



Optical Properties and Dynamic Extrinsic Chirality of Structured Monolayer Black Phosphorus

Mengke Sun, Ying Wang, Hui Hu, Hao Zhang, Wenjia Li*, Bo Lv, Zheng Zhu, Chunying Guan and Jinhui Shi*

College of Physics and Optoelectronic Engineering, Harbin Engineering University, Harbin, China

Chiral metamaterials have drawn increasing attention due to their strong chiral responses. Monolayer black phosphorus is a tunable two-dimensional material with anisotropy that plays an important role in a variety of fields such as chirality and polarization control. In this work, we propose a metamaterial with structured monolayer black phosphorus to manipulate the transmission properties of circularly polarized waves. The metamaterial exhibits strong circular dichroism and circular birefringence effects depending on oblique incidence of the circularly polarized wave and has a weaker circular conversion dichroism effect as well. Moreover, this work also investigates effects of different chiral phenomena of the metamaterial on various structural parameters as well as incident angles and the electron concentration. It has been proved that the electron concentration of monolayer black phosphorus can dynamically tune the chirality properties. Remarkably, the non-zero pure optical activity always occurs at one certain frequency regardless of the elevation angle and the azimuthal angle. The proposed framework provides opportunities for designing meta-devices with monolayer black phosphorus and practical potentials for novel and high-performance infrared metamaterials.

Keywords: metamaterial, black phosphorus, extrinsic chirality, circular dichroism (CD), circular birefringence

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*Correspondence:

Wenjia Li
liwenjia@hrbeu.edu.cn
Jinhui Shi
shijinhui@hrbeu.edu.cn

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INTRODUCTION

In the past years, many attempts have been made to realize the manipulation for phase, amplitude, and polarization of electromagnetic waves by exploiting metamaterials (Shi et al., 2014; Gansel et al., 2015; Scheuer, 2017; Yang et al., 2018). Metamaterials, a kind of artificial structures, can be elaborately designed to acquire extraordinary and unique properties including negative refraction, perfect lens, giant chiral response, and nonlinearity (Shelby et al., 2001). Metamaterials can be also extended into diverse fields, such as electromagnetism, acoustics, and mechanics (Deng et al., 2019; Fan et al., 2019; Thevamaran et al., 2019). Chiral objects lack mirror symmetry, and it cannot superimpose with its mirror image through planar translation or rotation. Chirality plays an indispensable role in chemistry, biology, and medicines (Lewis et al., 1999; Ćorić and List, 2012; Collins et al., 2017). However, since chiralities of natural materials are inherently limited, metamaterials with giant chiral responses have attracted considerable attention and shown importance in optics (Rogacheva et al., 2006; Decker et al., 2010; Zhou et al., 2012). Chirality has been classified into intrinsic chirality and extrinsic chirality. Extrinsic chirality refers to chiral response of achiral anisotropic planar metamaterials at oblique incidence (Plum et al., 2009a; Plum et al., 2009b). To achieve extrinsic chirality, two-dimensional (2D) materials, including graphene and black phosphorus (BP), have drawn

much attention over the past decade (Shi et al., 2018; Wei et al., 2020) because they have excellent electromagnetic and optical properties. With the comprehensive study of 2D materials, BP has gradually become one of the most popular materials for its unique optical properties including intrinsic anisotropy and has been explored for extensive applications such as polarization converters (Valagiannopoulos et al., 2017; Cheng et al., 2019) and plasmonic devices (Liu and Aydin, 2016; Hong et al., 2018).

