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Nanocrystal formation during thermal treatment of phosphoric-acid-swollen polybenzimidazole membranes

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The design of suitable and efficient membranes for fuel and electrolyzer cells has been a long standing problem that still seeks for technical improvements. The structural properties of proton conducting poly [2,2-(m-phenylene)-5,5-bibenzimidazole] (PBI) membranes doped with phosphoric acid as electrolyte have been investigated prior and after thermal treatment by small angle neutron scattering (SANS), neutron diffraction with polarization analysis and transmission electron microscopy (TEM). Our observations indicate that during the first heating attempt stable crystalline domains in the PBI are formed and they persist in all further thermal cycles typical for electrolyte membrane operation. The combination of real space TEM picture under cryogenic conditions and reciprocal space scattering pattern with enhanced contrast by partial deuteration allows obtaining unique insight into the structure of proton conducting membranes, i.e., characteristic distances associated with irreversible changes as well as volume fraction of the different phases of the sample.

KEYWORDS

HT-PEFC, proton conducting membranes, PBI, membrane structure, SANS, diffraction

1 Introduction

Different polymeric membranes are candidates for polymer electrolyte fuel cells (PEFC) Page et al. (2012); Wang et al. (2011). High temperature polymer electrolyte fuel cells (HT-PEFC) operating at elevated temperatures between 100°C and 200°C (mainly in the range 140°C–180°C) are an interesting class of PEFCs with advantages for certain applications Kreuer (2013); Asensio et al. (2004). The most common material studied for low temperature polymer electrolyte fuel cells (LT-PEFC), operating at temperatures below 100°C is Nafion[®] Page et al. (2012); Schmidt-Rohr and Chen (2008); Moilanen et al. (2008). Operating temperatures above 100°C in HT-PEFCs simplify significantly the water management during operation, since water is mainly present in the vapour phase. The CO tolerance is also better in HT-PEFCs compared to the LT-PEFC counterparts. The typical application range is at medium power (in the order of 5 kW) Liu et al. (2016). At the heart of a HT-PEFC is the

proton conducting polyelectrolyte membrane between anode and cathode of the cell. The properties in terms of ion conduction and chemical stability are key for the performance of a HT-PEFC and are widely studied Kreuer (2013); Li et al. (2009); Wannek et al. (2009a). The operation temperatures of HT-PEFC above 100°C in acidic environment require polymerswith excellent thermal and chemical stability, one promising material being poly (2,5benzimidazole)-based (PBI) membranes Wannek et al. (2009b). The aromatic backbone is the reason for its good thermal and chemical stability and its high glass transition temperature Kreuer (2013); Di et al. (2012). Due to the basic nature of PBI it can be doped with large amounts of phosphoric acid (PA) as an electrolyte, which has a high proton conductivity Vilciauskas et al. (2012) and provides in combination with the scaffolding PBI material a well-performing proton conducting membrane Maier et al. (2012). Conductivity data of phosphoric acid and PBI are presented in Ref. Korte et al. (2016). A recent review on PBI based membranes can be found in Ref. Aili et al. (2020).

The structure of the polymer membrane in PEFCs is an important parameter for the thermal stability and ageing of the membrane under the harsh operating conditions, and also has an influence on the conductivity due to the confinement of the electrolyte inside the porous membrane. Many studies of Nafion-based perfluorosulfonic acid (PFSA) membranes for LT-PEFCs use small angle scattering (SANS) for structural investigations Schmidt-Rohr and Chen (2008); Moilanen et al. (2008); Haubold et al. (2001); Elliot et al. (2000). Molecular dynamics (MD) simulations of a hydrated perfluoro sulfonyl imide (PFSI) model system address the question how nanostructure and proton transport are related Idupulapati et al. (2011). The relationship between structure and transport for similar systems was investigated recently in details Berrod et al. (2015). X-ray radiography Maier et al. (2012) and tomography Eberhardt et al. (2015) as well as neutron imaging Boillat et al. (2014); Arlt et al. (2015) served as tools to study HT-PEFC fuel cells or different part of its core component, the membrane electrode assembly (MEA). Structural length scales of a similar ABPBI membrane were investigated on a broad range of length scales with microscopic and scattering techniques Ivanova et al. (2017). The swelling of the PBI membrane upon doping with an alkali ion conductor has been investigated with SANS and prompt gamma activation analysis Babcock et al. (2019). Phosphoric acid diffusion in the swollen PBI membrane has been measured with quasielastic neutron scattering and analyzed by a fractal diffusion model Hopfenmüller et al. (2018). Aili et al., stated that the degree of crystallinity of PBI membranes increases under heating, which results in creating highly proton conducting pathways through the membrane cross-section by forcing phosphoric acid into the amorphous regions Aili et al. (2012). However, details of structural changes and its interpretation are lacking. This report aims of shedding light on the mesoscopic structural changes of the PBI membrane demonstrating irreversibility of these changes under thermal cycles from room temperature to typical operating temperatures of HT-PEFCs with neutron scattering and diffraction as well as with transmission electron microscopic techniques.

