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Role of salts on the electrical performance of ceria-based electrolytes: An overview

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This work provides an overview on established achievements and debatable findings involving Ca, Gd or Sm-doped ceria-based electrolytes, using Li₂CO₃, LiNO₃ and Na₂CO₃ as sintering aid or as second phase. The performance of these materials is discussed considering the characteristics of the oxides and of the salts or derived second(ary) phases (e.g., alkali metal oxides and hydroxides, eutectic mixtures), extensively surveyed to identify influential parameters with respect to processing and electrical performance (e.g., melting and boiling points, thermal decomposition, hydrolysis). The analysis of published data highlights the possible contribution of additional charge carriers to the total conductivity, besides oxide-ion vacancies. Claimed bulk and grain boundary conductivity enhancements are deeply discussed, as well as advantages and limitations of impedance spectroscopy as characterization tool. Irrespective of controversial reasons, reports on unusual improvements of grain boundary conductivity sustain the possibility of advanced grain boundary engineering to enhance the performance of these materials.

KEYWORDS

ceria, alkali metal salt, solid electrolyte, bulk conductivity, grain boundary conductivity, impedance spectroscopy, microstructure

Introduction

Solid oxide fuel cells (SOFCs) are devices able to convert chemical into electrical energy *via* electrochemical reactions, where electrolyte and electrodes play a key role. Lower processing and operating temperatures are desirable targets to decrease cost and improve long term performance. Ceria-based ceramics are promising electrolyte materials due to their high ionic conductivity with respect to conventional zirconia-based electrolytes, enabling lower operating temperatures (Mogensen et al., 2000; Kharton et al., 2004). Drawbacks include the partial reduction of Ce⁴⁺ to Ce³⁺ under reducing conditions, also the high sintering temperatures (>1500°C) needed to achieve almost full densification (Inaba and Tagawa, 1996; Kharton et al., 2001).

Chemical routes, yielding reactive ceria-based powders, were exploited to lower sintering temperatures (Herle et al., 1997; Chinarro et al., 2007; Ding et al., 2008;

Moure et al., 2009; Cela et al., 2011; Teoh and Chiang, 2012). Large attention was also given to sintering aids, including transition metal oxides (Kleinlogel and Gauckler, 2000; Fagg et al., 2002; Avila-Paredes and Kim, 2006; Pérez-Coll et al., 2006; Pikalova et al., 2007; Gao et al., 2010; Ge et al., 2011; Xu et al., 2011; Villas-boas et al., 2014; Taub et al., 2015; Neuhaus et al., 2018a, 2018b; Santos et al., 2018), and alkali metal salts (Nicholas and De Jonghe, 2007; Li et al., 2012; Le et al., 2013; Zhu et al., 2014; Maheshwari and Wiemhöfer, 2016). Sintering temperatures around 1000°C are easily reached using these admixtures, providing high densifications (in the 95–99% range), and regular grain size distributions.

This work exploits alternative explanations to claimed conductivity enhancements of selected materials using salts as sintering aids or second phases. Bulk conductivity augmentation is hardly explained by conventional effects involving dopant concentration and/or mobility of ionic defects (Figueiredo and Marques, 2013). Thus, emphasis is dedicated to grain boundary effects, large in most reports. Own expertise in similar materials provides a sound background for this purpose (Grilo et al., 2019, 2020; 2021a; 2021b).

Selected examples used in this work might be described as typical of debatable findings, deviating from a conventional review where extensive sets of data on similar oxides are collected and discussed. Here, emphasis is on data on distinct salt-derived phases that might coexist under specific conditions, yielding unusual performance effects. Proper consideration of second(ary) phase effects is almost absent from previous analyses in the literature. While new hypotheses are introduced to explain some findings, there is no intention to disregard previous explanations. The complexity of studied systems cannot be ignored. The presence of minor salt-derived phases is a clear source of uncertainty. Molten phases easily originate quasi steady state rather than true equilibrium conditions. Even in solid state systems, due to slow equilibration kinetics, exact phase boundaries are often uncertain (e.g., see (Duwez et al., 1952; Grain, 1967; Stubican, 1986; Duran et al., 1991)).

This work has a clear focus on ceria-based solid electrolytes, in contrast to composites involving one oxide and one molten phase (e.g., *see* (Benamira et al., 2007; Ferreira et al., 2011b; Rondão et al., 2013)). The borderline between these families of materials is established by the fully dominant solid-state condition. In oxide+molten salt composites, the molten phase corresponds to more than 10 vol%, and is conceived as provider of specific functionalities, like additional ionic pathway in fuel cells or CO_2 gas separation membranes (Li et al., 2009; Patrício et al., 2014).

The electrical characterization of single-phase solid electrolytes is often based on impedance spectroscopy measurements. This technique provides information on contributions within the bulk grains, across grain boundaries and even on the electrolyte/electrode interfacial processes of a TABLE 1 List of acronyms and selected physical properties of singlephase electrolytes (Duran et al., 1990; Zha et al., 2003; Maheshwari and Wiemhöfer, 2016).

Acronym	d_{th}^{a}	МР ^ь (°С)
5CDC	6.96	2445
GDC	7.23	
SDC	7.13	
	Acronym 5CDC GDC SDC	Acronym d _{th} ^a 5CDC 6.96 GDC 7.23 SDC 7.13

^ad_{th}-theoretical density.

^bMP, melting point for pure CeO₂.

polycrystalline ceramic electrolyte (Van Dijk and Burggraaf, 1981; Maier, 1986; Guo et al., 2002, 2003).

The presence of a second phase, even in minor amounts, can be the source of disturbance of impedance arcs. Porous electrolytes are one of the best studied examples, showing that pore size distribution and volume fraction influence the so-called bulk and grain boundary arcs (Roberts and Schwartz, 1985; Kleitz and Steil, 1997; Pérez-Coll et al., 2010; Marrero-López et al., 2014; El Khal et al., 2017; Cordier et al., 2019). For dispersed insulating phases within a conducting phase, the effective medium theory offers proper background for data analysis (Maxwell, 1904; Kleitz and Steil, 1997). The often-disregarded complexity of analysis in electrolytes with two potential conducting phases is hereby addressed, including a survey of impedance spectra characteristics and corresponding equivalent circuits.

Materials overview

Oxides

Materials considered in this work are based on ceria, with Ca, Gd or Sm as dopants. The first composition is interesting with respect to (low) cost, the remaining compositions are electrolytes with premium ionic conductivity. This selection is enough to cover a variety of situations where interaction between oxide and salt (or derived phases) might reveal distinct characteristics. Table 1 shows the acronyms and basic properties of these high-density ceramic refractory materials.

