



# Background Pressure Induced Structural and Chemical Change in NiV/B<sub>4</sub>C Multilayers Prepared by Magnetron Sputtering

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Wei Z, Zhang Z, Jiang L, Yang Y, Chang C, Feng Y, Qi R, Huang Q, Yan W, Xie C and Wang Z (2022) Background Pressure Induced Structural and Chemical Change in NIV/B<sub>4</sub>C Multilayers Prepared by Magnetron Sputtering. Front. Phys. 10:837819. doi: 10.3389/fphy.2022.837819 NiV/B<sub>4</sub>C multilayers with a small d-spacing are suitable for multilayer monochromator working at a photon energy region from 5 to 8 keV, or photon energy region from 10 to 100 keV. To investigate the influence of background pressure during fabrication, NiV/B<sub>4</sub>C multilayers with a d-spacing of 3.0 nm were fabricated by magnetron sputtering with different background pressures. The grazing incidence x-ray reflectivity (GIXR) and transmission electron microscopy (TEM) measurement illustrated the structural change that happened in NiV/B<sub>4</sub>C multilayers when background pressure is high. The electron dispersive x-ray spectroscopy (EDX) of NiV/B<sub>4</sub>C multilayer deposited with a high background pressure suggests a gradient distribution of oxygen, which corresponds to the gradient thickness change. The detailed x-ray absorption near edge spectroscopy (XANES) comparison of NiV/B<sub>4</sub>C multilayers, NiV coating, and B<sub>4</sub>C coating showed the chemical state change induced by background pressure. We concluded that during the deposition, vanadium oxide promoted the oxidation of boron. In order to fabricate a good performance of NiV/B<sub>4</sub>C multilayers, the background pressure needs lower than 1 × 10<sup>-4</sup> Pa.

Keywords: NiV/B<sub>4</sub>C, multilayer, magnetron sputtering, XANES, background pressure

# **1 INTRODUCTION**

Multilayer interference coating is an important optics in the hard x-ray region, which enables the reflection of the short-wavelength light beyond the total external reflection region, and chooses a narrow-band spectrum from the white light radiation. Thus, it is widely used in advanced imaging, spectroscopy, and monochromator systems for synchrotron radiation facilities and astronomical observation projects [1-3]. As photon energy requirement of synchrotron radiation facilities and space telescopes extends to several tens of kiloelectron volts, the multilayer working at such wavelength region requires an extremely small d-spacing which is less than 3.0 nm [4-7]. This makes the x-ray reflectance very sensitive to the interface and surface quality of the multilayer, which also needs the more precise control of fabrication and the more suitable selection of materials. As for fabrication, magnetron sputtering has been proved as an effective technique for fabricating small d-spacing multilayers, which has the good controllability of the process and the relatively low cost of the instrument [8]. For materials selection, more and more researchers began using a compound to replace the elementary substance [9, 10].

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In the hard x-ray region, a multilayer monochromator has the advantage of a large spectral bandwidth and high reflectance, which can provide one or two orders of magnitude higher integral flux than a crystal monochromator [11]. Nickel is one of the ideal absorption materials in the whole hard x-ray region, except for 8.3 keV (Ni K-edge) [12, 13], but nickel is a magnetic material, which is not compatible with magnetron sputtering for small d-spacing multilayers. For solving this problem, a nonmagnetic alloy was made up of 93% nickel, and 7% vanadium has been proposed (in the following article, we will use NiV instead of the more correct Ni<sub>93</sub>V<sub>7</sub>). At the beginning, NiV has been suggested by Danish Space Research Institute as a potential candidate for hard x-ray telescope. [14] Because compared with tungsten and platinum, there was no absorption line for NiV at the photo energy range from 20 to 200 keV [9]. Until 2017, NiV was first used as a multilayer monochromator in European Synchrotron Radiation Facility (ESRF) [15]. Given the higher reflectivity and integral flux, the NiV/B4C combination was thus chosen as the material of choice for a monochromator working at 8 keV and around 20 keV [16-18]. Apart from that, the detailed studies on NiV/B<sub>4</sub>C multilayers were really limited, especially on how to fabricate this promising X-ray multilayer mirror.

