



Dynamic Performance Detection of CFRP Composite Pipes based on Quasi-Distributed Optical Fiber Sensing Techniques

Hua-Ping Wang^{1,2*}, Si-Yuan Feng¹, Xiang-Shu Gong¹, Yan-Xin Guo¹, Ping Xiang^{3*}, Yu Fang¹ and Qi-Ming Li¹

¹School of Civil Engineering and Mechanics, Lanzhou University, Lanzhou, China, ²Key Lab of Mechanics on Disaster and Environment in Western China (Lanzhou University), Ministry of Education, Lanzhou, China, ³School of Civil Engineering, Central South University, Changsha, China

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*Correspondence:

Hua-Ping Wang
wanghuaping1128@sina.cn
hpwang@lzu.edu.cn
Ping Xiang
pxiang2-c@my.cityu.edu.hk

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For the high strength, corrosion resistance, and good stability, carbon fiber-reinforced polymer (CFRP) composites have been made into pipes to transfer gasses and oils in subsea environment. Structural performance of CFRP composite pipes is particularly important to sustain the regular operation of the delivery system. To obtain the in-field behavior of the CFRP composite pipes, quasi-distributed optical fiber sensing techniques are developed based on the multiple configuration of fiber Bragg grating (FBG) sensing elements. Theoretical investigation on the dynamic response of the pipes is performed. Experiments on cantilever CFRP pipes with surface-attached FBGs in series and packaged FBG sensors have been conducted to check the feasibility and effectiveness of the proposed sensing technique. Results validate the good measurement performance of the proposed sensors and the accuracy of the vibration analysis. The study can be adopted to instruct the establishment of the structural health monitoring system of CFRP composite pipes and assess the safety operation state of the pipe systems.

Keywords: CFRP composite pipe, dynamic performance, vibration response, quasi-distributed optical fiber sensing technique, dynamic strain measurement

INTRODUCTION

Pipe transportation has become the main channel for oil and gas, due to the advantages of high efficiency, low energy consumption, and reasonable cost. Submarine pipes with extremely high fabrication cost are prone to local buckling, cracking, perforation, and fatigue fracture damages induced by seawater corrosion, sea sand erosion, waves, and other multimedia cycle coupling action during the long-term operation process (Zhang et al., 2015; Ferras et al., 2017; Wu et al., 2020). The cumulative or sudden damage induced pipe defect or fracture can result in huge gas or oil leakage, which can inspire serious explosion and environmental pollution. Therefore, it is essential to develop effective, real-time, and long-term online monitoring technology to obtain the service status of submarine pipeline structures, conduct timely online damage diagnosis, and instruct the rapid rehabilitation and long-term safe operation and maintenance (Drumond et al., 2018; Psyras and Sextos, 2018).

Submarine pipelines have always been under the action of waves and currents. When the fluid flows bypass both sides of the suspended pipe, the change of pressure difference makes the boundary layers separate, and then the wake vortex falls off, leading to vibration of the pipe with certain frequency. When the excitation frequency that makes the vortex fall off is equal to or approaches the inherent frequency of the submarine pipe, the vibration locking phenomenon (resonance) occurs, and the amplitude of the pipe increases dramatically, which can induce fatigue failure of the pipe. Therefore, it is particularly important to explore the dynamic characteristics, dynamic failure mode, and dynamic damage monitoring methods of submarine pipelines, so as to develop preventive control measures for preventing the fatigue fracture damage inspired by vortex-induced vibration (Datta and Sarkar, 2016).

Currently, considerable attempts have contributed to analyzing the dynamic characteristics of pipe structures. Kershenbaum et al. (Kershenbaum et al., 2000) explored the seismic performance of pipelines with snake shape and restrained sections. Choi (Choi, 2001) established the dynamic model of offshore pipelines to discuss the natural frequencies and sensitive parameters of the free spans, and provide improved design suggestions regarding the vortex-induced vibration and free span analysis. Kouretzis et al. (Kouretzis et al., 2007) theoretically studied the surface point-source blast-induced ground wave propagation with radial attenuation and spherical front in pipelines and the distribution of strain responses. Giannaros et al. (Giannaros et al., 2016) adopted finite element simulation for understanding the blast response of composite pipelines. Tian et al. (Tian et al., 2019) used the Euler–Bernoulli foundation beam model to analyze the dynamic response of pipelines under impact load based on Fourier series. Talemi et al. (Talemi et al., 2019) conducted the 3D dynamic ductile fracture of steel pipeline based on coupled interaction simulation. Deng et al. (Deng et al., 2019) explored the dynamic scaling mechanism of glass fiber–reinforced plastics pipelines by experimental and numerical studies. Manolis et al. (Manolis et al., 2020) conducted the dynamic response of buried pipelines undergoing both axial and flexural vibrations based on Newmark assumption. Wang et al. (Wang et al., 2020) performed the experimental testing and numerical investigation on the blast-induced ground vibration of gas pipelines.

The mentioned analysis on pipes under various vibration actions based on the simplified theoretical solutions, finite element simulation (Zhu et al., 2020; Qu et al., 2021), and experimental testing can declare some dynamic features of submarine pipes. However, it is generally based on the simplified models and has difference with the practical pipe structures. Therefore, it is important to develop the dynamic testing technique for obtaining the vibration response of submarine pipeline, which can provide compensative and feasible supports for the vibration design, damage identification, and performance rehabilitation of submarine pipe structures.

Some scholars used piezoceramic transducers to detect the defect of pipelines (Datta and Sarkar, 2016; Tian et al., 2019). However, it is based on electrical input, which may lead to the

blast or fire of pipelines for oil and gas transportation. Due to the advantages of absolute measurement, anti-electromagnetic interference, high precision, multiplexing, and integration of sensing network over other sensors (Ansari, 2007; Wang et al., 2014; Wang et al., 2016; Wang et al., 2018a, Wang et al., 2018b; Wang et al., 2019a; Wang and Dai, 2019), optical fiber based sensors have thus been adopted to monitor the health condition and structural performance of pipe structures. For the excellent characteristics of low weight, corrosion resistance, and high strength (Dai et al., 2013; Gao et al., 2015; Wang et al., 2019b), carbon fiber–reinforced polymer (CFRP) composite pipes have been considered, and the exploration on the replacement of steel pipes still requires further investigation.

In view of this, the vibration characteristics of the CFRP composite pipe structures have been analyzed, and the displacements and strain responses of the pipe under bending vibration have been described from the perspective of continuum vibration differential equations. Experimental investigations on the vibration response characteristics of the CFRP composite pipes have been performed based on quasi-distributed optical fiber sensing technology. Furthermore, the effectiveness and applicability of the testing method have been demonstrated by the comparative analysis on calculated and measured dynamic strain responses.

VIBRATION ANALYSIS ON CFRP COMPOSITE PIPE

The aspect ratio of the CFRP pipe is larger than 50, which has the characteristics of the beam. Therefore, the infinite degree of freedom beam vibration model with uniform mass distribution is used to study the vibration of the CFRP pipe. The research objective is referred to submarine pipelines for transporting oil/gas, which belongs to slender beam and is dominated by low-order natural vibration. Therefore, the Bernoulli–Euler beam model can be used for analysis. The coordinates of the CFRP pipe follow **Figure 1A**, and an infinitesimal element is abstracted from the pipe, with force distribution displayed in **Figure 1B**. To simplify the theoretical derivation, the hollow-pipe cross section is equivalent to a “good” arbitrary cross section during the analysis (Dai et al., 2013; Wang et al., 2019b). $p(x, t)$ is the external vibration force. $f(x, t)$ is the inertial force. m is the mass per unit length of the pipe, and it is a constant for the beam model with uniform cross sections. EI is the bending stiffness of the pipe. $V(x, t)$ and $M(x, t)$ separately stand for the shear force and bending moment of the element, respectively.

The dynamic equilibrium condition in y direction gives

$$\frac{\partial V(x, t)}{\partial x} = p(x, t) - m \frac{\partial^2 u(x, t)}{\partial t^2}. \quad (1)$$

Ignoring the second-order small quantity, the balance condition of the rotation of the infinitesimal element around the centroid can be obtained.

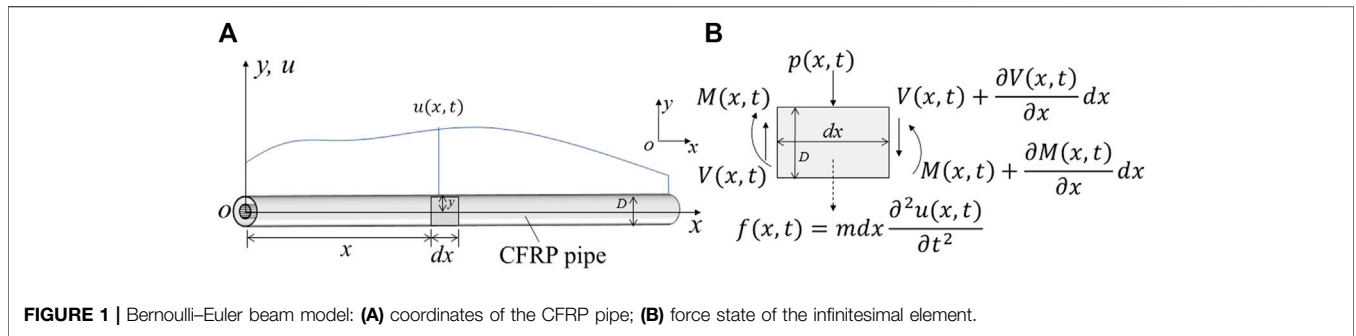


FIGURE 1 | Bernoulli-Euler beam model: **(A)** coordinates of the CFRP pipe; **(B)** force state of the infinitesimal element.

