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Optimization of entrainment and interfacial flow patterns in countercurrent air-water two-phase flow in vertical pipes

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This study investigates countercurrent air-water two-phase flow in vertical pipes with inner diameters of 26 mm and 44 mm and a height of 2000 mm, under controlled conditions to eliminate heat and mass transfer. Cutting-edge techniques were employed to measure the liquid film thickness (δ) and entrainment (e) within the annular flow pattern. The methodology involved a systematic comparative analysis of experimental results against established models, identifying the most accurate methods for predicting flow behavior. Specifically, the Schubring et al. correlation was found to most accurately predict e in 26 mm pipes, while the Wallis correlation was more accurate for 44 mm pipes. Additionally, interfacial shear stress was analyzed, confirming the high precision of the δ and e parameters. This research enhances the understanding of countercurrent air-water two-phase flow by providing reliable estimation methods for different pipe diameters and emphasizes the significance of accurately determining interfacial shear stress. Key findings include the identification of the most accurate models for different pipe sizes and addressing challenges in measuring δ and e under controlled conditions. The study's novelty lies in its comprehensive comparative analysis of existing models, leading to improved predictions of flow dynamics in vertical pipes, thereby contributing valuable insights into two-phase flow behavior in geosciences and environmental engineering.

KEYWORDS

entrainment, optimization, flow patterns, two-phase flow, interfacial, vertical pipe, geoenvironmental science

Highlights

- The study investigates countercurrent air-water two-phase flow in vertical pipes with 26 mm and 44 mm diameters.
- Liquid film thickness and entrainment were quantified using state-of-the-art techniques under controlled experimental conditions.
- The Oliemans correlation best predicted entrainment for 26 mm pipes, while the Wallis correlation was more accurate for 44 mm pipes.

- Precise estimation of interfacial shear stress and key parameters δ (film thickness) and e (entrainment fraction) was achieved, enhancing the understanding of flow behavior.
- A comprehensive analysis was conducted to identify the most reliable methods for predicting flow behavior, significantly advancing the understanding of two-phase flow dynamics.

Introduction

Two-phase flows are frequently encountered in various types of heat transfer equipment and industries, including ventilation systems, refrigeration systems, cooling towers, countercurrent flows inside vertical pipes in nuclear reactors and coal-fired power plants, the oil and gas industry, the petrochemical sector, and oil transmission pipelines (Choubineh et al., 2017; Ghorbani et al., 2019; Sorgun et al., 2015; Tang et al., 2023). Recent studies on two-phase flows have predominantly focused on concurrent flows (Akhtar et al., 2021; Azzopardi, 1986; Taitel et al., 1980; Wang et al., 2020a; Zhao et al., 2023). Despite the complexities in analyzing flow conditions that arise from the interactions of countercurrent flows of two phases, this paper aims to address these challenges. The countercurrent air/water interaction in cooling towers serves as a prominent example in this context. In countercurrent flow regimes, turbulence increases as each phase progresses, leading to blockage of the other phase's entrance (Huang et al., 2023; Shagiakhmetov et al., 2016; Taitel et al., 1980). Predicting flow behavior in such cases is difficult due to the existence of distinct flow regimes (Ashraf et al., 2021; Azzopardi, 1986; Jayanti et al., 1990; Li et al., 2018; Liao et al., 2020; Rajabi et al., 2022; Sawant et al., 2008). Some researchers have developed semi-analytical models for water flow (Ashraf et al., 2024; Zhang et al., 2024), gas flow (Allawi et al., 2024; Cao M. et al., 2024) and CO_2 -brine flow (Thanh et al., 2024; Wang et al., 2020a; Wang et al., 2020b).

The annular flow regime is one of the most prevalent types of two-phase air-water flows within pipes, characterized by a broad range of gas and liquid flow rates (Hurlburt et al., 2006; Li et al., 2019; Liao et al., 2021; Manzoor et al., 2024; Shedd and Newell, 1998). This pattern is characterized by a thin liquid layer along the walls and a gas core in the center of the flow (Ahmed and Srivastava, 2017; Cao et al., 2021; Haider et al., 2021; Hu et al., 2024a; Wang et al., 2023). Since two-phase air/water annular flow occurs frequently in nuclear, chemical, and oil industries, recognizing and identifying this pattern is essential. Hence, detecting the dominant flow regime is the initial step in analyzing two-phase flows (Fershtman et al., 2021; Hu et al., 2024b; Shen, et al., 2024; Imran et al., 2019; Joonaki et al., 2017; Mehmooda et al., 2021). Given the inherent complexities, knowledge of two-phase flow patterns heavily depends on experimental data and relevant empirical relationships. Pipe diameter plays a significant role in determining the type of two-phase flow pattern (Ashraf et al., 2020; Cao D. et al., 2024; Jo and Yoon, 2018; Peng et al., 2021; Ren et al., 2021). Various classification criteria have been proposed for categorizing pipes by size. For example, Taitel et al. (1980) provided the following Equation 1 (Taitel et al., 1980):

$$d \leq \frac{19}{\rho_L} \sqrt{\frac{\sigma(\rho_L - \rho_g)}{g}} \quad (1)$$

Where; d represents the inner diameter of the pipe, σ denotes the liquid surface tension at the interface, ρ_L and ρ_g show the respective densities of liquid and gas phases and g is the gravitational acceleration. If the values of σ , ρ_L and ρ_g are assumed to be 0.072N/m, 997 kg/m³ and 1.184 kg/m³, then $d \leq 5.16$ cm. Ts, if the inner diameter of the pipe is less than or equal to 5 cm, the pipe is considered a small to medium-sized pipe; for larger values, the pipe is classed as large pipe. Extensive studies have been conducted on flow regimes in the pipes of different sizes.

In 1997, Ghiaasiaan et al. presented empirical relations for flow regimes and volume fractions of two-phase flow in steep and vertical channels. Their study, conducted with air, water, milk, and paraffin in a channel with a height of 2 m and an inner diameter of 1.9 cm at various angles relative to the vertical, highlighted the influence of liquid viscosity on volume fraction and flow patterns (Ghiaasiaan et al., 1997). In 1983, Taitel and Barnea proposed a model for countercurrent horizontal pipes using pressure drop and flow patterns. They identified the probable formation of two regions: one without a regime and another with possible annular, slug, and bubble flow patterns (Taitel and Barnea, 1983). In 1994, Hasan et al. utilized the drift flux model to investigate liquid volume fractions in vertical pipes with a height of 5.5 m and a diameter of 127 mm, presenting a model capable of examining slug and bubble flow patterns in countercurrent flow (Hasan et al., 1994). In 2006, Chen et al. evaluated the effect of pipe inner diameter on different regimes using R134a refrigerant in vertical pipes with diameters ranging from 1.1 to 4.26 mm, concluding that pipe diameter significantly influences flow pattern formation (Chen et al., 2006). In 2016, Besagni et al. (2017) examined the hydrodynamic properties of a bubble column in a vertical pipe with countercurrent flow, suggesting a new correlation the water volume fraction and acknowledging the influential effect of pipe diameter on liquid volume fraction (Besagni et al., 2017). Li et al. (2023) investigated the mechanical properties of hydrate-bearing sand-well interfaces. They found that interfacial shear strength, cohesion, and dilatancy increase with higher interfacial roughness and hydrate saturation. After hydrate dissolution, these properties decrease, with the decrease being more significant at higher roughness and saturation levels. They concluded that their proposed mechanical model accurately simulates the stress-displacement relationship, aiding stability evaluations of hydrate-bearing reservoirs (Li et al., 2023). Wang et al. (2024) investigated cavitation in Venturi tubes using high-speed photography and particle image velocimetry. They found that cavitation occurs at a pressure ratio of 0.595, with increased cavitation length as the ratio decreases. The flow becomes choked at 0.280. They concluded that cavitation evolves periodically, with vortices forming and fragmenting near the wall (Wang et al., 2024).

The annular flow pattern occurs at high gas velocities and low to medium liquid velocities. In this pattern, liquid flows as a film along the pipe wall and is also drawn in the form of droplets towards the gas core. The wavy interface between the gas core and the liquid film causes the liquid droplets to be sprayed and deposited. The thickness of the liquid film and the distribution of the droplets in the gas phase further complicate the prediction problem. Under equilibrium conditions, steady entrainment (e) is achieved due to the equal rates of spraying and depositing of droplets. This parameter plays a vital role in understanding and modeling the behavior of the annular flow pattern, which is affected by various factors such as pipe

size, fluid velocity, and fluid properties. Several investigations have been conducted to predict this important parameter; among these studies, the following works are noteworthy:

In 1968, Wallis: Wallis proposed a new empirical correlation by substituting dynamic liquid viscosity for gas viscosity in the Paleev and Flipovich model. The results showed improved accuracy in predicting multiphase flow behaviors, concluding that liquid viscosity plays a crucial role in flow dynamics (Wallis, 1968). In 1986, Oliemans et al. (1986): This study introduced a correlation for vertical annular flows using the Harwell Database, spanning a range of pipe diameters and surface tensions. They achieved accurate predictions for various flow regimes, concluding that the correlation effectively covers a broad range of flow types (Oliemans et al., 1986). In 1989, Ishii and Mishima (1989): Ishii and Mishima developed a correlation based on the Reynolds number of the liquid phase and Weber number of the gas phase for annular flows. Their results showed good agreement with experimental data, concluding that these numbers are critical for accurate flow predictions (Ishii and Mishima, 1989). In 2002, Pan and Hanratty (2002): Pan and Hanratty established a correlation for vertical flow in pipes, focusing on the relationship between liquid droplet velocity and gas flow. They demonstrated that their correlation is reliable across a range of flow conditions, concluding that it effectively models droplet behavior in gas (Pan and Hanratty, 2002). In 2008, Schubring et al. (2008): Schubring et al. proposed a correlation accounting for δ in two-phase flows, achieving high accuracy. They identified challenges with large-scale experimental setups and highlighted the need for research on medium-sized pipes, emphasizing the impact of slope and phase velocity (Schubring et al., 2008). This study investigates several key aspects of two-phase flow patterns, focusing on countercurrent air-water flows in vertical pipes. It examines the comparative effects of pipe diameter on the accuracy of flow pattern predictions, specifically evaluating whether smaller (26 mm) or larger (44 mm) diameters yield more reliable results. Additionally, the research assesses the performance of established empirical correlations and models, such as those by Schubring et al. and Wallis, in predicting critical parameters like steady entrainment and liquid film thickness. The impact of interfacial shear stress on the formation and stability of annular flow patterns is explored to understand its role in flow stability. Furthermore, the study analyzes how varying fluid properties, such as viscosity and surface tension, and different flow conditions influence model accuracy. Finally, it aims to identify potential improvements to existing models to enhance their predictive capabilities for two-phase flows in small-to-medium diameter pipes.