During the magnetron sputtering process, there are a lot of factors that will influence the final performance of multilayers. For example, the pressure of working gas (typically argon) impacts stress and the layer density, and the speed of the spin motion of the substrate during deposition influences the uniformity of the coating. The background pressure is also one of the parameters that needs to be characterized and precisely controlled. It is reported that the background pressures affect the stress and layer morphology of layers [19, 20]. The residual gas in vacuum chamber will react with sputtering atoms and alter the refractive index contrast between multilayer materials, furthermore changing the film structure by the formation of compounds and impacting the interface quality [21, 22]. In order to understand the background pressure effect on the growth of the NiV/B<sub>4</sub>C multilayer, a comparative study of the NiV/B4C multilayers fabricated with different background pressures was performed in this paper. The multilayer structure, layer morphology, and chemical state of the multilayers were investigated and discussed systematically.

## 2 EXPERIMENTS

In order to realize the application of NiV/B<sub>4</sub>C multilayers in hard x-ray telescope or multilayer monochromator, the multilayer d-spacing in this research was set to d = 3.0 nm and the thickness ratio  $\gamma = d_{\text{NiV}}/d = 0.6$ . The number of bilayers was N = 60 to test the period drift of multilayer. The NiV/B<sub>4</sub>C multilayers were fabricated by direct current magnetron sputtering technique. The B<sub>4</sub>C (purity 99.5%) target was used at a constant power of 80 W, the NiV target (Ni:V = 93:7, purity 99.99%) at 20 W. Both targets have a diameter of 4 inches. The distances from substrate to target are 8.5 and 9.0 cm, respectively to NiV and B<sub>4</sub>C. During the deposition, high-purity argon (99.999%) was used as the working gas at a pressure of

0.16 Pa, for keeping a relatively high kinetic energy of the deposited atoms and, thus, to form smooth interfaces, and the flow rate for Ar is 8 sccm. The multilayers were deposited on superpolished Si (100) wafers (20 mm × 20 mm) with a root-meansquare roughness of 0.2 nm  $(1 \times 1 \,\mu\text{m}^2 \text{ atomic force microscope})$ scan). The individual layer thicknesses in the multilayer were, thus, adjusted by controlling the sputtering time independently over each target. The background pressures, which arise from residual air in the deposition chamber, were measured before inflating Ar. In each case, the chamber was pumped from the normal air condition to the background pressures of  $4 \times 10^{-4}$ ,  $2 \times 10^{-4}$ ,  $1 \times 10^{-4}$ ,  $8 \times 10^{-5}$ ,  $4 \times 10^{-5}$ , and  $2 \times 10^{-5}$  Pa, by using a combination of a turbomolecular pump and a mechanical pump. The fabricated multilayers are symbolized as  $S-4 \times 10^{-4}$ ,  $S-2 \times 10^{-4}$ ,  $S-1 \times 10^{-4}$ ,  $S-8 \times 10^{-5}$ ,  $S-4 \times 10^{-5}$ , and  $S-2 \times 10^{-5}$  Pa, respectively. Except for background pressure, all the deposition parameters were the same in all cases.

The grazing incidence x-ray reflectivity (GIXR) measurement used a laboratory-based diffractometer (Bede, Durham, UK) with the Cu K $\alpha$  line as the source ( $\lambda = 0.154$  nm). Fits of the GIXR curve were performed using the Bede Refs software [23]. The information of the individual layer thickness and interface width was determined from the fitting.