$$V(x, t) = \frac{\partial M(x, t)}{\partial x}. \tag{2}$$

According to the mechanics of materials,

$$M(x, t) = EI \frac{\partial^2 u(x, t)}{\partial x^2}. \tag{3}$$

Substituting Eqs 2, 3 into Eq. 1, the partial differential control equation is given by

$$m \frac{\partial^2 u(x, t)}{\partial t^2} + \frac{\partial^2}{\partial x^2} \left[EI \frac{\partial^2 u(x, t)}{\partial x^2} \right] = p(x, t). \tag{4}$$

For free vibration, the external excitation force $p(x, t)$ in Eq. 4 is 0, and the solution of this equation can be written as

$$u(x, t) = \phi(x)q(t), \tag{5}$$

then

$$\frac{\partial^2 u}{\partial t^2} = \phi(x)\ddot{q}(t) \quad \frac{\partial^2 u}{\partial x^2} = \phi''(x)q(t). \tag{6}$$

Substituting Eq. 6 into Eq. 4, it derives

$$m\phi(x)\ddot{q}(t) + q(t)[EI\phi''(x)]'' = 0. \tag{7}$$

The above formula can be rewritten as

$$\frac{-\ddot{q}(t)}{q(t)} = \frac{EI[\phi''(x)]''}{m\phi(x)}. \tag{8}$$

The premise that Eq. 8 is effective for all the x and t should be that the ratio of the left and right terms is a constant, denoted as ω^2 . Then, the partial differential equation $m \frac{\partial^2 u(x,t)}{\partial t^2} + \frac{\partial^2}{\partial x^2} [EI \frac{\partial^2 u(x,t)}{\partial x^2}] = 0$ is separated into two ordinary differential equations. One equation controls the time function $q(t)$, and the other controls the space function $\phi(x)$, which gives

$$\ddot{q}(t) + \omega^2 q(t) = 0, \tag{9a}$$

$$EI[\phi''(x)]'' - \omega^2 m\phi(x) = 0, \tag{9b}$$

and then Eq. 9b can be rewritten as

$$\phi^{(4)}(x) - \beta^4 \phi(x) = 0, \tag{10}$$

where $\beta^4 = \frac{\omega^2 m}{EI}$. Then, the general solution of Eq. 10 is

$$\phi(x) = C_1 \sin \beta x + C_2 \cos \beta x + C_3 \sinh \beta x + C_4 \cosh \beta x. \tag{11}$$

Considering the cantilever constraints, at the fixed end ($x = 0$), the displacement and inclination of the beam are equal to zero, then

$$u(0) = 0 \Rightarrow \phi(0) = 0 \Rightarrow C_2 + C_4 = 0 \Rightarrow C_4 = -C_2, \tag{12a}$$

$$\frac{\partial u(x, t)}{\partial x} \Big|_{x=0} = 0 \Rightarrow \phi'(0) = 0 \Rightarrow \beta(C_1 + C_3) = 0 \Rightarrow C_3 = -C_1. \tag{12b}$$

At the free end ($x = L$), both the bending moment $M(x, t)$ and the shear force $V(x, t)$ are zero. Therefore, according to Eqs 2, 3, and using Eq. 11, it yields

$$EI\phi''(L) = 0 \Rightarrow C_1(\sin \beta L + \sinh \beta L) + C_2(\cos \beta L + \cosh \beta L) = 0, \tag{13a}$$

$$EI\phi'''(L) = 0 \Rightarrow C_1(\cos \beta L + \cosh \beta L) + C_2(-\sin \beta L + \sinh \beta L) = 0. \tag{13b}$$

If both C_1 and C_2 are equal to zero, it will give a solution without vibration. To make at least one of the two coefficients not equal to zero, the frequency equation can be solved.

$$1 + \cos \beta L \cosh \beta L = 0. \tag{14}$$

To solve βL from Eq. 14, numerical methods can be used. That is, when $n = 1, 2, 3, 4, \dots$, we get $\beta_n L = 1.8751, 4.6941, 7.8548, 10.996, \dots$, where $\beta_n L \approx (2n - 1)\frac{\pi}{2}$. It then yields

$$\omega_n = \frac{(\beta_n L)^2}{L^2} \sqrt{\frac{EI}{m}} \quad n = 1, 2, 3, \dots \tag{15}$$

and then the natural vibration mode of each order corresponding to $\beta_n L$ is

$$\phi_n(x) = C_1 \left[\cosh \beta_n x - \cos \beta_n x - \frac{\cosh \beta_n L + \cos \beta_n L}{\sinh \beta_n L - \sin \beta_n L} (\sinh \beta_n x - \sin \beta_n x) \right], \tag{16}$$

where C_1 is the arbitrary constant.

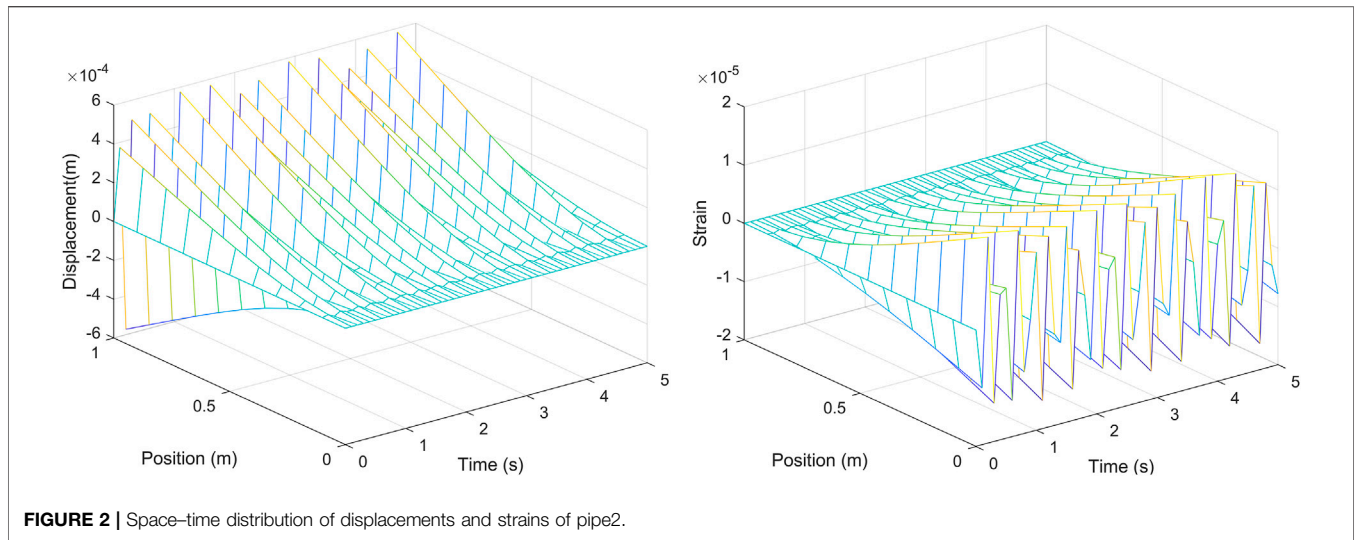


FIGURE 2 | Space-time distribution of displacements and strains of pipe2.

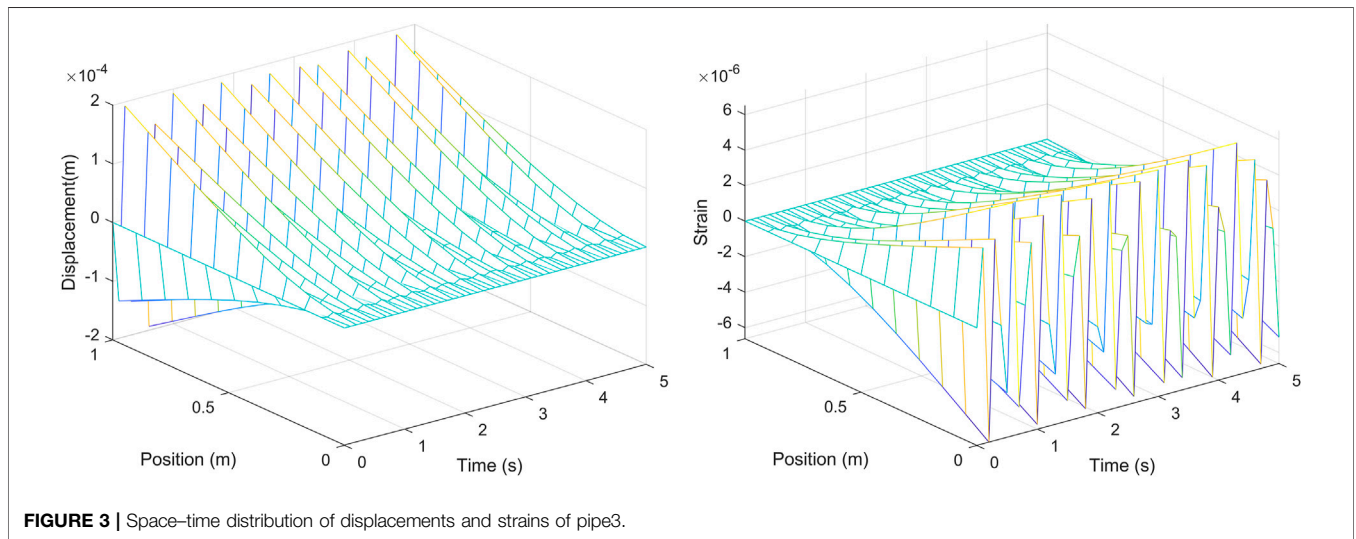


FIGURE 3 | Space-time distribution of displacements and strains of pipe3.