2 Experimental details

2.1 Materials

The membrane material poly [2,2-(m-phenylene)-5,5bibenzimidazole] (PBI) was obtained from Danish Power Systems or Fumatech and doped with phosphoric acid according to the following (standard) procedure: The pristine ~40 µm thick PBI membrane has been immersed for several hours (at $T = 110^{\circ}$ C) in a bath with 85 wt% solution of phosphoric acid (PA). The PA uptake has been calculated from the weight increase of the pristine membrane upon doping (for details see ref. Conti et al. (2012); Majerus (2014)). The resulting PA uptake varied from the sample to sample. Depending on the experiment requirements, different samples have been used. However, we ensured sufficient saturation of PA in the sample by using relatively high concentration of PA: 392,498 wt%, corresponding to about 10.512.5 molecules per PBI repeat unit for neutron small angle scattering experiment, two samples with similar and higher concentration of PA, i.e., with 377 and 558 wt%, (corresponding to ca. 8.7 and ca. 14 molecules per PBI repeat unit), for neutron difraction and a sample with 488 wt% (corresponding to ca. 12.2 molecules per PBI repeat unit) for TEM. A detailed description of the phosphoric acid uptake in PBI can be found in Ref. Korte et al. (2015). It is important to note that deuterated phosphoric acid has been used to provide a sufficient contrast for neutron experiments. Phosphoric acid (H₃PO₄, 85%), and deuterated phosphoric acid (D₃PO₄, 85%) were purchased from Sigma Aldrich and used as received. Information concerning the synthesis of poly (2,5-benzimidazole) for use as a fuel cell membrane can be found elsewhere in literature Asensio et al. (2004); Li et al. (2009); Kim et al. (2004).

Sample preparation for both neutron experiments was similar apart from the different sample holder geometry (standard sandwich cell with quartz windows for the small angle scattering experiment and flat aluminium container for diffraction) and sample environment (furnace and closed cycle cryofurnace, respectively). Several sheets of dPA-doped PBI membrane were stacked together under Argon atmosphere and hermetically sealed in a PFA foil sachet in order to achieve the necessary sample volume and to prevent losing deuterated PA (dPA) by pumping in the diffraction experiment. The resulting samples have been measured at several temperatures (from RT temperature to 180°C and back) with equalibration time of 30 min. The temperature range is chosen in order to cover effect of water loss (above 80°C) and change from phosphoric acid (H_3PO_4) to pyrophosphoric acid $(H_4P_2O_7)$ above 140°C. The heating stages and temperatures presented for SANS and neutron diffraction are explained in more detail in the corresponding Section 3, Results and Discussion.

2.2 Experimental techniques and models used for data analysis

2.2.1 Small angle neutron scattering

With small angle neutron scattering (SANS), the ensemble averaged structure of the sample on length scales of nm to μ m can

be probed in reciprocal space where the modulus of the scattering vector *Q* is linked to the real space distance *d*. The average distance between the particles can be estimated by $d \sim 2\pi/Q$ Teixeira (1988).