High-resolution transmission electron microscopy (TEM) measurements were made to compare the microstructure and layer morphology of the two samples deposited with different background pressures. The samples were prepared by focused ion beam and observed by FEI Talos F200X (FEI, Hillsboro, OR, United States). The elemental composition of the layers was further measured by energy dispersive x-ray spectroscopy (EDX) during the TEM measurements. Mapping and one-dimensional line scans across several bilayers were performed to measure the composition depth profiles of the different elements.

X-ray absorption near-edge structure (XANES) spectroscopy was commonly used to probe the local atomic structure around specific elements in films and often used as a fingerprint of the changes in the coordination numbers of elements [24–26]. XANES measurements were conducted after fabrication at the Beamline U12b at the National Synchrotron Radiation Laboratory (NSRL) in the total electron yield (TEY) mode by collecting sample drain currents under a vacuum, which was greater than  $5 \times 10^{-8}$  Pa. The beam from a bending magnet was monochromatized with a varied line-spacing plane grating and refocused using a toroidal mirror. An energy range from 100 to 1,000 eV was used with an energy resolution of approximately 0.2 eV. All XANES spectra were measured using the process of normalization by considering the low and high photon energy parts of the spectra far from the threshold [26, 27].

## **3 RESULTS AND DISCUSSION**

# 3.1 Grazing Incidence X-Ray Reflectivity Measurement Analysis

The GIXR measurements of as-deposited samples are shown in **Figure 1A**. The different colors represent different background pressures. As the background pressure becomes lower, the Bragg



peaks become narrower, which indicates that the multilayer sample has less d-spacing drift from the surface to the substrate. Besides the width of Bragg peaks, the positions of Bragg peaks also changed with different background pressures. With the background pressure decreasing from  $4 \times 10^{-4}$  to  $1 \times 10^{-4}$  Pa, the grazing angles of Bragg peaks kept increasing, which means the d-spacing of multilayer becomes thinner. When the background pressure reached  $8 \times 10^{-5}$  Pa and further lower, the grazing angle of Bragg peaks did not change, indicating that the d-spacings of different samples became consistent with varied background pressures. Figure 1B shows the fitted result of NiV/B4C multilayer deposited with background pressure of  $4 \times 10^{-5}$  Pa. The fitted curve matched very well with GIXR measurement result, which means the fitted model is in agreement with the deposited layer structure, and the bandwidth of the measured Bragg peaks is similar with the fitted curve, which indicates the good uniformity of 60 bilayers. The fitted values of two individual layers are also shown in this figure; the thicknesses of NiV and B<sub>4</sub>C were 1.65 and 1.10 nm, respectively. The average interface width for the two kinds of materials were 0.29 nm (NiV) and 0.31 nm (B<sub>4</sub>C).

Based on GIXR measurements, we can assume that the change in layer structure happened in NiV/B<sub>4</sub>C multilayers when the background pressure was higher than  $1 \times 10^{-4}$  Pa because all of the parameters for GIXR measurement and deposition are the same except for background pressure. There are two most possible reasons for broader Bragg peaks, one is the gradient d-spacing change in a multilayer, and the other one is the gradient refractive index change. How did the background pressure induce a gradient change in multilayer? We need further analysis to investigate how the change happened in NiV/B<sub>4</sub>C multilayers.

# **3.2 Transmission Electron Microscope Analysis**

The cross-sectional TEM images and selected area electron diffraction (SAED) patterns of NiV/B<sub>4</sub>C multilayer deposited

