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Dipole doping effect in MoS₂ field effect transistors based on phase transition of ferroelectric polymer dopant

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Molybdenum disulfide (MoS₂) has great potential for next-generation electronic devices. On the other hand, stable doping methods are required to adjust its physical properties so MoS₂ can be utilized in practical applications, such as transistors and photodetectors. On the other hand, a conventional doping method based on ion implantation is incompatible with 2D MoS₂ because of the damage to the lattice structures of MoS₂. This paper presents an n-type doping method for MoS₂ field-effect transistors (FETs) using a poly (vinylidene fluoride-co-trifluoroethylene) (P (VDF-TrFE)) and polar polymer. The dipole moment of P (VDF-TrFE) provides n-type doping on MoS₂ FETs. The polar phase formation in dopant films enhances the doping effects, and the relationship between phase transition and n-type doping states was investigated using optical and electrical characterization methods. Under the optimal doping conditions, the doped MoS₂ FET achieved an improved field effect mobility of 34.4 cm² V⁻¹s⁻¹, a negative shift in the threshold voltage by -25.6 V, and a high on-current of 21 μA compared to the pristine MoS₂ FET.

KEYWORDS

transition-metal dichalcogenides, molybdenum disulfide, polar polymer, dipole moment effects, field effect transistors

1 Introduction

Since the first application of graphene in an electronic device, two-dimensional (2D) materials, including transition-metal dichalcogenides (TMDs), hexagonal boron nitride (h-BN), and Mxenes, have been highlighted in various research fields. Among the 2D materials family, MoS₂ is a promising next-generation semiconducting material owing to its mechanical flexibility (Samy et al., 2021), appropriate energy band gap (Kim et al., 2020), high carrier mobility (Shen et al., 2020), and high transparency (Lee et al., 2013). Therefore, MoS₂ may have a wide range of applications, including displays (Choi et al., 2020; Hwangbo et al., 2022), chemical sensors (Pham et al., 2019; Kumar et al., 2020), optical detectors (Gant et al., 2019; Woo et al., 2021a), and integrated circuit devices (Lin et al., 2020; Wang et al., 2022).

Despite this potential, issues, such as variation in physical properties based on the number of layers and a high contact resistance caused by the high Schottky barrier between

the MoS₂ and metal electrodes, are a problem in the application field of MoS₂-based electronic devices (Yoo et al., 2018). In general, doping is the most common technique for controlling the properties of semiconductors. Ion implantation is the most common technique for silicon-based electronics because it allows for delicate control and good quality in the doping process of thick silicon. On the other hand, this approach is incompatible with 2D MoS₂ because it causes significant damage to the lattice structure and defects in atomically thin layers.

Various attempts have been reported to overcome the limitations of the above-described conventional doping technology. For example, the surface functionalization of MoS₂ can be considered an alternative approach. Depending on the donor or acceptor states, p-type or n-type doping effects are achieved on MoS₂ devices (Pak et al., 2019). In doped transistors, however, a large reduction in the subthreshold swing was often observed due to the formation of trap states during the functionalization. Substoichiometric metal oxide dopant films fabricated by atomic layer deposition (ALD) can also provide efficient doping effects on MoS₂ FETs (Woo et al., 2021b). Metal oxide dopant layers exhibit high stability without degradation, but the on/off current ratios or subthreshold swings of doped devices are typically reduced (McClellan et al., 2021). Furthermore, the ALD-based doping approach requires a complex and high-cost process. Doping results for small molecule dopants have a high degree of freedom regarding the charge carrier concentration and doping types (Wang et al., 2020). Nevertheless, this method has stability issues in dopant films (McClellan et al., 2021) in the viewpoints of air exposure or thermal degradation. To maintain a doping effect in practical environments, the stability of the dopant layer should be secured. In general, the fixed physical or chemical properties of dopant layers provide high chemical resistance. However, fixed charges or states of dopant films have a limitation in terms of controllable doping effect with the variation of dipole moments. A doping approach that simultaneously provides characteristics such as controllability, reliability, and stability, which are important for the practical use of TMD-based electronic devices, is still lacking.

Along this line, this paper presents a controllable, reliable, and thermally stable n-type doping approach for MoS₂ FETs by forming a polar P(VDF-TrFE) polymer dopant layer. P(VDF-TrFE) is one of the most promising piezoelectric materials with good chemical stability (Choi et al., 2015) and dielectric properties (Meng et al., 2020). The dopant film was fabricated using a simple spin-coating method without any functionalization, which reduces the possibilities of trap-state formation. The top dopant layer provides n-type doping effects to the bottom MoS₂ based on dipole moments, which can be controlled further by the crystalline phases *via* the melting and recrystallization process. As the concentration of polar phases in dopant films increased, the proposed device exhibited improved electrical performances, such as on-current, on/off ratio, and field effect mobility. Furthermore, the uniformity and reproducibility of the proposed approach were investigated with 8 different devices. The n-type doping mechanism on MoS₂ was examined according to the phase transition of dopant films using comprehensive electrical

characterization, Raman and photoluminescence (PL) spectroscopies.

