



Compressive Interferenceless Coded Aperture Correlation Holography With High Imaging Quality

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Interferenceless coded aperture correlation holography (I-COACH) provides an alternative way for the 3D imaging of spatial incoherent illuminated or fluorescent sample. However, the low imaging signal-to-noise ratio (SNR) is one of the bottlenecks that restrict the application of I-COACH. The limitation is mainly originated from the strong bias level that presents in the recorded holograms. Phase shifting methods were implemented in I-COACH to eliminate the background noise while the multiple-exposures recording mechanism significantly reduces the temporal resolution of the system. In this paper, we proposed a compressive I-COACH imaging method with high reconstruction quality and without the sacrifice of the imaging speed. The 3D holographic image reconstruction was implemented under compressive sensing framework while only one single-exposure object hologram and one point spread hologram are necessary. High quality reconstructions were obtained using the proposed method, even for the downsampled holograms. The imaging SNR of the I-COACH system was improved by a factor of more than 16.5% when comparing with the imaging SNR obtained by the conventional cross-correlation reconstruction method. The proposed method provides a fast and high-fidelity imaging method that can potentially benefit the imaging through scattering medium, partial aperture imaging, and other fields.

Keywords: incoherent holography, coded aperture imaging, 3D imaging, compressive sensing, high quality reconstruction

INTRODUCTION

Holography based on two-beam mutual-coherent and interference has unique advantages in the field of quantitative phase contrast imaging and 3D imaging because it can record and reconstruct not only the amplitude but also the phase of the light that emitted or reflected from the object. However, coherent sources are not always practically available, especially in some important applications such as astronomical imaging and fluorescence microscopy. In addition, the coherent speckle noise in the reconstructed images has significantly affected the imaging performance of the holographic system (Men et al., 2017). On the other hand, incoherent holography that first was proposed by Mertz and Young in 1961 has extended the application of holography to the field of spatial incoherent imaging (Rogers, 1950; Mertz et al., 1961). The basic idea of original incoherent holography is that the spatial incoherent objects can be 3D reconstructed from their Fresnel zone plate projection patterns (hologram). The incoherent holographic techniques were then further developed and improved by Lohmann, Stroke, and Cochran (Lohmann, 1965; Stroke and Restrick, 1965; Cochran, 1966). One significant difference between incoherent holography and conventional holography is that the

OPEN ACCESS

Edited by:

Yaping Zhang, Kunming University of Science and Technology, China

Reviewed by:

Jianglei Di, Guangdong University of Technology, China Shuming Jiao, Peng Cheng Laboratory, China

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Specialty section:

This article was submitted to Optical Information Processing and Holography, a section of the journal Frontiers in Photonics

Received: 31 December 2021 Accepted: 25 January 2022 Published: 04 March 2022

Citation:

Liu C, Man T and Wan Y (2022) Compressive Interferenceless Coded Aperture Correlation Holography With High Imaging Quality. Front. Photonics 3:846731. doi: 10.3389/fphot.2022.846731

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mutual coherence between any two different points on the sample is no longer necessary for holographic recording. Instead, the holographic recording scheme was based on the characteristics of spatial self-coherence by splitting the light that emitted from each point on the object into two beams and then recording the resulting interference patterns. However, incoherent holographic methods are normally failed in providing images with sufficient SNR. The problems are mainly arising from the strong bias in the recorded holograms as the results of the incoherent superposition of all the point holograms.