with different background pressures are shown in Figure 2. The bright layers are B<sub>4</sub>C, and the dark layers are NiV. Based on lowmagnification images, we found that the total thickness of sample  $S-4 \times 10^{-4}$  Pa was 47.7 nm thicker than that of sample S-4  $\times$  $10^{-5}$  Pa, which confirmed the structural change induced by background pressure. In Figures 2A,B, both samples have flat layers, but the interfaces of sample  $S-4 \times 10^{-5}$  Pa are shaper. In the high-resolution images (Figures 2C,D), it can be easily found that the individual d-spacing of S-4  $\times$  10<sup>-4</sup> Pa was different from S-4  $\times$  $10^{-5}$  Pa, indicating the different positions of Bragg peak in Figure 1A. Based on further analysis of high-resolution images, we found that the d-spacing of sample S-4  $\times$  10<sup>-4</sup> Pa changed gradually from the surface to the substrate, which range from 3.38 to 4.18 nm. Such a gradient d-spacing will cause the broadened Bragg peak as shown in Figure 1A. The thicker layers were close to the substrate, and the B<sub>4</sub>C layers changed much more than the NiV layers. As for sample S-4  $\times$  10<sup>-5</sup> Pa, the d-spacing was constant for the whole multilayer. The d-spacing was 2.79 nm measured from high-resolution images, which was close to the fitted result (shown in Figure 1B) of GIXR measurement and indicating that the fitted model was similar with the actual structure. SAED images were recorded with a large electron beam covering several tens of bilayers of the NiV/B<sub>4</sub>C multilayers. Figures 2E,F display the SAED images of samples S-4  $\times 10^{-4}$  and S-4  $\times 10^{-5}$  Pa, respectively. Both of them only have a diffraction pattern coming from the multilayer structure indicating an amorphous layer structure in these samples [7]. However, compared with  $S-4 \times 10^{-4}$  Pa, the electron diffraction pattern of S-4  $\times$  10<sup>-5</sup> Pa is more regular, which means that the layer structure in S-4  $\times$  10<sup>-5</sup> Pa is more uniform.

To our knowledge, the background pressure may cause the change in deposition rate for sputtering materials, but would not induce a gradient d-spacing in multilayers. Therefore, the only physical reason that can explain the formation of a gradient d-spacing is the element composition change in the multilayer due to the residual gas in the chamber [21, 22]. The element composition change will also cause the refractive index change.





Analysis of this hypothesis will be performed in the following section.

# **3.3 Energy-Dispersive X-Ray Spectroscopy** Analysis

The higher background pressure means more residual gas stayed in the chamber. During the deposition process, some gas would be activated by plasma environment, then forming compounds with sputtering particles. In order to analyze the elemental composition of samples S-4 × 10<sup>-4</sup> and S-4 × 10<sup>-5</sup> Pa, EDX mapping and onedimensional line scans across several bilayers were performed during the TEM measurement. In our work, the EDX measurements were performed to characterize elements of boron (B), carbon (C), nitrogen (N), oxygen (O), nickel (Ni), and vanadium (V).

The EDX measurement allowed to approximately determine the atomic concentration of different elements. In **Figures 3A,B**, the high angle annular dark field (HAADF) images are combined with element mapping images, for observing the element distribution in samples S-4  $\times$  10<sup>-4</sup> and S-4  $\times$  10<sup>-5</sup> Pa, respectively. Vanadium and nitrogen mapping images are not included in these figures because of the low atomic concentration (vanadium is only 7% in NiV target). In HAADF images, the bright layers are NiV and the dark layers are B<sub>4</sub>C. As shown in Figures 3A,B, the Ni and C element distributions are both in accordance with HAADF images; the main difference happened in O and B elements distributions. In sample S-4  $\times$  10<sup>-4</sup> Pa, oxygen distribution exhibits the layer structure, and these layers are in keeping with B<sub>4</sub>C layers, which means a lot of oxygen atoms stayed in the B<sub>4</sub>C layers. On the contrary, the oxygen distribution in sample  $S-4 \times 10^{-5}$  Pa does not have such an obvious layer structure. Because the sensitivity for boron detection in the EDX measurement is not good, the boron distributions in both samples are not as clear as the other elements, but in sample S-4  $\times 10^{-5}$  Pa, boron distribution seems to have a blurred layer structure compared with S-4  $\times$  10<sup>-4</sup> Pa.