2 Materials and methods

2.1 The fabrication process of MoS₂ FET

A 300 nm of silicon dioxide film was formed using a dry oxidation process on a heavily boron-doped silicon substrate with 500–550 μm of thickness (supplied from Namkang Hi-Tech Inc., South Korea). Heavily boron-doped Si/SiO₂ substrates were sonicated for 10 min with isopropyl alcohol, and ethanol, subsequently and then blown with nitrogen gas on Si/SiO₂ substrates to remove residue. Semiconducting MoS₂ layers were mechanically exfoliated from the MoS₂ crystals (supplied from HQ Graphene, Netherlands) using adhesive 3 M Scotch tapes and transferred onto heavily boron-doped Si/SiO₂ substrates. At this time, the average thickness of 7 pristine MoS₂ flakes is 25.6 nm ± 12.3 nm. The source and drain electrodes were patterned on as-transferred MoS₂ flakes using the conventional photolithography method (EVG 620, EV Group, Austria). To pattern the source and drain electrodes, the image-reverse photoresist AZ 5214E (supplied from Clariant Corp., United States) was spin-coated on Si/SiO₂ substrates with MoS₂ crystals transferred. Then, the solvent was removed from the spin-coated photoresist by annealing process on a hot plate at 110°C for 1 min. The exposure process was performed on photoresist-coated Si/SiO₂ substrates using a photomask with a mesh structure. To remove the source/drain-patterned photoresist, phase transition of the photoresist was performed by annealing on a hot plate at 120°C for 150 s. Subsequently, photoresist-patterned Si/SiO₂ substrates were dipped in AZ 300 MIF developer (supplied by Az Electronic Materials Co. Ltd., United States) for 19 s to remove the photoresist with the source/drain patterned. To deposit the source/drain electrodes, the photoresist-patterned Si/SiO₂ substrates were mounted in a vacuum evaporation equipment (SHE-8T-500, Samhan Vacuum Development, South Korea). The 20 nm of titanium (Ti) for the adhesion layer was deposited by electron-beam evaporation method in a vacuum condition of 3 × 10⁻⁶ mTorr. Then, subsequently, 100 nm of gold (Au) was deposited on the entire surface of photoresist-patterned Si/SiO₂ substrates using the thermal evaporation method. Ti/Au-deposited Si/SiO₂ substrates were immersed in 60°C of acetone and the source and drain electrodes were patterned using the lift-off method. At last, the fabricated MoS₂ FET was dried using nitrogen gas. The length and width of the fabricated MoS₂ FET were 10.7 and 12.9 μm, respectively. The proposed device consisted of a bottom-gate/top-contact structure.

2.2 Remote n-type doping process

The dopant solution was prepared by dissolving P(VDF-TrFE) in acetonitrile (anhydrous, 99.8%) at a concentration of 7 mg mL⁻¹. The as-prepared solution was aged for 24 h for stabilization. The dopant solution was spin-coated on MoS₂ FETs at 500 rpm and 30 s conditions. The doped MoS₂ FETs were annealed at 60°C on a hotplate for 10 min after spin coating to remove the residual solvent.

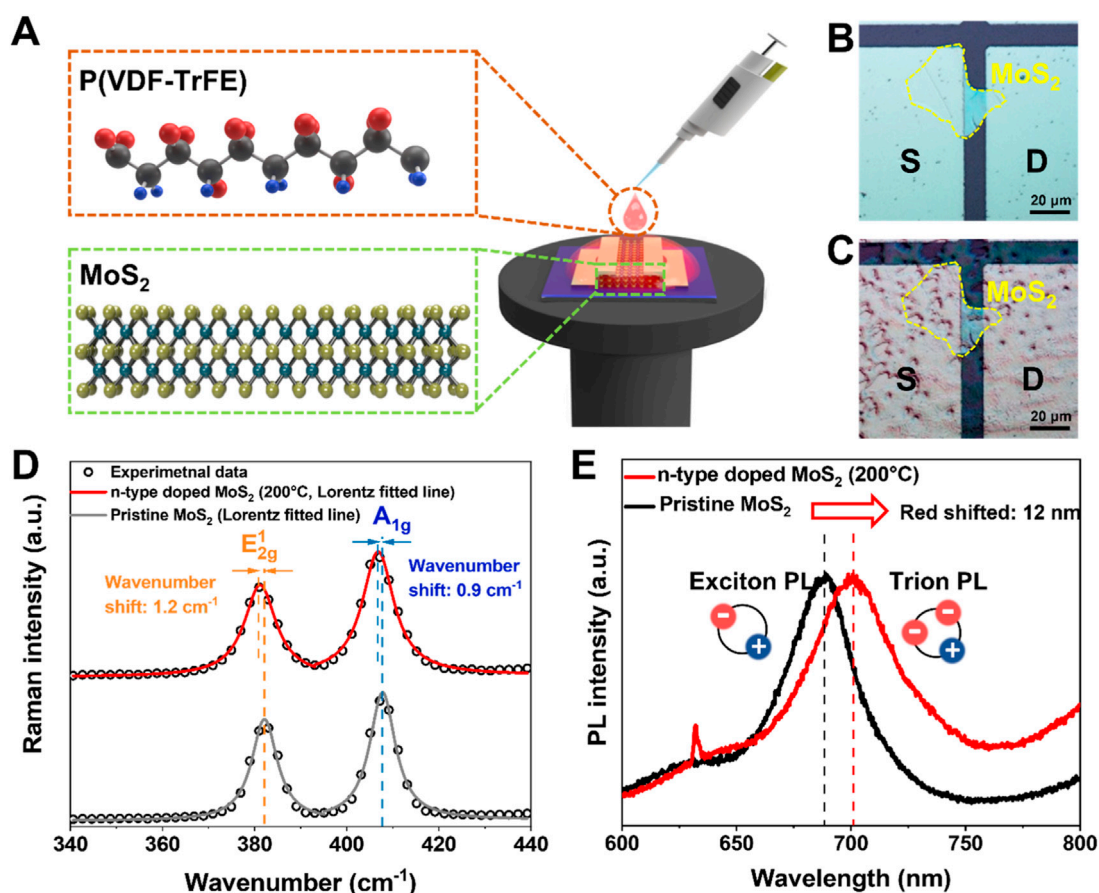


FIGURE 1

(A) Schematic diagram of dopant, MoS₂, and doping process. Optical micrographs of the MoS₂ FETs: (B) pristine MoS₂ FET and (C) P(VDF-TrFE)-doped MoS₂ FETs. (D) Raman spectroscopic analysis results of n-type doped and pristine MoS₂ flakes. The Lorentzian fits are represented with solid lines. (E) PL emission spectra of MoS₂ flakes according to the n-type doping states under 532 nm excitation.

The effects of P(VDF-TrFE) phase transition were examined by sequential annealing of the as-doped devices at 100, 200, and 300°C. In addition, the P(VDF-TrFE)-doped MoS₂ FETs with 10 min annealed had slowly cooling to 60°C.