Our SANS experiments presented in this paper were conducted at the KWS-2 beam line Radulescu et al. (2012) operated by Jülich Centre for Neutron Science at the FRM-II reactor, Heinz-Maier-Leibnitz Zentrum (MLZ), Garching. Two sample-to-detector distances of 2 and 8 m have been used to obtain a Q-range of $0.006 < Q < 0.3 \text{ Å}^{-1}$ with a wavelength of 4.97 Å and $8 \times 8 \text{ mm}^2$ sample aperture. The instrument had a velocity selector with a wavelength spread of $\delta\lambda/\lambda = 0.2$. SANS data were corrected with the usual procedure for instrumental background, sample transmission and detector sensitivity. The freely available software QtiKWS Pipich (2012) has been used for SANS data treatment.

2.2.2 SANS data model

SANS experiments allow to extract parameters such as the volume fraction of two different phases of the sample with a modelindependent approach by determining the scattering invariant Porod (1982). The scattering invariant *INV* reads:

$$INV = \int Q^2 I(Q) \, dQ = 2\pi (\Delta \rho)^2 \Phi (1 - \Phi) \tag{1}$$

with the difference in scattering length density between the two phases, $\Delta \rho$ and the volume fraction of one phase with respect to the other, Φ .

The *Q*-dependent scattering intensity has a power law decay of the form $I(Q) \propto Q^{-p}$ at low *Q* typical for a fractal structure at this range of length scales, with an exponent *p* which indicates the fractality of the structure. At larger *Q* a correlation peak appears which indicates a typical characteristic repeat distance within the sample. Developed for bicontinuous microemulsions, the Teubner-Strey formula Teubner and Strey (1987) describes such correlations with a typical repeat distance *d* and a correlation length in the sample, ξ :

$$I_{\rm TS}(Q) = \frac{I_0}{Q^4 - 2(Q_0^2 - \xi^{-2})Q^2 + (Q_0^2 + \xi^{-2})^2}$$
(2)

The characteristic length scale *d* can be simply obtained from Q_0 and the correlation length ξ is a measure how "crumpled" the structures are.

2.2.3 Neutron diffraction with polarization analysis

Neutron diffraction experiments that give access to the higher q-values compared to SANS data, have been conducted at the diffuse neutron scattering spectrometer with polarization analysis DNS Su et al. (2015); Schweika and Böni (2001) operated by Jülich Centre for Neutron Science at the FRM-II reactor, Heinz-Maier-Leibnitz Zentrum (MLZ), Garching. The measurements have been performed with an incident wavelength $\lambda = 4.2$ Å that allows to cover the Q-range from 0.2 to 2.5 Å⁻¹.

Despite using deuterated phosphoric acid, all our samples still contain a high amount of hydrogen. Due to its huge spin incoherent scattering cross-section, a strong shapeless incoherent background is found under the scattering from the structural correlations. In order to separate the coherent and incoherent differential cross sections in the DNS experiment we have used neutron polarization analysis. Neutron polarization was chosen to be perpendicular to the scattering plane. We have measured separately the intensities with and without flipping of the neutron polarization. Raw data have been corrected for detector efficiency and for the flipping ratio Stewart et al. (2009). Combining spin-flip and non-spin-flip Schärpf and Capelmann (1993); Stewart et al. (2009) we have separated the coherent and incoherent contributions. The coherent intensity provides information on the structural correlations. The incoherent signal is an independent quantification of the hydrogen amount in the sample. Below we present results of the diffraction experiments as a ratio of the coherent to incoherent scattering. In this representation the incoherent scattering term serves as an internal calibration factor for the coherent one. Such a calibration allows to get rid of the effects of sample size and geometry, as well as of influence of the inelastic scattering on the detector efficience Gaspar et al. (2010).

Two samples have been used: a fresh sample with 377% of D_3PO_4 , and the sample previously used in our quasielastic neutron scattering experiments Hopfenmüller et al. (2018) with 558% of D_3PO_4 . Background from the PFA foil has been measured and subtracted. Data reduction has been performed using the "dnsplot" interface of plot.py package Glavic.

2.2.4 Transmission electron microscopy (TEM)

In real space, TEM provides a view onto the PA-swollen PBI membrane on similar length scales as SANS (although without the possibility to perform *in situ* heat cycles). Therefore only the initial structure of the PBI membrane filled with phosphoric acid could be investigated.