The detailed one-dimensional line scans of samples  $S-4 \times 10^{-4}$ and  $S-4 \times 10^{-5}$  Pa are shown in **Figures 3C,D**. In sample S-4 ×



 $B_4C$  multilayers fabricated with background pressures of  $4 \times 10^{-4}$  Pa (A,C) and  $4 \times 10^{-5}$  Pa (B,D).

 $10^{-4}$  Pa (**Figure 3C**), except for boron, the other elements both have a periodic oscillation of concentration. The element with the most remarkable profile is oxygen, the average atomic concentration of oxygen in the whole scan area is more than 30%, and from the surface to substrate, the atomic concentration of oxygen keeps increasing, which is consistent with the d-spacing change in sample S-4  $\times$  10<sup>-4</sup> Pa. In sample S-4  $\times$  10<sup>-5</sup> Pa (Figure 3D), all measured elements have a periodic oscillation of concentration, which further illustrates that sample S-4  $\times$ 10<sup>-5</sup> Pa has a uniform layer structure.

According to the EDX measurement results, we find that oxygen is the main reason for the gradient d-spacing in NiV/ B<sub>4</sub>C multilayer. When the deposition process happened with a higher background pressure, the oxygen in residual gas would interact with boron and carbon atoms. Then the tri-element compound layer replaces B<sub>4</sub>C and causes the volume expansion. As the deposition time went by, the oxygen concentration in the chamber becomes lower, the tri-element compound layer also becomes thinner, which means the volume of the tri-element compound is proportion to the concentration of oxygen. Due to this phenomenon, the d-spacing of the multilayer is gradually thinning from the substrate to the surface. This gradient change is consistent with the results obtained from the GIXR and TEM measurements. Except for the volume change, the refractive index of this tri-element compound layer may also change with the

oxygen concentration. When the deposition process happened with a lower background pressure, the oxygen concentration also lower in the residual gas. Although the oxygen still has the opportunity to interact with sputtering particles (which has been proved in our previous work [4]), the lower concentration of oxygen would not cause the obvious thickness change in the boron carbide layer, which promote the uniformity between bilayers (as shown in Figures 1B, 3D).

# 3.4 X-Ray Absorption Near-Edge Structure **Spectroscopy Analysis**

In order to further explore the chemical change caused by different background pressures, XANES measurements have been implemented in several samples (S-4  $\times$  10<sup>-4</sup>, S-1  $\times$  10<sup>-4</sup>,  $S-8 \times 10^{-5}$ ,  $S-4 \times 10^{-5}$ ,  $S-2 \times 10^{-5}$  Pa). Except for multilayer samples, we also fabricated two single-layer samples for comparison. One is NiV coating marked as S-NiV in the following part, and this sample was deposited with a background pressure of 1  $\times$   $10^{-4}\,\text{Pa.}$  Another one is  $B_4C$ coating marked as S-B<sub>4</sub>C, which was deposited with a background pressure of  $6 \times 10^{-5}$  Pa. These two background pressure values were suggested by GIXR and TEM measurement results to mitigate the influence caused by residual gas. The thicknesses of S-NiV and S-B<sub>4</sub>C are 16 and



35 nm, respectively, which fitted the GIXR measurements. Because the purpose for this study is investigating the influence from background pressure, pre- and post-edge normalizations have not been carried out during the data analysis in this paper.

# 3.4.1 B K-Edge X-Ray Absorption Near-Edge Spectroscopy

The B K-edge absorption spectra of five NiV/B<sub>4</sub>C MLs are presented in **Figure 4A**. There are three prominent features in the B-K absorption spectra, which are marked as a, b, and c, respectively. Features a and b are both narrow peaks appearing at photon energies of ~192.8 and ~194.2 eV. Feature c looks like a broad bump in the high-energy regime (195–205 eV). The narrow edge peak a (192.8 eV) occurs due to the transition of a B 1s electron to the unoccupied B 2p  $\pi^*$  state. According to Jiménez' s research, this feature seems related to inter-icosahedral B–B bonds [28]. The other narrow peak b (194.2 eV) is due to transition of a B 1s electron to the unoccupied B 2p  $\pi^*$  state in B in the presence of oxygen [29, 30]. The feature c is due to the transition of B 1s electrons to the unoccupied B  $\sigma^*$  states, which corresponds to bonding within the C–B–C chain structure [31, 32].