2.3 Characterization methods

The phase transition of dopant P(VDF-TrFE) and lattice vibration of the MoS₂ flakes were observed by Raman and PL spectroscopy with an in-house-built, combined spectroscopy system at room temperature. Both Raman and PL spectroscopy measurements were performed using a SYCERITY 1024 × 256-OE TE cooled CCD detector (HORIBA Scientific, Japan) for UV-VIS-NIR regimes and a HORIBA Jobin Yvon NanoLog equipped with a laser excitation wavelength of 532 nm. All spectra were collected with an optical microscope (BX43-32P01HAL, Olympus, Japan) with the ×100 objective lens (N.A. 0.9). Water contact angles of P(VDF-TrFE) films annealed at different conditions were measured with the Contact angle meter (DSA100, KRÜSS Scientific Instruments Inc., Germany) All electrical characterization of MoS₂ FETs in this work was measured in air-ambient condition

(RH = 21% and temperature = 23.6°C) using a Keithley 4,200 parameter analyzer (Keithley Instruments Inc., United States) installed in a probe station (MS-Tech, South Korea). In addition, the thermal stability of the P(VDF-TrFE)-doped MoS₂ FET was measured using a probe station chamber connected to a Keithley 4,200 parameter. A hot chuck mounted in the probe station chamber applies a temperature of up to 100°C to the devices.

3 Result and discussion

3.1 Characterization of n-type doping effects on MoS₂ FETs

First, the changes in morphological and optical properties of the MoS₂ FET according to the remote n-type doping process were examined. Figure 1A illustrates the n-type doping process and molecular structure of the components. Optical microscopy photographs according to n-type doping states (Figures 1B, C) showed no observable decomposition in MoS₂ flakes, proving that the proposed solution coating method is an appropriate

n-type doping approach without causing significant damage to the thin MoS₂ semiconductor. The optical characterization of the MoS₂ region revealed the effect of the n-type doping process. The optical characterization of the MoS₂ region revealed the effect of the doping process with Raman spectroscopy (Figure 1D). To clearly observe the shape and shift in A_{1g} and E_{2g}¹ peaks, we have performed Lorentzian function fitting for each raw data point. As shown in Figure 1D, the A_{1g} and E_{2g}¹ modes provide information on the out-of-plane strain and in-plane vibration, respectively (Zhou et al., 2014). The A_{1g} and E_{2g}¹ peaks are both red-shifted in the n-type doped MoS₂ sample compared to the pristine MoS₂. The specific positions of A_{1g} peak in the fitted line were 406.8 cm⁻¹ and 407.7 cm⁻¹ at n-type doped MoS₂ and pristine MoS₂, respectively. The E_{2g}¹ peak positions were extracted with the values of 380.9 cm⁻¹ and 382.1 cm⁻¹ for n-type doped and bare MoS₂ films. As reported in previous studies, heavy n-type doping shifts the Fermi level near the conduction band, causing interactions between the continuum of electronic states and optical phonons (Sahoo and Kale, 2019; Schauble et al., 2020; Tanwar et al., 2022). The above electron-phonon interaction (i.e., Fano interaction) arises owing to appropriate matching between discrete phonon energy and electron continuum energy by heavy doping and photoexcited carriers. Resonance of the two energies leads to constructive or destructive interference and results in perturbed line shapes at Raman spectral parameters (Tanwar et al., 2022). In the case of A_{1g} mode, peak shift is the main indicator of the n-type doping of MoS₂ (Chakraborty et al., 2012; Miller et al., 2015; Manamel et al., 2020; Krishnan K et al., 2021). Furthermore, the increased asymmetric ratio in Raman peaks provide the quantitative information of doping induced Fano interaction (Sahoo and Kale, 2019; Schauble et al., 2020; Tanwar et al., 2022). To quantitatively compare the MoS₂ E_{2g}¹ peak asymmetric ratios depending on the n-type doping procedure, we extracted the peak parameters such as full width at half maximum (FWHM, Γ), Half-width in pre-maximum side (γ_L), and Half-width in post-maximum side (γ_R) by fitting the raw data points with bi-Gaussian function. As shown in Supplementary Figure S1, the asymmetric ratio (γ_L/γ_R) of pristine MoS₂ was 1.01, which is almost close to 1 and reveals the peak is symmetric. In contrast, the asymmetric ratio of the n-type doped MoS₂ was calculated to be 0.95, suggesting that the peak asymmetry increased through the electron-phonon interaction (i.e., Fano interaction) with the doped state. The above-mentioned asymmetry in E_{2g}¹ Raman peaks were further investigated through Fano fitting which defined by the following Eq. 1 (Schauble et al., 2020).

$$a(\omega) = a_0(q + 2(\omega - \omega_p)/\gamma)^2 / (1 + (2(\omega - \omega_p)/\gamma)^2) \quad (1)$$

Where a_0 is the prefactor, ω is frequency, ω_p is bare phonon frequency, γ is the linewidth, and q is the symmetry parameter depends on electron-phonon coupling strength. As previously reported in literatures, a lower absolute value $|q|$ indicates greater asymmetry and heavier doping (Nickel et al., 2000; Schauble et al., 2020). The calculated q values in the Fano profiles for pristine and n-type doped MoS₂ were 29.89 and 12.25, respectively (Supplementary Figure S2). The lower q value for the n-type doped MoS₂ suggests a greater asymmetry and a heavier doped state compared to the pristine MoS₂.

The changes in the electronic states of the MoS₂ region after n-type doping were also discovered in the PL spectrum (Figure 1E). The PL emission peak of doped MoS₂ was red-shifted by 12 nm compared to the pristine MoS₂. The PL emission of MoS₂ was attributed to exciton (X, peak at ~ 690 nm) and trion (X⁻¹, peak at ~ 702 nm) recombination (Brill et al., 2021). Because trions are more negatively charged than excitons, an increase in trion-related PL emission shows that the MoS₂ region is more negatively charged (n-doped). The PL peak was red-shifted towards the trion PL emission direction when the P(VDF-TrFE) film was fabricated on the MoS₂. This result confirms the n-type doping effect through dopant fabrication. The effect of n-type doping on MoS₂ observed in the PL measurements was verified through the electrical characterization of the device.