Using the mode superposition method to calculate the vibration of the beam, the vibration response of the beam can be expressed as

$$u(x, t) = \sum_{n=1}^{+\infty} \phi_n(x)q_n(t). \tag{17}$$

Substituting it into Eq. 4, and using the boundary conditions and the orthogonality of natural vibration modes, the principal coordinate differential equation can be obtained

$$M_n \ddot{q}_n(t) + K_n q_n(t) = P_n(t), \tag{18}$$

where $M_n = m \int_0^L [\phi_n(x)]^2 dx$, $K_n = EI \int_0^L [\phi_n''(x)]^2 dx$, and $K_n = \omega_n^2 M_n$. Since the simple harmonic excitation point acts on the end ($x = L$) of the cantilever pipe, substituting $p(x, t) = p_0 \sin \omega t \delta(x - \xi)$ ($\delta(x - \xi)$ a Dirac- δ function centered on ξ with

value equal to L in this case) into the generalized force $P_n(t) = \int_0^L p(x, t)\phi_n(x)dx$, it generates

$$P_n(t) = p_0 \sin \omega t \phi_n(\xi) \tag{19}$$

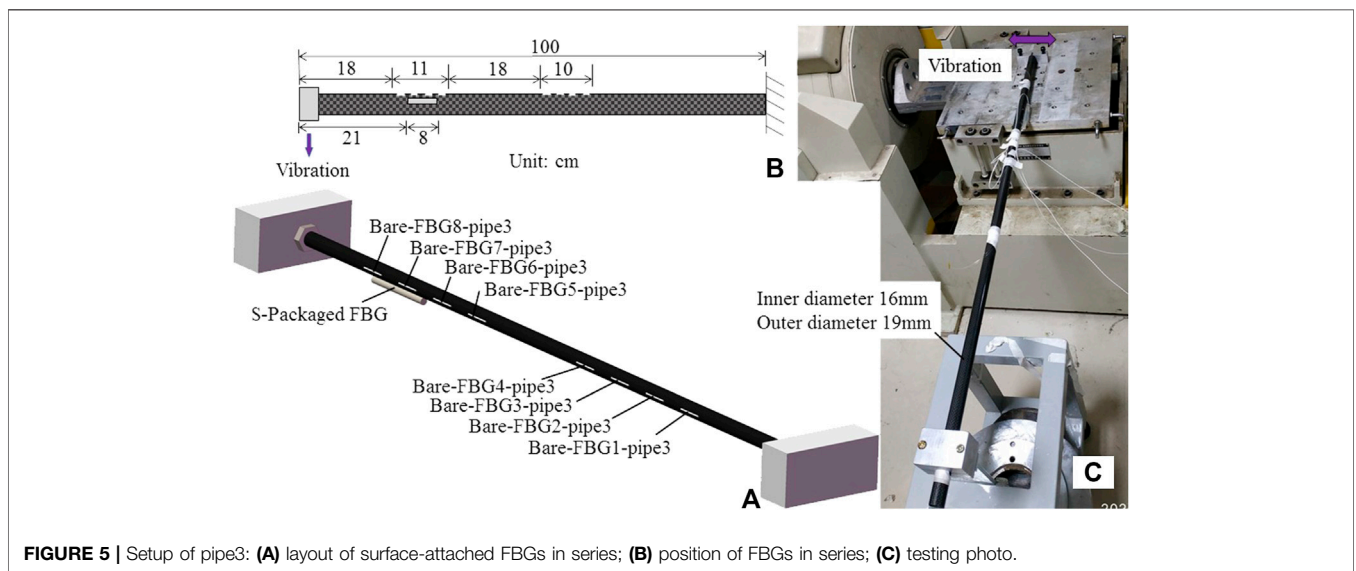
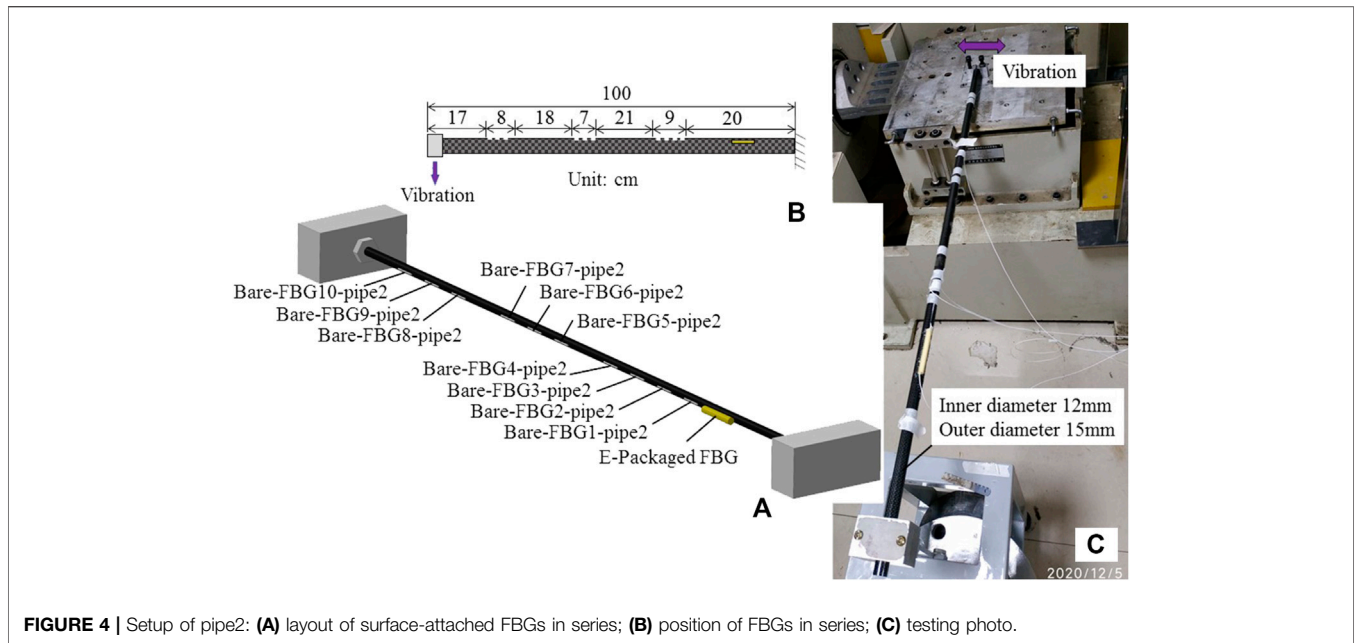
and then Eq. 18 can be rewritten as

$$\ddot{q}_n(t) + \omega_n^2 q_n(t) = \frac{\omega_n^2}{K_n} p_0 \sin \omega t \phi_n(\xi). \tag{20}$$

Introducing the boundary condition $q_n(0) = \dot{q}_n(0) = 0$, Eq. 20 is solved

$$q_n(t) = \frac{p_0 \phi_n(\xi)}{K_n} \frac{1}{1 - (\omega/\omega_n)^2} \left(\sin \omega t - \frac{\omega}{\omega_n} \sin \omega_n t \right). \tag{21}$$

Substituting $\xi = L$ into Eq. 21, and noticing that $\phi_n(x)$ has been solved from Eq. 16, then the vibration response of the beam can be given by



