



Topological Study of Phase-Separated Ag-Conducting Chalcogenide Glasses Using Peak Force Quantitative Nano-Mechanical Characterization

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Peak Force Quantitative Nano-Mechanical property mapping (PF-QNM) was used to explore the nanomechanical properties of binary $\text{Ge}_y\text{Se}_{100-y}$ glasses and phase-separated $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses. The indentation modulus decreases when the selenium content increases in the binary glasses and when the silver content increases in phase-separated ternary glasses as a consequence of the increasing flexibility of the network. At nanoscale level, PF-QNM measurements highlighted a higher indentation modulus for the Ag-rich phase than for the Ag-poor phase, together with a decrease of both local indentation moduli with the increase of silver content in the glass. A model to gain a further insight in the mean coordination numbers and compositions of the Ag-rich and Ag-poor phases, based upon the assumption of a decoupled effect of the variations in Ag and in Ge on the mechanical properties, is described in detail. In the framework of the model, the mean coordination numbers for the Ag-rich phase $\langle N_R \rangle$ and Ag-poor phase $\langle N_P \rangle$ keep on decreasing when the total silver content in the glasses increases. The Ge concentration differs in Ag-rich and Ag-poor phases. When the silver content in the glass increases, the Ge concentration decreases in both phases. The result, in turn, indicates that an increase in the conductivity of the Ag-rich phase when the total silver content increases, as shown previously by C-AFM, would not be only explained by an increase in silver content but also by a softening of the network and consequently by an increase of charge carrier mobility.

Keywords: Ag-doped chalcogenide glasses, phase separation, ionic conduction, near field microscopy, nano-mechanical mapping

INTRODUCTION

Owing to their remarkable properties (IR transparency, high ionic conduction, photo-diffusion, etc.), chalcogenide glasses aroused strong interest in both the academic community and private companies. Among potential applications, the Conductive-Bridging Random Access Memories (CB-RAM), susceptible to replace the current flash memories, take advantage of two properties

of chalcogenide films, their sensitivity to photo-diffusion and their high ionic conductivity when doped with silver. A CB-RAM memory typically comprises a Ag-photo-doped glassy thin film of composition $\sim\text{Ge}_{0.25}\text{Se}_{0.75}$ placed between two electrodes, a silver one and a nickel one for example. The conductivity of the film is reversibly changed by several orders of magnitude when a low voltage is applied ($\sim 0.2\text{V}$). When applied to nanometric devices (when the thickness of the glassy film is typically 20–30 nm), the phenomenon is characterized by a very short time for commutation (~ 10 ns) and a very high cyclability ($> 10^6$ cycles) (Kozicki et al., 2005).

In order to explain the operation of this memory and in particular its very short switching time hardly compatible with the mobility of Ag^+ ions, various hypotheses have been proposed and among them, the presence of pre-existing silver-rich zones within the amorphous film—which would be heterogeneous—and the creation of conduction paths between these zones. To confirm this hypothesis and to better understand the phenomena at the origin of the switching from a resistive state to a conductive state in CB-RAM memories, the electrical, structural, and thermal properties of bulk Ag-Ge-Se glasses have been widely studied (Dejus et al., 1988, 1991, 1992; Kawasaki et al., 1999; Mitkova et al., 1999; Iyetomi et al., 2000; Piarristeguy et al., 2000, 2003, 2007a, 2012; Ureña et al., 2003, 2005; Wang et al., 2003; Cervinka et al., 2005; Tafen et al., 2005; Cuello et al., 2007; Prasai and Drabold, 2011; Le Parc et al., 2013; Stellhorn et al., 2016a,b; Zeidler et al., 2016, 2018). It was shown that the conductivity of $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses with $1 < x < 30$ at. % increases suddenly by seven to eight orders of magnitude at 7–10 at. % Ag.

Insight into the microstructure of these glasses was clearly needed. However, owing to the ease of diffusion of silver in these materials (Kolobov and Elliott, 1991), the use of classical microscopy techniques to probe the structure of Ag-based glasses is tricky; a risk of altering the Ag-centered local structures exists if the light or electron beam energy is too high. As a matter of fact, few papers on the electrical properties of chalcogenide glasses report attempts to get insights into the homogeneity of the materials.