As mentioned in the introduction, certain organic materials, such as self-assembled monolayers (SAMs), and polar polymers, can provide n-type doping effects to control the electrical characteristics of MoS₂ FETs (Li et al., 2013; Kang et al., 2015; Maity et al., 2017). The dipole effect could change the electrical characteristics of a MoS₂ FET in the case of the P(VDF-TrFE) film coated with the surface of the MoS₂ semiconductor. Figures 2A, B present 3D schematic diagrams of pristine and MoS₂ FETs doped with a P(VDF-TrFE) solution. An as-coated P(VDF-TrFE) thin film was formed after 10-min annealing to evaporate the solvent. Subsequently, annealing was performed at 100, 200, and 300°C using a hotplate for 10 min. The electrical characteristics of the MoS₂ FETs were measured to assess the n-type doping effect of P(VDF-TrFE). Figure 2C shows the transfer curves of a pristine MoS₂ FET and P(VDF-TrFE)-doped MoS₂ FET. The drain voltage of 1 V and the gate voltage range of -40 V-40 V were applied to two types of MoS₂ FETs to collect the transfer curve. The pristine MoS₂ FET with typical n-type behavior showed an on-current of 8.7 μ A, and a field effect mobility of 26.9 cm² V⁻¹s⁻¹. In addition, the P(VDF-TrFE) doping effect was maximized according to the annealing temperature. The on-current and field effect mobility of the P(VDF-TrFE)-doped MoS₂ FET (100°C) increased significantly to 12 μ A and 30.7 cm² V⁻¹s⁻¹, respectively. The P(VDF-TrFE)-doped MoS₂ FET (200°C) with the most improved n-type operation has a 21 μ A on-current and 34.4 cm² V⁻¹s⁻¹ of mobility. The crystal structure of P(VDF-TrFE) has α , β , and γ phases and can be controlled by annealing temperature. The P(VDF-TrFE) annealed at 200°C has high alignment due to the formation of the β phase. The β phase of P(VDF-TrFE) induces a strong dipole moment and enhances the n-type operation of the MoS₂ FET. On the other hand, the electrical characteristics of the P(VDF-TrFE)-doped MoS₂ FET (300°C) are deteriorating, with an on-current of 1.6 μ A and a field effect mobility of 12.4 cm² V⁻¹s⁻¹, respectively. The dipole moment of P(VDF-TrFE) induced on the surface of the MoS₂ semiconductor was weakened because the alignment of the P(VDF-TrFE) was lower. Figure 2D shows the output curves of pristine MoS₂ FET and doped MoS₂ FET (200°C). The output curves of the MoS₂ FETs were measured by applying a drain voltage of 0 V-10 V and a gate voltage of -10 V-40 V. The P(VDF-TrFE) has the effect of increasing the on-current of MoS₂ FETs with increased drain current in the same gate voltage range. Figure 2E presents the contact resistance and channel resistance of a pristine MoS₂ FET and a P(VDF-TrFE) doped MoS₂ FET extracted by the G-function method (Liu et al., 2015; Jiang et al., 2019). The channel resistance