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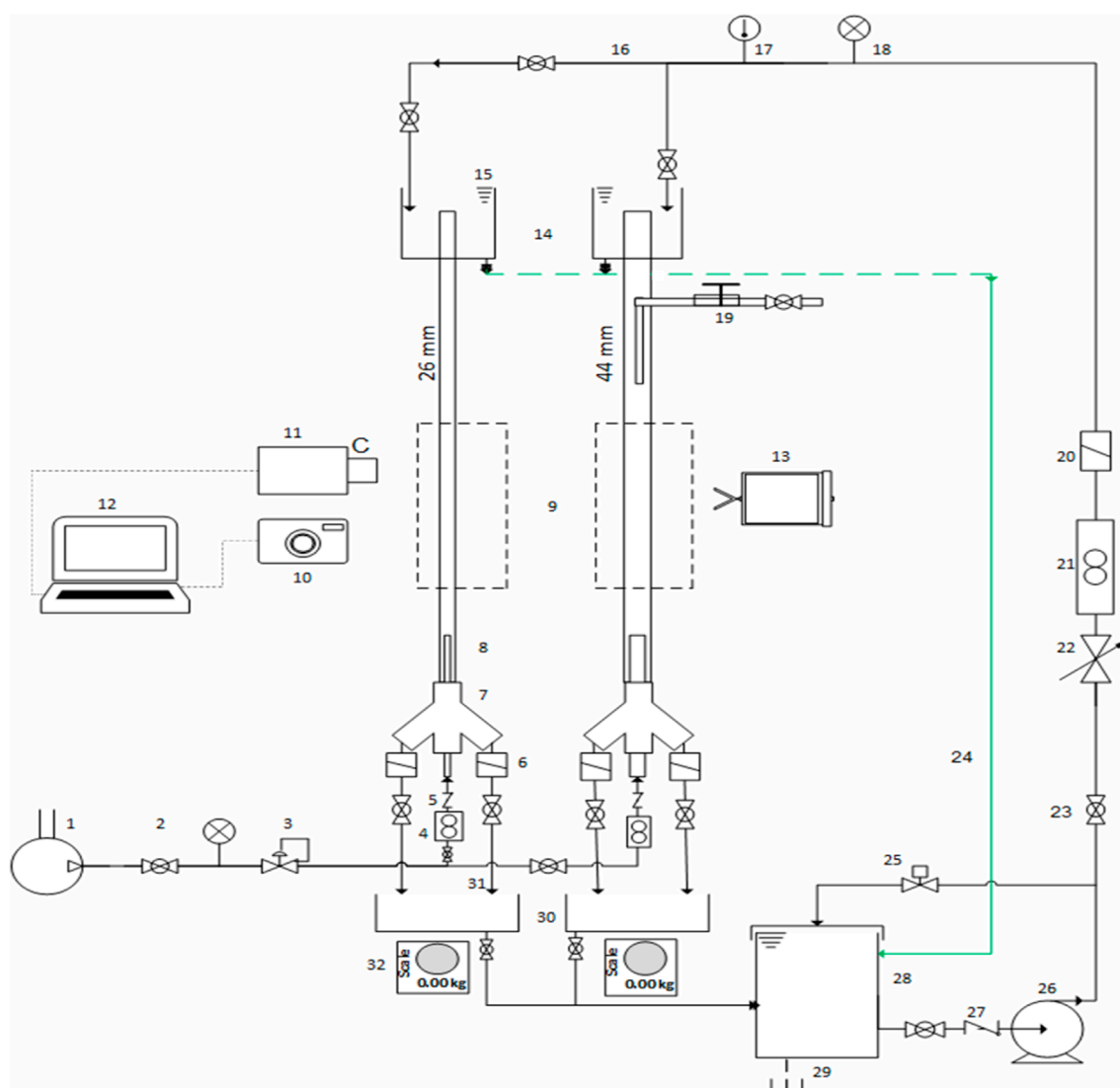
flow conditions influence model accuracy. Finally, the study aims to identify potential improvements to existing models to enhance their predictive capabilities for two-phase flows in small-to-medium diameter pipes.

This research offers a novel and significant contribution to the field of two-phase flow dynamics by providing a comprehensive and detailed analysis of countercurrent air-water flow in vertical pipes, specifically focusing on two different pipe diameters (26 mm and 44 mm). Unlike previous studies, this research meticulously investigates the entrainment (e) and liquid film thickness (δ) under controlled conditions that minimize heat and mass transfer, ensuring the accuracy of the results. The study employs advanced techniques and a comparative analysis of existing models, leading to the identification of the most reliable estimation methods for different pipe diameters, particularly highlighting the superior performance of the Schubring et al. model for the 26 mm pipe and the Wallis correlation for the 44 mm pipe. By analyzing the influence of interfacial shear stress and validating the proposed models against established data, the research addresses key challenges in accurately determining flow behavior parameters. The findings underscore the importance of removing insignificant variables in predictive models, thus refining the accuracy of two-phase flow predictions. This work significantly advances the understanding of flow dynamics in vertical pipes by providing a detailed evaluation of the predictive capabilities of existing models, thereby improving the reliability of flow behavior predictions and contributing to the development of more accurate and efficient engineering applications in geoenvironmental sciences.

Methodology

Experimental vertical setups

The schematics of the experimental setup, consisting of two pipes with heights of 2000 mm and inner diameters of 26 mm and 44 mm, are shown in Figure 1. For more convenient observation and analysis of two-phase flow regimes, transparent Plexiglas pipes were used. Moreover, while the scale of the experimental setup is representative of certain industrial scenarios, it may not fully reflect the complexities encountered in larger systems. In real-world applications, factors such as scale effects, material heterogeneity, and environmental conditions like temperature and pressure variations play significant roles in influencing flow dynamics. The smaller dimensions of the experimental pipes might not exhibit phenomena such as turbulence, phase separation, or pressure drops to the same extent as in full-scale systems. Therefore, while the controlled environment and specific material choices in the experiment provide clear insights into fundamental flow behaviors, further studies using materials and conditions more representative of actual geoenvironmental systems are necessary to fully validate and generalize these findings. Such studies would ensure that the models developed from these experiments are robust and reliable when applied to the design, operation, and optimization of industrial pipelines and energy extraction systems. Additionally, air and water were used as the working fluids. Under standard conditions, the density, dynamic viscosity, and surface tension of water were 998 kg/m³, 0.001 Pa s, and 0.073 N/m, respectively. A 7.5 kW compressor was used to



1. Air compressor	9. Test section	17. Thermo meter	25. By pass valve
2. Ball valve Gas	10. Digital camera	18. Pressure Gauge	26. Pump
3. Regulator Gas	11. Video camera	19. Iso kinetic probe	27. Check valve water
4. Flow meter (Gas)	12. Laptop	20. Electrical valve	28. Storage water Tank
5. Check valve Gas	13. LED Projector	21. Flow meter (water)	29. Drain
6. Electrical valve water	14. Water Open Vessel	22. Gate valve	30. Water open vessel
7. Three way junction	15. Water surface	23. Ball valve water	31. water outlet Line
8. Gas inlet	16. Water Line	24. Water Return line	32. Digital Scale

FIGURE 1
Schematic of the experimental setup for flow in vertical Plexiglas pipes.

supply the required compressed air flow rate. In the experiments, water and air were chosen as the liquid and gas phases.

According to Figure 1, an 80-L main tank, equipped with a pump, was employed for pumping purposes. The separated water droplets from the two pipes were collected in a small tank for weighing and calculating efficiency. Excess air and liquid droplets

were vented to the atmosphere at the outlet section of the pipe to reduce the back pressure of the system and increase the air flow rate. Air and water flow rates were measured through an air rotameter (model ACA 05-25 ZT) and a water rotameter (model Z-5008), both with accuracies of ± 0.1 LPM, respectively, at a maximum pressure of 10 bars. In accordance with the ANSI/ASME standard

(10), the maximum measurement uncertainties for the water and air flowmeters were $\pm 4\%$. Under these conditions, the measurement values remained accurate under any environmental conditions. The temperature measurements were taken using thermometers with a precision of $\pm 0.01^\circ\text{C}$ and an uncertainty of 3% for air/water flows before they entered the pipe containing the two-phase flow. For pressure measurements, barometers with an accuracy of ± 3 mbar and a maximum uncertainty of 1.38% were used.