A new type of incoherent digital holography, coined as coded aperture correlation holography (COACH), was proposed in 2016 (Vijayakumar et al., 2016), which originated from Fresnel incoherent correlation holography (FINCH) (Men et al., 2017). COACH possesses the advantages of higher axial-resolution and higher spectral-resolution when comparing with FINCH. In COACH, the pseudorandom coded phase mask (CPM) generated using Gerchberg-Saxton (GS) algorithm is loaded onto the spatial light modulator (SLM), and the self-interference holographic recording is realized based on the beam splitting using the polarization sensitive characteristics of the SLM (Vijayakumar et al., 2016). The 3D reconstruction of the object is realized by cross-correlating the object hologram (OH) with the point spread holograms (PSHs) that have been recorded at all possible axial planes. That the 3D information of an object can be encoded and decoded even without the use of self-interference recording was found by later research. The interferenceless coded aperture correlation holography (I-COACH) was proposed for 3D imaging without any two-wave interference (Vijayakumar and Rosen, 2017). Comparing with COACH, I-COACH has a much simpler optical setup and improved photon-energy efficiency. However, the temporal resolution of the system is sacrificed because multi-exposure phase shifting technique and image averaging method were used to suppress the background noise that dominates the reconstructed images. Different techniques such as phase filtering (PF) (RatnamRai et al., 2017), non-linear reconstruction (NLR) (Rai and Rosen, 2019; Liu et al., 2020), and adaptive reconstruction (Wan et al., 2021) were used to suppress the background noise and improve the temporal resolution of the I-COACH system. However, the increases of temporal resolution normally come with the decreases of reconstruction quality. On the other hand, compressive sensing (CS) theory has demonstrated that signals, assumed to be sparse in some basis and sampled by multiplex encoding systems, may be accurately inferred with high probability from much fewer measurements than suggested by Nyquist's sampling theorem. In digital holography, the CS technique provides a high-efficiency way for 3D object reconstruction from single real-valued (intensity) hologram. Specifically, the 3D sample is modulated by the optical system to form a 2D hologram, and the reconstruction process is implemented under the CS framework to achieve 3D tomographic imaging of the sample (Zhang et al., 2018). In compressive digital holography, the theoretical model of the holographic recording procedure of a 3D object was first established based on vector diffraction theorem. The forward and backward propagation models, with sparsity regularization imposed on the object, were then used to solve the so-called under-determined inverse problem that will be met when trying to reconstruct 3D objects from their

single-exposure digital hologram. As a result, the bias and twinimages can be suppressed (Brady et al., 2009; Lim et al., 2011; Brady et al., 2015; Zhang et al., 2018).

In this paper, inspired by the basic idea of compressive holography, we propose a compressive I-COACH 3D imaging method with high imaging quality and most importantly without the sacrifice of the system temporal resolution. In CS-I-COACH, only single-exposure PSH and OH are necessary for the 3D reconstruction of the object. The imaging model under the CS frame of the I-COACH was first discussed, followed by the preliminary experimental results on 2D and 3D objects that demonstrate the validity of the proposed method. The dependence of the CS reconstruction accuracy in I-COACH on the scattering degree of the CPM was analyzed. The implementation of CS high fidelity reconstruction for under-sampled holograms was also discussed. We believe the proposed method provides a powerful research tool and benefits the optical research fields such as 3D tracking, imaging through scattering medium, and partial aperture imaging.

THE SENSING MODEL OF THE I-COACH SYSTEM

Nyquist sampling frequency is a sufficient but not necessary condition for the high-fidelity retrieval of a discrete signal. CS theory demonstrates that if the signal is sparse itself or can be represented sparse in a certain domain (i.e., Fourier domain or wavelet domain), the original signal can be reconstructed with high probability using numerical optimization approaches, even when the sampling rate is much lower than the value required by the Nyquist criterion. CS can be used for image reconstruction in imaging systems because most natural images are sparse or have sparsity in some transform domains. Assume the image *f* can be expressed on the orthogonal base $\varphi = \{\varphi_1, \varphi_2, ..., \varphi_n\}$ as:

$$f = \sum_{i=1}^{n} \varphi_i \mu_i = \varphi \mu \tag{1}$$

where the coefficient $u = [u_1, u_2...u_n]^T \in \Re^{n \times n}$. If $\|\mu\|_0 \le k$, $(\|\cdot\|_0$ is defined as the l_0 -norm and k is the number of non-zero elements in $u, k \ll n$), then f is a k-sparse under matrix φ , while φ is a sparse operator. Given the measurement matrix $R \in \mathbb{C}^{m \times n}$, the matrix that describes the imaging procedure and details of object-image transformation, the measured image g can be expressed as:

$$g = Rf = R\sum_{i=1}^{n} \varphi_{i}\mu_{i} = R\varphi\mu = \Theta\mu$$
(2)

where $\Theta \in \mathbb{C}^{m \times n}$ is the sensing operator, $g \in \mathbb{C}^{m \times 1}$ is the measurement, and the dimension *m* of *g* is much smaller than the dimension *n* of *f*. Then an accurate reconstruction can be obtained with high probability by solving:

$$\hat{\mu} = \arg\min_{u} \left\| g - \Theta \mu \right\|_{l_1} \qquad s.t. \quad g = \Theta \mu \tag{3}$$

