscopy. *Phys. Rev. Lett.* (2006) 97:168102. doi:10.1103/physrevlett.97.168102
- Ou X., Horstmeyer R., Zheng G., Yang C. High Numerical Aperture Fourier Ptychography: Principle, Implementation and Characterization. *Opt. Express* (2015) 23:3472–91. doi:10.1364/OE.23.003472
- Ledwig P., Robles F. E. Quantitative 3D Refractive index Tomography of Opaque Samples in Epi-Mode. *Optica* (2021) 8:6–14. doi:10.1364/OPTICA.410135
- Zuo C., Li J., Sun J., Fan Y., Zhang J., Lu L, et al. Transport of Intensity Equation: a Tutorial. *Opt. Lasers Eng.* (2020) 135:106187. doi:10.1016/j.optlaseng.2020.106187
- Nguyen T. H., Kandel M. E., Rubessa M., Wheeler M. B., Popescu G. Gradient Light Interference Microscopy for 3D Imaging of Unlabeled Specimens. *Nat. Commun.* (2017) 8:1–9. doi:10.1038/s41467-017-00190-7
- Zernike F. Phase Contrast, a New Method for the Microscopic Observation of Transparent Objects Part II. *Physica* (1942) 9:974–986. doi:10.1016/S0031-8914(42)80035-X10.1016/s0031-8914(42)80079-8
- Webb K. F. Condenser-free Contrast Methods for Transmitted-light Microscopy. *J. Microsc.* (2015) 257:8–22. doi:10.1111/jmi.12181
- Wang Z., Millet L., Mir M., Ding H., Unarunotai S., Rogers J, et al. Spatial Light Interference Microscopy (SLIM). *Opt. Express* (2011) 19:1016–26. doi:10.1364/OE.19.001016
- Zhang J. K., He Y. R., Sobh N., Popescu G. Label-free Colorectal Cancer Screening Using Deep Learning and Spatial Light Interference Microscopy (SLIM). *APL Photon.* (2020) 5:040805. doi:10.1063/5.0004723
- Ma Y., Guo S., Pan Y., Fan R., Smith Z. J., Lane S, et al. Quantitative Phase Microscopy with Enhanced Contrast and Improved Resolution through Ultra-oblique Illumination (UO-QPM). *J. Biophotonics* (2019) 12:e201900011. doi:10.1002/jbio.201900011
- Guo S., Ma Y., Pan Y., Smith Z. J., Chu K. Organelle-specific Phase Contrast Microscopy Enables Gentle Monitoring and Analysis of Mitochondrial Network Dynamics. *Biomed. Opt. Express* (2021) 12:4363–4379. doi:10.1364/BOE.425848
- Ma Y., Li D., Smith Z. J., Li D., Chu K. Structured Illumination Microscopy with Interleaved Reconstruction (SIMLR). *J. Biophotonics* (2017) 11: e201700090. doi:10.1002/jbio.201700090
- Descloux A., Grusmayer K. S., Radenovic A. Parameter-free Image Resolution Estimation Based on Decorrelation Analysis. *Nat. Methods* (2019) 16:918–924. doi:10.1038/s41592-019-0515-7

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