Two resulting doped membranes were cut into pieces sized 3.5×10 mm. The obtained pieces were stacked and placed in a Leica FC7 cryo-ultra-microtome holder. The large-area sides were protected with Teflon tape and only the side perpendicular to the surface was cryo-sectioned: 100 nm sections of PBI polymer were produced by cryo-sectioning at -90°C in liquid nitrogen atmosphere using a glass knife with a clearance angle of 4°. The aim of using cryo-sectioning was to harden the specimen block. The sections were collected and deposited onto 100-200 mesh bare gold grids. Samples and grids were stored in Argon atmosphere when not used. The specimens were inserted into a Multi-Specimen Cryo Transfer Holder (Model 910, Gatan, Munich, Germany) and transferred to a JEM 2200 FS EFTEM instrument (JEOL, Tokyo, Japan). Experiments were carried out at temperatures around -180° C and a vacuum of 2.85×10^{-6} Pa. The transmission electron microscope was operated at an acceleration voltage of 200 kV. Zeroloss filtered images, enhancing the contrast, were recorded digitally by a bottom-mounted 16 bit CCD camera system (FastScan F214, TVIPS, Munich, Germany). Images of the PBI membrane slices (perpendicular to the surface) have been taken with the EMenu 4.0 image acquisition program (TVIPS, Munich, Germany) and processed with a free digital imaging processing system Image J Rasband (2012); Abràmoff et al. (2004); Schneider et al. (2012).

3 Results and discussion

The temperature dependence of the structure of PBI membranes doped with D_3PO_4 has been studied with SANS on nanometer scale prior to heating to the typical operation temperature of HT-PEFCs and and at elevated temperatures in the range of operating temperatures. Figure 1 (left) shows two curves of the small angle scattering data for the initial membrane (top) measured at a temperature of 80°C (which is identical to the measurement at lower temperature), and the same membrane heated for the first time to the 140°C (bottom). Any further cooling and heating of the membrane resulted in no further change. The structures developed during the first heating cycle and discussed in more detail below remained also after cooling down and in all further heating cycles without additional changes.

The SANS intensity resulting from the measurement of PBI membrane doped with dPA at 80°C is fitted with a model combining the power law decay and the Teubner-Strey description for the broad

shoulder at larger Q:

$$I_{\rm TS,LT}(Q) = I_{\rm TS}(Q) + I_1 Q^{-p} + bgr$$
(3)

The resulting parameters for the low temperature phase were $Q_0 = 0.03 \pm 0.001$ Å⁻¹, $d = 2\pi/Q_0 = 20.9 \pm 0.5$ nm and $\xi = 0.8 \pm 0.1$ nm.

The power law decay of the second term with I_1Q^{-p} describes the large scale fractal structure of the membrane at small Q-values. It is then followed by the broad shoulder, described with Eq. 2, which finally results in a Q^{-4} decay at large Q. The broad structure reflects the network of PBI strands in the swollen membrane.

As one can see from Figure 1 (left, bottom), for the measurement at 140° C the shoulder at large Q developed into a much more pronounced correlation peak. This has been accounted for by a



FIGURE 1

Left: Small angle neutron scattering data from phosphoric-acid-doped membrane: Top left: membrane measured at 80°C or 353 K and the same membrane heated to 140°C or 413 K (bottom left). The correlation peak at q~0.07 Å⁻¹ emerges above 80°C and reflects an additional repeat distance, which is attributed to the formation of crystalline domains induced by thermal treatment at typical operation temperatures. Right: Ratio of coherent to incoherent differential scattering cross section obtained from the neutron diffraction experiment plotted vs. momentum transfer Q. Top right: two measurements at the same temperature of 80°C (353 K) are shown - prior to heating to 140°C (413 K) in blue and after subsequent cooling down to room temperature in red. Distinct Bragg peaks emerge after the first heating of the sample, indicating the formation of crystalline regions. Bottom right: A previously heated sample measured at the same (353 K) and even higher temperature (453 K) as well as the base T (253 K) is shown to demonstrate the irreversible nature of the observed structural change to stable semi-crystalline structures. Arrows indicate positions of the peaks, corresponding to certain interatomic distances given in the text.