Different from other 2D materials (Chen et al., 2019), BP is the most stable allotrope of phosphorus with a stratified structure that breaks narrow bandgap limitation and low absorption efficiency of graphene. In addition, BP has tunable thickness-dependent bandgaps changing from 0.35 eV (bulk) to 2 eV (monolayer), filling up the “blank space” between graphene and transition-metal dichalcogenides (Li et al., 2014; Chen et al., 2015; Saito and Iwasa, 2015; Jia et al., 2020). Thus, BP has been widely studied and applied in some electronic devices, such as field effect transistors (Buscema et al., 2014a; Liu et al., 2014; Wang et al., 2014; Liu et al., 2016), photovoltaic devices (Buscema et al., 2014b), rechargeable batteries (Park and Sohn, 2007), solar cells (Zhou et al., 2016), and heterostructures (Shen et al., 2015; Niu et al., 2019). On the other hand, BP has a large tunable bandgap. Thus, the manipulation of electromagnetic waves with BP, such as absorbers (Lu et al., 2015; Niu et al., 2018; Kong et al., 2019) and reconfigurable color display (Jia et al., 2021), has aroused widespread interest. The monolayer BP is composed of hexagonal honeycomb lattice with a folded structure (Low et al., 2014; Hong et al., 2019). Due to strong in-plane anisotropy, BP is one of the most important 2D materials to manipulate the polarization state of light as well as polarization-dependent optical phenomena. At mid-infrared wavelength, polarization-dependent localized surface plasmon resonances have been studied and numerically verified in monolayer BP nanoribbon and nanopatch arrays (Liu and Aydin, 2016). The optical activity resulting from the extrinsic chirality has been proved theoretically and numerically in unpatterned monolayer BP, which can be adjusted by the mutual orientation of the in-plane anisotropic BP film and the incident light (Hong et al., 2019). The monolayer BP is promising to realize tunable polarization manipulation; however, chiral metamaterials with the whole BP layer have limited tunable properties and working frequency range.

In this work, we propose monolayer nanostructured BP metamaterial, patterned on a dielectric substrate and study its optical properties, such as transmittance, optical activity, and circular conversion dichroism. In addition, the dependence of optical properties in the BP metamaterial is presented on different parameters including the period, duty ratio, doping concentration levels of the monolayer BP, and angle of incidence. The different parameters differently impact the chiral responses; thus, nanostructured square BP is beneficial to realize diverse optical properties. Compared with the previously reported monolayer BP (Hong et al., 2019), the designed monolayer BP metamaterial has more flexibilities and

more freedoms to tailor tunable circular response. The proposed structure is promising to design various novel and high-performance infrared metamaterials.

DESIGN AND ANALYSIS OF THE TUNABLE METAMATERIAL

The proposed structure is schematically illustrated in **Figure 1A**. The armchair and zigzag directions of the monolayer BP are along x - and y -directions, respectively. The square BP marked by the black box is periodically patterned on a dielectric substrate with a relative dielectric constant of 4 and relative permeability of 1. Both periodicities along x and y directions are identical and marked by p . The doping concentration level of the monolayer BP is denoted as n . The BP thickness is assumed to be 1 nm, and the width is w . The parameter w is equal to the product of the period (p) and the duty ratio (k). The angle between the wave vector and the z -direction is defined as the elevation angle β , and the angle between the incidence plane and the x -direction is defined as the azimuthal angle α . The metamaterial will exhibit chiral response at oblique incidence of the circularly polarized wave. A frequency domain solver is adopted, and the periodic boundary conditions are applied along x and y directions in software Microwave Studio (CST).

To demonstrate the extrinsic chirality of the designed metamaterial, the electromagnetic responses of the square BP are simulated as shown from **Figure 1B** to **Figure 1E**, in which $n = 6 \times 10^{13} \text{ cm}^{-2}$, $p = 500 \text{ nm}$, $k = 0.95$, $\beta = 75^\circ$, and $\alpha = 45^\circ$. The circular dichroism (labeled as Δ) can be written by