Salts

Salts are often used only as sintering aids (Nicholas and De Jonghe, 2007), as those listed in Table 2. Carbonates or nitrates, the most common admixtures, might end up as oxides or hydroxides, after thermal processing or exposure to specific environmental species. Comparison with data from Table 1 shows the sharp difference in density and melting point between these compounds and previous oxides.

Compound	Acronym	d _{th} ^a	MP ^b (°C)	BP ^b (° C)	Ref	
Na ₂ CO ₃	NaC	2.53	851	1600 (d*)	Haynes, (2005)	
NaOH	_	2.13	318	1390 (d*)		
Na ₂ O	_	2.27	1132	1950		
LiNO ₃	LiN	2.38	264	600 (d*)		
Li ₂ CO ₃	LiC	2.11	723	1310 (d*)		
LiOH	—	1.46	450	924 (d*)		
Li ₂ O	—	2.01	1438	2600	Ortman and Larsen, (1983)	
Mixture		Acronym		ET ^c (°C)	Ref	
LiNO ₃ + 40 mol% LiOH		LiN+LiOH		180	Janz and Tomkins, (1983)	
Na ₂ CO ₃ + 90 wt% NaOH	ł	NaC+NaOH		283		
Li ₂ CO ₃ + 84 mol% LiOH		LiC+LiOH		433	Bale et al. (2016)	
Na ₂ CO ₃ + 52 mol% Li ₂ CO ₃		NLC		498	Jiang et al. (2017)	
Li ₂ CO ₃ + 13 wt% Li ₂ O		LiC+LiO	LiC+LiO		Kaplan et al. (2011)	

TABLE 2 List of acronyms and selected physical properties of salts, derived compounds and eutectic mixtures.

^ad_{th}, theoretical density.

 $^{\mathrm{b}}\mathrm{MP}$ and BP, melting and boiling points, respectively.

 $^{\rm c}{\rm ET},$ eutectic temperature; d*, decomposes, see Table 4.

TABLE 3 List of ceria-based solid electrolytes (from top to bottom: Maheshwari and Wiemhöfer, 2016; Zhu et al., 2014; Li et al., 2012; Grilo et al., 2020).

Material		Firing stages and temperatures (°C)		Grain size (μm)	Secondary phase			
						Nominal composition		vol%
Oxide	Salt	Acronym	1st	2nd		Mass	vol%	
5CDC	NaC	5CCC4:1	400	600 (20 h)	<1	20 wt% NaC	40.7	≤40.7 (NaC)
GDC	LiN	5LiGDC	600	1000	<1	5 cat% (Li ⁺)	6.0	$\leq 1.6^{a}$ (Li ₂ O)
SDC	LiC	SDC15Li	700	900	~5	15 cat% (Li ⁺)	11.3	\leq 5.1 ^a (Li ₂ O)
GDC	_	GDC	1500		2.6	_	_	nd
	NLC^{b}	GDC-5M ^c	1100		1.9	5 mol% NLC	7.8	nd
		GDC-5S ^c			2.6			

^aAssuming only thermal decomposition of LiN and LiC, and no porosity.

^bNLC, NaC+ 52 mol% LiC.

^cM, mechanical (ceramic) route.

^dS, chemical synthesis route; nd, not detected.

Oxide electrolytes with salt addition

Table 3 shows the exact compositions and firing conditions experienced by electrolytes considered in this work, grouped according to data sources rather than composition. The first three rows correspond to materials prepared by authors from distinct institutions. The last row includes three materials prepared in one single institution (own work). These materials will be used as reference.

Materials acronyms in Table 3 were partly introduced in Tables 1, 2. These follow closely those adopted in the source references. However, the range of combinations is wide, and a few

Reaction	T (°C)	Remarks	Ref
$Na_2CO_3 (l) \rightarrow Na_2O (s) + CO_2 (g)$	>1200	In Argon, weight loss w = $9.66 \times 10^2 \exp(-58600/RT)$ (g/cm ²)	Nakamura et al. (1980)
	>950	In Argon, very slow rate	Kim and Lee, (2001)
	650-700	Sharp weight loss (TGA)	Maheshwari and Wiemhöfer, (2016)
$2\rm NaOH \rightarrow \rm Na_2O+\rm H_2O$	475-750	Decomposition rate r = 50.5 × exp(-30230/T) (gh ⁻¹), vacuum = 13.3 Pa	Yurkinskii et al. (2005)
LiNO ₃ →Li ₂ O (Li ₂ O ₂)+NO _x	(500 _{start})-700 _{peak}	Based on TGA-DTA data	Sweeney, (1975)
	(485_{start}) -650 _{peak}	Based on TGA-DSC data (air: 40 vol%, He:60 vol%)	Ruiz et al. (2010)
$Li_2CO_3 \rightarrow Li_2O+CO_2$	>800	Under Argon	Kim and Lee, (2001)
	>710	Under Argon, complete at 900°C	Li et al. (2012)
	825-975	$P_{CO_2} = 2.33 \times 10^9 \times \frac{1 - X_{U_2O}}{X_{U_2O}} exp(-32500/T)$	Kaplan et al. (2011)
2LiOH (s) \rightarrow Li ₂ O (s) + H ₂ O (g)	290-405	Fraction of released water x = $1-k_1 \times exp(-k_2.t)$, vacuum = 4×10^{-4} Pa	Kudo, (1979)

TABLE 4 Selected reactions involving alkali metal-based materials.

examples are used to assist in the capture of the meaning of acronyms. 5CCC4:1 means a Ceria (5CDC) - Carbonate (NaC) -Composite with a 4:1 weight ratio. 5LiGDC means GDC with 5 cat% Li⁺. GDC-5S corresponds to GDC with 5 mol% NLC, prepared by a chemical synthesis (S) route. The extension (M) in GDC-5M designates a ceramic route, starting with mechanical milling. SDC15Li is SDC with 15 cat% Li⁺. When needed, the material acronym also includes a reference to the highest firing temperature (e.g., SDC15Li-900 means SDC15Li sintered at 900°C). Table 3 also provides information on grain size and vol% of distinct phases, with impact on electrical performance.

The first obvious conclusion from these data is that processing temperatures are below salt melting (5CCC4:1), exceed the salt melting and boiling points (5LiGDC), or surpass only the melting point of precursor salts (SDC15Li-900).