FIGURE 2

Cryo-transmission electron microscopy images (cryo-TEM) of PA-doped PBI domains on the same location at 500–1000 ms intervals each. Beam heating led to the structural contrast from the initially homogeneous sample. The middle image reveals the spots where the phosphoric acid is trapped prior the massive beam damage depicted in the third image.

second contribution of the form of Eq. 2. The high temperature SANS intensity $I_{TS,HT}(Q)$ is fitted with the following function:

$$I_{\text{TS,HT}}(Q) = I_{\text{TS}}(Q) + I_{\text{TS,LT}}(Q)$$
(4)

which results in $Q_{0,HT} = 0.073 \pm 0.001$ Å⁻¹, $d_{HT} = 2\pi/Q_{0,HT} = 8.6 \pm 0.1$ nm and $\xi_{HT} = 3.2 \pm 0.3$ nm.

The second Teubner-Strey contribution indicates that an additional structure with a characteristic average distance between the emerged objects, d_{HT} , has been created.

The scattering invariant gives a hint on the volume fraction occupied by one phase in a two phase system. If we take only the fitted component from $I_{\rm TS}(Q)$ and integrate it over Q, this leads to an indication of the volume fraction of the part of additionally emerged structure after the first heating cycle. Here, we find that the resulting volume fraction of the additional phase after heating is $\Phi = 0.04$. We assumed here that the difference in scattering length density between the amorphous phase and the crystalline phase of PBI is approximately 0.5×10^{-6} Å⁻².

Crystalline polymers are typically denser than amorphous ones Kholodovych and Welsh (2007), thus we assume here an increase in density by a factor typical for polymers of about ×1.16 resulting in the estimate of the scattering length densities above. From the weight measurements during sample preparation the ratio between polymer and phosphoric acid is approximately 1:4, i.e., one may assume that only 20% of the sample consists of polymer. The part exhibiting a clear correlation peak from the Teubner-Strey analysis of the heated sample however can be attributed to 4% only, which we interprete as the formation of small crystalline PBI domains within the strands of amorphous PBI matrix. Following this estimation with the crystal volume fraction 0.04 and the polymer fraction 0.17, we have $\Phi_{polymer}-\Phi_{cryst}$ = 0.13 left for amorphous polymer parts connecting and surrounding the crystals. The corresponding size qua diameter of the crystals is similar to the correlation length 8.6 nm from the Teubner-Strey analysis. Assuming spherical particles with $r = \xi/2$, this leads to a number of crystalline particles of $N=0.04/V_{\rm sphere}=2.3\times10^{21}$ particles per cubic centimeter. With an average distance between the crystals of d = 8.6 nm, the maximum thickness of the connecting amorphous polymer strands can then be estimated by distributing the remaining 16% of polymer volume to thin cylinders, resulting in $NV_{\text{PBIcyl}} = N(\pi R_{\text{PBIcyl}}^2 d) = 0.16$ and therefore $R_{\text{PBIcyl}} = (0.16/(N\pi d))^{1/2} = 1.6$ nm with the radius of the polymer strand R_{PBIcyl} . The diameter of a polymer strand connecting the crystalline clusters is then approximately $2R_{\text{PBIcyl}} = 3.2$ nm. Since the polymer strands obtained from the geometrical considerations above are smaller than the crystalline parts, a necklace structure where the crystal "beads" are connected with amorphous strands in between them is suggested.

The Porod part of the scattering $(I_1Q^{-p} \text{ in Eq. 3} \text{ with } p = 3.2 \text{ and } 3.6 \text{ for the initial and tempered membrane) can be connected to large scale inhomogeneities. It can be seen in a context of gels where large scale inhomogeneities are described by a Debye-Bueche scattering function (although we do not see the Guinier region here). It is characterized by a typical length scale of the inhomogeneity,$ *R*, and can be obtained from the slope at low Q of the SANS data. Since I (Q) of SANS is measured in absolute units it can provide the relevant length scale. Just from the Porod behaviour by taking the lowest-*Q* $data point and assuming an ideal <math>Q^{-4}$ dependence, the following relation is used:

$$I(Q_{\text{small}}) = 2\pi\phi(\Delta\rho)^2 \frac{3}{R} \frac{1}{Q_{\text{small}}^4}$$
(5)