The experiments consisted of two main components: 1) the rapid closure solenoid section, and 2) the isokinetic sampling section. Due to the short length of the pipe, the isokinetic sampling was not affected by the inlet, and the annular flow was fully developed in the measuring sections. To determine the structure and flow pattern of the two-phase water-air mixture at different velocities, images were captured by a Nikon D72 high-speed camera with a video recording speed of 240 f/s and a Casio camera (EX-ZR1200) with a video recording speed of 240 f/s and a resolution of 1.16 mp. The cameras were fixed at a specified height from the bottom of the pipe. The distance between the camera and pipes varied between 15–20 cm, depending on the flow conditions inside the pipe, to cover 15 cm of the pipe length. Additionally, a white light was shone on opaque white plates placed at a suitable distance behind the test pipe to diffuse the light over a wider area inside the pipe. As a result, the camera lens was focused on the flow structure, enhancing the quality of the captured images. In this setup, a one-way valve and a solenoid valve were embedded before the entrance of gas to the test pipe to largely prevent the penetration and return of water into the air path. Air and water entered the test pipe through two inlets with inner diameters of 20 mm.

Determination of flow pattern

Determining the flow pattern for countercurrent two-phase flow in vertical pipes with diameters of 26 mm and 44 mm, and evaluating the effect of pipe diameter on the flow pattern boundaries, are additional objectives of this study. The simplest procedure to develop a flow pattern map is to characterize the flow regimes by the superficial velocities of gas (u_{sg}) and liquid (u_{sl}) and to establish a relationship between them.

Determination of flow pattern for 26 mm pipe

For this pipe, the ratio of length to inner diameter (L/D) was 76.92, and the experiments were conducted for 64 cases corresponding to the desired superficial velocities of water and air, in the ranges of 0.03–0.09 m/s and 1.3–5.2 m/s, respectively. No two-phase flow pattern appeared in areas with low superficial air velocity due to the air's inability to enter, as blocked by the water column. With the gradual increase in air superficial velocity, slug, turbulent, annular, and eruptive patterns were observed. Initially, the slug regime was noted, where water containing numerous small bubbles trailing the Taylor bubble was detected. The Taylor bubble formed an umbrella-shaped dome, as shown in Figure 2A. Then, with the acceleration of the gas phase, the Taylor bubbles began to break down, and an unstable, transitional flow developed between

the slug and annular flows. In this transition regime, the water was drawn upward due to air momentum, but then pulled downward again due to its higher density relative to air. The repetition of this process led to a continuous oscillating movement upward and downward (Figure 2B).

With a further increase in gas velocity, a combination of turbulent and annular regimes emerged, followed by a highly turbulent and oscillating flow for a period of time. After some amount of water exited the pipe, the flow pattern transitioned into an annular flow. This transition recurred as water accumulated again (with a U_{sg} fluctuation magnitude of 5.2 m/s). The annular flow then developed in the middle of the pipe, with a layer of water around the pipe, forming the shape of an empty cylinder. The annular flow was defined as a gas core surrounded by a continuous liquid layer (Figure 2C). Eventually, in the return flow, the air velocity surpassed the water velocity, causing the air to prevent the water from flowing, which resulted in the upward movement of the water flow. As a result, the countercurrent flow transitioned into a concurrent upward flow. The observed flow patterns in the 26 mm pipe are depicted in Figure 3. In this figure, obtained from Design Expert software, codes 1, 2, and 3 represent slug, annular, and flooding flow (Figure 2D), respectively.

Flooding, which involves the upward transfer of annular waves, occurs in 26 mm pipes. This mechanism is observed in pipes with diameters of 26 mm and smaller. At low gas flow rates, the liquid film flow remains undisturbed. Shortly after the separation point, the waves return to the liquid and do not reach the upper part of the test section. The gas stream carries the waves upward for a short distance. As the gas flow increases further, the amplitude of the circular waves becomes so large that the available cross-sectional area of the gas stream is restricted. At this point, the subsequent waves are carried upward beyond the liquid inlet by the gas and may even occupy the entire tube. By maintaining a constant gas and liquid flow rate, waves are periodically formed near the liquid outlet and carried upwards by the gas from the liquid inlet. Flooding occurs at this stage. Large amounts of bubbles are present in the waves moving upwards. When flooding happens, the pressure gradient increases sharply due to the small diameter. According to this figure, the annular flow pattern was the dominant one observed in the experiments.

Determination of flow pattern in 44 mm pipe

For this pipe, the ratio of length to inner diameter (L/D) was 45.45. The experiments were conducted for 64 cases, with the superficial velocities of water and air ranging from 0.01037 to 0.2075 m/s and 3.6555–9.138 m/s, respectively. The main flow patterns observed included slug-churn, annular, and reverse flow patterns. Compared to the 26 mm pipe, the length of the slug regime was slightly shorter, and the observed Taylor bubbles were relatively slower and smaller. This phenomenon was attributed to the increase in the water phase volume and the corresponding pressure (Figure 4A). Subsequently, a churn regime developed between the slug and annular regimes, covering a larger area than in the 26 mm pipe. On one hand, the faster breakdown of Taylor bubbles caused part of the slug regime in the smaller pipe to transition into churn flow. On the other hand, the increased volume of water inside the

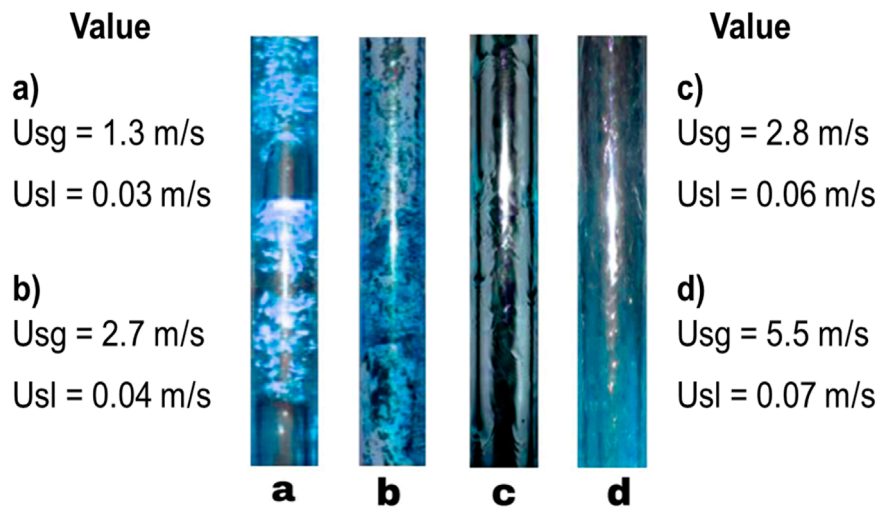


FIGURE 2 Observed flow patterns for the pipe with an inner diameter of 26 mm (A) slug, (B) churn, (C) annular, (D) flooding.