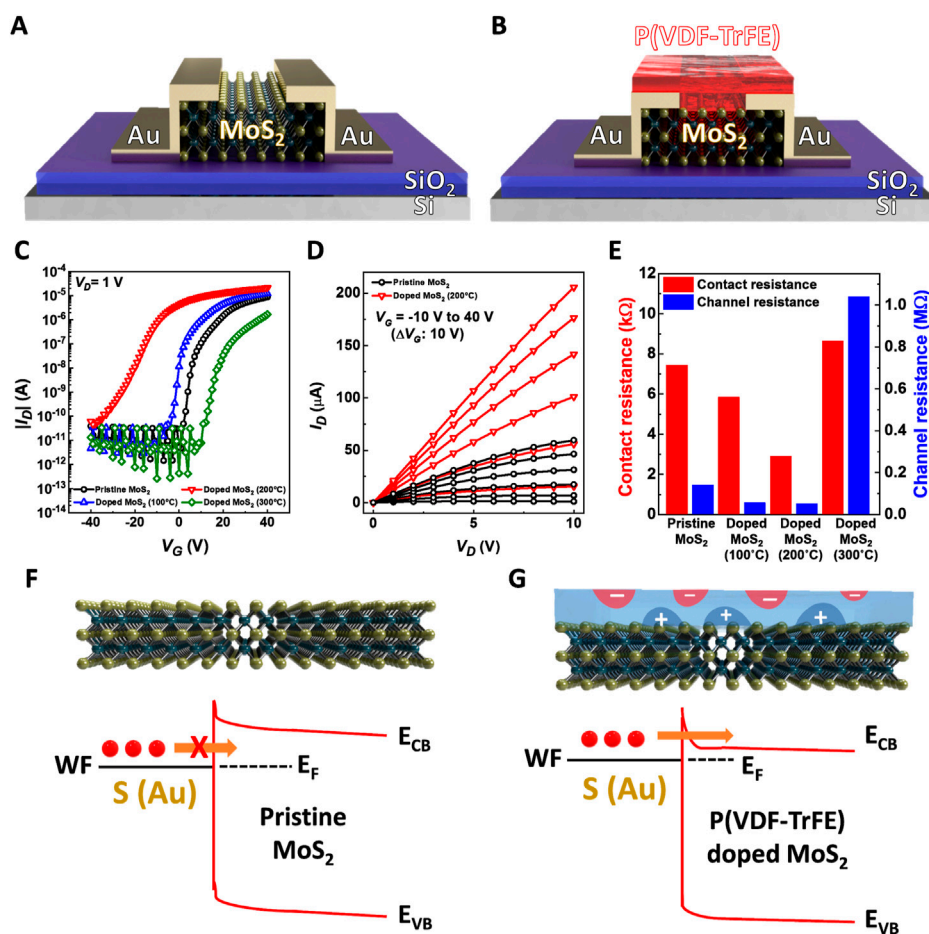


FIGURE 2 3D schematic diagram of the (A) pristine and (B) P(VDF-TrFE)-doped MoS₂ FET. (C) Transfer curves of the pristine MoS₂ FET and the P(VDF-TrFE)-doped MoS₂ FETs (100, 200, and 300°C). (D) Output curves of the pristine MoS₂ FET and the P(VDF-TrFE)-doped MoS₂ FET (200°C). (E) Contact resistance and channel resistance of the pristine MoS₂ FET and the P(VDF-TrFE)-doped MoS₂ FETs (100, 200, and 300°C) extracted by the G-function method. (F) Energy band diagram of the pristine MoS₂ FET. (G) Energy band diagram of the P(VDF-TrFE)-doped MoS₂ FETs.

and contact resistance of the extracted MoS₂ FET decrease as the annealing temperature increases. The channel resistance and contact resistances of the pristine MoS₂ FET are 0.14 MΩ and 7.4 kΩ, respectively. When MoS₂ FET is doped with P(VDF-TrFE) and subsequently to an annealing process at 100°C, the channel resistance (0.55 MΩ) and contact resistance (5.8 kΩ) are significantly reduced. The P(VDF-TrFE)-doped MoS₂ FET (200°C) exhibits the lowest channel resistance (0.51 MΩ) and contact resistance (2.9 kΩ). Low channel resistance and contact resistance increase the field effect mobility, on-current, and lower the operating voltage of MoS₂ FET. On the other hand, the channel resistance (1.04 MΩ) and contact resistance (8.6 kΩ) of the P(VDF-TrFE)-doped MoS₂ FET (300°C) increased compared to pristine MoS₂ FET. The P(VDF-TrFE) doping effect lowers the Schottky barrier in the MoS₂ FET and induces a low contact resistance. Owing to the dipole effect of P(VDF-TrFE), the n-type behavior of the doped MoS₂ FET was enhanced. Figure 2F shows the energy band diagram of a pristine MoS₂ FET. The energy band of pristine MoS₂ makes it difficult to tunnel electrons into the electrode. On the other hand, P(VDF-

TrFE) brings the conduction band and Fermi level of MoS₂ closer (Figure 2G). Hence, the dipole moment of P(VDF-TrFE) shifts the energy band of MoS₂ downward, facilitating electron tunneling.

The electrical parameters were extracted to investigate the n-type doping effect of P(VDF-TrFE) induced on MoS₂ FETs. The uniform n-type doping effects of the proposed approach were observed in terms of field effect mobility, on-current, and negatively shifted threshold voltage with 8 different MoS₂ FETs. Figure 3A presents the average and maximum field effect mobility of pristine MoS₂ FETs and P(VDF-TrFE)-doped MoS₂ FETs. The average field effect mobility of 8 MoS₂ FETs was extracted as a histogram and an error bar of standard error was added. In this work, the transfer curve of 8 MoS₂ FETs for extracting electrical parameters was reported in Supplementary Figure S2. The electrical characteristics of MoS₂ FETs are significantly impacted by various process factors. Chemical vapor deposition (CVD) and mechanical exfoliation with two common synthesis methods influenced the quality of MoS₂ due to process variations (Kim et al., 2014, 2016; Choi et al., 2022). In addition, the contact resistance between the electrodes and channel of MoS₂ FETs is changed based on the

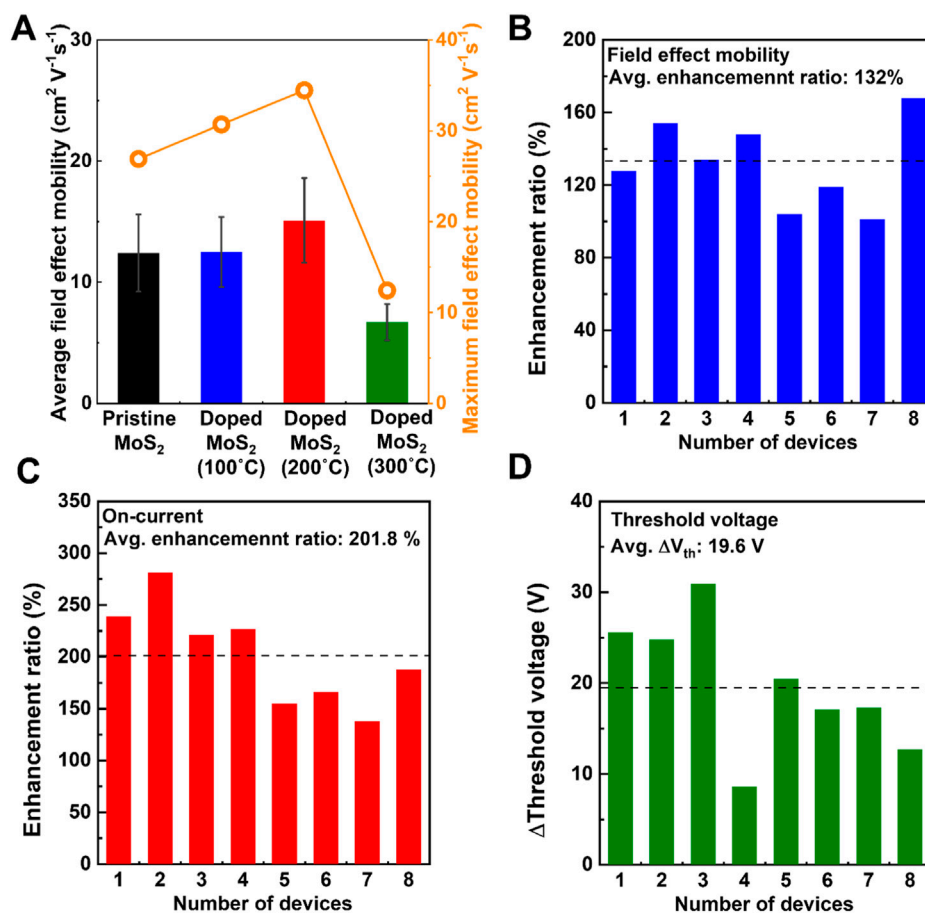


FIGURE 3

Electrical parameters of the pristine MoS_2 FET and the P(VDF-TrFE)-doped MoS_2 FETs (100, 200, and 300°C) (A) Average and maximum field-effect mobilities. The enhancement ratio of 8 different P(VDF-TrFE)-doped MoS_2 FETs (200°C): (B) Field-effect mobility, (C) on-current, and (D) changes in the threshold voltage values compared to pristine devices.