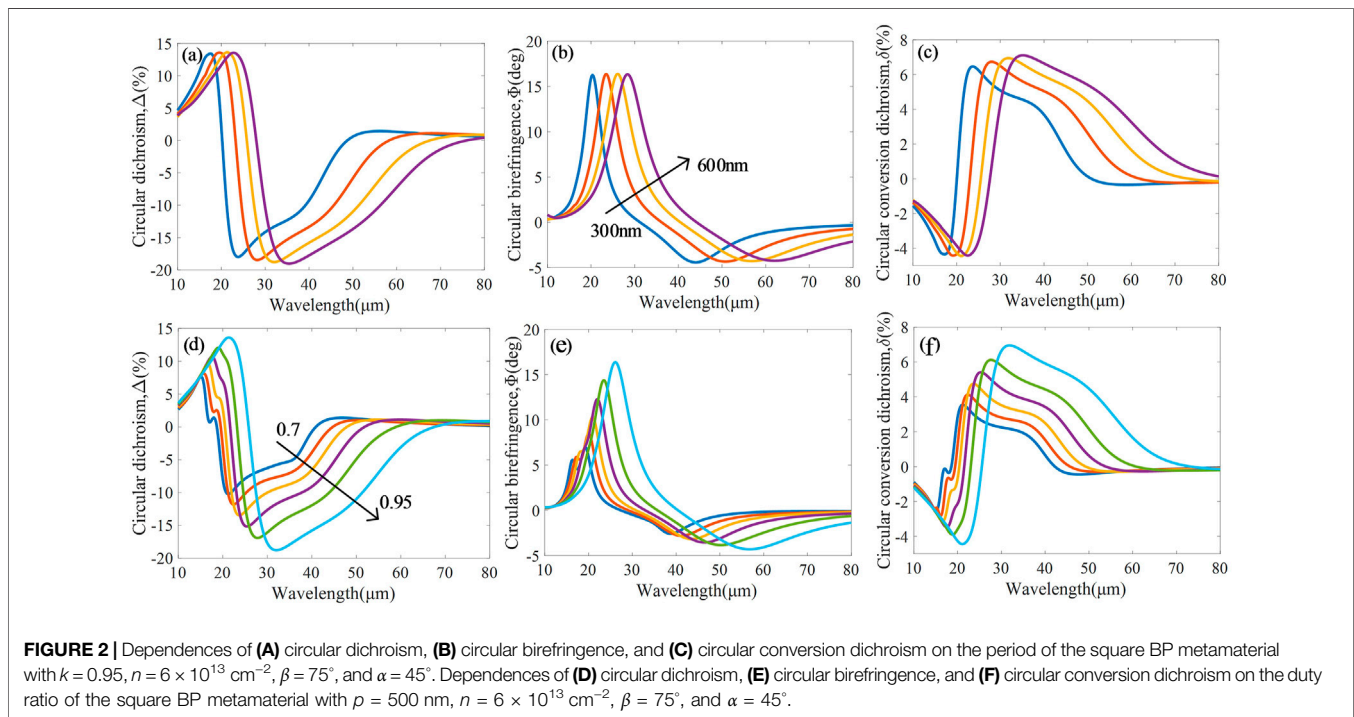
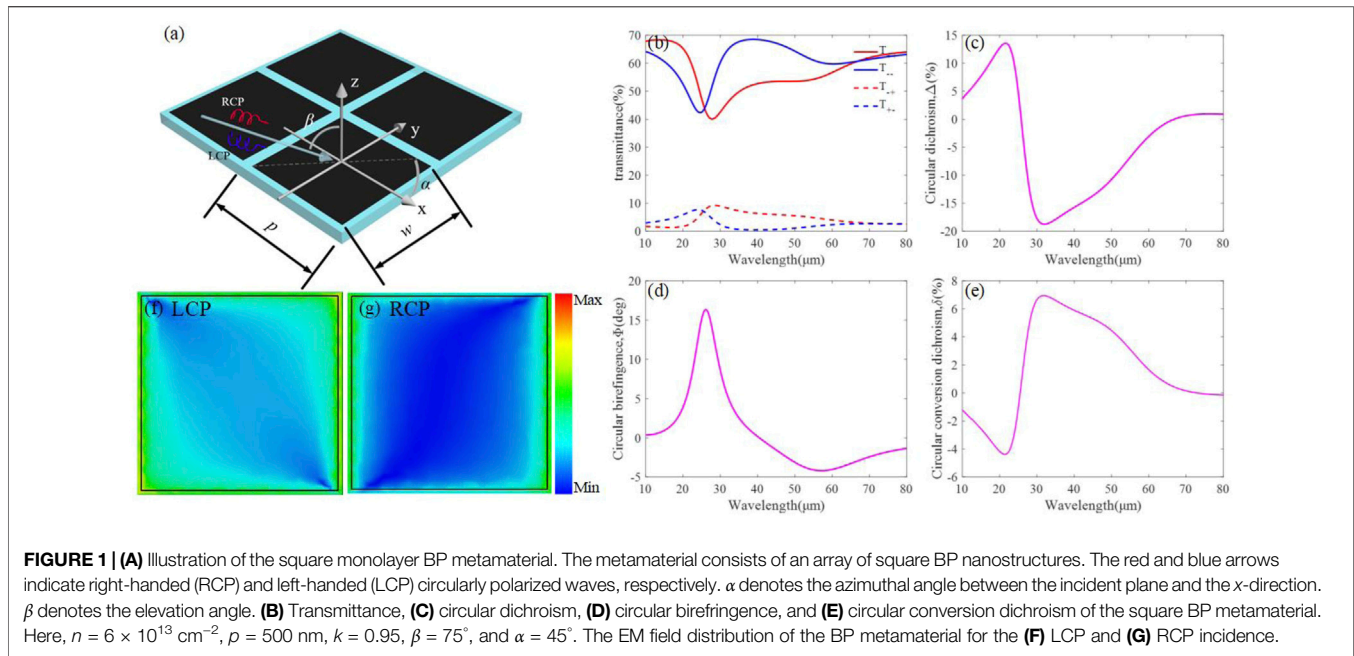
$$\Delta = T_{++} - T_{--}, \quad (1)$$