Gutenev et al. (1991) and Wang et al. (2003) indicated for the first time that the glasses were phase separated according to optical microscopy and modulated differential scanning calorimetry (MDSC) experiments, respectively. However, these experiments cannot provide insight on, for example, the evolution of the composition or the conductivity of each phase with increasing silver content. Several years later, the ability and the complementarity of near field microscopy techniques provided insight into the intrinsic electrical properties of each phase present in bulk $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses and their composition dependence (Balan et al., 2006; Piarristeguy et al., 2007b, 2010a,b). Electrostatic force microscopy (EFM) was indeed used to get images of electrical contrasts (dielectric permittivity) in phase separated glasses. Combined with electrical conductivity measurements using complex impedance spectroscopy, EFM helped in understanding the jump in conductivity at 7–10 at% in $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses. While a Ag-rich phase is embedded in a Ag-poor phase for the low conductivity region ($x < 7$ –10 at. % Ag), the reverse is true for

high conductivity region ($x > 7$ –10 at. % Ag) with the Ag-rich phase controlling the silver diffusion throughout the material. The percolation of the Ag-rich phase was at the origin of a sudden jump in conductivity (Balan et al., 2006; Piarristeguy et al., 2007b).

Recently another near field microscopy technique, i.e., the Contact Resonance Atomic Force Microscopy (CR-AFM) in combination with Raman mapping was used to get additional information on the phase compositions in bulk $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses (Piarristeguy et al., 2018). The results suggested a different Ge/Se ratio in Ag-poor and Ag-rich phases, the first one being richer in selenium than the second one.

In this work, an innovative and complementary technique, the Peak Force Quantitative Nano-Mechanical mapping (PF-QNM) was used to explore the nano-mechanical properties of $\text{Ge}_y\text{Se}_{100-y}$ and $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses and get information on the topology of these phase-separated glasses. A model, based upon the assumption of a decoupled effect of the variations in Ag and in Ge on the mechanical properties, helped to gain a further insight in the mean coordination numbers and compositions of the Ag-rich and Ag-poor phases.

MATERIALS AND METHODS

Sample Preparation and Characterization

Bulk $\text{Ge}_y\text{Se}_{100-y}$ glasses ($y = 15, 20, 25, 30,$ and 33 at. %) and $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses ($x = 0, 7, 15,$ and 20 at. %) were synthesized from a mixture of high-purity starting elements [Ag (Aldrich, 99.9 %), Ge (Aldrich, 99.999 %) and Se (Aldrich, 99.99 %)] Three grams of the stoichiometric powders were placed in a cylindrical quartz ampoule (11 mm inner diameter, 1 mm thick) sealed under secondary vacuum (10^{-5} mbar). The tube was subsequently heated with a low heating rate of $6^\circ\text{C}/\text{h}$ up to 600°C for $\text{Ge}_{15}\text{Se}_{85}$, 700°C for $\text{Ge}_{20}\text{Se}_{80}$ and 850°C for $\text{Ge}_{30}\text{Se}_{70}$ and $\text{Ge}_{33}\text{Se}_{67}$, held at this temperature for 12 h, and finally quenched in air. For the $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses, the powder was heated up to 950°C with a $10^\circ\text{C}/\text{h}$ heating rate and kept at this temperature for 8 h for homogenization before being quenched in a salt-ice-water mixture. The samples were named Ag_x according to the concentration of Ag (in at. %) and the nomenclature R-Ag_x and P-Ag_x was used for the Ag-rich and Ag-poor phases, respectively.

The amorphous nature of the ingots was confirmed by performing X-Ray diffraction using a PANalytical X'Pert diffractometer. A Cu ($K\alpha$) source ($\lambda = 1.5406\text{\AA}$) was used with an operating voltage of 40 kV and a beam current of 30–40 mA.

Peak Force Quantitative Nano-Mechanical Property Mapping (PF-QNM)

The mechanical characterization was performed with a Multimode AFM instrument (from Bruker Corporation, USA) upgraded with Nanoscope V using PF-QNM imaging mode and a special RTESPA-525-30 (BatchA050 wafer1 Box9 probe#3) (Bruker) probe with a lever spring constant of 167 N/m and a tip radius of 34 nm. A FSILICA-12M sample from Bruker Instruments was first used to calibrate sensitivity deflection and sync distance QNM. The calibration was checked by

measuring a sample (given by Arnould et al., 2017) consisting in a mixture of resin and fiber of K28 grade Kevlar, providing indentation moduli of around $4 \text{ GPa} \pm 0.4 \text{ GPa}$ and $23 \text{ GPa} \pm 4 \text{ GPa}$, respectively, in agreement with the data of Arnould et al. (2017). The applied maximum load was set at 200 nN for all the measurements. Different mechanical properties, i.e., adhesion force, elastic modulus, deformation, energy dissipation, were obtained by a real time analysis of the “force versus separation” curves. The well-known Derjaguin-Muller-Toporov (DMT) model (Derjaguin et al., 1975) was used to fit the retract curves and calculate the indentation moduli (Poisson coefficient set at 0).

Measurements were performed on the surfaces of freshly fractured glass to avoid any problem due to surface oxidation. The experiments were carried out within the following hours.