For the inhomogeneities we assume a contrast obtained from crystalline vs. amorphous PBI. For the concentration of crystalline PBI we assume $\phi = 0.2 \times 0.04$ from considerations above. So we arrive at inhomogeneities of the size of either approx. $2 \mu m$ or 400 nm before and after the first heating to elevated temperatures. The latter value also seems to be motivated by the large scale brightness change of Figure 2.

The same effect of crystallite formation upon first heating of the membrane has been revealed at higher Q-values (sub-nm distances in real space) in the diffraction experiment. As one can see in Figure 1 (top right), the broad diffuse pattern at 1.6 Å⁻¹ corresponding to 3.9 Å transforms upon heating, thus leading to



Illustration of the suggested structure of phosphoric-acid-doped PBI membrane before and after first heating. The initial PBI scaffold is shown left (PA areas are represented by the voids in polymeric structure). After the first thermal cycle crystallization most probably takes place in small domains within the initially mainly amorphous scaffold (right, with additional crystallization domains).

the formation of a number of relatively well-defined peaks at the following Q-values: 0.5, 1, 1.25 and 1.75 Å⁻¹. On length scales of characteristic molecular distances this indicates the higher degree of morphological ordering. This sharpening of the Bragg peaks in the sample after the first heating to 140°C reflects the formation of small PBI crystals, which are also responsible for the additional contrast in the SANS experiment with the pronounced additional correlation peak from the average distance between those semi-crystal particles. A previously heated sample doped with a higher amount of phosphoric acid is shown in Figure 1 bottom right. It demonstrates the irreversibility of the produced semi-crystalline structures. One may notice the different ratio of the coherent and incoherent contributions for these samples. However, the lower ratio in Figure 1 (top right) indicates the presence of high amount of hydrogen in the system, which is consistent with a moderate PA concentration of the sample. The Q-values of the peak positions indicated in the figures can be related to the following distances: 13.4 Å, 8.9 Å, 6.4 Å, 4.9 Å, 4.2 Å, and 3.6 Å. These distances correspond to certain inter- and intra-chain correlations, i.e., the inter-segmental distance between polymer chains as well as interatomic distances between phosphoric acid molecules appearing during dehydration of phosphoric acid (formation of pyro-phosphoric acid and water above 140°C) in accordance with the available literature information [see, e.g., References Trisno et al. (2022); Persson et al. (2018)]. The changes in the structure slow down dramatically after the first heating and main changes remains after multiple subsequent cooling and heating cycles, as shown in Figure 1 (bottom right) for a sample with an even higher content of PA, where we have observed sharper peaks at similar distances after several heating and cooling cycles. This can be considered as an evidence of irreversible thermally induced formation of a (semi-)crystalline structure during operation at elevated temperatures.

Transmission electron microscopy (TEM) has been used to highlight the structure of a swollen membrane. In real space, TEM provides a view onto the PA-swollen PBI membrane on similar length scales as SANS. However it lacks the possibility to perform thermal cycles during the measurements. Therefore only the initial structure of the PBI membrane filled with phosphoric acid could be investigated. Figure 2 shows cryo-TEM images on the very same location (although in two different magnifications) at fewmilliseconds intervals. Initially there is no contrast between the PA and PBI domains and only a fully unstructured image is visible. The structural features of the membrane emerge then in subsequent images. It is obvious from this series of images that the phosphoricacid-doped membrane inside the cryo-TEM is very sensitive to electron irradiation. At the low pressure (of 2.85×10^{-6} Pa), the vapor pressure of PA or remaining water contents might result in an evaporation at the outer parts of the membrane, making the remaining thermally more stable polymer structure appear. Pure PBI membranes show almost no structure in the TEM, as has been shown, e.g., in Ref. Holderer et al. (2015) for a similar (ABPBI) membrane.