where the subscripts “+” and “-” denote right-handed circularly polarized light (RCP) and left-handed circularly polarized light (LCP), respectively, and the order of the subscripts corresponds to transmission and incidence, respectively. The transmittance is expressed as $T_{ij} = |t_{ij}|^2$, where t represents the magnitude of electromagnetic waves and “ i ” as well as “ j ” represents the transmitted and incident electric components, respectively. The circular birefringence (labeled as Φ) and circular conversion dichroism (labeled as δ) can be given by

$$\Phi = \arg(t_{++}) - \arg(t_{--}); \quad (2)$$

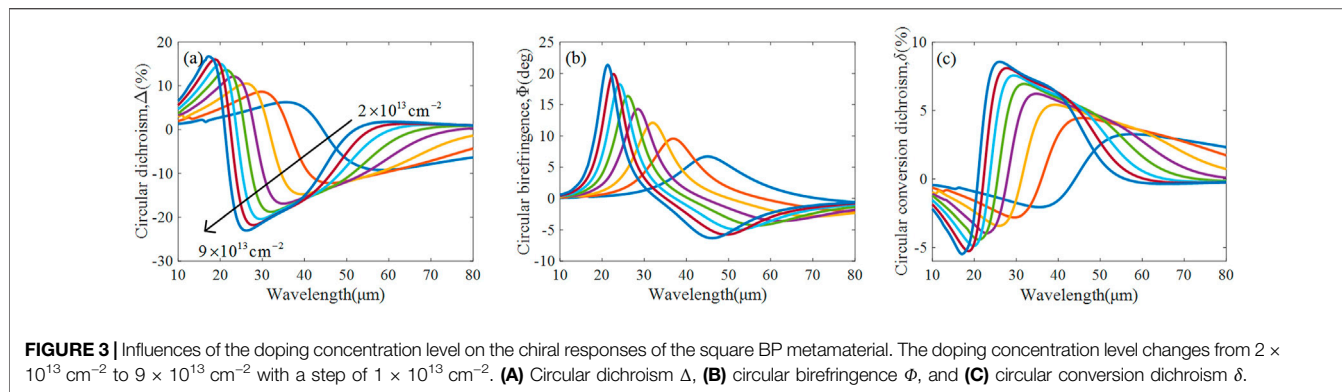
$$\delta = T_{+-} - T_{-+}. \quad (3)$$

It is obvious in **Figure 1B** that the square BP metamaterial distinctly responds to the incident RCP and LCP waves in the nearly whole wavelength range. Both the transmittance spectra T_{++} and T_{--} exhibit resonant dips at the wavelength ranging between 20–30 μm . The transmittance T_{++} and T_{--} are different except for the cross point at about 26 μm ; thus, the square BP metamaterial reveals extrinsic chirality at oblique incidence since it is intrinsically achiral. The relatively weak circular polarization conversion is shown in **Figure 1B**. The co-polarization transmittance is several times higher than the cross-polarization transmittance. Similar to the co-polarization transmittance spectra, there is a cross point between two cross-polarization transmittance spectra at about 26 μm , and they are different somewhere else. The circular dichroism of the



