



Individual Roles of Network and Sheet Solids in Architected Porous Materials

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To understand the different roles played by sheet solids and network solids in complex porous biomaterials/native tissues, we designed a new kind of nature-inspired structure comprising these two solids by using triply periodic minimal surfaces and Voronoi diagrams using the CAD method and compared them with the previously reported Poisson-Voronoi (PV) solids that only comprise network solids. Here, we show that the sheet solids contribute greater stiffness and solid/void interface than the network solids, and our TPMS-Voronoi solids can improve/tune both the elastic moduli and specific surfaces even at fixed solid volume fraction and can be stiffer than the Poisson-Voronoi solids. This can directly guide the porous materials design for use in tissue engineering and aerospace.

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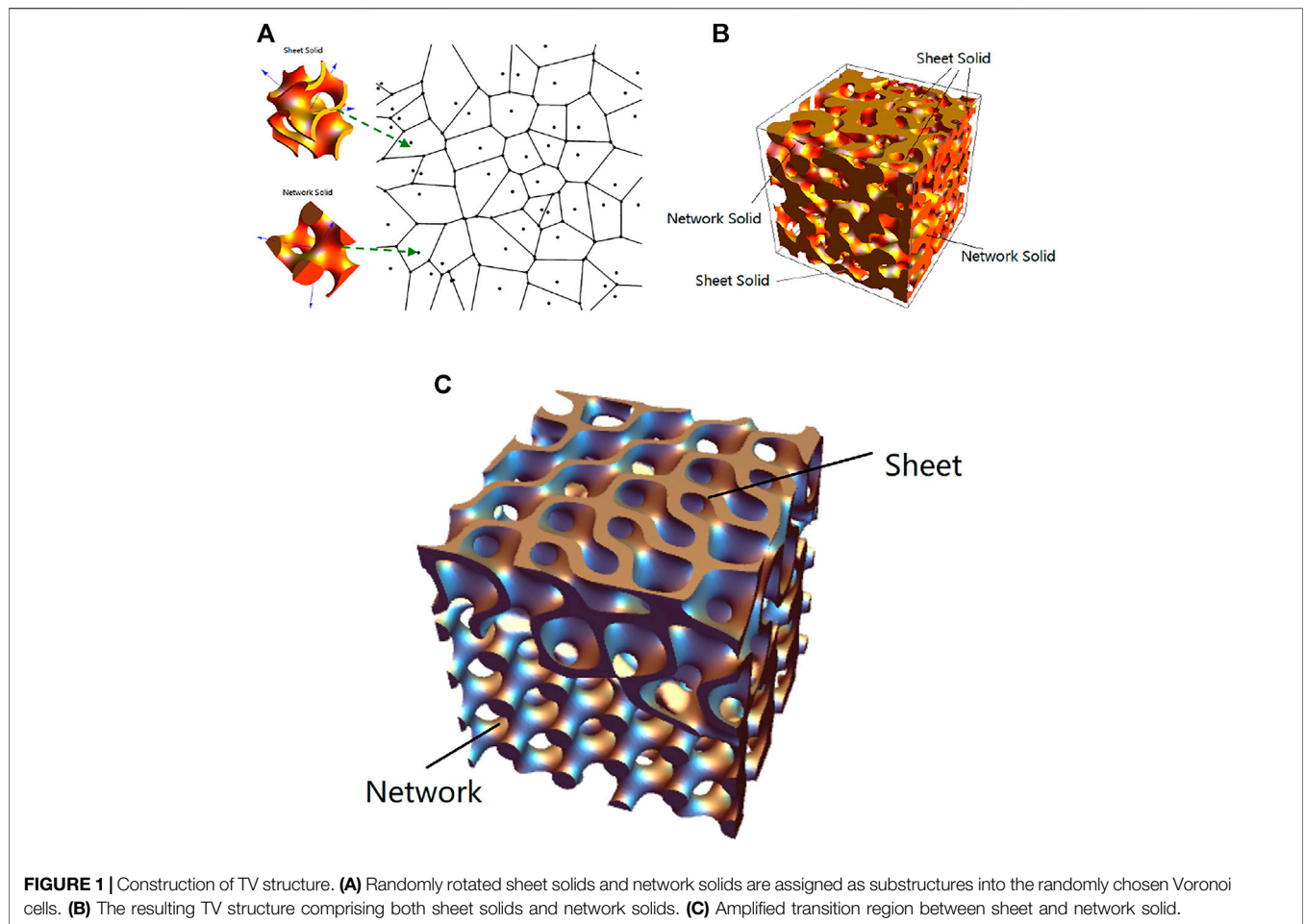
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INTRODUCTION