RESULTS

Ge-Se Glasses

PF-QNM experiments allowed getting information about the mechanical properties of glasses. The indentation moduli, reported for glasses for the first time, were estimated from data recorded over $5 \times 5 \mu\text{m}$ images. **Figure 1** shows the indentation moduli for the different $\text{Ge}_y\text{Se}_{100-y}$ glasses. A monotonous increase in modulus with increasing germanium concentration (y) is observed, which is in agreement with Vickers and Meyer’s hardness and Young’s modulus measured by conventional mechanical techniques (Guin et al., 2002). The insert in **Figure 1** shows both the indentation moduli obtained by PF-QNM measurements and the Young’s moduli for the same glassy compositions (Guin et al., 2002). Similar trends are observed, which allows to validate the PF-QNM technique to evaluate realistic glass indentation moduli.

Because of similar Se-Se and Se-Ge bond strengths (223 and 230 kJ/mol, respectively) and only slightly weaker Ge-Ge bond strength (186 kJ/mol), Guin et al. (2002) claimed that the source for the increase of the hardness has to be topological in nature. Consequently, the continuous increase of hardness is of the same origin as the increase of the elastic moduli, and it results from the increasing number of strong covalent bonds with increasing Ge concentration.

Ag-Ge-Se Glasses

Figure 2 shows the topography and the indentation modulus images (PF-QNM) for the $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses. These images (**Figures 2b,d,f**) clearly show heterogeneous domains, which indicates that the PF-QNM technique is sensitive to modulation of mechanical properties at the local stage. The heterogeneous domains are similar to those already reported for this family of glasses using Electrostatic Force Microscopy (EFM) and Conductive-Atomic Force Microscopy (C-AFM) (Balan et al., 2006; Piarristeguy et al., 2007b, 2010a,b). For glassy compositions belonging to the low-conductivity region ($x < 10$ at. % Ag, **Figure 2b**), nodules of a Ag-rich (R-Agx) phase (light region) are embedded in a connected Ag-poor (P-Agx) phase (dark region). At the opposite, for the glassy compositions belonging to the high conductivity region