deposition method for metal, annealing temperature, and metal/ MoS_2 interface, resulting in variations in carrier mobility. Also, the performance of MoS_2 FET is affected by the photolithography conditions that affected the surface morphology and quality of MoS_2 (Plechinger et al., 2014; Wang et al., 2019; Yang et al., 2021). The average field effect mobility of the 8 pristine MoS_2 FETs was $12.4 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and increased slightly to $12.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ when the P(VDF-TrFE) annealing temperature was 100°C. The average field mobility of MoS_2 FETs in the optimized P(VDF-TrFE) annealing (200°C) process was $15.1 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$. On the other hand, the field effect mobility of P(VDF-TrFE)-doped MoS_2 FETs (300°C) was $6.7 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, which is lower than pristine MoS_2 FET. In addition, the MoS_2 FET with the highest mobility was $26.9 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for the pristine sample, while it was $34.4 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ under optimized annealing conditions. Also, the negative shifts in threshold voltage were observed at n-type doped MoS_2 FETs. The dipole effect of top P(VDF-TrFE) region induces the conduction band and Fermi level of MoS_2 closer. The downward energy band promotes rapid inversion due to enhanced electron tunneling, resulting in the threshold voltage shifting negatively. Supplementary Figure S3 shows the threshold voltages of pristine MoS_2 FET and P(VDF-

TrFE)-doped MoS_2 FETs. The pristine MoS_2 FET has a threshold voltage of 16.7 V. The P(VDF-TrFE)-doped MoS_2 FET (100°C) has a negatively shifted threshold voltage of 9.8 V due to relatively weak n-type doping effect from the large fraction of non-polar α -phase P(VDF-TrFE). The P(VDF-TrFE)-doped MoS_2 FET (200°C) under the optimized conditions with strong n-type doping effect had -8.9 V of threshold voltage due to the enlarged fraction of the polar β phase P(VDF-TrFE). However, the threshold voltage of P(VDF-TrFE)-doped MoS_2 FET (300°C) was shifted in the positive gate voltage direction ($V_{th} = 30.6$ V). This phenomenon is related with the Mo-O bonds generated in MoS_2 during the high-temperature annealing. The Mo-O bonding created by adsorbed oxygen and water molecules causes a p-type doping effect in MoS_2 (Wei et al., 2014). As a result, the n-type doping effect on P(VDF-TrFE)-doped MoS_2 FET (300°C) was reduced owing to the Mo-O bonding induced p-type doping effect. Statistical analysis of 8 P(VDF-TrFE)-doped MoS_2 FETs (200°C) was performed to investigate the uniformity of the P(VDF-TrFE) doping effect. The field effect mobility of the P(VDF-TrFE)-doped MoS_2 FETs (200°C) increased by an average of 132% compared to the pristine MoS_2 FET (Figure 3B). Figure 3C shows the enhancement ratio of on-current with the P(VDF-TrFE)-doped MoS_2 FETs (200°C) compared to the

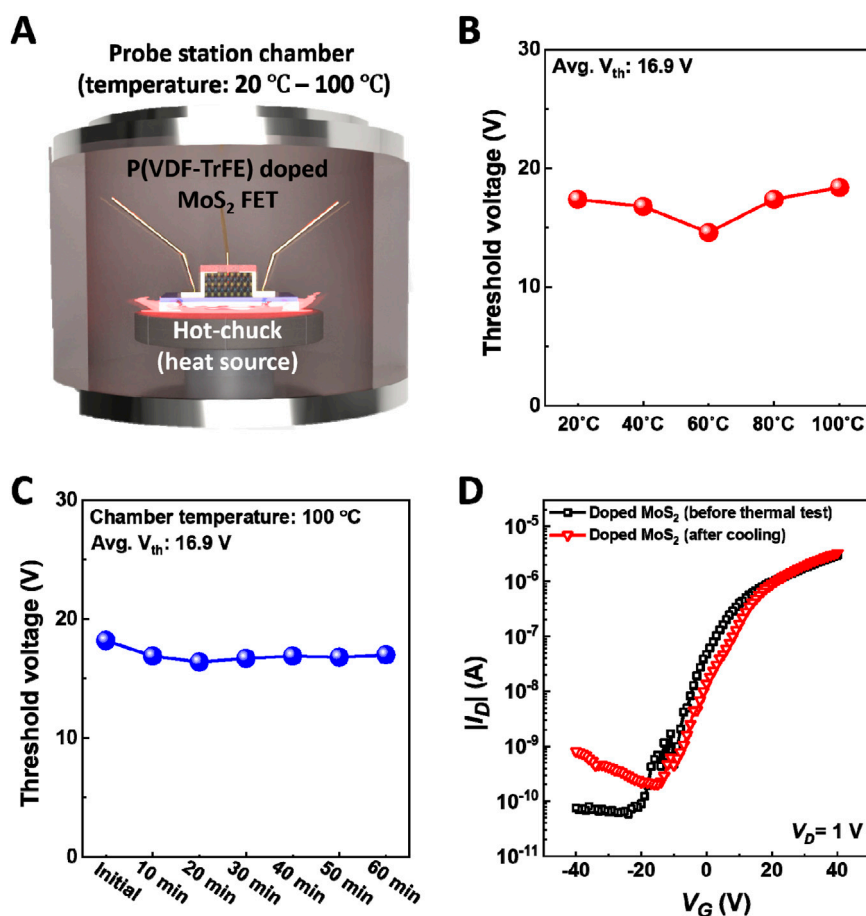


FIGURE 4

(A) 3D schematic description of the thermal stability characterization environment. Threshold voltage change according to the experimental conditions: (B) during the chamber temperature incrementation from 20 to 100°C, (C) continuous 100 °C environment condition for 1 h. (D) Transfer curves of the P(VDF-TrFE)-doped MoS₂ FET (200°C) before and after the thermal stability characterization.