Salt phase changes

Salts are also likely to suffer total or partial decomposition during thermal processing. Table 4 lists typical examples of known reactions, conversion levels and conditions. NaC (used in 5CCC4:1) is expected to be preserved after firing at 600°C. Melting is only at 851°C, and thermal decomposition is significant only above the sintering temperature (600°C). LiN (used in 5LiGDC) might easily boil but also partly decompose, being converted into Li₂O. LiC (used in SDC15Li) is expected to be partly converted to Li₂O. NLC (used in GDC-5S and GDC-5M), the eutectic mixture of NaC and LiC, melts at around 500°C. The stability of this phase is high at moderate temperatures (e.g., 650°C), but volatilization and even decomposition to oxides are likely at higher temperatures. These comments explain the values listed in the last column of Table 3, on predictable higher limits of oxides, from thermal decomposition of salts.

Standard "lab working conditions," involving ambient air during measurements (with moisture and CO_2), and unknown/ diverse storage conditions of samples, are additional sources of considerable uncertainty. As example, an investigation on the hydrolysis of lithium oxide showed that the starting material (nominally pure Li₂O) consists of particles with a Li₂O core covered by a LiOH protective shell (Weber et al., 2018). Furthermore, the loss of water from alkali hydroxides is not fast or easy, requiring high temperatures and long periods of time (*see* Table 4).

5CDC with NaC should have a considerable volume fraction of this phase (30–40 vol%), the most stable with respect to the adopted thermal processing (600°C). SDC with LiC should have a moderate volume fraction of Li₂O and/or LiOH (up to 5.1 vol%). The oxide is the most stable phase with respect to the thermal processing steps (up to 900°C), but the hydroxide should form easily from exposure to moisture. GDC with LiN should have a residual volume fraction of Li₂O (up to 1.6 vol%) and/or LiOH, since LiN starts decomposing before boiling and the oxide and hydroxide are the most stable after the adopted processing temperature (1000°C). GDC-5S and GDC-5M might have residual amounts of alkali-metal carbonates, oxides and/or hydroxides.

The likely presence of alkali oxides and/or hydroxides in these materials is source of additional phase interactions. Table 2 lists eutectic compositions involving these compounds. The corresponding eutectic temperatures, starting below 200°C, are commonly used in electrical measurements. Partial decomposition of carbonates and nitrates, besides melting and boiling, were considered in the literature only as part of the TABLE 5 Literature data on bulk (σ_b), grain boundary (σ_{gb}) and total conductivity (σ_t) of electrolytes without and with salt additions (Li et al., 2012; Chen et al., 2014; Maheshwari and Wiemhöfer, 2016). Conductivity ratios (σ_{ad} - material with admixture; σ_f - admixture free) highlight conductivity enhancements (see Tables 1–3 for acronyms).

Material ^a	T (°C)	σ _b (S/cm)	σ_{ad}/σ_{f}	σ _{gb} (S/cm)	σ_{ad}/σ_f	T (°C)	σ _t (S/cm)	σ_{ad}/σ_{f}
5CDC	350	7.4×10^{-5}	_	3.5×10^{-6}	_	350	3.3×10^{-6}	19
5CCC4:1		—		—			$6.4 imes 10^{-5}$	
GDC	450	5.0×10^{-3}	3.0	7.0×10^{-3}	143	550	1.0×10^{-2}	4.0
5LiGDC		1.5×10^{-2}		1.0			4.0×10^{-2}	
SDC	370	7.0×10^{-4}	1.7	$3.5 imes 10^{-4}$	1.7	370	$2.3 imes 10^{-4}$	1.7
SDC15Li		1.2×10^{-3}		5.8×10^{-4}			3.9×10^{-4}	

^aPublished data (figures and tables) treated to provide a reasonable basis for comparison. Minor inaccuracy in values estimated from curves in plots might be present (few %).



processing route (Stern and Weise, 1969; Nakamura et al., 1980; Kim and Lee, 2001; Singh and Singh, 2007; Ruiz et al., 2010; Olivares, 2012; Lee et al., 2013). The presence of hydroxides, confirmed in different situations (Lapa et al., 2010; Ferreira et al., 2011a; Xing et al., 2015), is mostly disregarded.

As summary, in ceria-based solid electrolytes with salts, the nominal and actual phase compositions can be quite distinct because of the effects of multiple processing parameters (e.g., gas phase composition, weight losses during firing, weight gain during storage or measurements). Since the kinetics of some reactions is also slow, studied materials include metastable phases, with their proportions changing during measurements under diverse conditions. As example, conductivity hysteresis was found only during a first heating/cooling cycle in Na-doped ceria (Pearce and Thangadurai, 2009).



The presence of Li and/or amorphous phases in these electrolytes decreases the efficacy of energy-dispersive X-ray spectroscopy (EDS) to properly map chemical compositions, or X-ray diffraction (XRD) for complete phase identification. These limitations should be understood only as background information on potential scenarios behind published reports.

Electrical performance

Materials overview

Table 5 includes data obtained with Na2CO3, Li2CO3 and LiNO3 salt additions, and ceria-based electrolytes (CDC, GDC

and SDC) (Li et al., 2012; Zhu et al., 2014; Maheshwari and Wiemhöfer, 2016).

These data show impressive, claimed conductivity enhancements for salt containing materials with respect to pure (salt free) oxides, reaching factors of 3 in bulk conductivity, 143 in grain boundary conductivity, and 19 in total conductivity. While all microstructures are slightly distinct, grain size is between 1 and 5 μ m (Table 3), and densifications exceed 95%. Sintering temperature, duration, and sintering aids, originate distinct grain size and densification effects, preventing fully comparable sets of sintering parameters and final microstructures.

The microstructural characteristics of these materials can be obtained from a combination of information on grain size (in the order of 1 μ m for 5CCC4:1 and 5LiGDC and 5 μ m for SDC15Li (Li et al., 2012; Chen et al., 2014; Maheshwari and Wiemhöfer, 2016)) and of volume percentage of the likely residual second phase, after sintering (NaC or Li₂O, *see* Table 3). Assuming a conventional brick-layer model (*see* Section 5 for details), volume fractions can be used to relate grain size (*D*) and grain boundary thickness (δ_{gb}), as depicted in Figure 1.

Typical values of δ_{gb} in ionic conductors range between a few nm and a few dozens of nm (Kingery, 1974b, 1974a; Christie and Van Berkel, 1996). Increasing the grain boundary thickness, the volume fraction of this region easily slides to values in excess of 10 vol%, which is already composite territory. With more than 30 vol% of a second phase, percolation is almost inevitable (Jiang and Wagner, 1995). Materials under consideration here include thin grain boundaries (5LiGDC, $\delta_{gb} \approx 5$ nm), moderately thick grain boundaries (SDC15Li, $\delta_{gb} \approx 90$ nm), and others clearly seating in the composites category (5CCC4:1), with a "nominal grain boundary thickness" around 190 nm.