Porous materials have many applications. In engineering, porous metamaterials can realize sound absorption and insulation [1–3]. In medicine, porous structures can be used as implants and promote the growth of tissues and cells [4–7]. Natural porous structures, such as trabecular bones [4–7], consist of two kinds of solids: the network solid and the sheet solid. What are the different roles of these two solids? What are the benefits with this arrangement? Previous reports showed that at the same solid volume fraction (SVF), pure sheet solids tended to be stiffer than pure network solids [8, 9], but their roles in a hybrid complex structure still remain unclear. With nature materials, we cannot flexibly control the fraction of each solid. Thus, we designed a new kind of nature-inspired structure comprising both sheet and network solids by using triply periodic minimal surfaces (TPMS) and Voronoi diagrams [10–15]. Here we show that our TPMS-Voronoi (TV) solids allow us to flexibly vary the fraction of the sheet/network solids to obtain various elastic moduli and specific surfaces.

TPMS architectures are attractive candidates for designing innovative scaffold structures because they overcome the limitations of existing design and fabrication strategies, which require complex changes between different substructures during design and path planning for fabrication [16]. Even earlier, such triply periodic materials were discussed in soft materials, including biomorphic synthetic materials and biological structures (e.g., butterflies) [17, 18]. TPMS allows for scaffold structures to be simply represented using a single mathematical inequality and to be readily fabricated using additive manufacturing (AM) techniques [19–22]. Yoo proposed an effective method for arranging porosities for a porous scaffold using a thin-plate radial basis function based on distance field (DF)- and TPMS-based functions [23]. Melchels et al. [24] analyzed the porous structures and mechanical properties of conventional gyroid and diamond architectures. They also added a linear term in a conventional TPMS-based function to obtain gradient porosities



[10, 25]. Kapfer et al. [8] proposed TPMS-based sheet solids that could provide a relatively high stiffness with a high porosity. This structure could also combine a high specific surface area with a high flux and the desired pore accessibility [26]. Pasko et al. [27] proposed lattice porous structures based on trigonometric functions and arranged their pore distribution using a DF method. Fryazinov et al. [28] proposed a cellular structure based on periodic functions and obtained metamorphosis inside a cellular structure, which resulted in gradient pores. However, the capabilities of these current strategies are too limited to construct complicated random structures. Here, we broaden the capability for TPMS architectures to flexibly generate complex porous structures and investigate the elastic moduli and specific surfaces of the resulting structures.

MATERIALS AND METHODS

Construction of TV Solids

First, N Voronoi cells were obtained based on N random seed points in a region of interest with dimensions of $w \times w \times w$. N_s sheet solids and N_n network solids (where $N = N_s + N_n$) were randomly rotated, and each solid was then placed in a randomly chosen Voronoi cell as shown in **Figure 1A**. But the substructures

may connect poorly on the Voronoi boundaries. Using the boundary strengthening and substructures' smooth combination techniques for strong connections, all the substructures were integrated to form a hybrid structure as shown in **Figure 1B**. The sheet solids and network solids with SVF (ϕ) were combined to yield a hybrid solid with the same SVF value. The sheet solid fraction (α) is defined by $\alpha = N_s/N$.

We choose Gyroid TPMS for each network solid because this yields the stiffest network solid among the known minimal surfaces. An unrotated network solid (the solid phase) is defined by the 3D region $stru_{network}(x, y, z) \geq 0$ [11], according to the following:

$$stru_{network}(x, y, z) = \cos(cx) \sin(cy) + \cos(cy) \sin(cz) + \cos(cz) \sin(cx) + d \tag{1}$$

Similarly, an unrotated sheet solid is defined by $stru_{sheet}(x, y, z) \geq 0$, according to the following:

$$stru_{sheet}(x, y, z) = d - |\cos(cx) \sin(cy) + \cos(cy) \sin(cz) + \cos(cz) \sin(cx)| \tag{2}$$

where $c = 2p\pi/q$ and implies that p pores are arranged within a distance q in one direction, thus controlling the pore density/size,

TABLE 1 | Relationship between d and SVF for the Gyroid TPMS sheet solid and network solid.