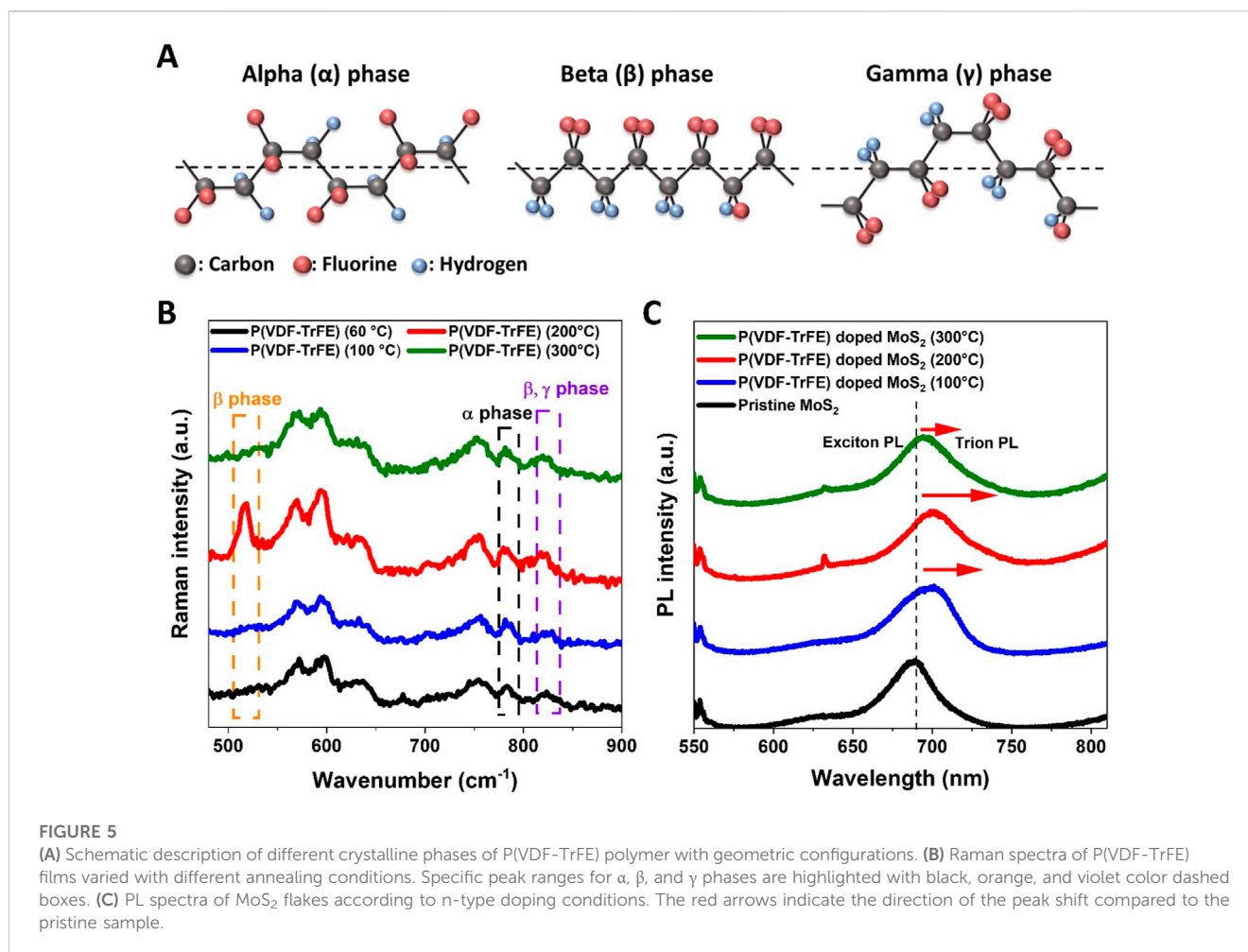
pristine MoS₂ FET. The n-type doping effect of P(VDF-TrFE) on MoS₂ FETs increased the on-current of the 8 devices by an average of 201.8%. [Supplementary Figure S4](#) represents the output curves of 8 P(VDF-TrFE)-doped MoS₂ FETs (200°C). As shown in [Supplementary Figure S4](#), 8P (VDF-TrFE)-doped MoS₂ FETs (200°C) achieved increased drain current compared to pristine MoS₂ FETs. In addition, the average threshold voltage of 19.6 V has been shifted negatively ([Figure 3D](#)).

We further validated the robust n-type doping effect of the proposed method through a thermal stability characterization. [Figure 4A](#) shows a 3D schematic description of the experimental environment with P(VDF-TrFE)-doped MoS₂ FET (200°C). More specifically, the P(VDF-TrFE)-doped MoS₂ FET (200°C) was placed in the hot-chuck mounted probe station chamber. The temperature was increased from room temperature (RT, 20°C) to 100°C and maintained for 1 h. After 1 h, the chamber was naturally cooled to 20°C. The electrical characterization of the device was measured during the whole thermal sequence at specific points. [Figure 4B](#) displays the threshold voltage variation of the P(VDF-TrFE)-doped MoS₂ FETs (200°C) during the heating-up condition range of 20–100°C. The threshold voltage before the thermal stability test was measured to 17.4 V (at 20°C condition). While increasing the

thermal condition of the chamber up to 100°C, the threshold voltage values were varied with the lowest value of 14.6 V (at 60°C) and the highest value of 18.4 V (at 100°C). The average threshold voltage measured during the 5 steps is 16.9 V, which represents a negligible difference of 0.5 V compared to the pristine value (17.4 V). Furthermore, the above-mentioned stable threshold voltage retention was demonstrated even under continuous thermal conditions (100°C for 1 h). As shown in [Figure 4C](#), the threshold voltage values measured at each of the 10 min cycles were almost maintained without any sudden degradation or fluctuation, which proves the good thermal stability of the proposed device at 100°C environments. [Figure 4D](#) shows the comparison of transfer curves before and after the thermal stability test. Compared to the initial characteristics, only slight changes were observed with an increase in off-current and on-current with the values of 14.7 nA and 28.9 μA.