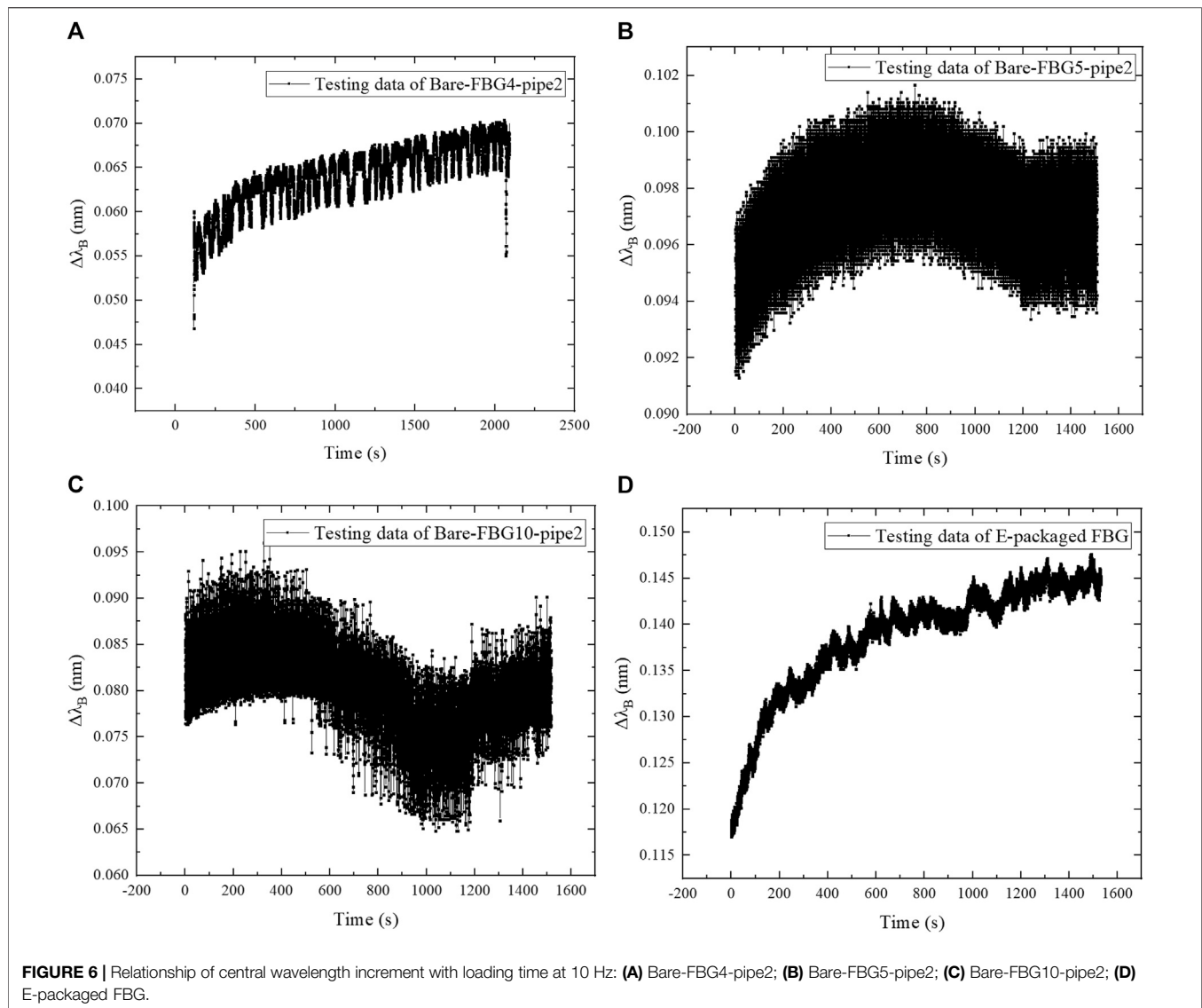
$$u(x, t) = \sum_{n=1}^{\infty} C_1 \left[\cosh \beta_n x - \cos \beta_n x - \frac{\cosh \beta_n L + \cos \beta_n L}{\sinh \beta_n L - \sin \beta_n L} (\sinh \beta_n x - \sin \beta_n x) \right] \cdot \frac{p_0 \phi_n(L)}{K_n} \frac{1}{1 - (\omega/\omega_n)^2} \left(\sin \omega t - \frac{\omega}{\omega_n} \sin \omega_n t \right) \quad (22)$$

According to the bending theory of beams, the axial strain response caused by bending can be obtained

$$\epsilon_x(x, t) = \frac{\partial^2 u(x, t)}{\partial x^2} y = \sum_{n=1}^{+\infty} \phi_n''(x) q_n(t) y, \quad (23)$$

where y is the distance from the strain measured point on the cross section to the neutral axis of the beam along the y -direction, as shown in **Figure 1A**. The expression of the strain response can be obtained by substituting the second derivative of $\phi_n(x)$ into **Eq. 23**.

Given the theoretical analysis above, a related numerical calculation program has been compiled to calculate the displacement and strain response of the cantilever CFRP pipe under the action of a unit harmonic force with an excitation frequency of 10 Hz. **Figures 2, 3** are corresponded to CFRP pipes with inner and outer diameters of 12–15 mm and 16–19 mm, respectively. It can be seen that the displacement of the CFRP pipe is largest at the free end (i.e., the loading point), and the



displacement undergoes periodic vibration changes over time. The strain of the CFRP pipe reaches the maximum at the fixed end, which also undergoes periodic vibration changes over time. Under the same load, since the section modulus of pipe3 is larger than that of pipe2, the strain and displacement responses of pipe3 are smaller than those of pipe2. The above results are in accordance with the laws of mechanics, indicating the effectiveness of the above solution.

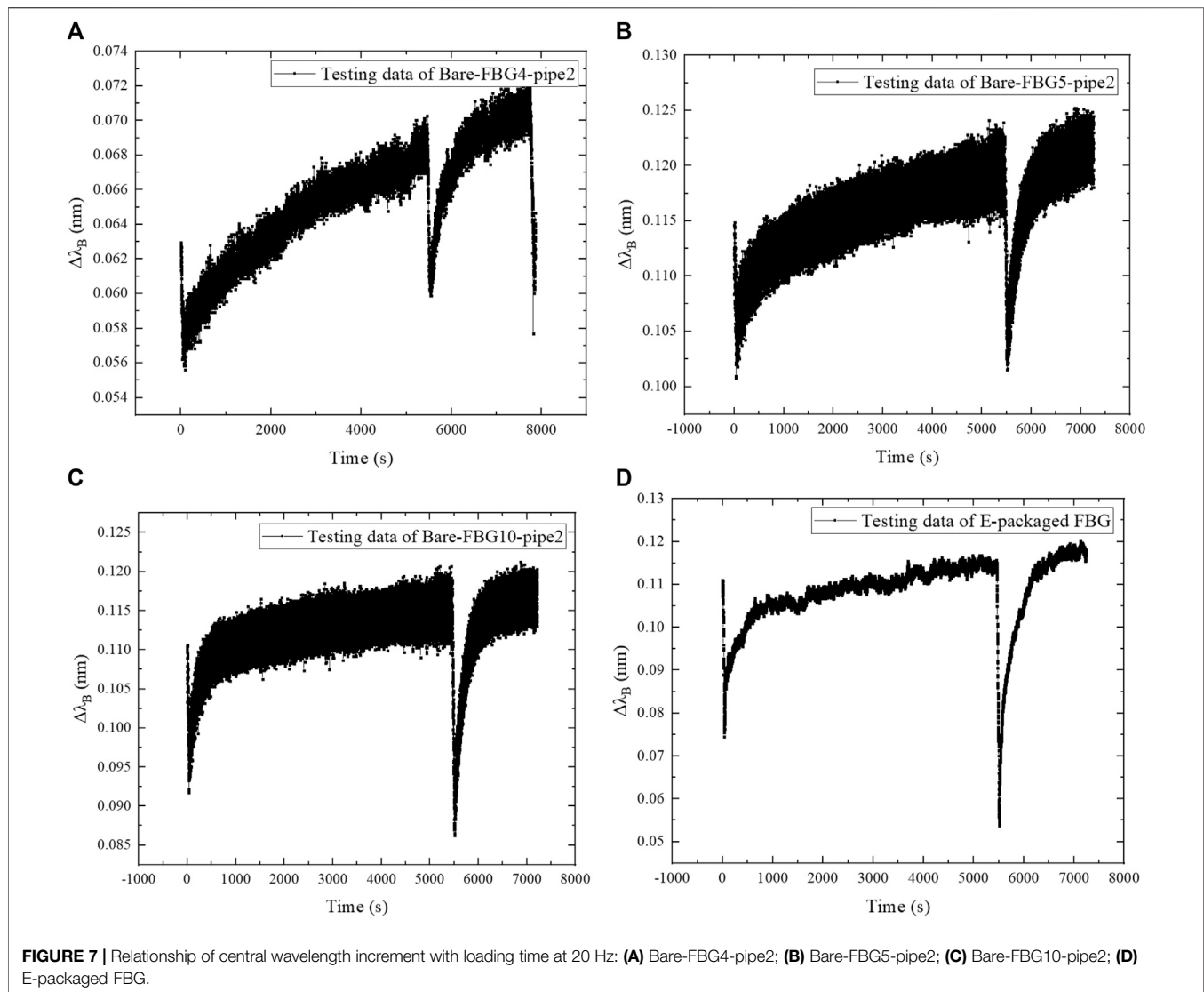
EXPERIMENTAL INVESTIGATION

To verify the feasibility and effectiveness of the installed FBGs in series to measure the vibration response of the pipe structures, experiments on vibration testing of the CFRP tube have been carried out. Surface-attached FBGs in series have been designed to measure the distributed strains in the axial direction of the pipe. The Optical system 200 has been

used to demodulate FBG signal information, and the sampling frequency is 100 Hz.

Testing Description

Two kinds of CFRP pipes have been used in the cantilever vibration experiments, as shown in **Figures 4, 5**. The relevant size information is as follows: the CFRP pipe with a length of 1 m, and with inner and outer diameters of 12 and 15 mm, respectively, is marked as pipe2; the CFRP pipe with a length of 1 m, and with inner and outer diameters of 16 and 19 mm, respectively, is marked as pipe3. Three FBGs in series have been arranged on pipe2: two series are composed of three FBGs, and the other one is composed of four FBGs. The related numbers are marked in **Figure 4A**. In addition, near the constraint end, an epoxy resin packaged FBG sensor has also been placed, marked as E-packaged FBG, which is mainly used to detect the influence of packaging measure on the vibration deformation measurement of FBGs. The similar information