	SVF	11%	21%	30%	50%	61%	72%	83%
D	Sheet solid	0.18	0.33	0.5	0.735	0.876	1.09	1.24
Value	Network solid	-1.15	-0.89	-0.55	0	0.4	0.6	1

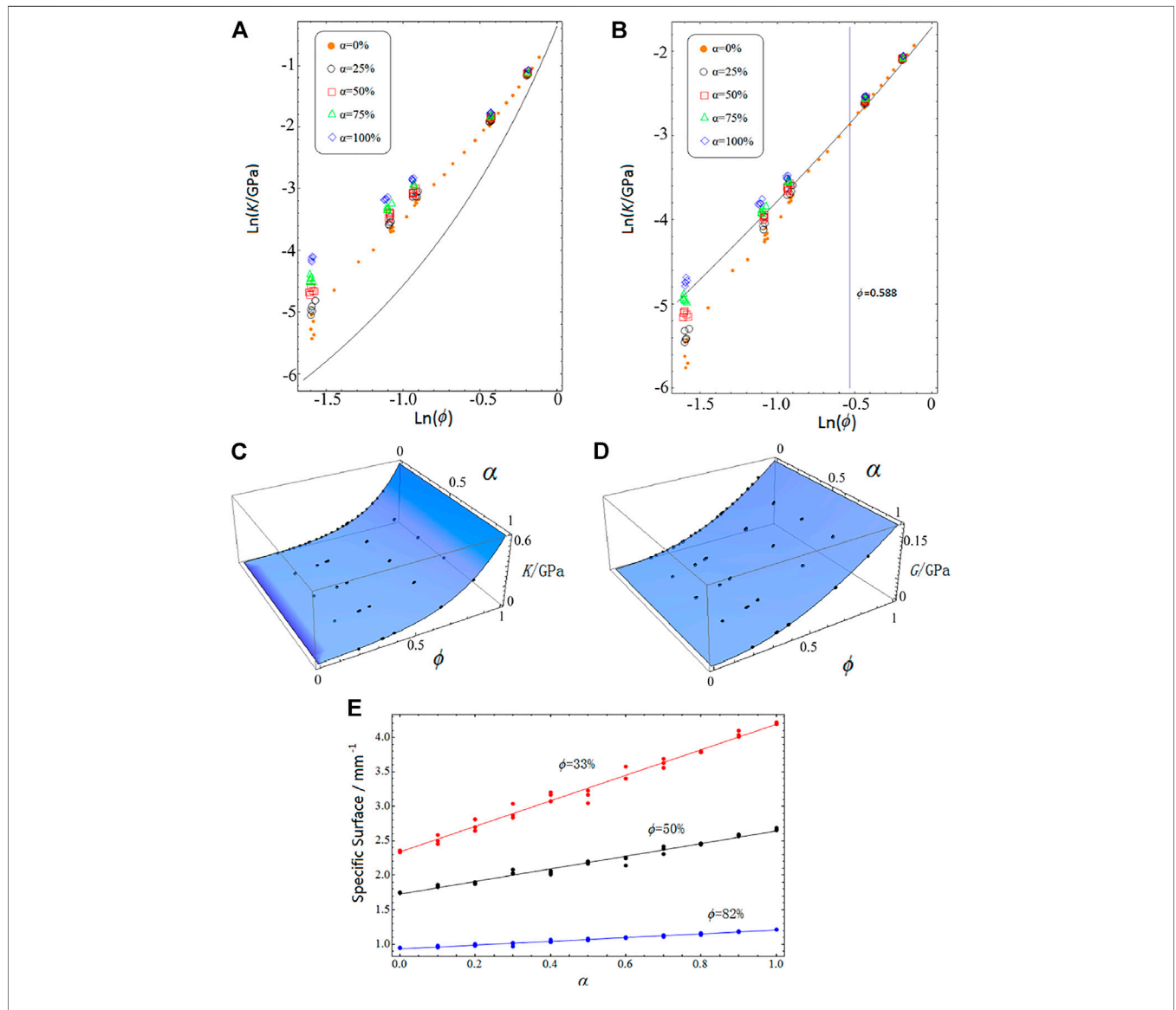


FIGURE 2 | Comparisons of the logarithm of bulk moduli (A) and shear moduli (B) versus the logarithm of SVF (ϕ) for TV solids and PV solids, where the points represent the results for TV solids when $\alpha = 0$ (orange dots), 25% (black circles), 50% (red squares), 75% (green triangles), and 100% (blue rhombi), and the solid lines represent the results for PV solids. The bulk moduli of TV solids are greater than those of the PV solids, but the shear moduli of TV solids without the sheet solids ($\alpha = 0$) are lower than the PV solids for $\phi < 0.588$ but exceed the PV solid values at higher ϕ . Moreover, for low SVF, the moduli of TV solids increase as the sheet solid fraction (α) increases at a given SVF; however, for high SVF, α has only a slight effect on the moduli. The blue surfaces (C) and (D) are the best fits of bulk modulus data and shear modulus data for TV solids (i.e., the data represented by dots, circles, squares, triangles, and rhombi in Panel 2A,B), respectively. Five structures were constructed for each α with each ϕ . (E) The relationship between the α and specific surface of TV solids for different ϕ , that is, $\phi = 33\%$ (red), 50% (black), and 82% (blue). The specific surface increases with increased α for all ϕ , but the slope of each line decreases as ϕ increases. Three structures were built for each α with each ϕ .

d controls the SVF for each solid (Table 1), and x , y , and z are the three-dimensional coordinates. Since the elastic and shear modulus are not changed with the sizes, the structural models

for modulus calculation are dimensionless. See **Supplementary Material** for the transition and connection between the sheet and network solids.

Elastic Moduli

Using the voxel-based finite element method [29], the elastic moduli of our TV solids were calculated as the minimized sum of the elastic energy stored in all the voxels with a given applied strain. The 3D region of interest containing the solid was discretized by 200^3 voxels. The TV solids were constructed as collagen. The bulk modulus and shear modulus of collagen were set to $K_s = 0.69$ GPa and $G_s = 0.18$ Pa [29], where the subscript s denotes the non-porous solid.