Ca-doped ceria (CDC)

For CDC-based electrolytes, NaC is present in large quantity (about 40 vol%), percolating throughout the material. Figure 2 shows the electrical conductivity of 5CDC, 5CCC4:1, and likely secondary phases. Irrespective of this overpopulated plot, naked eye analysis immediately shows that several secondary phases exceed the conductivity of 5CDC (\Box) or even 5CCC4:1 (Δ).

With respect to salts, pure NaC (in CO_2 and air) and Sr or Ca-doped NaC are all considered. The trends are for higher conductivity of NaC in air with respect to CO_2 (*see* NaC curves), also higher conductivity of Ca-doped NaC with respect to pure NaC (*see* CaDNaC curve in Figure 2). 5CCC4:1 shows conductivity levels within the range of values reported for pure NaC in air or CaDNaC in CO_2 .

The consideration of Ca or Sr-doped NaC deserves a short comment. Doped NaC has a conductivity higher than pure NaC. This is a consequence of the formation of mobile and negatively charged Na⁺ vacancies (V_{Na}' , in Kröger-Vink notation) to



compensate the presence of positively charged M^{2+} ions (M = Ca^{2+} or Sr^{2+}) in normal Na⁺ lattice positions (M_{Na}^{\bullet}). The simplified electroneutrality condition can be expressed as:

$$[\mathbf{M}_{\mathrm{Na}}^{\bullet}] = \left| \mathbf{V}_{\mathrm{Na}}^{\prime} \right| \tag{1}$$

The Na⁺ vacancies are the dominant charge carriers in Cadoped NaC, and the effect of Ca-doping was demonstrated measuring the electrical conductivity of materials with dopant concentrations increasing up to 0.093 mol% (Figure 2 (Cerisier and Roux, 1978)). The 5CDC electrolyte could easily be the source of minor amounts of Ca as dopant of NaC, where only 0.1 mol% is a meaningful level.

The "fictional" curve suggested for CaDNaC in air exceeds the values in CO_2 by one order of magnitude, taking into consideration that these relative magnitudes were observed for pure NaC in air and in CO_2 . To corroborate this approach, data for SrCO₃ (1 mol%)-doped NaC, in air, is also plotted (Guth et al., 1987). These results are impressively close or even exceed those observed for 5CCC4:1. Overall, the high-volume fraction (about 40 vol%) of NaC in these composites would justify the consideration of the exact role of this phase in the composite material.

If, irrespective of cautious thermal processing, some NaC partly decomposed to Na_2O , and converted into NaOH, the presence of liquid phase (acronym extension "m" for molten) could be expected at temperatures as low as 283°C. Alkali salts or hydroxides show extremely high values of electrical conductivity above melting, in the range 0.1–1 Scm⁻¹ (Li et al., 2007). The electrical conductivity of the NaC+NaOH eutectic composition is also plotted in Figure 2 to highlight this fact. The presence of a



molten phase (NaC+NaOH (m)), even a few vol%, could explain or contribute to the observed high conductivity. Molten phases easily spread and percolate in between grains.

If present in significant amounts, a solid to liquid transition often appears as a sharp change in conductivity with temperature (elongated z-shape in Arrhenius plots). The observed conductivity trend for 5CCC4:1 indeed shows a slight upwards bending close to the NaC+NaOH eutectic temperature (NaC+NaOH et), highlighted in Figure 2 (vertical dashed line). This is the opposite of the downwards bending observed at high temperature in ceramic oxide-ion conductors. This is not a demonstration of the presence of a molten phase but justifies due consideration of this possibility.

Gd-doped ceria (GDC)

LiN was added in a small quantity (5 cat% Li⁺) to GDC, sintered at a temperature (1000°C) where melting, boiling and partial salt thermal decomposition are expected. If either LiN, or Li₂O, or LiOH were present, molten phases could be found starting at temperatures as low as 264°C (melting of LiN), 480°C (melting of LiOH) or even 180°C (LiN+LiOH eutectic temperature). However, the total equivalent amount of Li_2O is below 1.6 vol% in the studied electrolyte. From a microstructural point of view, the material has the features of a polycrystalline ceramic with a thin grain boundary.

Figure 3 shows multiple sets of data on the electrical conductivity of distinct materials. Unlike the previous case of NaC+CDC, the reported solid LiN conductivity values are all below the conductivity of 5LiGDC1000. Pure LiN has a phase change with some hysteresis on heating/cooling, with the breaking region shown within a dashed circle (Figure 3), separating data for the two solid LiN phases (s1 and s2, (Claes and Glibert, 1985; Yamashita et al., 1995; Skobelev et al., 1996; Zhu et al., 2014)).

The electrical conductivity of molten LiN and molten LiOH (acronyms with extension "m") highlights the enormous impact that minor amounts of these phases would have on conductivity. Again, a small upwards bending in the conductivity curve of 5LiGDC1000 (*see* arrow in Figure 3) coincides with the melting point (mp) of LiN (signaled with a vertical dashed line). Accordingly, the potential presence of molten LiN could be considered as a possible source of unusual conductivity contributions.





Recently, the addition of NaC+LiC (NLC) to GDC, was assessed in own work, to decrease the GDC sintering temperature to 1100°C (Grilo et al., 2020). This is above the eutectic point of NLC (500°C). We recall that materials with a small amount of NLC (5 mol%) were prepared using a ceramic route starting with mechanical milling (M) or a wet chemical synthesis route (S). There was no evidence for localized spots of Na in any sample, and lattice parameters also did not show evidence for solubility of alkali-metal ions in the ceria lattice. Figure 4 depicts low temperature impedance spectra for these materials in air. An introduction to this type of plots is skipped here since these are rather standard in the literature on solid electrolytes.

Figure 4 shows that the same sintering admixture might affect the performance of these materials in distinct manners, depending on the exact processing route. We invoke these materials here since we can fully compare couples of materials versus performance (e.g., "S" versus "M" routes, or pure GDC versus "M" route). All materials were prepared using the same chemicals and shaping routes, and final grain size (around $2{-}2.5\,\mu m)$ and densifications (close to 100%) are comparable. Fluctuations often found in data obtained in distinct labs, are ruled out here.

The impedance spectra in Figure 4 include two arcs at high and intermediate frequency usually attributed to the contributions of bulk (higher frequency range) and grain boundary (intermediate frequency range). An incipient third arc, in the low frequency range, indicates the contribution of electrode processes. Figure 4B shows that there is no meaningful bulk impedance effect, only a slight data scatter. However, when obtained by wet chemistry (GDC-5S), materials present a much lower grain boundary resistivity (Figure 4A). We could add extra evidence to show that materials with high densification, show slight changes in bulk conductivity, but grain boundary conductivity is highly sensitive to contact with salts (e.g., *see* (Grilo et al., 2021b)).