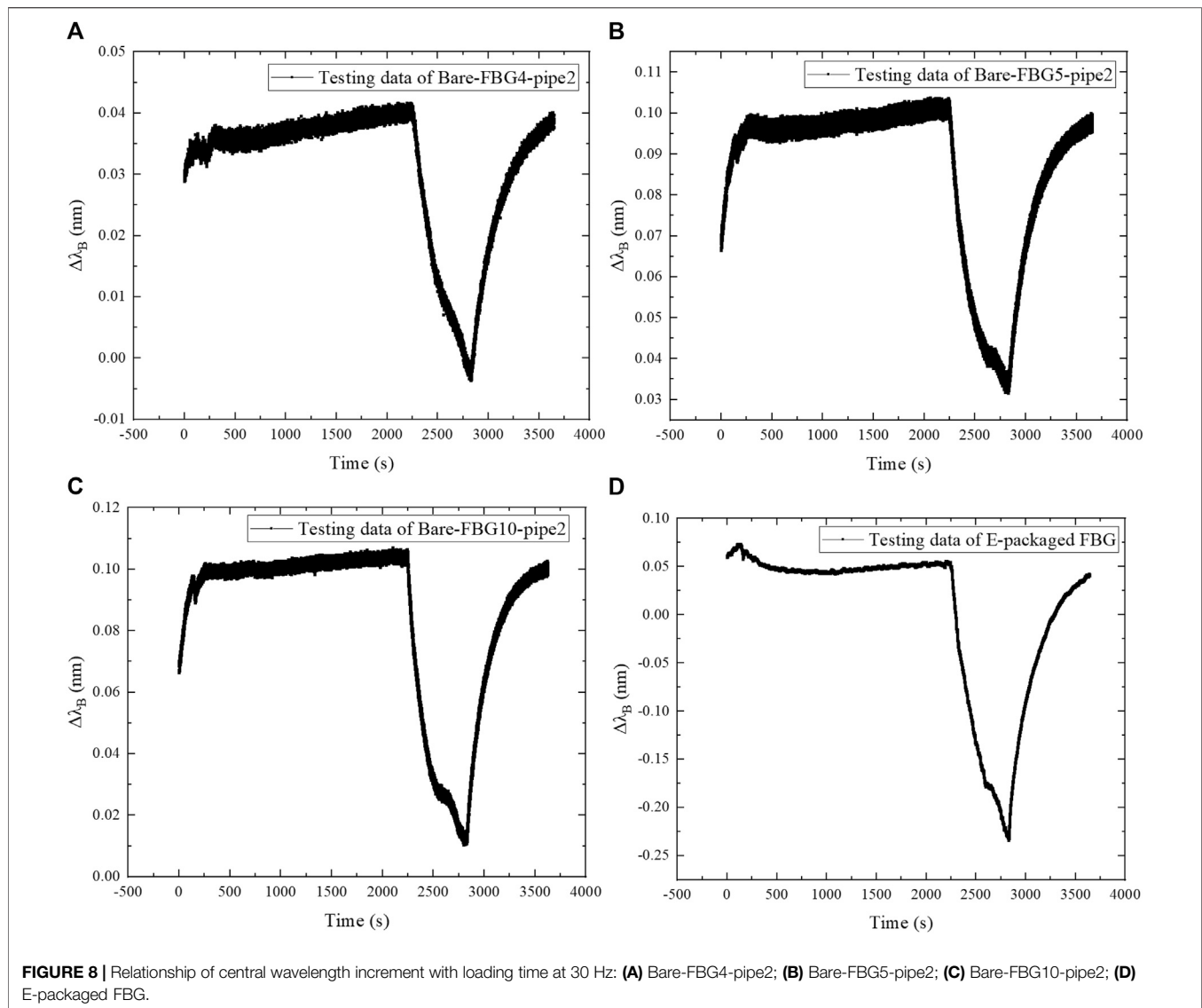
about the FBGs in series on pipe3 follows **Figure 5A**. Silicone rubber packaged FBG sensor has also been installed close to the vibration point.

Discussion on the Results

In the vibration experiment, the effect of the frequency change on the response information of the CFRP pipes and the effectiveness of the FBG measurement has been checked by changing the excitation frequency. The excitation frequencies acting on pipe2 include 10, 20, 30, and 50 Hz. Three bare FBGs in series attached on pipe2 are separately connected to channel 4, channel 5, and channel 6 of the demodulation device, and the E-packaged FBG is connected to channel 8. A simple harmonic force with a frequency of 50 Hz has been applied to the end of pipe2, and the loading time lasts about 2,000 s. After the completion of the first load, the excitation frequency is adjusted to 30 Hz, and the loading time of pipe2 lasts about 4,000 s. After that, the third load has been carried out, and the excitation frequency is 20 Hz with

the load time about 8,000 s. Finally, the fourth load with an excitation frequency of 10 Hz has been conducted, and the load time lasts about 2,000 s. An FBG from each channel has been taken out for analysis. The changes of FBG central wavelength increments with excitation frequencies are shown in **Figures 6–9**.

Figure 6 shows that the wavelength increment of the four FBG sensors has obvious consistency at the excitation frequency of 10 Hz, and they all increase with small fluctuations. That is, the internal deformation of the pipe accumulates after the action of applied simple harmonic excitation. Due to the disturbance in the loading process, the center wavelength of Bare-FBG5-pipe2 and Bare-FBG10-pipe2 has certain drop, but the overall trend is increasing. The change of FBG wavelength near the fixed end is significantly greater than that near the free end. It means that the vibration strain at the fixed end is larger than that at the free end, which validates the effective measurement. The strain fluctuation measured by E-packaged FBG near the fixed end is



small, indicating that epoxy resin packaging measure improves the stability of FBG in measuring dynamic strain.

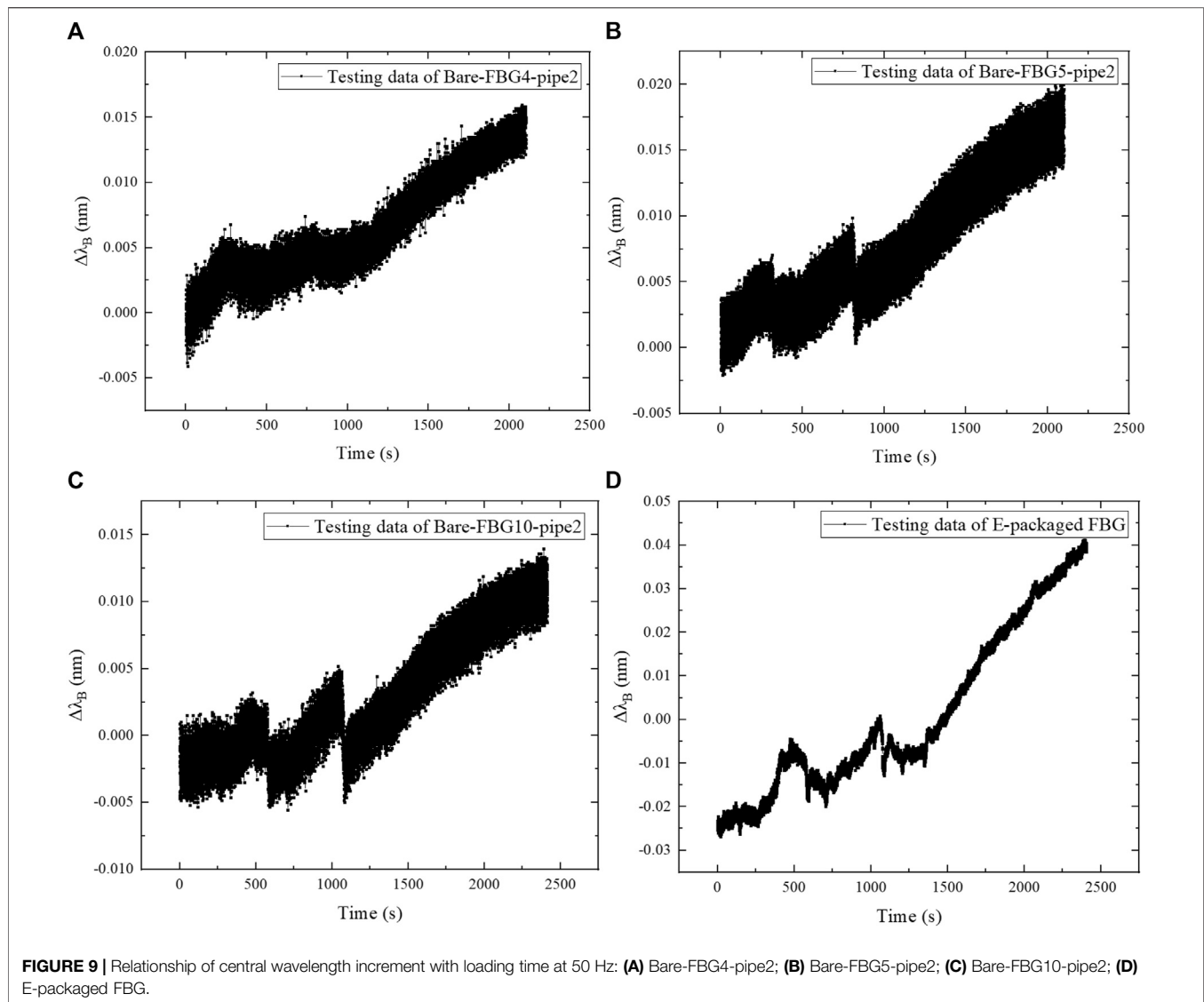
Figure 7 indicates that the variations of wavelength increments of the four FBG sensors at the excitation frequency of 20 Hz have high consistency. The wavelength of the FBG increases with fluctuation. That is, after harmonic excitation is applied, the pipe has a certain degree of deformation accumulation, and the degree is greater than that at 10 Hz. Due to the high sensitivity of FBG, the large “V” shape fluctuations in the figure can be attributed to the disturbance of the pipe. The overall deformation is relatively stable and has a certain rise. The vibration strain at the fixed end is greater than that at the free end, and the strain measured by the packaged FBG close to the fixed end has small fluctuation.

Figure 8 shows that the wavelength increments of the four FBG sensors under 30 Hz vibration have high consistency, and the wavelength of FBG fluctuates slightly (less than 0.005 nm). The small disturbance on the pipe produces the

large “V” shape fluctuation, which has been detected by the installed FBG sensors. The overall deformation is relatively stable, and the strain at the fixed end is greater than that at the free end.