With the background pressure changing from  $4 \times 10^{-4}$  to  $2 \times$  $10^{-5}$  Pa, the relative intensities of features a and b become lower, but the intensity of feature c becomes higher. This result indicates the sputtered boron atoms tend to bond with carbon atoms to form a C-B-C chain with a lower background pressure in a vacuum chamber. These C-B-C chains mainly stay in boron carbide layers and make the boron carbide layer have a regular layer structure, which is proven by boron distribution in Figure 3D, but with a higher background pressure, the oxygen concentration becomes higher, the sputtered boron atoms seem to prefer to combine with oxygen atoms. Therefore, the oxygen atoms destroyed parts of the C-B-C chains and formed a lot of B-O bonds. These B-O bonds stayed in boron carbide layers and formed a tri-element compound layer, which is consistent with the oxygen distribution in Figure 3C. The increased B-B bonds may come from some dissociated boron in NiV/B<sub>4</sub>C ML [33],

which results in an aperiodic boron distribution shown in EDX measurements (Figure 3C).

In **Figure 4B**, the B K-edge absorption spectra of NiV/B<sub>4</sub>C ML and B<sub>4</sub>C coating are presented. The main difference between these two curves is the feature b, which is marked by an orange line. The background pressure during the deposition of S-B<sub>4</sub>C (6  $\times 10^{-5}$  Pa) is higher than S6 (2  $\times 10^{-5}$  Pa), but S6 has a higher intensity at feature b, which means the boron in NiV/B<sub>4</sub>C ML is easier to bond with oxygen. We suspect in NiV/B<sub>4</sub>C ML that boron was excited by other elements then tend to combine with oxygen. There is a material worked as catalyst [34] during the deposition of NiV/B<sub>4</sub>C ML.

## 3.4.2 V L-Edge and O K-Edge X-Ray Absorption Near-Edge Spectroscopy

The V L-edge and O K-edge XANES were measured to investigate the catalysis of the boron oxidation that happened in NiV/  $B_4C$  ML. The V L-edge and O K-edge absorption spectra of samples S-NiV, S-1 × 10<sup>-4</sup> Pa and S-B<sub>4</sub>C are presented in Figure 5A. S-NiV and S-1 × 10<sup>-4</sup> Pa were deposited with the same background pressure. There are much more features in V L-edge and O K-edge XANES compared with B K-edge. Each feature is labeled with different letters shown in Figure 5A.

In the spectrum of S-NiV, the most distinct dual-peak feature (marked as e) at the region from 519.6 eV (left peak; known as  $L_3$ -edge) to 525.7 eV (right peak; known as  $L_2$ -edge), which correspond to the electronic transitions from the V  $2p_{3/2}$  and V  $2p_{1/2}$  core levels, respectively, to the unoccupied 3d levels [35]. Based on previous researches, the spectrum of S-NiV is similar with the V<sub>2</sub>O<sub>5</sub>, especially that they both have the small peak (peak d) besides the V  $L_3$ -edge [36]. This small peak (centered at 513.0 eV) can be assigned to the splitting of the unoccupied V 3d orbitals, which also indicates that the V<sub>2</sub>O<sub>5</sub>-liked structured formed in the S-NiV sample during the deposition. The features located at the energy region higher than 530.0 eV originated from the electronic transitions from O 1s to O 2p hybridized V 3d orbitals and are known to provide valuable information on the crystal structure, local symmetry of metal ions, and ligand-field