Sm-doped ceria (SDC)

LiC was added in considerable amounts (15 cat% Li⁺) to SDC. Sintering of these samples was performed at temperatures where melting and even thermal decomposition of LiC would be expected. The upper value of about 5 vol% of a secondary phase after sintering (as Li₂O) shows that this material is in the transition from a typical polycrystalline ceramic to a composite. If LiC, Li₂O, or LiOH, were present, molten phases could be found starting at temperatures as low as 723°C (melting of LiC), 480°C (melting of LiOH) or 705°C (LiC+ Li₂O eutectic temperature). Figure 5 shows multiple sets of data on the electrical conductivity of SDC and likely secondary phases. The reported pure solid salt (LiC) conductivity is anisotropic (parallel or perpendicular to plane (002)), with distinct values depending on direction. In all cases, the performance of solid LiC is below the performance of SDC15Li-900.

The electrical conductivity of molten LiOH is also shown in Figure 5 to highlight again the enormous impact that minor amounts of a molten phase may have on conductivity. The most controversial aspect in this case is related to the claimed temperature dependence of bulk and grain boundary conductivity, which deviates considerably from common trends for SDC1400 and SDC15Li-900. Activation energies for electrical conductivity in clear excess of 1 eV are unusual in ceria-based electrolytes (*see* SDC-Y and SDC trend lines, Yamashita et al., 1995). Deeper analysis of source data might explain these findings.

Impedance spectra reported at temperatures below 400°C show two arcs and a spike at low frequency, schematically reproduced in Figure 5, together with the adopted equivalent circuit. The lower frequency arc is interpreted as typical of grain boundary performance and the higher frequency arc is related to the bulk performance (Li et al., 2012). The high frequency intercept of this arc with the Z' axis (real part of impedance)

is distant from the origin. To account for this situation the authors introduced a resistor ($R_1 > 100$ ohm) in series with bulk ($R_2||CPE_1$) and grain boundary ($R_3||CPE_2$) standard parallel R||CPE circuits (resistance||constant phase element). However, no physical meaning was provided for R_1 . Standard setups (external electrical cables and cell wiring/electrodes) hardly show a resistance compatible with the reported value of R_1 . The analysis of these data is halted here. Additional comments on impedance spectroscopy measurements are provided in the last section of this work.

Considerations on conductivity enhancement

Bulk effects

Ceria-based electrolytes are oxide-ion conductors where the prevailing defects are acceptor dopants (e.g., Ca, Gd or Sm, negative defects as substitutional dopants) and oxide-ion vacancies (positive defects). Enhanced bulk conductivity must rely on enhanced concentration and/or mobility of oxide-ion vacancies.

The presence of alkali metal ions (Na⁺, Li⁺) in solid solution could imply an enhanced concentration of oxideion vacancies to balance the total concentration of dopants. The dimensions of these ions for 8-fold coordination (1.18, 0.92 and 0.97 Å for Na⁺, Li⁺ and Ce⁴⁺, respectively (Shannon, 1976)), seem consistent with the possibility of partial replacement. The solubility of up to 7 cat% Na⁺ in ceria was assumed from XRD patterns (Pearce and Thangadurai, 2009), but without consensus (Grilo et al., 2020). The solubility of Li+ in ceria also seems controversial. This possibility was raised due to variable lattice parameter of GDC after distinct firing conditions (Zhu et al., 2014). However, in the case of SDC, the XRD patterns showed no changes and this was interpreted as sign of no solubility of Li⁺ in the lattice (Li et al., 2012). In any case, Li was lost from the SDC15Li samples with increasingly high sintering temperatures, with only 1.5 cat% remaining in SDCLi15 fired at 800°C (Li et al., 2012).

The plausible simplified electroneutrality conditions, assuming that alkali metal ions are dissolved in ceria, should be:

$$2\left[Ca_{C_{e}}^{"}\right] + 3\left[Na_{C_{e}}^{"'}\right] = 2\left[V_{o}^{\bullet}\right]$$
(2)

for Na^+

as additional dopant in calcia-doped ceria, and

$$\left[\mathrm{RE}_{\mathrm{Ce}}^{'}\right] + 3\left[\mathrm{Li}_{\mathrm{Ce}}^{'''}\right] = 2\left[\mathrm{V}_{\mathrm{o}}^{\bullet\bullet}\right] \tag{3}$$

for Li⁺ as additional dopant in rare earth (RE)-doped ceria, with RE = Gd^{3+} or Sm^{3+} . The total concentration of oxide-ion vacancies ([V₀^{••}]) should depend on the concentrations of

alkali metal dopants ([Na_{Ce}^{''}] or [Li_{Ce}^{''}]) besides the alkalineearth or rare-earth (RE) concentrations ([Ca_{Ce}^{''}] and [RE_{Ce}[']], respectively).

Dissolved Na⁺ would imply an enhanced concentration of oxide-ion vacancies but impedance spectroscopy showed an extremely poor electrical conductivity for 7 cat% Na⁺, about two orders of magnitude lower than reported for pure SDC (Pearce and Thangadurai, 2009). From simple charge balance effects, 7 cat% Na⁺ as dopant should be identical to 21 cat% Gd³⁺ or Sm³⁺. A conductivity two orders of magnitude lower with respect to SDC must imply vestigial solubility of Na⁺ or deep trapping of oxide-ion vacancies due to strong defect association. The latter effect would be expected for heavily (triply) charged acceptor dopant defects (Na_{Ce}^{'''}) with strong tendency to associate with positively (double) charged oxideion vacancies. Bulk oxide-ion conductivity enhancements, without substantial increase in the activation energy of ionic conduction, are inconsistent with this picture.

Following a similar reasoning, in the case of Li⁺ as dopant, the reported low solubility (negligible?) cannot justify substantially higher concentration of oxide-ion vacancies. A negative impact on the mobility of oxide-ion vacancies should also be considered since dopants smaller than the host cation originate local lattice strain effects (Mogensen et al., 2000; Skinner and Kilner, 2003). There is no apparent rationale for a reported bulk conductivity enhancement (up to three times higher, *see* Table 5) when salts with Li are used as sintering admixtures. Both, concentration and mobility of oxide-ion vacancies, are unable to explain this claim. Enhanced bulk conductivity effects exceed what might be considered the state-of-the-art knowledge on the defect chemistry of these systems.

Grain boundary effects

Unlike for bulk conductivity, grain boundary effects may easily surpass one order of magnitude (Table 5). However, significant grain boundary improvements are claimed in situations ranging from 1 to about 40 vol% of secondary phases. This diversity of situations justifies distinct attempts to explain grain boundary effects.