metamaterial is shown in **Figure 1C**. The signs of the two circular dichroism peaks are opposite due to the respective resonances of two circularly polarized waves. The peak value is 14.0% at the wavelength of 21.3 μm , indicating that the RCP transmittance is greater than that of LCP. Moreover, the peak value of the circular dichroism is -18.8% at the wavelength of 32.3 μm , indicating that the LCP transmittance is higher. **Figure 1D** shows that the maximum circular birefringence is 16.4° at the wavelength of $\sim 26.1 \mu\text{m}$, while the circular dichroism is

0 at the wavelength of $\sim 26 \mu\text{m}$; as a result, the metamaterial exhibits a pure optical activity, and the transmitted wave is linearly polarized. **Figure 1E** displays the circular conversion dichroism. The circular conversion dichroism is up to 7% at the wavelength of $\sim 30 \mu\text{m}$ that is twice as large as that in the reported metamaterial with a whole BP layer (Hong et al., 2019). **Figures 1F,G** depict the field patterns of the proposed metamaterial with the LCP and RCP incidence, respectively, which is used to tell the optical difference between



LCP and RCP waves (Cao et al., 2013; Cao et al., 2015; Cao et al., 2017). The circular conversion dichroism spectrum of the designed BP structure varies a lot with the wavelength. Compared to the previous literature (Hong et al., 2019), the square BP metamaterial offers versatile chiral responses and abundant resonances for easily tailoring the polarization properties.

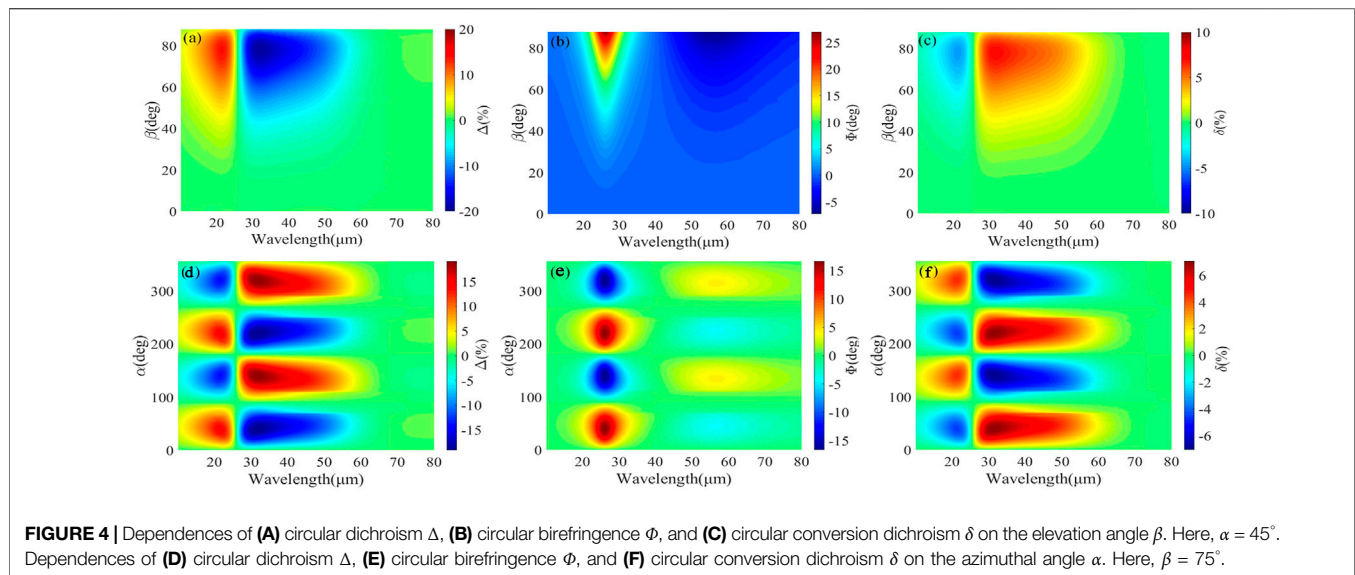
To fully explore multiple parameters-tailored chiral responses, we investigate how the responses of the square BP metamaterial depend on various structural parameters including period, duty ratio, doping concentration level, and angle of incidence. The conditions of $n = 6 \times 10^{13} \text{ cm}^{-2}$, $\beta = 75^\circ$, and $\alpha = 45^\circ$ are kept unchanged in **Figure 2**. The effects of the period and the duty ratio on the chiral responses are studied. **Figures 2A–C** show the dependence of the circular dichroism, circular birefringence, and circular conversion dichroism on the period ranging from 300 to 600 nm at an interval of 100 nm, in which $k = 0.95$. In fact, the period and the width of the square BP scale up or down together due to a fixed duty ratio. The peak values of all the responses remain unchanged with different periods, while the corresponding resonant wavelengths largely have red shifts with the increasing period. During the period change, the resonant wavelengths shift more than 30% (the shift is larger than 10 μm). The chiral responses can be well engineered *via* parameter study.