effects [37, 38]. In an octahedral environment, the metal  $e_g (d_z^2, d_{x-y}^2)$  orbitals are pointed directly toward the O ligands to form strongly bonded  $\sigma$  bonds with the O  $p_z$  orbital. On the other hand, the metal  $t_{2g} (d_{xy}, d_{xz}, d_{yz})$  orbitals point between the O ligands and make weak  $\pi$  bonds with O  $p_x$  and  $p_y$  [38]. In **Figure 5A**, peak f (531.7 eV) and peak g (534.1 eV) correspond to  $e_g$  and  $t_{2g}$ , respectively. The intensity of  $e_g$  is lower than  $t_{2g}$ , indicating that the V<sub>2</sub>O<sub>5</sub>-liked structured formed in S-NiV is closer to tetrahedrally coordinated compounds [38, 39]. The broad feature i is due to the transition of O 1s electrons to the unoccupied O  $\sigma^*$  states [38, 40].

Compared with S-NiV, the absorption spectrum of S-1 × 10<sup>-4</sup> Pa shows less features. The dual-peak feature has a blue shift, L<sub>2</sub>-edge centered at 517.0 eV, and L<sub>3</sub>-edge changed to 523.9 eV, which means that the oxidation state of vanadium changed [41]. The position of features g and i are consistent with S-NiV, indicating that the coordinated structure around O atoms is similar, but the features d and f become absent, which means that the vanadium oxide in S-1 × 10<sup>-4</sup> Pa is amorphous or highly disordered [36]. The O K-edge absorption spectrum of S-B<sub>4</sub>C only has two features, one looks like a shoulder around 535.9 eV (labeled as h), which is also shown in the curve of S-1 × 10<sup>-4</sup> Pa. The other one is also a broad feature, but peak has a little bit of shift compared with feature i in S-1 × 10<sup>-4</sup> Pa and S-NiV. This shift is because the local atomic structure around O atoms was changed in different samples.

In **Figure 5B**, we compared NiV/B<sub>4</sub>C MLs deposited with different background pressures. There are four orange dash lines corresponding to features e, g, h, and i (shown in **Figure 5A**), respectively. With the background pressure becoming lower, the L<sub>3</sub>-edge has a gradual blue shift, which means that the oxidation state of vanadium keep changing toward a higher oxidation state. All the five curves both have feature h, and this feature is located at the same position, but the relative intensity of feature h also becomes lower as the background pressure changed. This is the same with feature b shown in B K-edge XANES because these two features are both related to B–O bonds. As for features g and i, the curve of S-4 × 10<sup>-4</sup> Pa is totally different with other samples. The feature g in the curve of S-4 × 10<sup>-4</sup> Pa is absent, and feature i has a

shift compared with others. This shift is also shown in **Figure 5A**. Except for  $S-4 \times 10^{-4}$  Pa, in the other samples, features g and i are both located at the same position. The relative intensity of feature i has a slight change with the changed concentration of oxygen in a vacuum chamber. On the contrary, the relative intensity of feature g seems to become higher when oxygen atoms become fewer, which means the structure of vanadium oxide tends to become ordered with less oxidation of boron.

Based on the comparison of the V L-edge and O K-edge XANES of different samples, we can conclude that several facts had happened during the magnetron sputtering deposition process: 1) Vanadium will form vanadium oxide during the deposition even with a background pressure of  $2 \times 10^{-5}$  Pa, and vanadium oxide is a well-known catalyst [34], which enhanced the combination of boron and oxygen. 2) As the combination of boron and oxygen has been catalyzed, the boron will grab more and more oxygen atoms, which will prevent the vanadium oxide from going to its highest oxidation state. 3) When the background pressure is lower than  $1 \times 10^{-4}$  Pa, the local atomic structure around O atoms does not change a lot, which means the chemical state and the layer structure are relatively stable in NiV/B<sub>4</sub>C ML.