An I-COACH optical setup is shown in Figure 1. The object or pinhole is critically illuminated with spatial incoherent light, the



diffracted light is collected and collimated by the lens L, and then modulated by the phase mask which comprises pseudorandom CPM and a quadratic phase mask (QPM) with focal length f_s . The light scattered by the phase mask is then projected to the sensor plane by the QPM. The hologram of the object is recorded by the image sensor.

In the system as show in **Figure 1**, for a point object at the position of (x_s, y_s, z_s) , the intensity distribution I_{psh} on the CCD plane can be calculated and given as (Wan et al., 2021):

$$I_{psh}\left(\overline{r_{h}}; \overline{r_{s}}, z_{s}\right)$$

$$== \left| C \sqrt{A_{s}\left(\overline{r_{s}}; z_{s}\right)} L\left(\frac{\overline{r_{s}}}{z_{s}}\right) Q\left(\frac{1}{z}\right) \exp\left[i\varphi\left(\overline{r}\right)\right] \otimes Q\left(\frac{1}{z_{h}}\right) \right|^{2}$$

$$= I_{psh}\left(\overline{r_{h}} - \frac{z_{h}}{z_{s}}\overline{r_{s}}; z_{s}\right)$$
(4)

where the sign \otimes is 2D convolution, $z = z_s f_L f_s / (f_L f + z_s f_L - z_s f_s)$, and $\sqrt{A_s(\overline{r_s}; z_s)}$ is the amplitude of the object point at (x_s, y_s, z_s) . C is a complex constant. $\overline{r_s} = (x_s, y_s)$ is the transverse location vector. $\overline{r_h} = (x_h, y_h)$ is the transverse location vector on the CCD. z_h/z_s is the transverse magnification. The function $\varphi(\overline{r})$ represents the phase of the pseudorandom CPM calculated by the GS algorithm. L and Q are linear and quadratic phase functions, given by $L(\overline{r_s}/z) = \exp[i2\pi(b_x x + b_y y)/\lambda z]$ and $Q(b) = \exp[-i\pi b(x^2 + y^2)/\lambda]$, respectively. In the I-COACH system, the point spread function of the imaging system is actually the point spread hologram (PSH) of the point object.

The system of I-COACH is treated as a linear space-invariant imaging system (Wan et al., 2021). Assume that the intensity response of the *K*-th axial position plane of the 3D object at $z_s = z_k$ to the CCD plane is the convolution of the system point spread function and the 2D object plane. Therefore, the intensity pattern of the object at a certain depth level on the CCD and the hologram reconstruction with pure phase filtering can be given as **Eqs. 5**, **6**, respectively, which show the sensing model under the CS frame of I-COACH.

$$I_{oh}\left(\overline{r_{h}}; z_{k}\right) = O\left(\overline{r_{s}}; z_{k}\right) \otimes I_{psh}\left(\overline{r_{h}}; z_{k}\right)$$
$$= \sum_{j}^{n} a_{j}\delta\left(\overline{r_{s}} - \overline{r_{j}}\right) \otimes I_{psh}\left(\overline{r_{h}} - \frac{z_{h}}{z_{s}}\overline{r_{s}}; z_{k}\right)$$
$$= \sum_{j}^{n} a_{j}I_{psh}\left(\overline{r_{h}} - \left(1 + \frac{z_{h}}{z_{s}}\right)\overline{r}_{j}; z_{k}\right)$$
(5)

$$O'\left(\overline{r_s}; z_k\right) = I_{oh}\left(\overline{r_h}; z_k\right) \otimes I^*_{psh}\left(\overline{r_h}; z_k\right)$$
$$= \mathfrak{P}_{2D}^{-1}\left\{\mathfrak{P}_{2D}\left(I_{oh}\right) \cdot \mathfrak{P}_{2D}\left(I^*_{psh}\right)\right\}$$
$$= \left[\mathfrak{P}_{2D}^{-1} E \mathfrak{P}_{2D}\right] \mu = \Theta \mu \tag{6}$$