Early work in the field of ceramic electrolytes identified clear correlations between total conductivity and temperature. In Arrhenius-type plots of $\ln(\sigma_t)$ versus 1/T, the low temperature activation energy often exceeds the high temperature activation energy, and the effect is enhanced in materials with low purity or small grain size. With the adoption of impedance spectroscopy as characterization tool, two distinct contributions could be clearly identified, corresponding to the bulk and grain boundary (Bauerle, 1969). Defect association is also noticed in the low temperature range (Mogensen et al., 2000; Skinner and Kilner, 2003), but this will be ignored here.



Grain boundary schemes: (A) traditional description, including crystallographic orientation mismatch between grains, impurities in the grain boundary region and insulating particles in triple contact points (Bauerle, 1969); (B) electrical grain boundary, with core, space charge layers and concentration profiles of acceptor dopant ($[A_m']$) and oxide-ion vacancies ($[V_O^{\bullet^*}]$) (Guo et al., 2003; Guo and Waser, 2006). In the lower insert, levelling of concentration profiles due to an extrinsic species acting as local scavenger (Zhu et al., 2014); (C) combined oxide-ion transport mechanisms, where several species might assist O^{2-} transport in between grains or along the salt phase; (D) single (salt) phase and dual phase O^{2-} pathways (Maheshwari and Wiemhöfer, 2016). See text for details.

The classical description used to explain grain boundary effects is shown as scheme in Figure 6A). Grain boundaries are understood as a crystallographic feature, where crystal misorientation and presence of impurities implies excess energy requirements for ionic migration, and interfacial polarization effects.

Secondary phases, namely impurities, might be wetting and covering the entire grain boundary region (blocking ion transport), or might be non-wetting, forming isolated particles in triple joints between multiple grains. The discovery of grain boundary "scavenging" agents, able to combine chemically with impurities/glassy phases, forming isolated particles in joints between grains, like addition of alumina to zirconia-based electrolytes, is a good example of this process. This description, even if more than 40 years old, is still used in many cases to explain grain boundary conductivity improvements (Beekmans and Heyne, 1976; Butler and Drennan, 1982; Guo et al., 1995; Zhang et al., 2004; Sudarsan and Krishnamoorthy, 2018).

Several authors introduced a distinct grain boundary concept, often named as electrical grain boundary (Guo et al., 2003; Guo and Waser, 2006). In acceptor doped oxide-ion conductors, the defect concentrations in the grain boundary region deviate from bulk values due to segregation of dopant. Their negative charge equalizes a positively charged grain boundary core, which in turn implies the depletion of oxide-ion vacancies in the same region. The overall consequence is the formation of space charge layers around the grain boundary core, decreasing locally the oxide-ion conductivity due to an enhanced electrical potential barrier and depletion of mobile defects. This model is able to explain why ions are partly blocked in high purity materials. The physical scheme used to describe these situations is shown in Figure 6B.

The presence of salts in considerable amounts introduced the need to consider the role of this phase in such materials. As first remark, the concept of grain boundary conductivity seems inaccurate in this case. In the presence of large concentrations of both phases, the oxide grains are fully covered by the salt phase. This means that the often measured "grain boundary" conductivity corresponds to a designation "freely" adopted from the field of polycrystalline solid electrolytes, where impedance spectra can be deconvoluted into bulk, grain boundary and electrode contributions, as introduced earlier.

Having in mind this remark, an enhanced "grain boundary" performance was explained in 5CCC4:1 composites by the presence of a highly conductive phase in the inter-grain



region, where several species provide net transport of oxide ions (Maheshwari and Wiemhöfer, 2016). The firstly suggested cooperating mechanism involved the net transport of O_2^- (superoxide), *via* NaO₂* (intermediate species) and Na₂O₂ (peroxide), as drafted in Figures 6C,D.

The idea was revised after consideration of the possible role of CO₂, involving several elementary reactions:

$$O^{2-} + 2Na^+ \leftrightarrow Na_2O$$
 (dissolved in Na_2CO_3) (4)

and

$$Na_2O + CO_2 \leftrightarrow Na_2CO_3$$
 (5)

inside Na_2CO_3 , where CO_2 moves *via* pores of Na_2CO_3 . Other mechanisms can be envisaged, involving water molecules and OH^- ions:

$$O^{2-} + H_2O \rightarrow 2OH^-$$
 (oxide/molten hydroxide interface) (6)

$$2OH^{-} \rightarrow O^{2-} + H_2O(\text{molten hydroxide/oxide interface})$$
 (7)

Water vapor is present in ambient air and water molecules are known to easily combine with Na₂O or Li₂O. Fast transport of OH⁻ and fast diffusion of water molecules is expected in molten phases. In fact, the diffusion coefficient of water molecules in molten hydroxides (or salts) is one order of magnitude higher than for most ions (Janz and Bansa, 1982). Also, water molecules are smaller than CO₂ molecules and should diffuse faster (Bergmann et al., 2007; Talesh et al., 2010). In this case, no porosity is required, unlike when CO₂ is considered.

For materials with minor concentrations of sintering aids, a distinct line of reasoning was introduced, invoking modified electrical characteristics of the grain boundary region, with leveling of acceptor dopant and oxide-ion vacancy concentrations nearby the grain boundary core, (Figure 6B, lower insert). This possibility relies on experimental evidence of changes in acceptor dopant profiles close to the grain boundary core (Zhu et al., 2014). This might be named

advanced grain boundary engineering. Own experiments, designed to confine the interaction between oxide and salt exclusively to the grain boundary region of sintered GDC ceramics, found consistent grain boundary conductivity improvements (Grilo et al., 2021b).

Figure 7 shows Arrhenius-type plots in the 200–300°C temperature range, where deconvolution of impedance spectra into bulk and grain boundary arcs is quite accurate, for total (σ_t), bulk (σ_b), and grain boundary conductivity (σ_{gb}) of pure GDC sintered at 1500°C/4 h, and GDC+NLC ceramics prepared using (M) and (S) processing routes. The corresponding activation energies (E_a) are depicted in the σ_b and σ_{gb} plots.

The total conductivity of all materials is within a small range of values (Figure 7A). There is no evidence for any impact of molten phases as possible contributors to the total conductivity. Bulk conductivities of all samples showed in Figure 7B are nearly identical, confirming that bulk transport is poorly sensitive to the effect of NLC additions. On the contrary, σ_{gb} values are somewhat distinct, with samples prepared using chemical synthesis showing higher conductivity, even higher than observed for standard GDC (Figure 7C).