Figure 9 shows that the wavelength increments of the four FBG sensors under 50 Hz vibration have high consistency. The wavelengths of FBGs fluctuate and increase cumulatively with the loading time, and the increments are smaller than that at 30, 20, and 10 Hz. Compared with the bare FBGs, the strain fluctuation measured by E-packaged FBG is smaller, indicating that the treatment of epoxy resin package protection is beneficial to the dynamic measurement of FBG.

In the vibration test, the excitation frequencies of pipe3 were 20, 30, and 50 Hz. The two bare FBGs in series attached on pipe3 have been connected to channel 2 and channel 3 of the demodulation device, and the S-packaged FBG has been connected to channel 8. A simple harmonic force with frequency 50 Hz has first been applied to the end of pipe3,



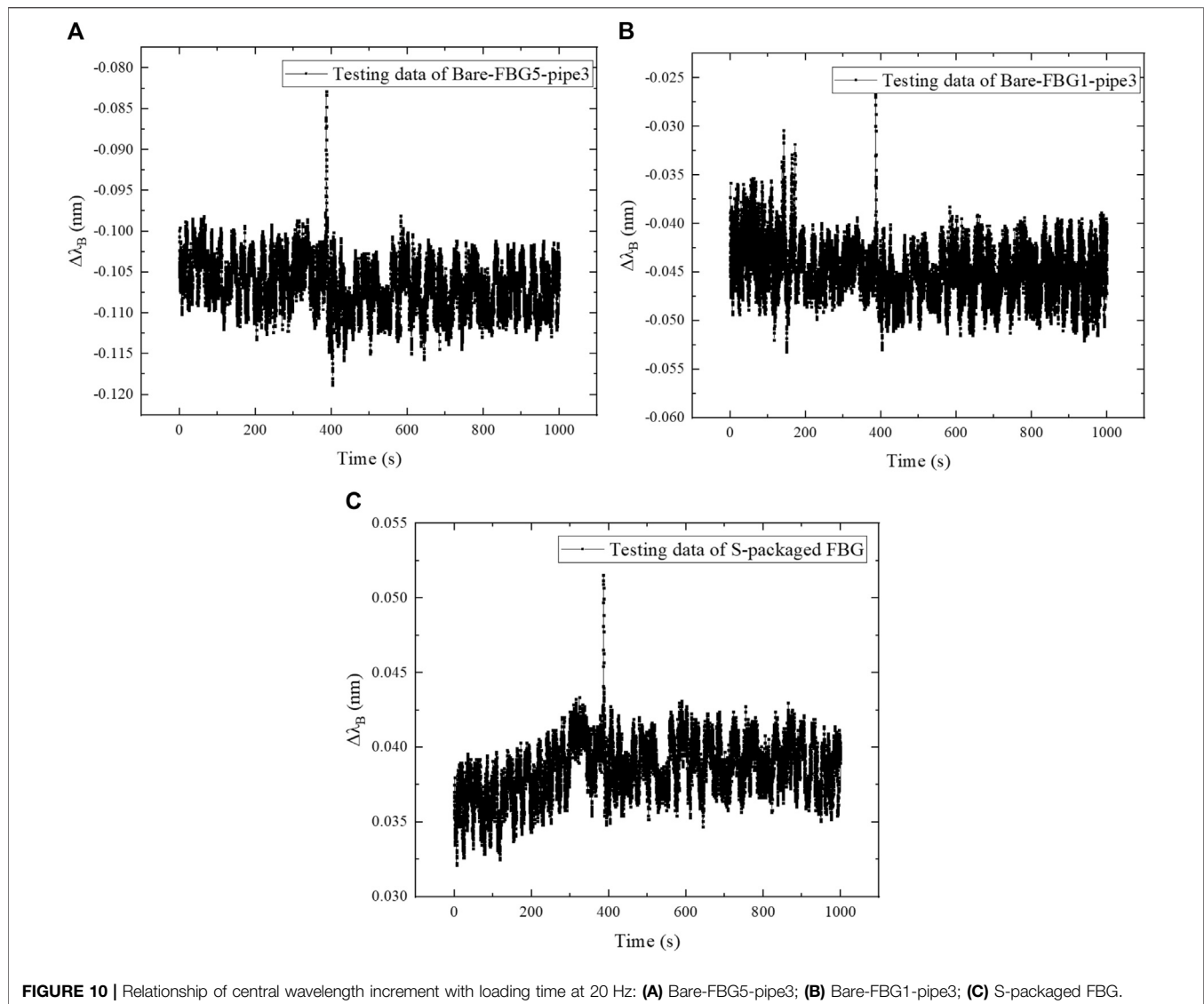
and the loading time lasts about 8,000 s. Second, adjusting the excitation frequency to 30 Hz, the load time on pipe3 lasts about 2,000 s. After that, the excitation frequency has changed to 20 Hz, and the loading time lasts about 1,000 s. An FBG from each channel was taken out for analysis. The FBG wavelength increment changes corresponding to the three excitation frequencies of 20, 30, and 50 Hz are shown in **Figures 10–12**.

Figure 10 shows that the wavelength increment of FBG sensors has very high consistency under the simple harmonic force with the excitation frequency 20 Hz, and the wavelength of FBG fluctuates steadily. The vibration strain at the fixed end is $0.15 \mu\epsilon$ greater than $0.1 \mu\epsilon$ at the free end. The measured strain of S-packaged FBG is smaller than that of bare FBG at the same position.

Figure 11 indicates that the wavelength increment of FBG sensors has very high consistency under the simple

harmonic force with 30 Hz. The wavelength of the FBG fluctuates steadily, and the vibration strain at the fixed end is larger than that at the far end. The strain of FBG with silicone rubber encapsulation is smaller than that of the bare FBG in the same position, which declares the obvious desensitization effect induced by the flexible package protection.

Figure 12 shows that the wavelength increment of FBG sensors has very high consistency under the simple harmonic force with 50 Hz. The wavelength of FBG fluctuates smoothly and steadily. The vibration strain at the fixed end is greater than that at the far end. The strain measured by silicone rubber encapsulated FBG is smaller than that measured by bare FBG at the same position, with relatively large wavelength increase. It indicates that large deformation accumulation occurs in the flexible silicone rubber material during the energy transferring process at high vibration frequency.



COMPARISON AND DISCUSSIONS

According to the aforementioned analysis on forced vibration of the cantilever beam, when a simple harmonic excitation is applied at the end of the beam, the strain response of the beam section at any point x and any time t can be described by Eq. 23. For the material parameters of the CFRP pipe, the modulus of elasticity E is 178.5 GPa obtained from the basic tensile testing, and the density of the pipe ρ is 1800 kg/m³. The simple harmonic excitation force $f = F \sin(\omega t)$ acts on the free end of the pipe. The first three-order steady-state responses of the pipe under the action of a single excitation force are considered in the calculation. The calculated values of vibration strains at $x = 0.75\text{m}$ and $x = 0.2\text{m}$, and the points located at adjacent FBGs are extracted from the programmed calculation and compared with the strains measured by FBGs. The distribution of the calculated and measured vibration strains with time is given in Figures 12, 13.

Figures 13A,B,C are separately corresponding to the strains measured by Bare-FBG2-pipe2, Bare-FBG3-pipe2, and Bare-FBG4-pipe2 attached on positions close to the fixed end of pipe2. It can be seen that the strain increases from the free end to the fixed end, and the average values of the amplitudes of the strains are 4.761, 4.15, and 3.907 $\mu\epsilon$ in order, showing a very uniform strain variation. The vibration frequencies and amplitudes of the measured and calculated strain responses are very close, and the amplitude error is around 0.7 $\mu\epsilon$.

Figures 13D,E,F are separately corresponding to the strains measured by Bare-FBG5-pipe2, Bare-FBG6-pipe2, and Bare-FBG7-pipe2 attached on positions close to the middle of pipe2. As can be seen from the figures, the strain amplitude closer to the fixed end is greater, and the values with relatively uniform variation are 3.866, 4.374, and 4.683 $\mu\epsilon$.

Figures 14A,B,C are separately corresponding to the strains measured by Bare-FBG2-pipe3, Bare-FBG3-pipe3, and Bare-FBG4-pipe3 attached on positions close to the fixed end of