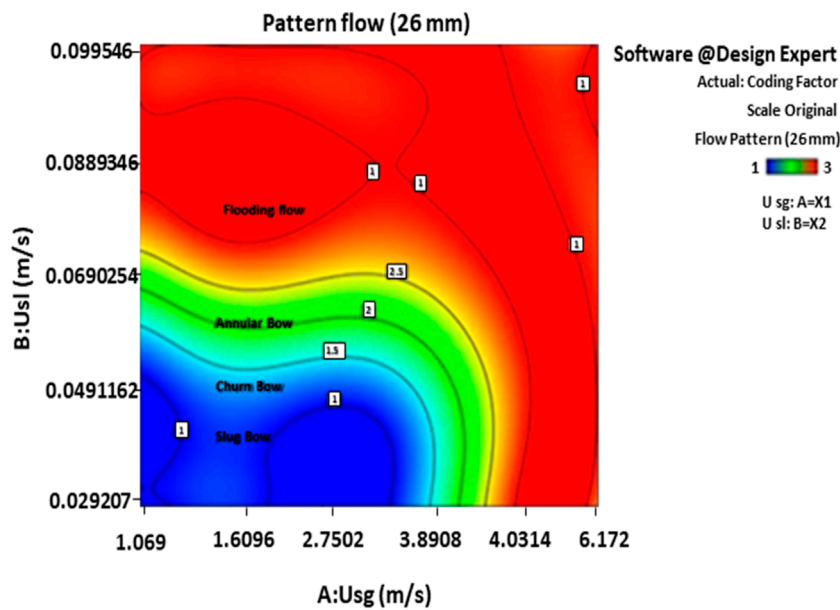
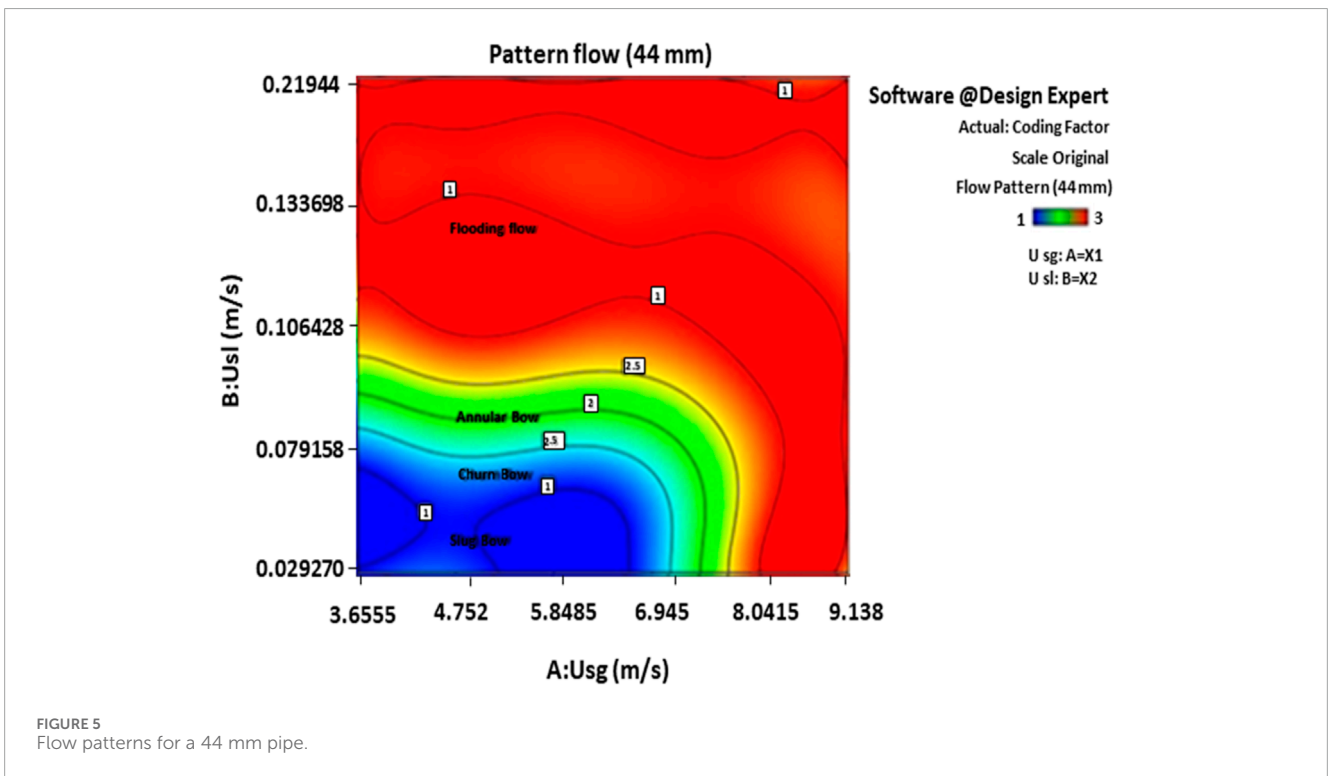
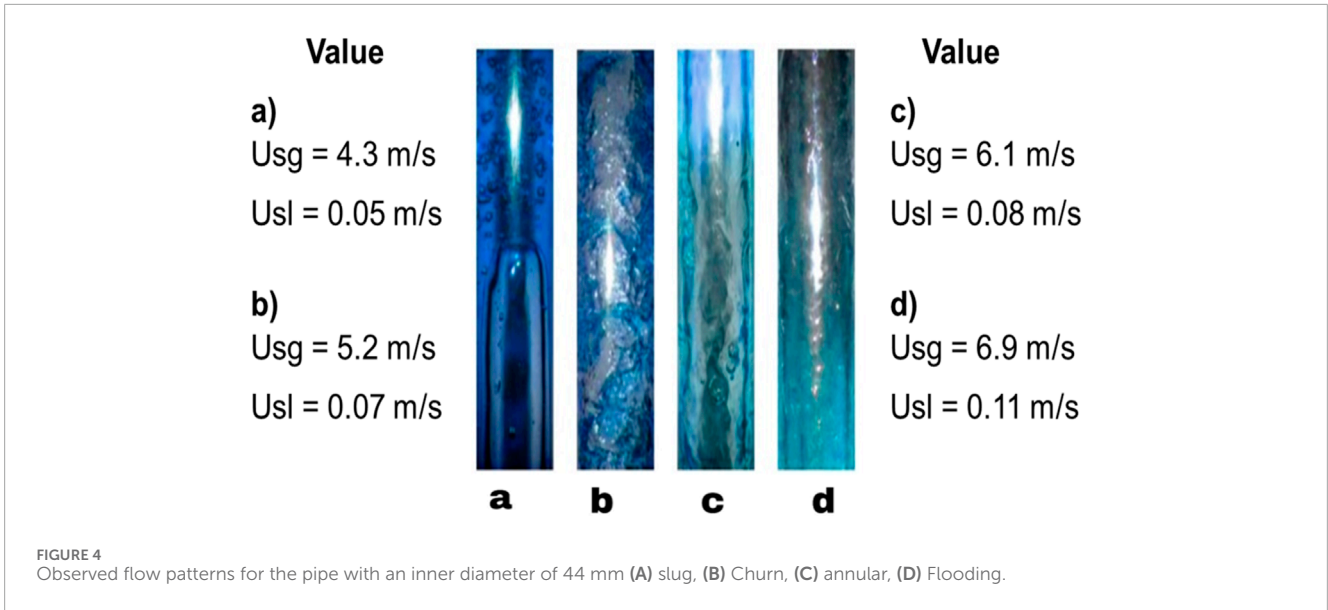


FIGURE 3 Flow patterns for a 26 mm pipe.

pipe reduced the air’s ability to penetrate the pipe’s center, leading to the annular regime transitioning into the churn regime (Figure 4B). Following this, a combination of churn and annular regimes was observed (Figure 4C). The flow then transformed into the annular regime due to the inability of the gas phase to displace the water, leading to the formation of a flooding flow, unlike in the 26 mm pipe. Consequently, part of the corresponding return flow in the smaller pipe transitioned into the annular regime. As previously mentioned, the extended turbulent regime delayed the onset of the annular regime (Figure 4C). Finally, the failure of the reverse flow

to fully expel the water from the pipe resulted in a balance between the interaction of gas momentum and gravitational force at a cross-section of the pipe (Figure 4D). The observed flow patterns in the 44 mm pipe are illustrated in Figure 5. In this figure, generated by Design Expert software, Code 1, Code 2, and Code three represent the slug, annular, and flooding flows, respectively. Flooding, as characterized by “e,” is observed in the 44 mm pipe. This mechanism occurs in pipes with a diameter of 44 mm and larger under low liquid flow conditions. When the gas flow rate increases, small droplets form near the liquid outlet, but these droplets return to the



liquid film in the same area. As the gas flow rate increases further, the number of droplets increases, and in the upper parts of the cross-section, droplets become visible above the liquid inlet to the gas outlet. They do not move upwards, but in some cases, waves are seen only in the lower half of the pipe. However, this occurrence of flooding through this mechanism is inaccurate. The dominance of the annular flow in the experiment can be inferred from this figure.

Shear stress theory in the interface of two-phase flow

In annular flow, unevenness occurs at the liquid surface, causing the annular air flow to be positioned in the middle of the pipe, surrounded by liquid. This is due to the slower velocity of the water compared to the gas, which results in some droplets separating from the liquid and moving with the gas. These separated droplets are

referred to as “ e ” and symbolized by ‘ e ,’ which ranges from $0 \leq e \leq 1$. The upper and lower limits are defined as:

If $e = 0$: no droplet was separated from the liquid.

If $e = 1$: all droplets were separated from the liquid.

This phenomenon increases the density of gas (ρ_{mix}) relative to the pure gas density. Consider a control volume with a length of Δx , through which \dot{m}_g flows and where some liquid droplets separate and enter the gas flow. As a result, the mass flow rate of this element becomes \dot{m}_{mix} (Equations 2, 3).

$$e = \frac{\dot{m}_e}{\dot{m}_l} = \frac{\dot{m}_l - \dot{m}_{fl}}{\dot{m}_l} \rightarrow \dot{m}_{fl} = \dot{m}_l - e\dot{m}_l = \dot{m}_l(1 - e) \quad (2)$$

$$\dot{m}_{\text{mix}} = \dot{m}_g + e\dot{m}_l \quad (3)$$

Where e represents entrainment, and m represents mass flow. It is worth noting that the isokinetic sampling method is used for the experimental calculation of e . Therefore, the two-phase shear stress (Equation 4) and the friction coefficient (Equation 5) of the interface can be calculated using the following relations (Equations 6–10):

$$\tau_i = \left[\left(-\frac{dp}{dz} \right)_{TP} + (\rho_{\text{mix}} \times g) \right] \times \left(\frac{D - 2\delta}{4} \right), \text{range} \left(\frac{dp}{dz} \right) = 0.26 - 0.76 \quad (4)$$

$$f_i = \frac{2 \times \tau_i}{\rho_g \times \left(u_{sg(\frac{m}{s})} - u_{sl(\frac{m}{s})} \right)^2} \quad (5)$$

$$\dot{m}_{fl} = \rho_l \times u_{fl} \times A_{fl} \quad (6)$$

$$u_{fl} = \frac{\dot{m}_{fl}}{\rho_l \times A_{fl}} \quad (7)$$

$$A_{fl} = \pi \times (D\delta(\text{mm}) - \delta^2(\text{mm}^2)) \quad (8)$$

$$u_{fl} = \frac{\dot{m}_{fl}}{\rho_l \times \pi \times (D\delta(\text{mm}) - \delta^2(\text{mm}^2))} = \frac{\dot{m}_l(1 - e)}{\rho_l \times \pi \times (\delta(\text{mm}) - \delta^2(\text{mm}^2))} \quad (9)$$

$$\rho_{\text{mix}} = \rho_g \left[1 + e \frac{\dot{m}_l}{\dot{m}_g} \right] \quad (10)$$

Where τ_i represents the interfacial shear stress, influenced by pressure gradient ($-dp/dz$), mixture density (ρ_{mix}), gravitational acceleration (g), pipe diameter (D), and liquid film thickness (δ). The interfacial friction factor (f_i) is calculated using interfacial shear stress (τ_i), gas density (ρ_g), superficial gas velocity ($u_{sg(m/s)}$), and superficial liquid velocity ($u_{sl(m/s)}$). The liquid film mass flow rate (\dot{m}_{fl}) is determined by liquid density (ρ_l), liquid film velocity (u_{fl}), and liquid film cross-sectional area (A_{fl}). Equations for u_{fl} involve \dot{m}_{fl} , ρ_l , and A_{fl} , while the liquid film cross-sectional area (A_{fl}) is calculated using pipe diameter (D), liquid film thickness (δ), and its square (δ^2). Finally, the mixture density (ρ_{mix}) is determined by gas density (ρ_g), entrained liquid fraction (e), liquid mass flow rate (\dot{m}_l), and gas mass flow rate (\dot{m}_g).