3.2 Effect of phase transition on n-type doping states

Next, to understand the optimized n-type doping condition (200°C annealing temperature), we investigated the change in the



dopant P(VDF-TrFE) film according to the annealing variation. As discussed previously, the main mechanism of n-type doping on MoS_2 by P(VDF-TrFE) is based on the dipole moment in the dopant film. Depending on the molecular arrangement, P(VDF-TrFE) has three different crystalline structures: α , β , and γ phases. Among them, the β and γ phases exhibited polar dipole moments owing to their geometric configuration, as shown in Figure 5A (Shepelin et al., 2019; Arrigoni et al., 2020). The crystalline phases of P(VDF-TrFE) mentioned above can be obtained through a sequential melting and recrystallization process (García-Gutiérrez et al., 2013; Arifin and Ruan, 2018). The non-polar α phase, however, is the most thermodynamically stable form, so the P(VDF-TrFE) polymer crystallizes preferentially into it during melting and recrystallization. Based on this phenomenon, it was predicted that different ratios of the polar phases (β and γ) and non-polar (α) phase would be formed according to the annealing conditions. As a result of the phase variation, the higher polar phase ratio leads to strong n-type doping effects on MoS_2 FETs observed during electrical characterization.

Raman spectroscopy was performed to observe the crystalline phases of the P(VDF-TrFE) films varied with annealing conditions (Figure 5B). Specific peaks according to the α , β , and γ phases of P(VDF-TrFE) were identified. The non-polar α phase appeared in all samples with the peak range at 775 cm^{-1} to 795 cm^{-1} (black dashed box)

(Riosbaas et al., 2014). The peak indicated a high polar β phase (510 cm^{-1} , orange dashed box) (Viswanath and Yoshimura, 2019) was observed only in the sample (annealed at 200°C), which exhibited the most n-doped characteristic. Furthermore, the intensity of peaks at 810 cm^{-1} to 840 cm^{-1} (violet dashed box, mixed phases of polar β and γ) (Singh et al., 2014) tended to increase in the order of efficiency of the n-type doping effect in electrical characterization results (Figures 2C, 3A). The most n-doped 200°C annealed sample exhibited a clear polar β phase. The intensity ratio of α phase/ β and γ mixed phase were reduced after 300°C annealing. The polar phase increase according to the recrystallizing condition in P(VDF-TrFE) film was cross-validated through contact angle measurements. As reported in previous studies, the enrichment of crystallization domains on the surface of the PVDF-based polymer films also can promote the hydrophilicity (low water contact angle) (Chen and Hong, 2002; Zhao et al., 2012). The increased polar crystalline phases enhance the surface polarity of the films, and results in the reduced water contact angle. As shown in Supplementary Figure S6, the water contact angle change according to the annealing condition is consistent with the β phase intensity ratio change from Raman analysis. The contact angle showed the lowest value at 200°C annealing condition where the β phase peak intensity and n-type doping effect on MoS_2 FET were maximum, and slightly increased at 300°C when the n-type doping effect was weakened. The above results are well-matched with the relationship

between polar phase formation and n-type doping efficiency. PL characterization supports the variation of n-type doping effects according to the dopant film phases. As shown in Figure 5C, the PL emission peaks were red-shifted for all doped samples, representing increased trion-related emissions in n-doped states. Nevertheless, the shifted values varied with the n-type doping conditions. The peak shifts were 10, 12, and 6 nm for the 100, 200, and 300°C annealed samples, which agree well with the n-type doping efficiency trends discussed above. Through the above characterization, it was confirmed that the P(VDF-TrFE) dopant film provides a strong n-type doping effect on MoS₂ FETs by forming polar crystalline phases.

4 Conclusion

In summary, we proposed the dipole-induced n-type doping approach for MoS₂ FET with the P(VDF-TrFE) polymer dopant. The n-type doping effects on MoS₂ FETs were varied with the phase transition of the top P(VDF-TrFE) region by controlling recrystallization conditions. The controllable doping states according to the phase transition of P(VDF-TrFE) layers were investigated with Raman and PL spectroscopies. The increased polar phases of P(VDF-TrFE) dopant layer provide the enhanced n-type operation to MoS₂ FET by reducing the contact resistance. At the optimized condition, the field effect mobility, and the on-current values of the P(VDF-TrFE) doped MoS₂ FET were achieved 34.4 cm² V⁻¹s⁻¹ and 21 μA, respectively. The n-type doping effects observed at 8 different MoS₂ FETs and operation in a harsh environment imply that the proposed doping methods stable and reliable. We hope that the proposed approach could be utilized for further applications in the field of 2D semiconductor-based transistors and related research fields.

Data availability statement

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

Author contributions

DHL, TP, and HY designed the conception of the article, analysis of data, and work on manuscript preparation. DHL and TP performed

experiments and collect the data. TJ, YJ, JP, NJ, and UW contribute to the interpretation of data and analysis of results. All authors listed have made a substantial, direct, and intellectual contribution to the work. DHL, and TP equally contribute this work.

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Conflict of interest

Authors TJ, YJ, JP, NJ, and UW were employed by the company Hyundai Motor Group.

The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Supplementary material

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fmats.2023.1139954/full#supplementary-material>

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