where $\mathfrak{F}_{2D}\mu$ corresponds to the calculation of $\mathfrak{F}_{2D}(I_{oh})$, \mathfrak{F}_{2D} represents the 2D discrete Fourier transform matrix, and \mathfrak{F}_{2D}^{-1} represents the inverse 2D discrete Fourier transform matrix. The sensing matrix of CS-I-COACH can be obtained from Eq. 6:

$$\Theta = \mathfrak{T}_{2D}^{-1} E \mathfrak{T}_{2D} = \mathfrak{T}_{2D}^{-1} \exp\left(-i \arg\left(\bar{I}_{psh}\right)\right) \mathfrak{T}_{2D}$$
(7)

The matrix $\Theta = \mathfrak{T}_{2D}^{-1} E \mathfrak{T}_{2D}$ denotes the I-COACH reconstruction process of the hologram. The object reconstruction is done with the pure phase filtering of *E* given by $E = \exp[-i\arg(\mathfrak{T}_{2D}(I_{psh}(\overline{r_h}; z_k)))]$. The compressed reconstruction problem of I-COACH can be denoted as an optimal solution problem, which is solved by the TwIST algorithm (Bioucas-Dias and Figueiredo, 2007). Further considering that the fluctuation of high frequency noise in the reconstructed image is small in the space domain, the total variable (TV) algorithm is chosen to suppress the reconstructed image noise:

$$\hat{u} = \operatorname*{arg\,min}_{u} \left\| g - \Theta u \right\|_{2}^{2} + \tau \| u \|_{TV} \tag{8}$$

As described above, the CS model of I-COACH and the optimized solution algorithm are established, combined with the sensing matrix to suppress the influence of the bias term on the reconstructed image, which is expected to obtain a high quality reconstructed image quickly by single exposure of PSH and OH.

EXPERIMENTS AND RESULTS

Preliminary experiments were carried out to demonstrate the validity and imaging performance of the proposed CS-I-COACH system. In the first experiment, the transmitted object was imaged, and the imaging performances of the CS-I-COACH were quantitatively evaluated and compared with conventional I-COACH methods. The next experiment verified the high SNR imaging capabilities of CS-I-COACH for under-sampled holograms.

As shown in **Figure 2**, the pinhole or object was illuminated by the incoherent light emitting diode (LED) (Thorlabs LED 625L4, 700 mW, center wavelength of $\lambda = 625$ nm, $\Delta \lambda = 17$ nm) placed at the front focal plane of the lens L₀ with focal length $f_0 = 150$ mm. The light from the beam splitter BS₁ was collimated by the lens L₀ and passed through a polarizer P. The polarizer P polarizes the light along the orientation of the active axis of the spatial light modulator (SLM, Holoeye PLUTO, 1,080 × 1920 pixels, 8 µm pixel pitch, phase-only modulation) located at a distance of 55 mm from the L₀. On the SLM, a phase mask is displayed whereas its phase is the combination of a CPM and a QPM with a focal length of $f_s = 150$ mm. The light modulated by the SLM was

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monochrome) located at a distance of $z_h = 171 \text{ mm}$ from the SLM. In the first experiment, the ability of CS-I-COACH to image two-dimensional transmissive objects was demonstrated. Similar with our work before (Wan et al., 2021; Liu et al., 2022), the scattering degree $\sigma = 0.167$ and the annular width ω = 1 pixel were chosen as the optimization parameters of the CPM in the modified GS algorithm. The PSH was recorded by

collected by a Charge Coupled Device (CCD, Thorlabs

CS235MU,1,200 × 1920 pixels, 5.86 µm pixel pitch, and

(a)

hologram.