Results and discussion

Statistical evaluation of results obtained from software

The statistical models were evaluated by comparing the correlation coefficient, mean squared error (MSE), F-value, p -value,

and ANOVA model, along with visual and graphical observations of the independent variables. The values of A, B, and AB shown in Tables 1–3 represent the superficial velocity of the gas, superficial velocity of the liquid, and the product of these two superficial velocities. Table 1 presents the results of δ and e values for the 26 mm pipe. Based on the provided p -value and F-value, the significant effect of A, B, AB, and A^2 input variables on determining the thickness of the liquid film in the 26 mm pipe can be inferred. However, the B^2 term (U_{sl}^2) was found to be insignificant and should be eliminated. Furthermore, A, B, AB, and B^2 were significant variables in determining e for a 26 mm pipe, but the term A^2 (U_{sg}^2) was insignificant and should be removed from the calculations. A more accurate estimation of the output variable can be achieved for the 26 mm pipe by eliminating insignificant variables.

The ANOVA analysis for the 44 mm pipe is also provided in Table 2. As can be inferred from this table, A, B, AB, and A^2 variables were significant for determining the thickness of the liquid film in the 44 mm pipe, whereas B^2 was insignificant. The significant variables for the determination of e in this pipe were A, B, AB, and B^2 , while A^2 was considered insignificant.

For the parameters δ and e , the optimized values were defined by a quadratic model. The correlation coefficients for δ and e of the 26-mm pipe were 0.9966 and 0.9806, respectively. These values were 0.9957 and 0.9866 for the 44-mm pipe. Accordingly, the accuracy of the proposed models for the determination of δ and e in 26-mm and 44-mm pipes was established. The diagrams of δ versus superficial gas velocity and estimated values against real measurements for the 26-mm and 44-mm pipes are depicted in Figures 6, 8, respectively. When comparing the diagrams, the higher performance accuracy of the 26-mm pipe diagrams compared to the 44-mm pipe diagrams was noted. The corresponding diagrams for the e variable in the 26-mm and 44-mm pipes are illustrated in Figures 6A–D, respectively. By comparing the diagrams of these figures along with the experimental and engineering perspectives, the adequacy of the obtained design of experiments (DOE) correlation model for predicting δ and e was concluded.

For the purpose of comparing and evaluating the experimental results with empirical relations from other studies, various statistical indices were employed, including percentage error (PE), mean squared error (MSE), mean relative error (MRE), root mean squared error (RMSE), mean absolute relative error (MARE), coefficient of determination (R^2), and standard deviation (STD) (Abad et al., 2022; Pan and Hanratty, 2002). The comparison of the obtained experimental data for determining the δ in 26 mm and 44 mm pipes with the results of Schubring et al. (2008) (Schubring et al., 2008) and Rahman et al. (2017) (Rahman et al., 2017) through the statistical indices is reported in Table 3. According to this table, good agreement was found in these references for the liquid film estimation in the 26 mm pipe; however, the accuracy of the Schubring et al. model was slightly higher. Conversely, the results for the 44 mm pipe revealed that the Schubring et al. results were more accurate.

The comparison between the calculated δ and the measured values is shown in Figure 7. As indicated in Table 3; Figure 7, the results from Schubring et al. demonstrate consistency with the experimental data and exhibit higher accuracy for both 26 mm and 44 mm pipes. In contrast, the model by Rahman et al. (2017) demonstrated a significant loss in precision when applied to 44 mm

TABLE 1 Analysis of ANOVA model to define influential input variables in determining the δ and e for 26 mm pipe.

Type of output	Source	Sum of squares	df	Mean square error	F-value	p-value
Liquid Film Thickness	Model	0.3016	5	0.0603	1100.78	<0.0001
	A-Usg	0.1760	1	0.1760	3211.46	<0.0001
	B-Usl	0.0781	1	0.0781	1425.49	<0.0001
	AB	0.0090	1	0.0090	165.10	<0.0001
	A ²	0.0120	1	0.0120	218.36	<0.0001
	B ²	4.256E-06	1	4.256E-06	0.0776	0.7835
	Residual	0.0010	19	0.0001		
	Cor Total	0.3027	24			
entrainment	Model	0.0630	4	0.0158	252.39	<0.0001
	A-Usg	0.0154	1	0.0154	246.85	<0.0001
	B-Usl	0.0454	1	0.0454	727.42	<0.0001
	AB	0.0010	1	0.0010	15.94	0.0007
	A ²	0.0000	1	0.0000	0.0000	1.0000
	B ²	0.0024	1	0.0024	38.11	<0.0001
	Residual	0.0012	20			
	Cor Total	0.0643	24			

pipes. Consequently, it can be concluded that the Rahman et al. model is inadequate for predicting the behavior of pipes with diameters greater than 26 mm.

For the purpose of comparing and evaluating the prediction in 26 mm and 44 mm pipes with previous studies, the results presented in Table 4 were utilized. As observed, the highest performance accuracy for determining the in the 26 mm pipe was associated with the Oliemans et al. (1986) model, while the lowest performance accuracy was attributed to the Ishii and Mishima (1989) model. Additionally, for the 44 mm pipe, the highest and lowest performance accuracy were linked to the Wallis (1968) and Pan and Hanratty (2002) models, respectively.

Models integrate experimental data by using the data to calibrate and validate their predictions, ensuring that the theoretical assumptions and mathematical formulations accurately reflect real-world behaviors. Certain models perform better under specific conditions because they are tailored to account for the particular characteristics of those conditions, such as fluid properties, flow regimes, or pipe dimensions. For instance, a model designed to capture the nuances of liquid film thickness in small-diameter pipes may incorporate specific terms or correction factors that enhance its accuracy for that scenario, whereas it might be less precise for larger pipes where different flow dynamics dominate. The cross-plot diagrams of calculated values of e versus measured ones for

26 mm and 44 mm pipes are represented in Figures 8, 9, respectively, for comparison with the previously proposed correlations. The coefficient of determination for the models of other researchers is provided in these diagrams. Based on these diagrams along with the results in Table 4, the order of performance accuracies for the previously proposed models is as follows:

For the 26 mm pipe: Schubring et al. (2008); Oliemans et al. (1986); Rahman et al. (2017); Pan and Hanratty (2002); Wallis (1968); Ishii and Mishima (1989).

For the 44 mm pipe: Schubring et al. (2008); Wallis (1968); Rahman et al. (2017); Ishii and Mishima (1989); Oliemans et al. (1986); Pan and Hanratty (2002).

Consequently, the inadequacy of the Ishii and Mishima (1989) model to determine the e for 26 mm pipes was inferred. The best model for this pipe size was the Oliemans et al. (1986) model. Moreover, the best model for the 44 mm pipe was the Wallis (1968) model, while the Pan and Hanratty (2002) model was unsuitable for determining the e in this size.

Validation of interfacial shear stress for experimental data

To validate the interfacial shear stress proposed in Equation 5, twenty experimental data points were utilized. Figure 10 illustrate

TABLE 2 Analysis of ANOVA model to define influential input parameters in determining the thickness of liquid film and e for 44 mm pipe.

Type of output	Source	Sum of squares	df	Mean square error	F-value	p-value
Liquid Film Thickness	Model	0.3619	5	0.0724	885.91	<0.0001
	A-Usg	0.1968	1	0.1968	2409.20	<0.0001
	B-Usl	0.1061	1	0.1061	1298.48	<0.0001
	AB	0.0141	1	0.0141	172.82	<0.0001
	A ²	0.0144	1	0.0144	175.88	<0.0001
	B ²	4.388E-06	1	4.388E-06	0.0537	0.8192
	Residual	0.0016	19	0.0001		
	Cor Total	0.3635	24			
entrainment	Model	0.7039	5	0.1408	278.75	<0.0001
	A-Usg	0.1045	1	0.1045	206.94	<0.0001
	B-Usl	0.5975	1	0.5975	1182.99	<0.0001
	AB	0.0104	1	0.0104	220.59	<0.0001
	A ²	0.0000	1	0.0000	0.0000	1.0000
	B ²	0.0182	1	0.0182	116.28	<0.0001
	Residual	0.0096	20	0.0005		
	Cor Total	0.7135	24			