# 3.4.3 Nickel L-Edge X-Ray Absorption Near-Edge Spectroscopy

The Ni L-edge x-ray absorption near-edge structure spectra were measured to analyze the change in local atomic structure around the nickel atoms. The Ni L-edge absorption spectra of five NiV/ $B_4C$  MLs are presented in **Figure 6A**. There are two features in the measurement results, marked as j and k. The feature j centered at 857.3 eV (assigned at Ni L<sub>3</sub>-edge) and the feature k located at 875.0 eV (assigned at Ni L<sub>2</sub>-edge), which correspond to the electronic transitions from the Ni  $2p_{3/2}$  and  $2p_{1/2}$  ground states, respectively, to the unoccupied Ni 3d final states [42, 43]. As shown in **Figure 6A**, the relative intensities and positions of the different samples are the same, which means the background pressure does not cause any change in the Ni element in NiV/B<sub>4</sub>C ML.



In **Figure 6B**, we compared the Ni L-edge absorption spectra of S-NiV and S- $1 \times 10^{-4}$  Pa. We found that features j and k both have a blue shift (~1.9 eV) in the curve of S-NiV. In the NiV coating, only vanadium and oxygen atoms are around Ni atoms, but in a multilayer, Ni atoms also adhere to the interface between NiV and B<sub>4</sub>C layers. We suspected that the presence of interface caused the change in local atomic structure around Ni atoms, which caused the shift of Ni L-edge.

# **4 CONCLUSION**

NiV/B<sub>4</sub>C multilayers fabricated with different background pressures were studied comparatively. When the background pressure was higher than  $1 \times 10^{-4}$  Pa, an obvious structural change has been found in multilayers. This structural change is mainly shown in the thickness change in the B<sub>4</sub>C layer. According to the EDX measurement, the thickness change is directly proportionate to the oxygen concentration in a vacuum chamber. Therefore, a gradient d-spacing change happened in a sample deposited with a high background pressure. Such a change will make the reflectivity spectrum broadened, which causes a severe performance loss when using an NiV/B4C combination for a multilayer monochromator. XANES has been applied to further investigate the chemical change that happened in NiV/B<sub>4</sub>C multilayers deposited with different background pressures. The oxidation of vanadium and boron has been found in all samples, which means the oxidation of active elements is inevitable during magnetron sputtering process, as long as the oxygen atoms stay in the vacuum chamber. After comparing the NiV/B<sub>4</sub>C multilayer, NiV coating, and B<sub>4</sub>C coating, an interesting phenomenon has been found. When the oxygen concentration is lower in the vacuum chamber, the vanadium oxidates work as a catalyst, which promotes the combination of boron and oxygen, but in the higher oxygen concentration environment, boron will "rob" the majority of oxygen atoms, which prevents the oxidation state of vanadium oxide toward its highest oxidation state. Besides the NiV/B4C multilayers, a similar catalysis has been found in the Pd/

B<sub>4</sub>C multilayers [27]. Due to the catalysis, the ultrathin Pd/B<sub>4</sub>C multilayers degraded rapidly when stored in air. Therefore, the lifetime and stability of NiV/B<sub>4</sub>C multilayers is worth to be further studied. In this study, the structural and chemical changes both suggest that the NiV/B<sub>4</sub>C multilayers need to fabricate with a background pressure lower than  $1 \times 10^{-4}$  Pa. Besides, the chemical state change found in this work can also provide useful guidance for further development and fabrication of NiV/B<sub>4</sub>C multilayers.

# DATA AVAILABILITY STATEMENT

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

# AUTHOR CONTRIBUTIONS

ZZ and ZWa conceived and designed the research work. QH and RQ carried out the simulations. ZWe, YY and LJ prepared the samples. ZWe, ZZ, YF and CC analyzed the data. ZWe and ZZ wrote the paper. WY, CX and ZWa contributed to the helpful discussions for the research and modifications to the article. All authors have read and agreed to the published version of the manuscript.

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