2022), the scattering degree $\sigma = 0.167$ and the annular width $\omega = 1$ pixel were chosen as the optimization parameters of the CPM in the modified GS algorithm. The PSH was recorded by placing in Channel-1 a 20 µm diameter pinhole at the front focal plane of the input lens L₀. Channel-2 was blocked during the recording of PSH. Then with Channel-1 blocked, the object hologram OH was recorded with a United States Air Force resolution chart (USAF 1951 RES-1, Newport, placed at the front focal plane of the L₀, line-groups of 4 and 5 have been illuminated) as the 2D object. The synthesized coded phase mask (SCPM), the hologram of a 20 µm diameter pinhole (PSH), and the hologram of a USAF target chart (OH) are shown in **Figure 3**, respectively.

The reconstructed images of the single-exposure PSH and OH were obtained using different traditional reconstruction methods and the CS-I-COACH reconstruction method,

respectively, and the results are shown in Figure 4. The reconstructed images from the single-exposure PSH and OH as shown in Figures 3B,C were obtained using different methods of cross-correlation (CC), pure phase filtering (PF), non-linear reconstruction (NLR), and CS-I-COACH, respectively. The results are shown in Figure 4. For the CS-I-COACH, Eq. 8 was solved using the TwIST algorithm with $\lambda_{TwIST} = 0.001$ (iteration = 50). Regularization parameter of $\tau_{TV} = 0.01$ was used. For comparison, the direct imaging result obtained in the same system (by using the SLM as a lens) is shown in Figure 4A. The strong background noise that presents in the CC results (Figure 4B) was suppressed partially by the PF (Figure 4C) and more effectively by the NLR (Figure 4D) reconstruction method. However, the imaging performances of the I-COACH system, especially the background noise level, were still much worse than the direct imaging (Figure 4A). The CS-based methods, on the other hand, provide results with much better quality (Figure 4E) than their conventional reconstruction counterparts. The improvements on the imaging quality can be seen more clearly from the inserted line profiles and the peak-signal-to-noise-ratio (PSNR) (Bioucas-Dias and Figueiredo, 2007) data of the reconstructions (Figure 4F).



(c)

PSH

(b)

FIGURE 3 | The SCPM used in hologram recording and experimentally recorded holograms. (A) Phase of the SCPM; (B) the Point spread hologram; (C) the object

2π



FIGURE 4 Comparison of the results of different reconstruction methods of PSH and OH for single exposure. (A) Lens imaging; (B) cross-correlation; (C) phase filtered cross-correlation; (D) non-linear reconstruction with parameters o = 0.8 and p = -0.2; (E) our proposed CS-I-COACH compression reconstruction method; (F) the PSNR of the reconstructed images from different methods.

	0.04629	0.07407	0.10185	0.13889	0.16667	0.18519	0.21296	0.24074	0.26852	0.28704	0.32407
1e-2	5.9353	5.7955	5.8269	6.3524	5.8459	6.1087	5.8695	6.3039	6.0154	6.1574	6.4073
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Among those, one can notice that the CS-I-COACH reconstruction has better SNR than all the conventional I-COACH reconstruction methods and direct imaging results. It turned out to be that for all the different methods, CS-based I-COACH reconstructions can always have an improved PSNR by a factor of 45.8% (CC), 42% (PF), and 16.5% (NLR). Since additional processing is required in NLR methods to find the optimal modulation

parameters (o and p in Eq. 7) (Liu et al., 2020), the more time-efficient phase-filtering-based methods (PF and CS-I-COACH) that provide similar imaging performance were used in the following experiments on 3D objects.

The CS reconstruction accuracy is highly correlated with the incoherence between the sparse operator (i.e., the sparsity transform matrix that has been applied on the object) and the imaging modal of the system (i.e., the matrix that describes the





transform relationship between the object and the measurement). The incoherence can be described theoretically by the number of the non-zero off-diagonal elements in the autocorrelation of the sensing matrix. In the