Radiation induced damage can be exploited in cryo electron microscopy to highlight structures otherwise not visible. For example, in polymer science, Talmon's group has used "radiation induced [...] formation of a high contrast network in the specimen" due to the presence of ice layer on organic film, Talmon et al. (1986) and references therein. In the case of polymer blends, the different sensitivities of the polymeric components can result in a stronger mass loss in one component than in another, yielding a development of contrast at the start of irradiation in the electron microscope. See Michler (2008) and references therein. According to the pressuretemperature phase diagram of H₃PO₄, the boiling point of the phosphoric acid is at -124°C under a vacuum of 3×10^{-6} Pa which is the vacuum condition in the TEM column during a cryo-TEM experiment. Although the cryo-TEM experiment is performed at ca. -180°C, the area exposed to the electron beam may be warmer locally, inducing the beam damage in the H3PO4 sites in the PBI membrane. Thus, electron beam focussing reveals the structure of the film, presumably by PA or residual water evaporation due to local heating. The resulting larger areas where the structure surrounding the created holes emerges (see Figure 2) might be interpreted as the area where the crystalline regions can be formed in the first thermal cycle.

Figure 3 illustrates the schematic picture from these considerations for a swollen PBI membrane (left), initially in a

fully amorphous state (left part of Figure 3), and after the first heating cycle with crystalline domains (right part of Figure 3). From our SANS results the microscopic picture of the PA-swollen PBI membrane can be represented by a continuous network of PBI strands with a diameter of about 3.2 nm, which is filled with PA. After the first heating cycles, about 1/5 of the PBI material forms crystalline clusters, which have an average distance from each other of about 8.6 nm and a diameter of roughly 3.2 nm. This is depicted by the sub-structures in Figure 3 (right). Neutron diffraction confirmed the increase of crystalline domains by a decrease of the width of the Bragg peaks. Also TEM images of swollen PBI membranes are consistent with this picture and the suggested structure of a swollen PBI network. The formation of crystalline domains seems to be an irreversible process, once the correlation peak is formed as shown by our neutron scattering data it stays for subsequent heating/cooling cycles.

4 Conclusion

Structural properties of phosphoric-acid-doped PBI membranes before the first use and after the first heating to operating temperatures have been investigated on length scales from atomic distances to micrometres with different neutron scattering techniques and electron microscopy. Scattering experiments allowed to observe the nano-structure of phosphoric-aciddoped PBI membranes under temperature conditions similar to operating conditions of the high-temperature fuel cells. Deuterated phosphoric acid has been used to achieve suitable contrast.

We found out that the first heating to the operating temperature of 140°C leads to a formation of a correlation peak significantly stronger than the shoulder-like broad slope of the initial sample. We interpret this as the formation of crystalline clusters in the thicker connecting points of the PBI network, which leads to an additional contrast with a characteristic repeat distance. Thermal cycles (not all curves are shown) indicate that only during the first heating of a PBI membrane, additional crystalline structures are created, most probably in small domains of the highly swollen PBI network. This higher crystalline ordering after heating the PBI structure for the first time is an irreversible process, and it can be speculated in which way such structures influence the mechanical stability of the membrane. Analysis of the scattering invariant together with peak position and correlation length of the additional peak allowed to deduce geometrical parameters of the PBI scatfold.

By combining small angle neutron scattering and neutron diffraction in reciprocal space with cryo-TEM images in real space a consistent picture of phosphoric-acid-swollen PBI membranes for fuel cell applications could be drawn. The overall picture obtained from cryo-TEM seems to be consistent with the initial stage of the

References

phosphoric-acid-doped PBI membrane. Additionally, electron beam damage has been used to observe the structure of the chemically and thermally stable PBI membrane under cryogenic conditions. Further studies would be required to shed a light on the mechanical influence of the formation of the crystalline domains in phosphoricacid-doped PBI membranes for high-temperature fuel cells.

Data availability statement

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

Author contributions

Conception and idea from OI, OH, and WL. Experiments by OI, KN, NS, and M-SA. Data evaluation and analysis by OI, OH, KN, M-SA, NS, VP, RZ, and MM. Writing of the paper by OH, OI, HF, KN, and RZ. All authors contributed to the article and approved the submitted version.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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