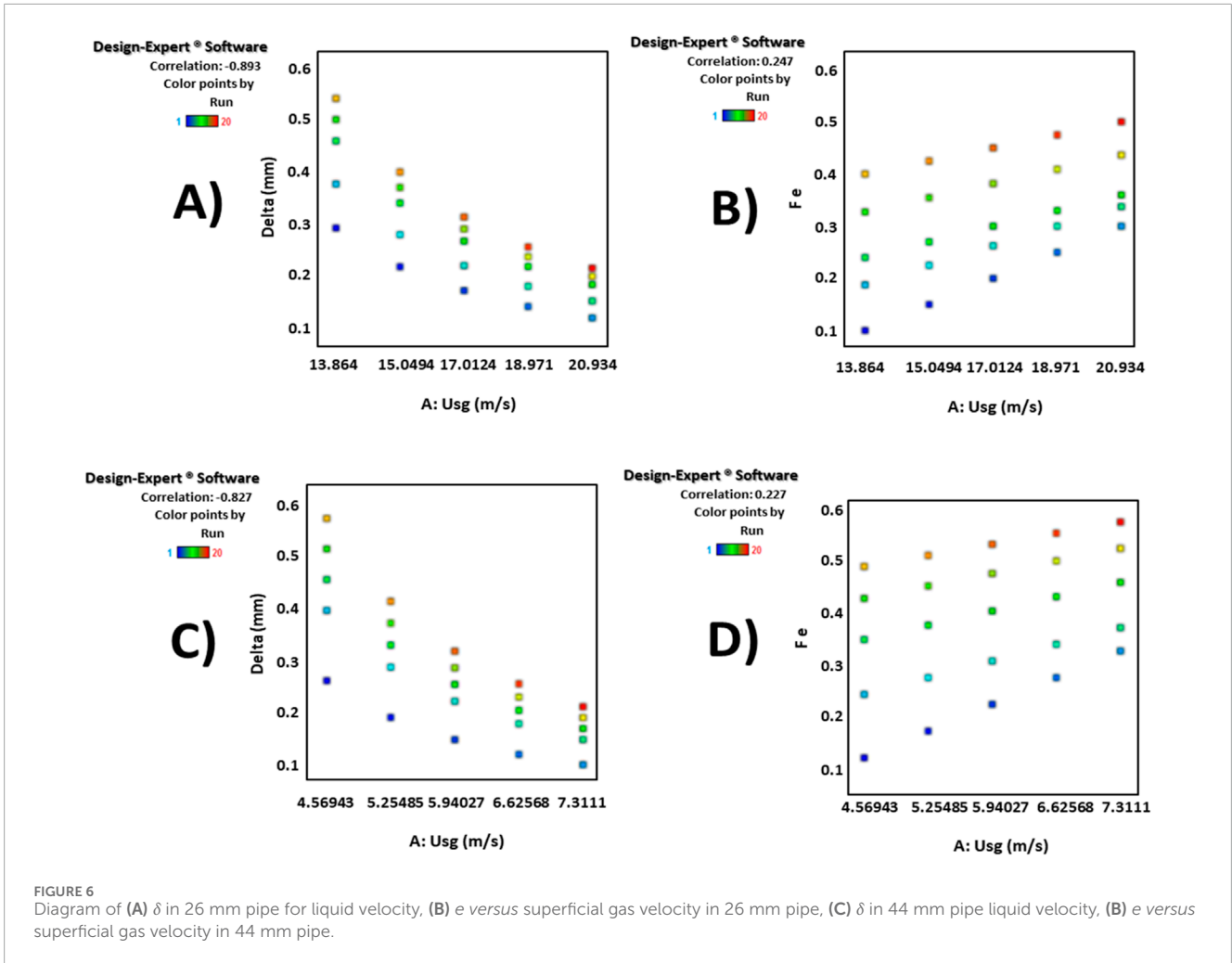
TABLE 3 Statistical results of prediction accuracy for δ in two-phase flow of vertical pipe.

Pipe size	Models	ARE	AARE	STD	MSE	RMSE	R2
mm	-	(%)	(%)	(mm)	(mm)	(mm)	-
26 mm	Schubring et al. (2008)	10.342	10.342	0.006	4.693E-04	0.02166	0.9957
	Rahman et al. (2017)	-13.380	15.064	0.022	1.210E-03	0.03479	0.9949
44 mm	Schubring et al. (2008)	-7.132	9.528	0.025	7.996E-04	0.02828	0.9608
	Rahman et al. (2017)	30.185	30.185	0.052	1.080E-02	0.10391	0.8192

the interfacial shear stress against superficial gas velocity for the 26 mm pipe (left side) and the 44 mm pipe (right side). Upon careful examination of these figures, lower values of interfacial shear stress were recorded for lower liquid superficial velocities, which increased with the enhancement of the liquid superficial velocity. This finding supports the experimental results of the conducted tests. As depicted in these figures, the highest interfacial shear stress value for the 26 mm pipe was associated with a liquid superficial velocity of 0.0747 m/s, while it was equal to 0.0498 m/s for the 44 mm pipe.

Measure the liquid film thickness

In this article, the height of the liquid surface is measured using graduated rulers. The rulers are placed in several positions in the test section of the pipe and are arranged along a vertical path across the entire width of the pipe section. The rulers are transparent, ensuring that there is no visual disturbance during the filming and imaging of the two-phase flow. Accurate calculation of liquid height is one of the most important parameters in measuring interface shear stress. Research indicates that when the apparent velocity of the gas is high,



the value of hL decreases; conversely, when the apparent velocity of the gas is low, the value of hL increases.

Uncertainty, limitations and its calculation

All major factors affecting measurement accuracy should be considered, along with correction steps to minimize measurement uncertainty. When the accuracy of the measuring equipment is known for each parameter (e.g., pipe diameter, pressure difference, temperature), the uncertainty of measuring the flow rate associated with each parameter can be predicted. This prediction method is based on the theoretical relationship of each parameter with fluid flow. Numerous standards are available for evaluating and estimating the contribution of measurement uncertainty for each measurement parameter (Ansi/Asme, 1986).

Uncertainty means the possible estimate of the error. Since we do not always have the exact amount of the error, we report the error with uncertainty. The reasons for this include: variations due to repeated measurements, undetectable fluctuations in environmental conditions such as temperature, misalignment of connections, wear

in components, reading errors, and small changes in the position of the measuring device.

The main purpose of a measurement is to determine the value of a specific quantity. Due to the errors created during the measurement process, the exact value cannot be determined, and only an estimate can be obtained. This estimate is complete when accompanied by a statement of uncertainty. The apparent velocities of water and air phases are obtained from Equation 11 (Holman, 2012).

$$u_{si} = \frac{4 \times Q}{\pi d^2} \tag{11}$$

Where, Q is the flow rate of the phases, u_{sl} is the apparent velocity of the air and water, and d is the inner diameter of the test pipe. The uncertainty of the apparent velocities of the air and water in all three tubes and at all test points is obtained using Equation (12) (Holman, 2012).

$$\omega_T = \left[\left(\frac{\partial I}{\partial Q} \omega_Q \right)^2 + \left(\frac{\partial I}{\partial d} \omega_d \right)^2 \right] \tag{12}$$

Where, ω_T the uncertainty of the measured value is stated, ω_Q and ω_d along with the uncertainty of the flow meters and pipe diameters, whose values are reported in Table 5.

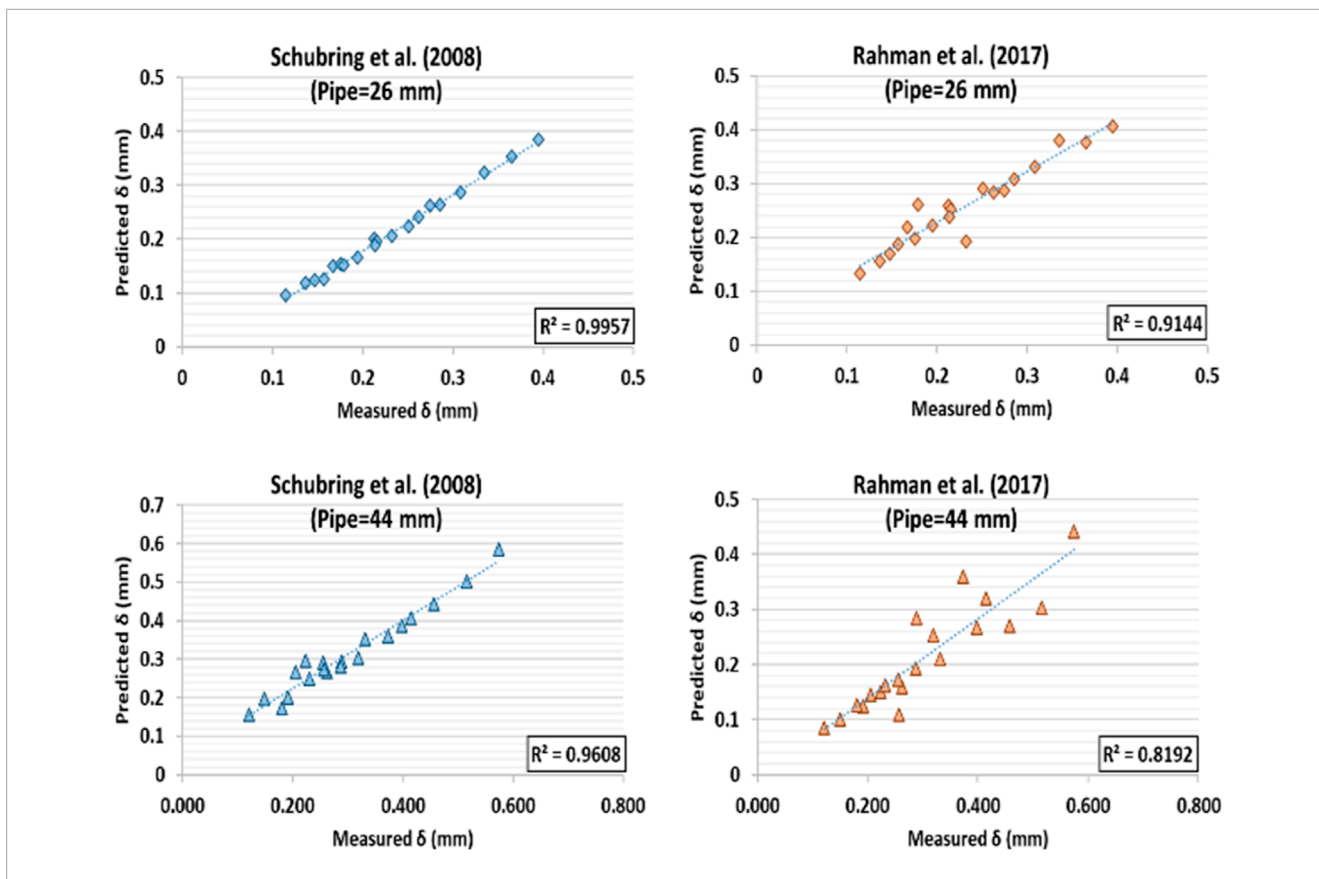


FIGURE 7 Cross plot diagram of calculated δ against the measured values (26 mm and 44 mm pipes).

TABLE 4 Statistical results of prediction accuracy for e in two-phase flow of vertical pipe.