I-COACH system, the incoherence can be controlled by using different scattering degrees σ of the CPM. To obtain the optimal scattering degree σ in the modified GS algorithm (Liu et al., 2020) corresponding to the best imaging quality

of the CS-I-COACH system, the entropy merit was used to quantitatively evaluate the quality of the reconstructed images in this paper (Men et al., 2017). Entropy is a measure of the disorder in a system, which in this case is maximized when the magnitude of the reconstructed image is distributed over the entire image plane and minimized when the entire image points are accumulated in the smallest area as possible. Therefore, for the relatively simple objects that have been used in the following experiments, a minimum entropy is expected to yield the reconstruction with the optimal SNR. From the intensity distribution function of the reconstructed image O'(m, n), the entropy can be calculated as:

$$\Re_{en} = -\sum_{m} \sum_{n} \psi(m, n) \log[\psi(m, n)]$$
(9)

where *m* and *n* are the pixel coordinates of the image and $\psi(m, n) = \frac{|Q'(m,n)|}{\sum_{m} \sum_{n} Q'(m,n)}$. The dependence of the CS

reconstruction accuracy (evaluated by the entropy of the reconstructed images, inserted blue-colored numbers) on the scattering degree σ of the CPM is shown in **Figure 5**. It was expected that higher accuracy would be obtained under higher σ . However, we found that this is not always true in our experiments and there exists an optimal value of σ that provides the reconstructed image with the relative optimal quality (indicated by the red rectangle in **Figure 5**). These can be explained by the fact that the imaging performances of the CS-I-COACH system are not disrupted only by the above-mentioned mathematical coherence, but also by physical factors such as the read-out noise of the camera and numerical factors in the reconstruction algorithm such as the coefficient of the regulator and iteration times.

We further demonstrated in our system one of the most interesting applications of CS in holography, i.e., the highfidelity image retrieval from under-sampled holograms. The recorded holograms of the PSH and OH were down-sampled randomly in the spatial frequency domain, and then used to reconstruct the object image. The dependence of the reconstruction accuracy (again evaluated by the entropy of the image, inserted blue-colored values) on the sampling rate S is shown in **Figure 6**. That the CS reconstructed images still preserve a reasonable quality until S = 10% was found in the results. Meanwhile, under the sample condition, the reconstructions obtained using conventional reconstruction methods were dominated by the strong background noise.

In the next experiment the proposed method was used for 3D imaging. The 20 μ m diameter pinhole was moved within a range of 12 mm (6 mm above and 6 mm below) around the front focal plane of L₀, with 1 mm step interval. The PSHs library at 13 different axial positions were then recorded. To simulate a 3D object, two transmission resolution targets (NBS, 1963A 1X R2L2S1N, and USAF 1951 1X, USAF GO Edmund optics), with an axial distance of 6 mm, were placed at the front focal plane of L₀, in Channel-1 and Channel-2, respectively. The line-groups of 4, 5 of the USAF targets, and the line-groups of 36 and

40 of the NBS target, were illuminated by the LED source. Initially, two targets were placed in the front focal plane of the lens L_0 . Then, the two targets in Channels 1 and 2 are moved simultaneously in reverse directions at 1 mm step, and the corresponding 7 object holograms are recorded to simulate 13 different depth planes of the 3D objects. The reconstructed images obtained using different reconstruction methods are shown in **Figure 7**. It can be seen from the results that 3D imaging performances of I-COACH can be improved by the proposed method since the defocus images attenuate more quickly with the increase of the axial space in the CS reconstructions when comparing it with PF and NLR results.

CONCLUSION

In this paper, we have successfully suppressed the background noise and improved the reconstruction quality of the conventional I-COACH methods without sacrificing the imaging speed. This was achieved by establishing the compressive sensing model and then by numerically reconstructing the holograms under the CS framework. Our experimental results have demonstrated that the PSNR of the reconstructed images can be improved by a factor of 16.5% via the proposed method. Meanwhile, CS reconstructions can still preserve reasonable quality for a 10% under-sampled hologram. In our method, the 3D imaging performance of I-COACH was improved by suppressing the effects of out-offocus images that are present in the reconstructions. With those improved system imaging performances, we believe the proposed CS-I-COACH method will benefit the research areas such as imaging through scattering medium and partial aperture imaging.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/Supplementary Material, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

YW: Conceptualization, Methodology, Supervision, Writing-Reviewing and Editing, Project administration, Funding acquisition. CL: Software, Writing-Original draft preparation, Methodology, Investigation. TM: Conceptualization, Methodology, Software, Investigation.

FUNDING

National Natural Science Foundation of China (61575009); Natural Science Foundation of Beijing Municipality (4182016).