Next, the dependences of the circular dichroism, the circular birefringence, and the circular conversion dichroism are investigated on the duty ratio shown in **Figures 2D–F**. The duty ratio changes from 0.7 to 0.95 with a step of 0.05 while $p = 500 \text{ nm}$, $n = 6 \times 10^{13} \text{ cm}^{-2}$, $\beta = 75^\circ$, and $\alpha = 45^\circ$ remain constant. The effect of the duty ratio is quite different from the aforementioned one. With the increase of the duty ratio in the range of 0.7–0.95, the peak values of all the chiral responses increase, and all the spectra move toward longer wavelengths. On the other hand, in **Figures 2A–F**, the tendency of the curves is alike with different duty ratios. Thus, the positions and amplitudes of the resonant peaks are changed simultaneously with different duty ratios. The resonances of the circular dichroism, circular birefringence, and circular conversion dichroism of the metamaterial could be moved to the required wavelength by changing the period and the duty ratio.

To well understand tunable chiral responses of the square BP metamaterial, it is necessary to study the effect of the doping concentration level of monolayer BP. **Figures 3A–C** depict the

changes of circular dichroism, circular birefringence, and circular conversion dichroism with different doping concentration levels. The doping concentration level varies from $2 \times 10^{13} \text{ cm}^{-2}$ to $9 \times 10^{13} \text{ cm}^{-2}$ with a step of $1 \times 10^{13} \text{ cm}^{-2}$ while $p = 500 \text{ nm}$, $k = 0.95$, $\beta = 75^\circ$, and $\alpha = 45^\circ$ are kept unchanged. The two resonant values of the circular dichroism spectrum vary from 6.2 to 16.7% and from -9.2% to -23% , respectively. The corresponding wavelengths alter from 36 to 17.4 μm and from 57.5 to 26 μm , respectively. **Figure 3A** implies that the metamaterial has a giant extrinsic chirality and reveals tunability in a broad frequency range. With the increase of the doping concentration level, the resonant peaks of the circular birefringence spectrum increase and the resonant wavelengths have a large blue shift. When the doping concentration level is $2 \times 10^{13} \text{ cm}^{-2}$, the metamaterial has a weak chirality effect, while the maximum value of the circular birefringence can reach 21.3° as the doping concentration level is $9 \times 10^{13} \text{ cm}^{-2}$. In addition, the circular conversion dichroism becomes stronger when the doping concentration level increases. The tuning ranges of three chiral responses are as high as 50%, for instance, the resonance of the circular dichroism spectrum moves from 36 to 17.4 μm . Therefore, the tunable extrinsic chirality of the BP metamaterial is strongly related to the doping concentration level. The doping concentration level is one of the most important parameters to actively modulate the chiral response of the metamaterial. In the proposed metamaterial, the chirality could be controlled by the doping concentration level of the square BP *via* the electric field, pressure, temperature, and light, especially it can act as a switch in the wavelength from 20 to 30 μm , which can modulate the chirality of the metamaterial by changing the doping concentration level of monolayer BP. Obviously, the structure can be optimized before the metamaterial is manufactured just by choosing appropriate parameters including period. Moreover, when the metamaterial is designed, BP as a tunable material can also achieve dynamic optical phenomena such as the increase or decrease of chiral properties through various doping concentration levels instead of redesigning the structure.