Pipe size	Models	ARE	AARE	STD	MSE	RMSE	R2
mm	-	(%)	(%)	(mm)	(mm)	(mm)	-
26 mm	Wallis (1968)	-8.266	17.518	0.042	1.848E-03	0.0438	0.4045
	Oliemans et al. (1986)	-2.073	4.960	0.018	3.595E-04	0.0190	0.8603
	Ishii and Mishima (1989)	1.857	12.265	0.070	5.591E-03	0.0740	0.3872
	Pan and Hanratty (2002)	-0.152	8.991	0.033	1.072E-03	0.0327	0.6240
	Schubring et al. (2008)	8.020	8.020	0.006	4.693E-04	0.02166	0.9957
	Rahman et al. (2017)	-10.12	11.805	0.022	1.210E-03	0.03479	0.9144
44 mm	Wallis (1968)	-3.114	7.606	0.056	1.167E-03	0.0342	0.9373
	Oliemans et al. (1986)	-3.114	7.606	0.032	5.111E-03	0.0715	0.7890
	Ishii and Mishima (1989)	6.886	9.665	0.049	3.709E-03	0.0609	0.8348
	Pan and Hanratty (2002)	-3.201	16.387	0.072	5.163E-03	0.0719	0.7123
	Schubring et al. (2008)	-7.132	9.528	0.025	7.996E-04	0.02828	0.9608
	Rahman et al. (2017)	30.185	30.185	0.052	1.080E-02	0.10391	0.8192

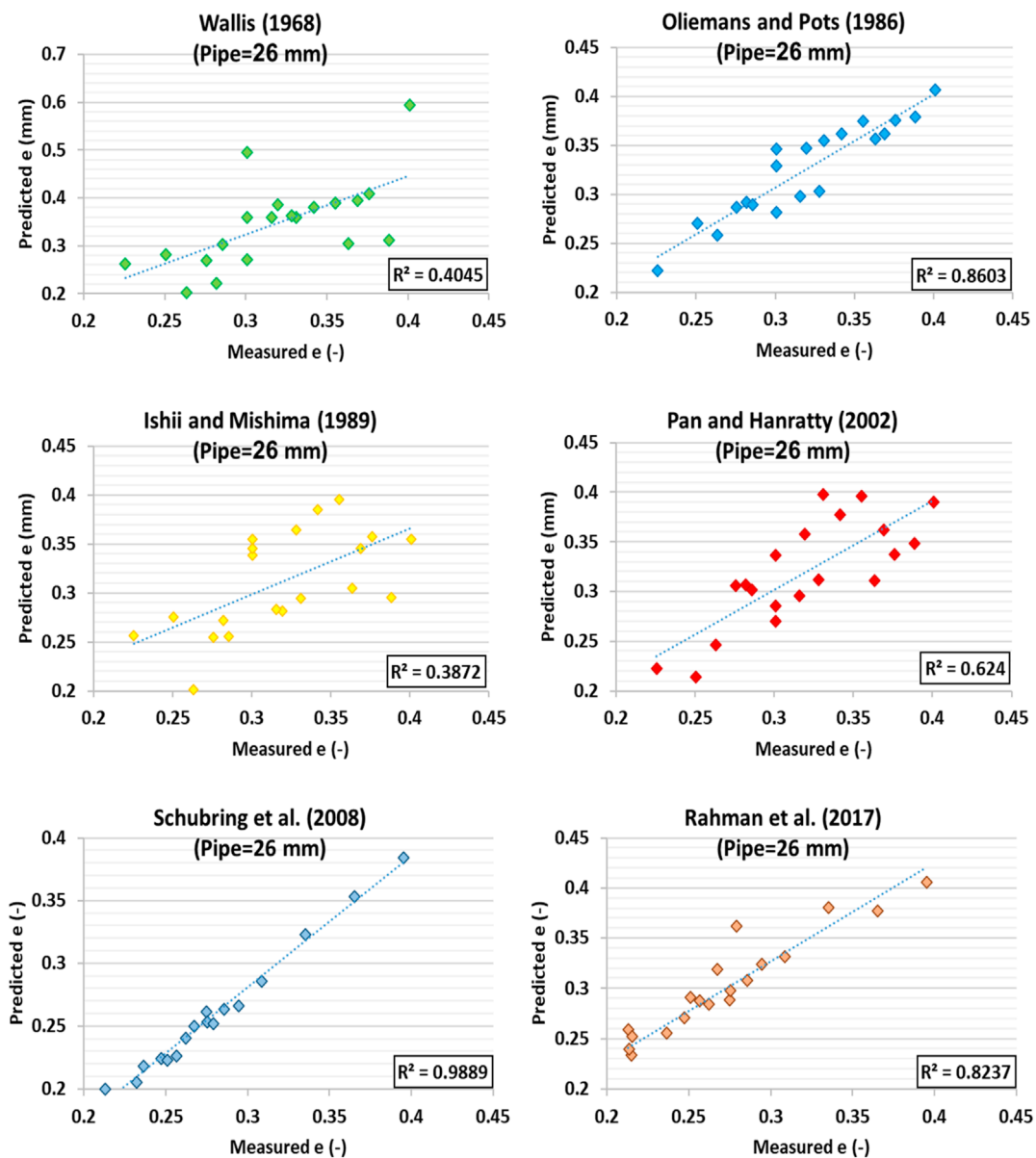
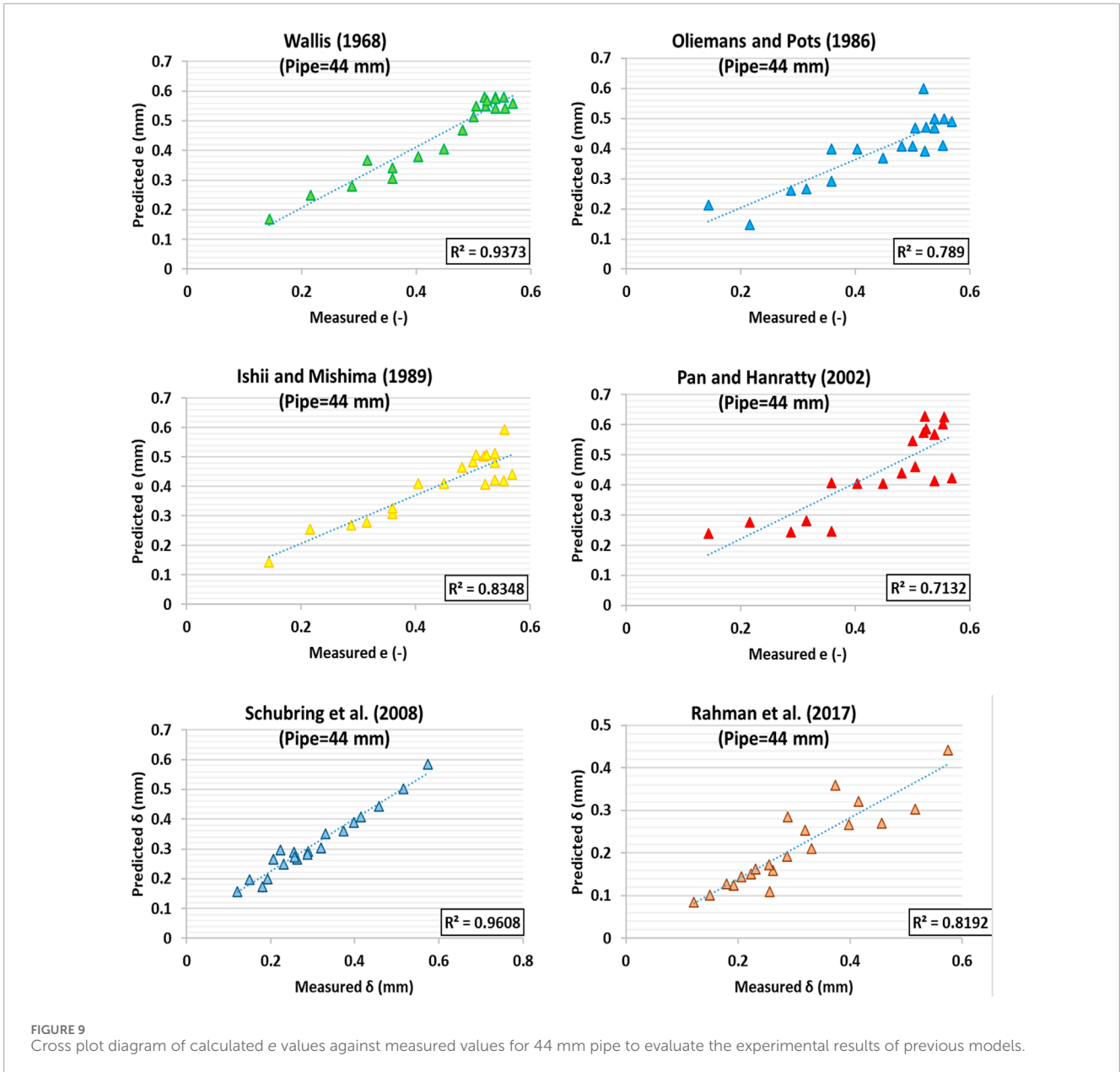


FIGURE 8
Cross plot diagram of calculated values against measured values for 26 mm pipe to evaluate the experimental results of previous models.

Tables 6, 7 respectively show the average uncertainties of air and water apparent velocities in the rectangular channel. The percent error is calculated based on the comparison of these values with the corresponding average apparent velocities. These tables indicate that the results show a low percentage of test uncertainties, demonstrating the high accuracy of the measuring equipment.

Despite the significant advancements in understanding countercurrent air-water two-phase flow in vertical pipes, several inherent limitations remain in this study. One primary limitation is the controlled nature of the experimental conditions, which does not account for real-world environmental variables, such as temperature variations and fluid impurities, that could affect flow dynamics and interfacial shear stress. Additionally, the study

examines only two specific pipe diameters (26 mm and 44 mm), which may limit the generalizability of the findings to pipes of other sizes. The reliance on existing models and correlations, while validated through comparative analysis, may still introduce biases or inaccuracies if these models fail to fully capture the complex interactions across all flow regimes. Furthermore, the experimental setup's restriction against heat and mass transfer may not accurately reflect practical scenarios where such transfers are unavoidable. Finally, the statistical evaluation, although robust, inherently assumes that the selected models and tested conditions are comprehensive, which may not always hold true, thus potentially limiting the applicability of the results to a broader range of operational conditions and pipe configurations.