REFERENCES

- Bioucas-Dias, J. M., and Figueiredo, M. A. T. (2007). A New TwIST: Two-step Iterative Shrinkage/Thresholding Algorithms for Image Restoration. *IEEE Trans. Image Process.* 16 (12), 2992–3004. doi:10.1109/tip.2007.909319
- Brady, D. J., Choi, K., Marks, D. L., Horisaki, R., and Lim, S. (2009). Compressive Holography. Opt. Express 17 (15), 13040–13049. doi:10.1364/oe.17.013040
- Brady, D. J., Mrozack, A., MacCabe, K., and Llull, P. (2015). Compressive Tomography. Adv. Opt. Photon. 7, 756–813. doi:10.1364/aop.7.000756
- Cochran, G. (1966). New Method of Making Fresnel Transforms with Incoherent Light*. J. Opt. Soc. Am. 56 (11), 1513-1517. doi:10.1364/josa.56.001513
- Lim, S. D., Marks, D. L., and Brady, D. J. (2011). Sampling and Processing for Compressive Holography [Invited]Sampling and Processing for Compressive Holography. Appl. Opt. 50 (34), H75–H86. doi:10.1364/ao.50.000h75
- Liu, C., Man, T. L., and Wan, Y. H. (2020). Optimized Reconstruction with Noise Suppression for Interferenceless Coded Aperture Correlation Holography. *Appl. Opt.* 59 (6), 1769–1774.
- Liu, C., Man, T. L., and Wan, Y. H. (2022). High-quality Interferenceless Coded Aperture Correlation Holography with Optimized High SNR Holograms. *Appl. Opt.* 61 (3), 661–668. doi:10.1364/ao.444783
- Lohmann, A. W. (1965). Wavefront Reconstruction for Incoherent Objects. J. Opt. Soc. Am. 55 (11), 1555–1556. doi:10.1364/josa.55.1555_1
- Men, G. F., Pascal, P., and Wang, H. Y. (2017). Research Progress of Fresnel Incoherent Correlation Digital Holography. *Imaging Sci. Photochem.* 35 (2), 199–207.
- Mertz, L., and Young, N. O. (1961). "Fresnel Transform of Images,". Editor K. J. Habell.Proceedings of the International Conference on Optical Instruments and Techniques.London: Chapman Hall, 305.
- Rai, M. R., and Rosen, J. (2019). Noise Suppression by Controlling the Sparsity of the point Spread Function in Interferenceless Coded Aperture Correlation Holography (I-COACH). Opt. Express 27 (17), 24311–24323. doi:10.1364/oe.27.024311
- Ratnam Rai, M., Vijayakumar, A., and Rosen, J. (2017). Single Camera Shot Interferenceless Coded Aperture Correlation Holography. *Opt. Lett.* 42 (19), 3992–3995. doi:10.1364/ol.42.003992

- Rogers, G. L. (1950). .Gabor Diffraction Microscopy: the Hologram as a Generalized Zone-Plate. *Nature*, 166,237. doi:10.1038/166237a0
- Stroke, G. W., and Restrick, R. C. (1965). Holography with Spatially Noncoherent Light. Appl. Phys. Lett. 7 (9), 229–231. doi:10.1063/1.1754392
- Vijayakumar, A., Kashter, Y., Kelner, R., and Rosen, J. (2016). Coded Aperture Correlation Holography-A New Type of Incoherent Digital Holograms. *Opt. Express* 24 (11), 12430–12441. doi:10.1364/oe.24.012430
- Vijayakumar, A., and Rosen, J. (2017). Interferenceless Coded Aperture Correlation Holography-A New Technique for Recording Incoherent Digital Holograms without Two-Wave Interference. *Opt. Express* 25 (12), 13883–13896. doi:10.1364/oe.25.013883
- Wan, Y., Liu, C., Ma, T., Qin, Y., and Lv, S. (2021). Incoherent Coded Aperture Correlation Holographic Imaging with Fast Adaptive and Noise-Suppressed Reconstruction. *Opt. Express* 29, 8064–8075. doi:10.1364/oe. 418918
- Zhang, H., Cao, L. C., and Jin, G. F. (2018). Applications and Challenges of Compressed Imaging. Sci. & Technology Rev. 36 (10), 20–29.

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