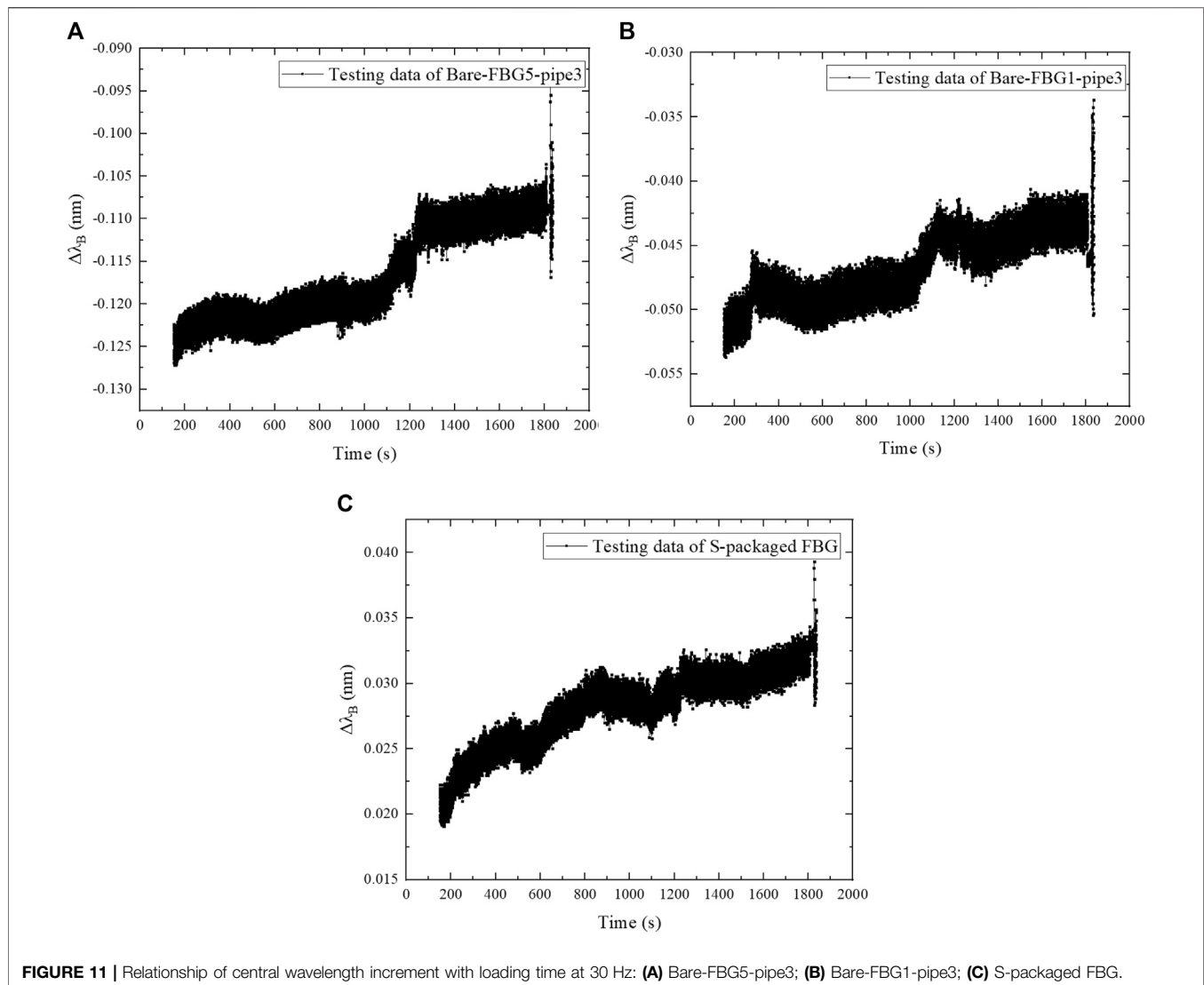


FIGURE 11 | Relationship of central wavelength increment with loading time at 30 Hz: **(A)** Bare-FBG5-pipe3; **(B)** Bare-FBG1-pipe3; **(C)** S-packaged FBG.

pipe3. The vibration strains vary uniformly, with the average amplitudes separately equal to 10.475, 9.975, and 10.375 $\mu\epsilon$. It can be seen from the figure that the strains present uniform change, and the vibration frequency and amplitude are very close, with the amplitude error about 0.7 $\mu\epsilon$. The difference from the previous ones is that the experimentally measured vibration curve shows a fluctuating change, which has a certain deviation from the theoretical calculation value but with close strain amplitudes. This phenomenon may be caused by insufficient clamping of the fixed end and large vibration energy at the loading end.

Figures 14D,E,F are separately corresponding to the strains measured by Bare-FBG5-pipe3, Bare-FBG6-pipe3, and Bare-FBG7-pipe3 attached on positions close to the free end of pipe3. It can be seen from the figure that the strain amplitudes are relatively uniform, with respective values of 8.443, 9.05, and 8.542 $\mu\epsilon$, which may be due to the close

distance between the FBG measuring points (the spacing is about 1 cm).

Statistical analysis on the data presented in **Figures 13E, 14D** has been displayed in **Table 1**. The output parameters indicate the high measurement accuracy and precision of the two sensing points. Similar analysis can also be extended to other points at different moments. It should be noted that for the degradation of the vibration device, the measurement accuracy and precision may be different. However, the good data indicate the effectiveness of the proposed technique.

Through comparison analysis, it can be seen that the error between the peak values of the theoretical vibration strain time-history curve and the FBG-measured curve is very small. That is, the peak value of the measured curve can much accurately reflect the vibration deformation response of the submarine pipes, and it is effective for adopting the proposed FBG sensors to measure the vibration strain. Based on the study,

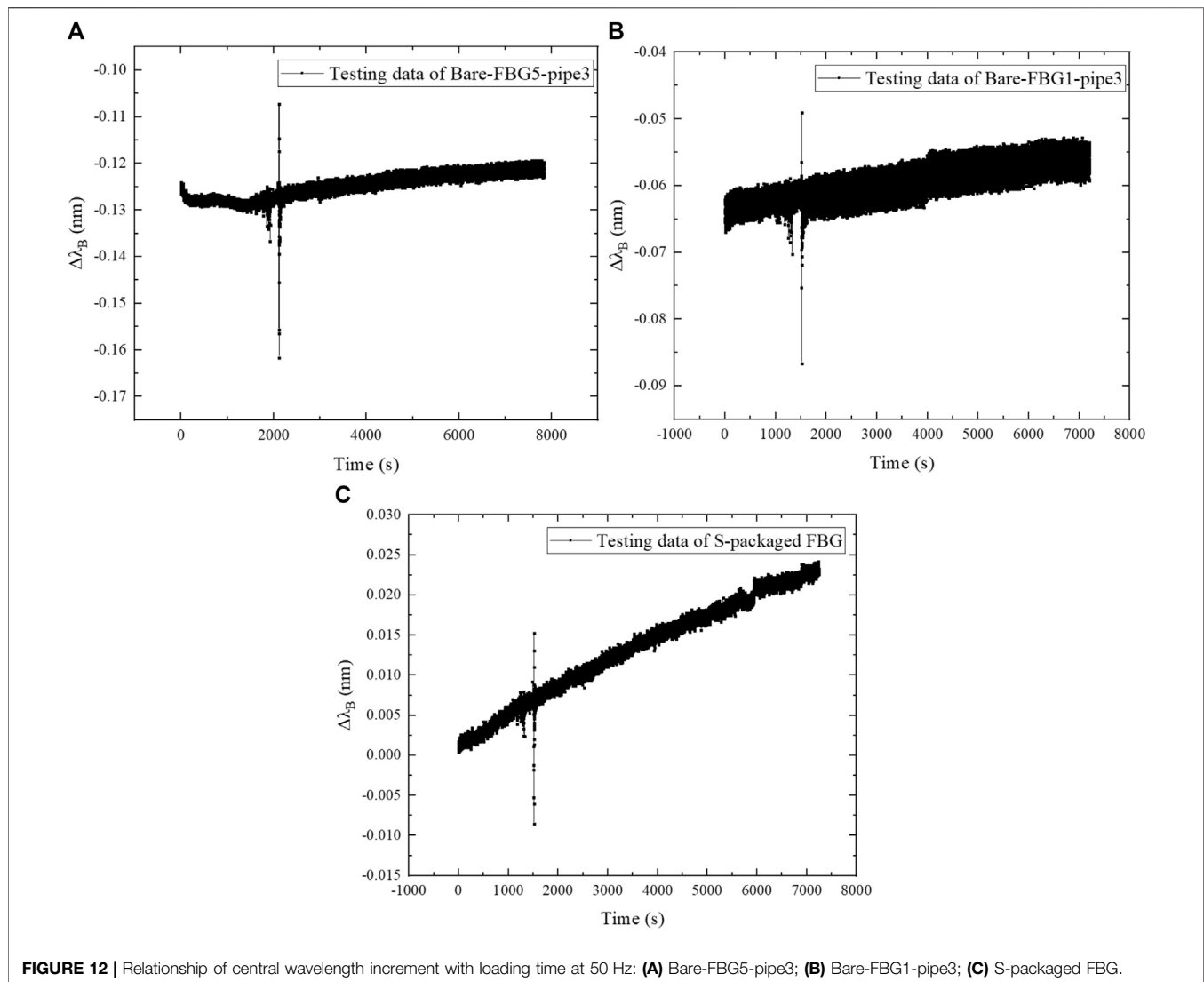


FIGURE 12 | Relationship of central wavelength increment with loading time at 50 Hz: **(A)** Bare-FBG5-pipe3; **(B)** Bare-FBG1-pipe3; **(C)** S-packaged FBG.

further work will be conducted to consider the random vibration-induced dynamic response of CFRP pipes based on the proposed FBGs in series and packaged FBG sensors. Furthermore, it is feasible to describe the dynamic deformation state of submarine pipelines under complex working conditions according to the information measured by FBGs.

Based on the fundamental study, the proposed sensing techniques can be further used to characterize the dynamic response of CFRP pipes under unknown excitation. It can be a challenging issue to configure the load spectrum and assess the structural performance of the pipes based on the effective data measured by FBG sensors.