The criterion for the number of test repetitions

Performing repeated tests is one way to increase the accuracy of test results. The confidence level is a parameter that can be used to control the number of repetitions of experiments to achieve the desired accuracy. Given that one of the goals of this research is to classify and present flow patterns, a 95% confidence level is the optimal value to achieve the objectives of this research. This level of confidence can be obtained using Equation 13 (Holman, 2012).

$$\alpha = \frac{t \times \omega}{\sqrt{n}} \tag{13}$$

Where, n is the number of experiments, t is a value corresponding to the desired confidence level, and based on $n = 3$ and a 95% confidence level, it is equal to 4.303. α is the uncertainty

resulting from the repetition of experiments. By replacing the values ω from Tables 6, 7, the uncertainty caused by repeating the experiment three times for the apparent velocities of air and water is presented in Table 8.

The role of interfacial shear stress in geoenery

Interfacial shear stress is a critical parameter in two-phase flow systems because it governs the interaction between different fluid phases, such as liquid and gas, within pipelines, geothermal wells, and CO₂ sequestration sites. Accurate determination of interfacial shear stress is essential because it influences the distribution of phases, the stability of flow regimes, and the efficiency of heat and mass transfer. Inaccurate estimations can lead to issues such as

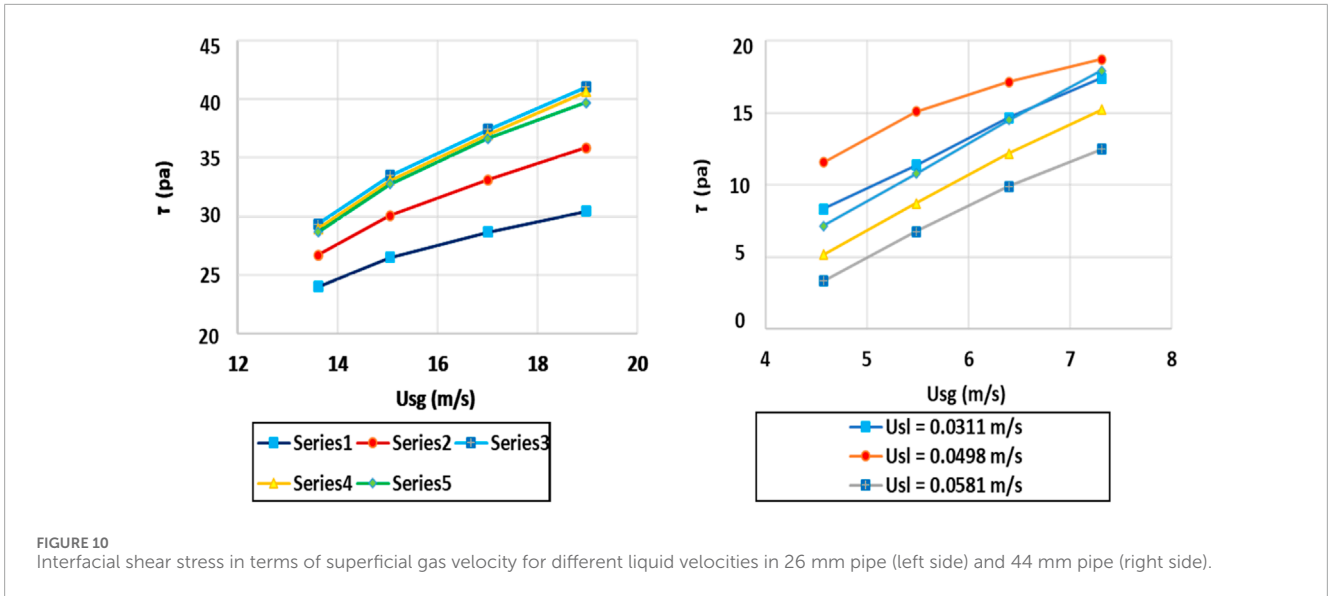


TABLE 5 Uncertainty of test equipment.

Measuring device	Uncertainty of measurement	Measurement range
Air flow meter	±4%	20–100 m ³ /h
Water flow meter	±4%	0.5–5 GPM
Ruler	±0.5 mm	0–200 mm

TABLE 6 Average and percentage of uncertainties of Usg.

Hydraulic diameter (mm)	Average measurement uncertainty	Uncertainty (%)
57.75	0.04	1

TABLE 7 Average and percentage of uncertainties of Usl.

Hydraulic diameter (mm)	Average measurement uncertainty	Uncertainty (%)
57.75	5.5×10^{-4}	0.47

TABLE 8 Uncertainty calculated from repeating experiments.

Hydraulic diameter (mm)	Average measurement uncertainty	Uncertainty (%)
57.75	±10 ⁻³	±0.09

excessive pressure drops, uneven phase distribution, and potential flow blockages, which compromise the performance and safety of geenergy systems. For example, in oil and gas pipelines, incorrect shear stress predictions can result in pipeline erosion or failure, while in geothermal systems, they can lead to inefficient energy extraction and increased operational costs. Thus, precise modeling of interfacial shear stress is crucial for optimizing the design, operation, and safety of geenergy infrastructure. Furthermore, the physical significance of these results lies in their implications for accurately modeling two-phase flow systems, which are critical in various geenergy applications. The accuracy of the liquid film thickness and entrainment models directly impacts the reliability of predictions related to fluid dynamics within pipelines, geothermal wells, and CO₂ sequestration sites. In real-world systems, incorrect predictions of liquid film behavior can lead to suboptimal design and operation, such as insufficient cooling in geothermal extraction, increased risk of pipeline corrosion due to uneven liquid distribution, or inefficient CO₂ trapping in sequestration projects. Therefore, understanding the limitations and accuracy of these models is crucial for optimizing the performance, safety, and longevity of geenergy systems.

Conclusion

Determining the key parameters of two-phase countercurrent liquid/gas flow in vertical pipes made of transparent Plexiglas for medium and small-sized pipes is fundamental for mechanical engineers. Hence, this paper aimed to define the influential parameters of two-phase flow, including entrainment (e), liquid film thickness (δ), and interfacial shear stress, and propose the corresponding experimental models. The main findings of this paper were as follows.

- Examining the ANOVA model for e in 26 mm and 44 mm pipes revealed the insignificance of the A2 parameter (square

power of U_{sg}). This variable should be eliminated from the calculations and presented models for better predictions.

- Examining the ANOVA model for δ in 26 mm and 44 mm pipes disclosed the insignificance of the B2 parameter (square power of U_{sl}), which should be eliminated from calculations and presented models.
- The correlation coefficients of δ and e for the 26 mm pipe were 0.9966 and 0.9806, respectively. The respective values for the 44 mm pipe were equal to 0.9957 and 0.9866. These results express the high accuracy of the proposed models for determining the δ and e in 26 mm and 44 mm pipes.
- Comparison of the experimental data for the determination of δ in 24 mm and 44 mm pipes with the results of Schubring et al. (2008) and Rahman et al. (2017) revealed the closeness of the results of the two references for the 26 mm pipe. However, the accuracy of the Schubring et al. model was slightly higher. For the 44 mm pipe, the Schubring et al. model provided better performance and offered more acceptable results.
- According to the presented results for determining the e in the 26 mm pipe, the best model with the highest performance accuracy was the model of Schubring et al. (2008), while the lowest performance accuracy was attributed to the Ishii and Mishima (1989) model. Additionally, for the 44 mm pipe, the models with the highest and lowest performance accuracies in the determination of e were associated with the Schubring et al. (2008) and Pan and Hanratty (2002) models, respectively.
- The related results to the flow patterns showed the dominance of the annular flow pattern in 26 mm and 44 mm pipes.
- The high validity of the prediction correlations for δ and e models was inferred from the results of interfacial shear stress diagrams for experimental data. Furthermore, the highest interfacial shear stress for the 26 mm pipe was related to the superficial liquid velocity of 0.0747 m/s, while this value was equal to 0.0498 m/s for the 44 mm pipe.

Data availability statement

The raw data supporting the conclusions of this article will be made available by the authors, without undue reservation.

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Author contributions

YW: Conceptualization, Formal Analysis, Resources, Software, Writing—original draft, Writing—review and editing. FL: Conceptualization, Data curation, Formal Analysis, Funding acquisition, Validation, Writing—original draft, Writing—review and editing. ZZ: Investigation, Methodology, Resources, Writing—original draft, Writing—review and editing. RL: Data curation, Investigation, Methodology, Validation, Visualization, Writing—original draft, Writing—review and editing. MS: Conceptualization, Formal Analysis, Investigation, Resources, Software, Visualization, Writing—original draft, Writing—review and editing.

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Conflict of interest

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Nomenclature

A_{fl}	liquid film cross-sectional area	ρ_g	Density of gas phases
u_{fl}	liquid film velocity	ρ_l	Density of liquid phases
m_{fl}	liquid film mass flow rate	δ	liquid film thickness
τ_i	interfacial shear stress	m_g	gas mass flow rate
f_i	interfacial friction factor	m_l	liquid mass flow rate
D	pipe diameter	n	number of experiments
ρ_{mix}	mixture density	Q	flow rate
$-dp/dz$	pressure gradient	ω_T	uncertainty of the measured value
u_{sl}	superficial velocities of liquid flow	ω_Q	uncertainty of the flow meters
u_{sg}	superficial velocities of gas flow	ω_d	uncertainty of the pipe diameters
g	gravitational acceleration	e	entrainment