For materials prepared using chemical synthesis, dense ceramics were obtained at 1100 °C. The presence of NLC changes the local chemical environment. At 1100 °C, migration of Gd^{3+} cations to the grain boundaries is restrained, unlike in pure GDC on cooling from high sintering temperatures. With lower concentration of acceptor cations, the grain boundary potential barrier decreases, and the oxide-ion vacancy concentration profile in the space charge layers are smoother (Lei et al., 2002; Guo and Waser, 2006). A specific mechanism might be considered, where Gd^{3+} cations combine with Li⁺ to form GdLiO₂, scavenging selectively the concentration of the acceptor dopant. The chemical reaction is known to occur at moderate temperatures (Yamauchi et al., 2008), but there is no report yet on a conclusive observation of GdLiO₂ in GDC+NLC grain boundaries.



In the case of materials prepared using mechanical milling, the effect seems the opposite, with $\sigma_{\rm gb}$ values lower than for all other materials. While the chemical synthesis route is expected to originate homogeneous materials at ionic level during wet chemistry steps, in the ceramic route, with low processing temperatures, heterogeneities are more likely. Any residual salt particles, with dimensions below the analytical capability of SEM/EDS, may act as potential oxide-ion blocking phases present in the grain boundary region (Cho et al., 2007, 2008; Lin et al., 2015).

The wide range of effects previously introduced is able to explain reports from enhanced to decreasing grain boundary conductivity. Previous comments also indicate that likely phenomena are local, within the few nm distance with respect to the grain boundary core. Only characterization tools with high spatial and chemical resolution (for light elements) can disclose the exact grain boundary characteristics. Such reports are scarce, and most analyses are based on indirect signs. In any case, experimental evidence opens the possibility of advanced grain boundary engineering solutions to circumvent drawbacks observed in conventional ceramics.

Impedance spectroscopy

Typical impedance spectra obtained at low temperature for solid electrolytes consist of two arcs, associated to bulk and grain boundary contributions (*see* examples in Figure 4). The grain boundary may represent a secondary phase distributed along intergranular regions, namely low melting point phases. The reported tendencies of the bulk and intergranular contributions are based on fitting of impedance spectroscopy data using a common equivalent circuit with a series association of resistors to account for the bulk (R_b) and grain boundary ($R_{gb,s}$) resistances, each having constant phase elements (CPE_b and CPE_{gb,s}) in parallel, as schematized by the bottom branch in Figure 8A. For simplicity, CPEs are replaced by simple capacitors (C_b , C_{gb}).

The so-called brick-layer model describes a material with cross-sectional area A and thickness L, composed of cubic grains with average grain size *D*, surrounded by grain boundaries of thickness δ_{gb} (Figure 8B) (Guo and Waser, 2006; Gomes et al., 2009). Considering only the bottom branch of the equivalent circuit to describe a polycrystalline ceramic with large average grain size, thus neglecting transport along parallel grain boundaries with huge electrical resistance due to obvious geometric reasons, the specific grain boundary conductivity (σ_{sp}) depends on the microscopic geometric factor of the grain boundaries $\delta_{gb}D^{-1}$. The latter is given by $\delta_{gb}D^{-1} = C_bC_{gb}^{-1}$, assuming equivalent dielectric constants for bulk and grain boundaries ($\varepsilon_b \approx \varepsilon_{gb}$). Therefore, δ_{gb} can be estimated from:

$$\delta_{\rm gb} = DC_{\rm b}C_{\rm gb}^{-1} \tag{8}$$

and σ_{sp} from:

$$\sigma_{\rm sp} = LC_{\rm b} \left(AR_{\rm gb}C_{\rm gb} \right)^{-1} \tag{9}$$

All values needed for Eqs 8, 9 are easily obtained from microstructures (*D*), impedance plots ($C_{\rm b}$ and $C_{\rm gb}$), and macroscopic sample dimensions. Often, authors try to correlate the reduction of the amplitude of the grain boundary arc with enhanced grain boundary conductivity, resulting from changes of the core potential or localized effects of sodium or lithium salts (*see* e.g. (Zhu et al., 2014)). This is a reasonable approach if the grain boundary thickness is small (typically of few nm), or the amount of second phases corresponds to that of an additive (1–2 vol% maximum), and not typical of composites (well above 10 vol%).

The possibility of a parallel branch to this series association was never considered in the analysis of impedance spectra of this kind of electrolytes. However, in a classical brick-layer model, grain boundaries offer two distinct pathways with respect to electrical transport. In between consecutive grains along the current direction the grain boundary is perpendicular to the ionic flow. In between sides of grains, the grain boundary offers a parallel pathway (Figure 8A). The usually assumed high grain boundary resistivity (with respect to the bulk grain) explains the tendency to discard the second branch of the equivalent circuit.

In the following discussion the full circuit with two parallel branches is considered (Figure 8A). The parallel, upper branch (circuit elements with subscript p) might correspond to a secondary phase or a parallel grain boundary. In the latter



FIGURE 9

Simulated impedance spectra based on the equivalent circuit of Figure 8A, where Z_s is the impedance of the series (bottom) circuit, Z_p is the impedance of the parallel (upper) branch and Z_t is the total impedance of the circuit. The spectra in Figures (A-C) were simulated with R_b, R_{qb,s} and $R_{gb,p}$ values given in the legend, and D=1 μ m and δ_{gb} =5 nm. Spectra in (D) obtained assuming a constraint on the electrical properties that scale with a variable δ_{qb} (see text for details).

case, the electrical properties of the perpendicular and parallel grain boundaries are interlinked (same conductivity and dielectric constant). In the former case they can be manipulated independently.

Other circuit elements besides these are neglected. This includes possible contact/electrical setup impedances and electrode impedances. A remark is needed to emphasize that if the upper branch involves the transport of the same ionic species the electrode impedance should be common to both branches (Figure 8A), case i). However, if a distinct ionic species is involved, a distinct electrode branch should thus be included (Figure 8A), case ii).

Figure 9 shows a set of model impedance spectra where the circuit (s) series and (p) parallel branches are both considered. The order of magnitude of resistances and capacitances is selected to generate spectra with the graphical features commonly observed in published data (clear bulk and grain boundary arcs), but there is no attempt to reproduce exact sets of published data. Arbitrary units are assumed in all examples discussed below.

In Figure 9A the bulk resistance ($R_b = 1000$, in the bottom series branch) is only one fifth of the perpendicular/series grain boundary resistance ($R_{gb,s} = 5000$), and both much smaller than the upper parallel resistance ($R_{\rm gb,p} = 180,000$). The much larger impedance of the parallel grain boundary branch (Z_p) prevents the observation in this plot of the impedance of the series branch (Z_s) and of the total circuit impedance (Z_t) , both detailed in

Figure 9B. The spectra for Z_s and Z_t are almost identical because of the large Z_p . In this case, the upper parallel branch is almost electrically inactive.