CONCLUSION

CFRP composite pipes have been considered as the replacement of steel pipes in the transportation of oil and gas under deep submarine environment. Thus, the vibration characteristics of the

CFRP pipes have been investigated by theoretical analysis and experimental testing. Measurement on vibration responses of the CFRP pipes have been performed based on proposed quasi-distributed optical fiber sensing technology. The following conclusions can be drawn from the study:

- 1) The dynamic response of the pipe beam model is derived from the continuum vibration differential equations, and the comparison analysis validates the accuracy of this simplified theoretical model to describe the bending vibration-induced response.
- 2) The designed FBGs in series attached on the CFRP pipes have measured the dynamic strain responses with much high sensitivity and accuracy, which indicates the feasibility of the proposed monitoring technique for detecting the dynamic performance of composite pipes under arbitrary excitation.
- 3) The sensing performance of two packaged FBG sensors has also been checked, and the reasonable outputs indicate the

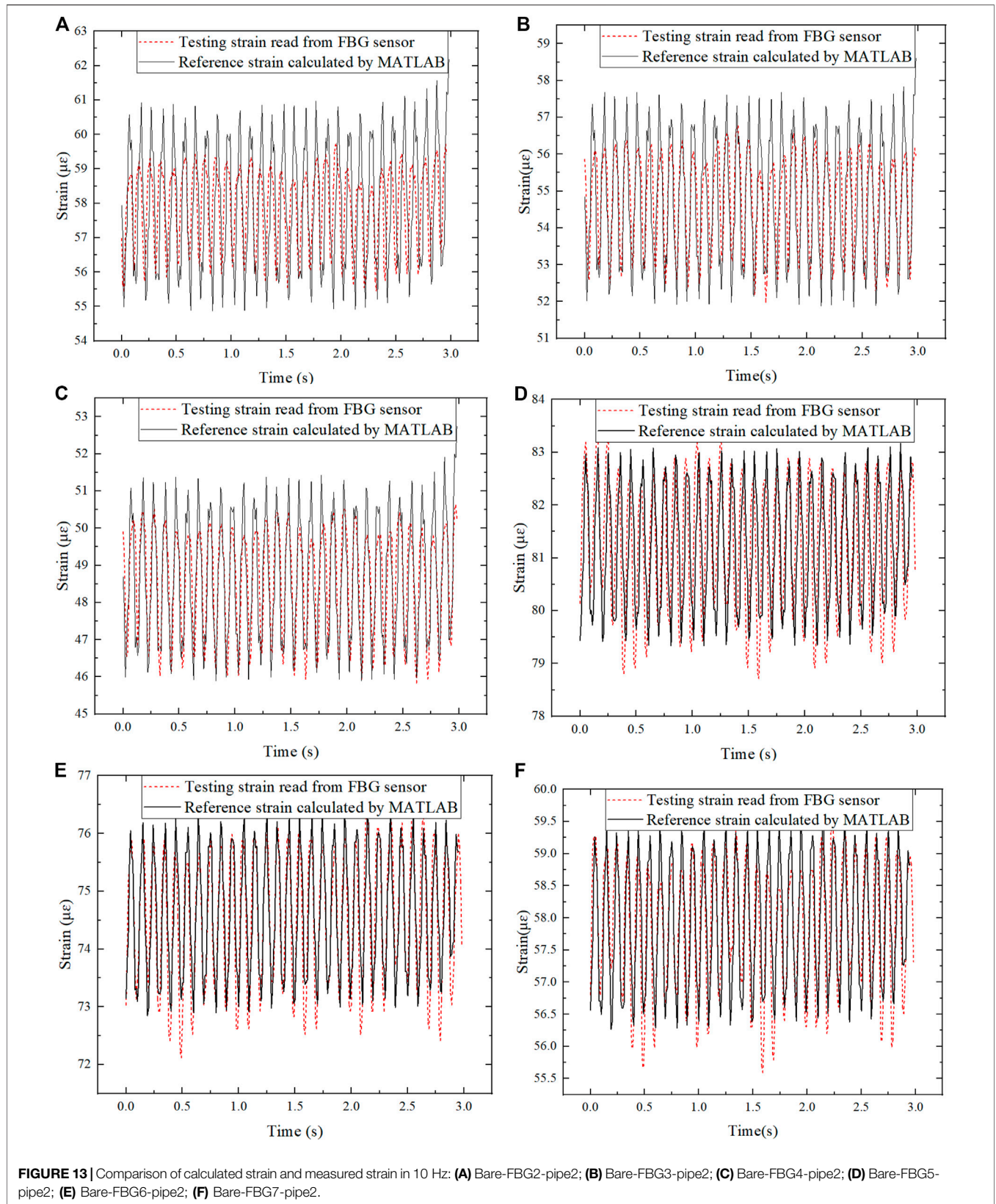


FIGURE 13 | Comparison of calculated strain and measured strain in 10 Hz: **(A)** Bare-FBG2-pipe2; **(B)** Bare-FBG3-pipe2; **(C)** Bare-FBG4-pipe2; **(D)** Bare-FBG5-pipe2; **(E)** Bare-FBG6-pipe2; **(F)** Bare-FBG7-pipe2.

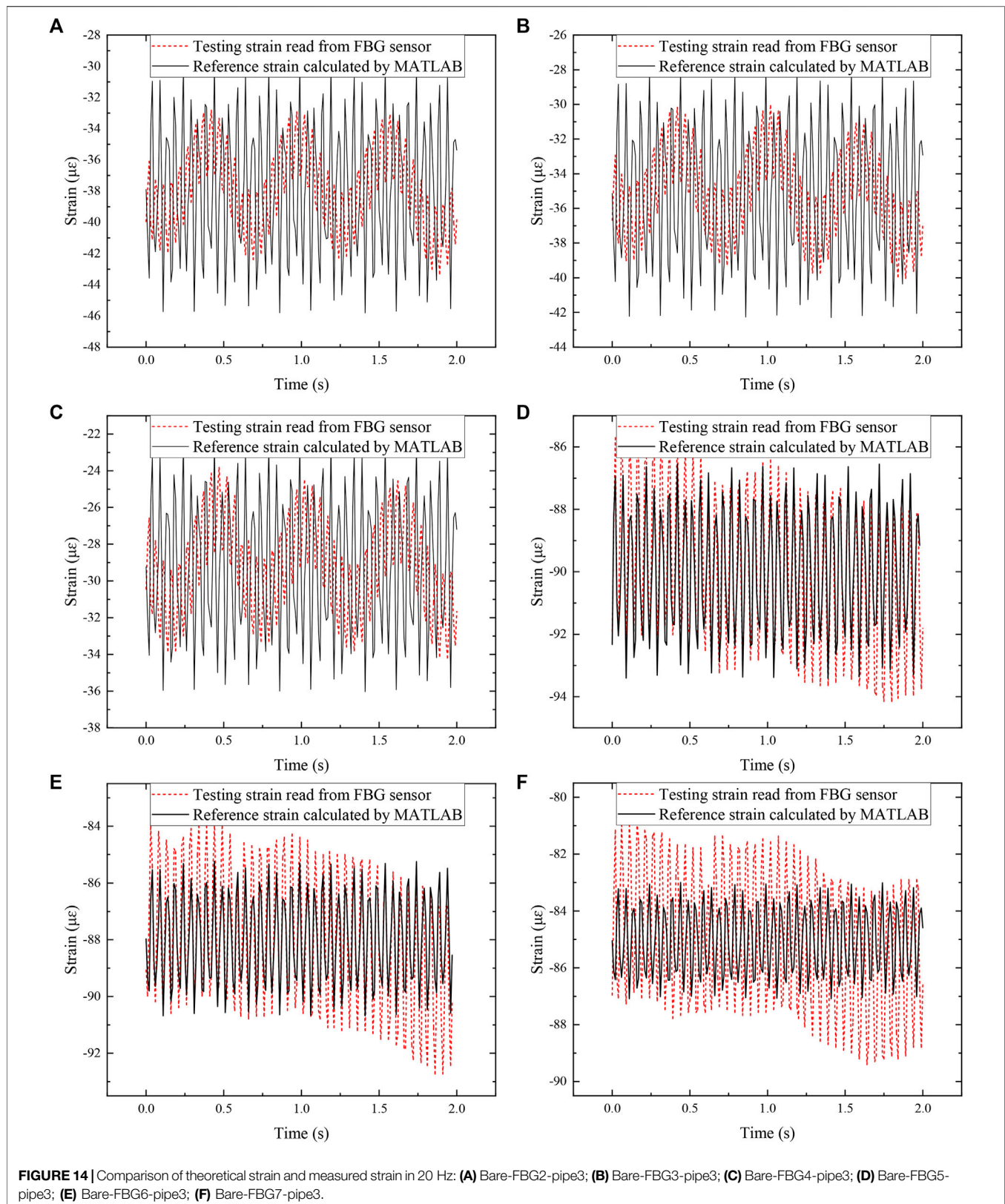


TABLE 1 | Statistical analysis on the measured data.

Item	Bare-FBG6-pipe2		Bare-FBG5-pipe3	
	Testing value	Theoretical value	Testing value	Theoretical value
Mean value	74.55	74.64	-89.96	-89.98
Standard error	0.07	0.06	0.17	0.14
Median	74.76	74.64	-90.08	-89.92
Standard deviation	1.16	1.05	2.39	1.96
Variance	1.35	1.11	5.73	3.86
Kurtosis	-1.34	-1.37	-1.23	-1.20
Min value	72.12	72.85	-94.15	-93.42
Max value	76.38	76.36	-85.71	-86.54
Sum	22,067.73	22,093.54	-17902.42	-17905.58
Number	296	296	199	199

effectiveness of the measurement. The results also indicate that the epoxy resin packaged FBG sensor installed on the CFRP composite pipes provides good measurement stability, which can be considered for the detection of submarine pipes under large dynamic deformation.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/Supplementary Material; further inquiries can be directed to the corresponding authors.

AUTHOR CONTRIBUTIONS

H-PW and S-YF contributed to the main part of the manuscript. X-SG, Y-XG, and YF contributed to the experimental study. PX checked the manuscript.

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