For the proposed structure, there is no chirality for the incident circularly polarized wave at normal incidence. It is necessary to study the angular dependence of the extrinsic chiral responses in the square BP metamaterial, while all other geometrical and physical parameters are unchanged as those in **Figures 1B–E**. The angular dependences of the metamaterial's responses are shown in **Figure 4** when the elevation angle β and the azimuthal angle α are individually



changed. All circular dichroism, circular birefringence, and circular conversion dichroism are mostly dictated by the elevation angle β in **Figures 4A–C**, in which $\alpha = 45^\circ$. When β is small, the metamaterial does not exhibit any chiral responses including circular dichroism, circular birefringence, and circular conversion dichroism due to the symmetry of the metamaterial. With increasing β , the extrinsic chirality occurs, and the chiral responses manifest their maxima at large angle of incidence. The circular dichroism, circular birefringence, and circular conversion dichroism reach their maxima of -18.8% , 28° , and 7% , respectively. Remarkably, the non-zero pure optical activity phenomenon always occurs at $\sim 26\ \mu\text{m}$ regardless of the elevation angle β . It can be seen from **Figures 4A–C** that there is a significant characteristic in the vicinity of the wavelength $26\ \mu\text{m}$, in which the circular dichroism and the circular conversion dichroism are opposite at the left and right sides of this wavelength, and the circular birefringence always reach its resonance. Here, $\beta = 75^\circ$. The chiral properties of the square BP metamaterial are shown in **Figure 4D–Figure 4F** when the azimuthal angle α varies in the full range of 360° . All spectra are periodically produced at every 180° , and their signs are reversed at every 90° . When the azimuthal angle α is 0° , 90° , 180° , and 270° , the whole system combined with the proposed structure and the incident circularly polarized waves is symmetric; thus, the chirality vanishes in the square BP metamaterial. The pure optical activity at $26\ \mu\text{m}$ does not shift with the azimuthal angle. Thus, the incident angles α and β provide a way to engineer the extrinsic chirality in the square BP metamaterial.

CONCLUSION

In conclusion, we have proposed a metamaterial with monolayer BP patches to demonstrate a giant tunability of the extrinsic chirality. The circular dichroism, circular birefringence, and circular conversion dichroism of the metamaterial have been numerically investigated. The extrinsic chirality can be well engineered to work in a broad

frequency range via studying the parameters such as period, duty ratio, doping concentration level, and angle of incidence. The circular conversion dichroism is up to 7% at the wavelength of $\sim 30\ \mu\text{m}$, which is twice as large as that in the whole BP layer. The doping concentration level of the square BP metamaterial leads to large tuning ranges of three chiral responses up to 50% . Remarkably, the non-zero pure optical activity always occurs at one certain frequency regardless of the elevation angle and the azimuthal angle, and the circular birefringence reaches 28° . Compared to the whole BP layer in the previous literatures, the square BP metamaterial offers versatile chiral responses and abundant resonances for easily tailoring the polarization properties. The proposed metamaterial with BP patches is promising to realize active polarization manipulation as well as chiral switching and biological chiral sensing.

DATA AVAILABILITY STATEMENT

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

AUTHOR CONTRIBUTIONS

MS, YW, and HH contributed to the conception and design of the study. MS performed the design simulation and wrote the first draft of the manuscript. YW and HH performed the statistical analysis. YW, WL, and JS wrote the sections of the manuscript. WL and JS supervised the project. All authors contributed to manuscript revision and read, and approved the submitted version.

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