Figure 9C shows the result of a much lower Z_p ($R_p = 1200$), while the remaining circuit parameters were preserved. Zt now is completely different from Z_s, due to the role of the low parallel resistance. Simple naked eye analysis of these spectra suggests that the (perpendicular) grain boundary impedance is only about 1/10 of the original value while the bulk impedance shrank to about 1/2 of the original value. In fact, they were fully preserved but they appear as if the corresponding conductivities increased. The situation just described fits quite well with previous reports on enormous enhancement of grain boundary conductivity and moderate bulk conductivity improvement. Migration from spectra in Figures 9B,C was obtained with simple consideration of a fastionic parallel pathway. This can be easily done assuming distinct phases and properties for the parallel and perpendicular grain boundaries.

Figure 9D repeats the previous exercise but with a constraint on the electrical properties of parallel and perpendicular grain boundaries, by assuming the same specific conductivity (σ_{sp} , see Eq. 9) and dielectric constant (ε_r) for the series and the parallel grain boundaries. The initial values of $R_{\rm b} = 1000$, $R_{\rm gb,s} = 300$, $C_{\rm b} =$ 10^{-9} and $C_{\rm gb,s} = 2 \times 10^{-7}$ corresponding to $\delta_{\rm gb} = 5$ nm and D =1 µm, determine the other relevant parameters, which scale with δ_{gb} according to

$$R_{\rm gb,p} = R_{\rm gb,s} D^2 \left(D + \delta_{\rm gb} \right) \cdot \left[\delta_{\rm gb}^2 \left(2D + \delta_{\rm gb} \right) \right]^{-1}$$
(10)

$$R_{\rm b} = \left[\sigma_{\rm b} \left(D - \delta_{\rm gb}\right)\right]^{-1} \tag{11}$$

$$R_{\rm gh,s} = \sigma_{\rm sp}^{-1} D^{-2} \delta_{\rm gh} \tag{12}$$

$$C_{\rm gb,s} = C_{\rm b} D \delta_{\rm gb}^{-1} \tag{13}$$

$$C_{\rm gb,p} = C_{\rm gb,s} R_{\rm b} R_{\rm gb,p}^{-1} \tag{14}$$

The spectra in Figure 9D show the impact of increasing δ_{gb} on the total impedance. Starting from the initial geometry corresponding to $\delta_{gb} = 5$ nm, one observes a significant increase of the total impedance up to $\delta_{gb} = 50$ nm, which is essentially determined by the increase of $R_{gb,s}$ (Eq. 12). The other parameter in the series branch, R_{b} , also increases with increasing δ_{gb} (Eq. 11), but much less. Comparison of the two spectra for $\delta_{gb} = 5$ nm and $\delta_{gb} =$ 50 nm visually confirms this trend since the high frequency arc (on the left, for low impedance) remains nearly unchanged, whereas the low frequency contribution is much larger for $\delta_{gb} = 50$ nm. Further increase of δ_{gb} reverses the trend (the impedance for $\delta_{gb} = 200$ nm is already smaller than for $\delta_{gb} = 50$ nm), and the two arcs lose some of their distinctive features and start to overlap.

At $\delta_{gb} = 400 \text{ nm}$ one observes that the total impedance is already (slightly) smaller than the initial condition ($\delta_{gb} = 5 \text{ nm}$), and the spectrum appears as a single, highly depressed arc. This behavior is explained by the accentuated decrease of the parallel branch resistance $R_{gb,p}$, which, as shown by Eq. 10, scales down with δ_{gb} to the power of -3, whereas $R_{gb,s}$ scales up with a simple unity exponent. Therefore, significant changes in the total impedance are only observed for an exorbitant δ_{gb} of 400 nm, requiring a large volume fraction of a secondary phase. This condition is typical of a composite but not of a ceramic with a minor secondary phase (partly) in the grain boundary region.

Overall, Figure 9 shows that the usually reported enormous enhancement of the grain boundary conductivity and moderate improvement of the bulk conductivity are not compatible with a simply enhanced grain boundary conductivity, even with due consideration of the so-called parallel pathway. On the contrary, consideration of a secondary phase with high conductivity (Figure 9C) can easily generate the reported changes in the spectra with sintering admixtures. From known properties of solid and molten ionic conductors, the latter are closer to fulfil the high conductivity requirements. From a microstructural point of view, even minor amounts of molten phases can easily percolate throughout contacts between grains, still bonded *via* common grain boundaries.

The presence of a secondary charge transport pathway and carrier (e.g., alkali metal ions) should show up under direct current (dc) conditions (and long tests), since typical electrodes cannot act as sink and source of alkali metal ions. However, this situation is not necessarily detected under alternate current (ac) analysis, as with impedance spectroscopy. A report on dc experiments (Hebb-Wagner method (Neuhaus et al., 2018b)) suggests peculiar constraints while testing Ca-doped ceria with Na_2CO_3 as second phase, namely with respect to the (limited) range of allowed applied voltages, with respect to the situation observed with pure Ca-doped ceria. Whether or not this is a confirmation of the presence of multiple species is obviously unclear.

Conclusion

The assumption of exclusive oxide-ion conductivity in ceriabased electrolytes prepared with salts of alkali metals neglects the possible contribution of other species and compounds with significant conductivity, situation aggravated when even minor amounts of molten phases might be present. A small fraction of a fast-ionic conductor can easily explain claimed trends in (moderate) bulk and (enormous) grain boundary conductivity enhancements.

With (ac) impedance spectroscopy all mobile species can contribute to the total conductivity, preventing the identification of separate contributions. Tests performed under dc conditions and complementary techniques able to detect alternative molten/ amorphous phases (e.g., FTIR to detect hydroxides) are needed to fully discard alternative explanations for reported conductivity enhancements. Furthermore, confirmation of exact mechanisms behind grain boundary conductivity enhancements, require advanced analytical tools to confirm most hypotheses under discussion. Advanced grain boundary engineering is feasible using admixtures, opening possibilities up to now rarely exploited. This involves manipulation of the local concentration of species at low temperature, preventing the accumulation of dopants, often observed after high temperature processing. This solution is of particular interest in materials like ceria-based electrolytes. Long term stability of these materials under typical operating conditions also needs proper assessment.

Author contributions

JG: writing—original draft, formal analysis; AJ and MS: writing—review; FF: conceptualization, writing—review and editing, supervision; FM: conceptualization, writing—review and editing, funding acquisition, supervision.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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