Specific Surfaces

The specific surface was defined as the area of the solid/void interface divided by the solid volume in the region of interest, calculated using Rhinoceros software.

RESULTS AND DISCUSSION

TV Solids

An example of TV solid is shown in **Figure 1B**, where the fraction of sheet solids is about 50%. The locations of sheet solids (plate elements) and network solids (rod elements) can be randomly assigned in a Voronoi diagram as shown in **Figure 1A**, and smooth transitions occur between these substructures. The numbers of sheet solids and network solids can be indicated by the designer to further study the properties of the resulting structures. Such a model can be exported as an STL file and then fabricated using the AM method for further geometrical measurement and mechanical test.

Linear-Elastic Moduli

Here, we analyze the effect of α on the elastic moduli of our TV solids and compare their moduli to that of the previously reported PV solid [29].

The elastic moduli of the TV solids are compared with those of the PV solids using the same calculation scheme and the same settings (where the PV solids were constructed based on Voronoi diagram alone, without TPMSs). The data for the moduli of our TV solids are shown by points in **Figures 2A,B** with the curved lines representing the PV solid moduli. **Figure 2A** shows that the TV solid bulk moduli are greater than those of the PV solids for all ϕ and increase with α at a given ϕ . **Figure 2B** shows that when $\alpha = 0$, most of the TV solid shear moduli are lower than those of the PV solids for $\phi < 0.588$, but exceed the PV solid values at the higher ϕ values. However, even for $\phi < 0.588$, the TV solid shear moduli can exceed the PV solid values when sheet solids are fully introduced into the structure ($\alpha = 100\%$). While the bulk and shear moduli can be significantly adjusted at low SVF by changing α , the moduli vary only slightly with α at high ϕ values. Although the sheet solids tend to be stiffer than the network solids at low SVF, their moduli become similar at high SVF. This indicates that even at a specified SVF, the use of TV solids gives us flexibility in the design of biomaterials with different mechanical properties.

Many porous structures exhibit a power-law relationship between the modulus and the SVF [29, 32]. Here, we use the following:

$$K = (u_k \alpha + v_k) \phi^{n_k} \left(\frac{K_s}{u_k \alpha + v_k} \right)^\phi \quad (3)$$

to fit the bulk modulus data with three free parameters (u_k , v_k , and n_k), and **Eq. 3** can ensure that $K \rightarrow 0$ for $\phi \rightarrow 0$, and $K \rightarrow K_s$ for $\phi \rightarrow 1$. An analogous form is used for the shear modulus expression. The fitting results are $n_k \approx 0.5$, $u_k \approx 0.0075$ GPa, $v_k \approx 0.0132$ GPa and $n_g \approx 2.29$, $u_g \approx 0.083$ GPa, $v_g \approx 0.22$ GPa for the bulk and shear moduli, respectively (where the subscripts k and g denote the parameters for bulk and shear modulus, respectively). The best fit surfaces for the bulk and shear modulus data are shown in **Figures 2C,D**. Comparison to the PV solids ($n_k \approx 1.0$ and $n_g \approx 1.69$ [29]) shows that the values of shear and bulk moduli of our TV solids are more and less sensitive to the changes in SVF, respectively.

For the overlapped results, we can arrange the sheet and network solids with the same α and ϕ as shown in **Figure 1A**, but we can change their locations, thus slightly changing their moduli.

Specific Surface

We can tune the specific surface value of the TV solids by changing α at a given SVF. As shown in **Figure 2E**, the specific surface linearly increases with α for all ϕ (33, 50, or 82%); this is because a sheet solid has a bicontinuous surface that has two parallel interfaces with the void phases. Generally, the total area of the two interfaces is larger than that of the single interface in a network solid [8]. However, when ϕ is large, the specific surface can only be tuned slightly as indicated by the slopes of the best fit lines.

DISCUSSION

Once the fabricated material, structural type, and SVF are fixed, it is not easy to tune the elastic moduli and specific surfaces for TPMS lattices. However, such properties can be adjusted by changing the numbers of sheet solids and network solids in a hybrid structure while keeping the fabricated material, structural type, and SVF constant. In other words, the structural property can be decoupled with the SVF, which is useful in the structural design process to independently achieve one property with fixing others.

Also, we can understand some mechanisms in the natural materials. Interestingly, similar geometries of the hybrid structures also occur elsewhere in biochemistry/biology. Such transition structures are commonly observed in the lipid analogues, for example, the coexistence between Diamond and Primitive morphologies in cubosomes [30], with similar dynamic transition structures described for lipid mesophase transitions [31]. But in these systems, there are much more rigorous requirements for how two domains of different types and orientation merge at their common interface, as the solid/void interface is determined by the optimization of surface curvature.

It shows that surface curvature (mean curvature H and Gaussian curvature K) is a key factor for guiding growth of cells and tissues [7]. For an individual cell, on a flat substrate ($H=K=0$), the cell nucleus tends to be flat with a small

compressive force by the perinuclear actin. On a convex spherical surface ($H > 0$ and $K > 0$), the nucleus is flatter as that compressive force becomes larger. On a concave spherical surface ($H > 0$ and $K < 0$), the nucleus is lifted off and tends to be a spherical shape [6]. Compared to the hemispherical cap, a cell on the cylindrical surface is not likely to cause the osteogenic formation, as the fibers align to the longitudinal orientation and the force that compresses the nucleus is relatively small [33]. For collective cells, the tissue growth rate linearly increasing where the negative curvature occurs is used in the tissue deposition model and agrees well with the experimental results [34]. Inspired by this, the proposed method can integrate substructures with different surface curvatures for biomimetic design purpose and further study the curvature-guiding mechanism for cell and tissue growth.

By controlling the fraction of sheet solids based on our model, we can tune the elastic modulus with fixed SVF and understand how lower moduli of trabecular bone in postmenopausal women are associated with a reduced percentage of plate elements over all elements [4, 5].

CONCLUSION

Here, we show that morphology and structure are very important to determine elastic modulus and specific surface. To design porous biomaterials with desired moduli and specific surfaces, we can make them tunable by combining sheet solids with network solids with different fractions even at a specified porosity. On the other hand, we can understand the individual roles played by these two solids in natural porous materials based on the compound porous models. Not limited to the types of sheet and network solids chosen in this study, we can combine other

mathematically defined structures using our method for porous scaffold design.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/**Supplementary Material**, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

NY: conceptualization, methodology, writing—original draft preparation, and supervision. SF, JZ, YZ, and ZQ: simulation.

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SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fphy.2022.892525/full#supplementary-material>

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