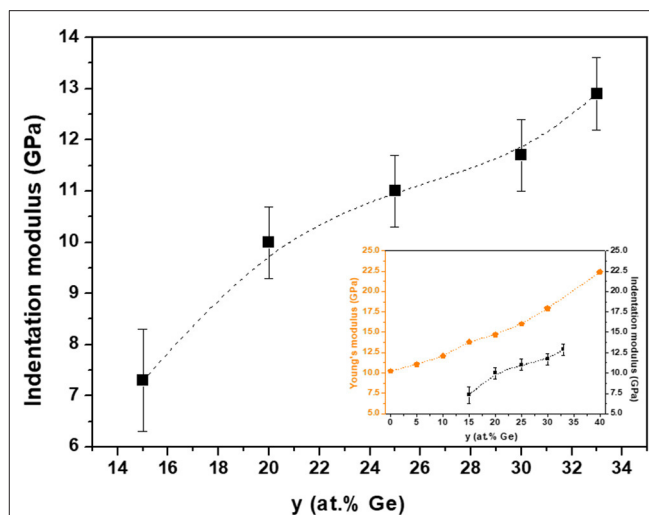
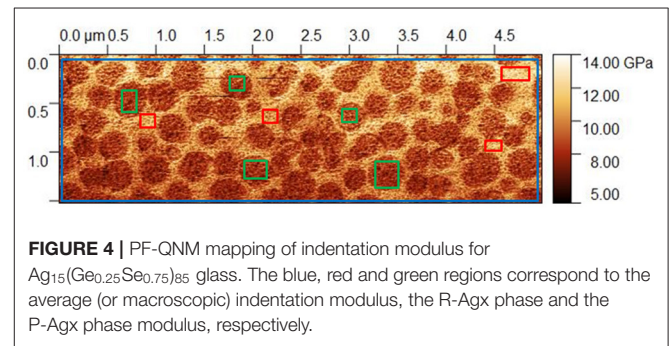
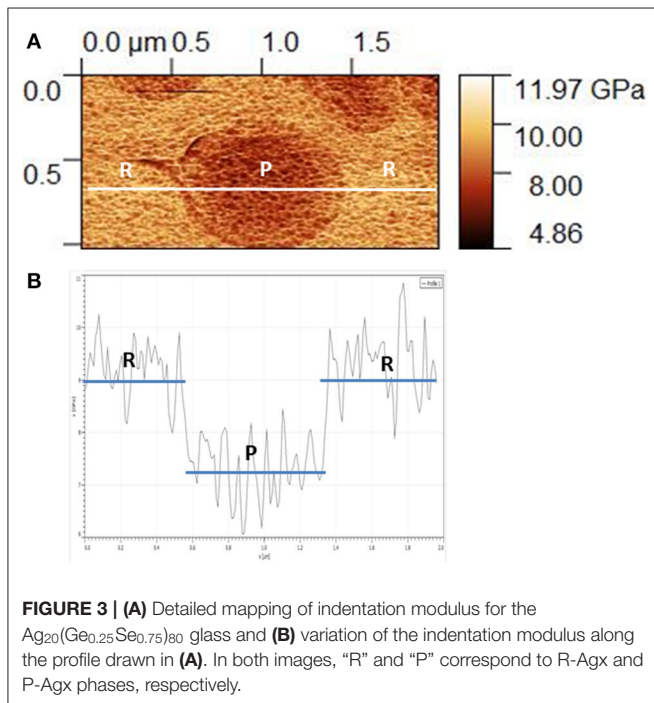
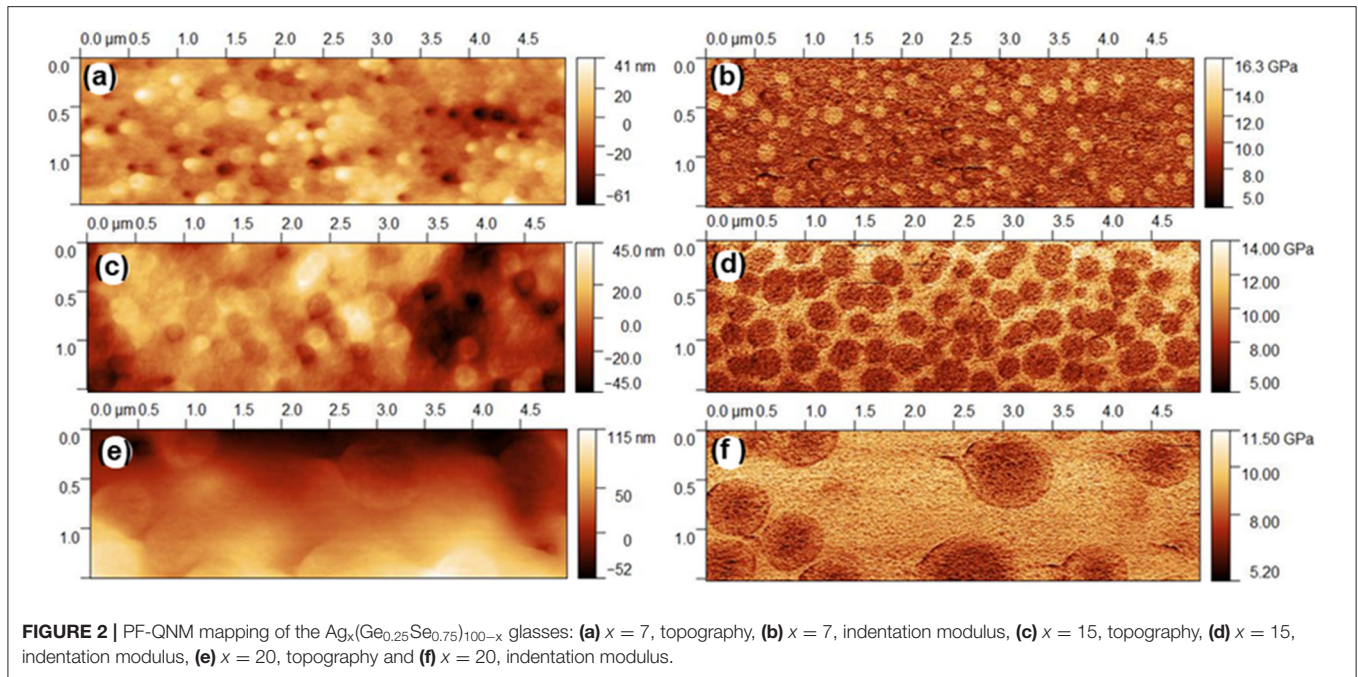


FIGURE 1 | Indentation moduli for bulk $\text{Ge}_y\text{Se}_{100-y}$ glasses. Insert shows the comparison between Young’s modulus (Guin et al., 2002) and indentation modulus obtained by PF-QNM measurements. The broken lines are drawn as guides for the eye. The indentation modulus values and the error bars were calculated from an average of three PF-QNM measurements and standard deviation evaluation.

($x = 15$ and 20), a connected R-Agx phase contains nodules of P-Agx phase (**Figures 2d,f**). Depending upon the quenching conditions, spinodal or nucleation/growth phase separation can occur in these glasses (Piarristeguy et al., 2018). In this work, the images show a morphology typical of nucleation/growth phase separation.

As an example, the analysis of a small zone of the indentation modulus mapping for a typical phase-separated $\text{Ag}_{20}(\text{Ge}_{0.25}\text{Se}_{0.75})_{80}$ glass is detailed (**Figure 3**). This glass presents nodules of about $0.7 \mu\text{m}$ (P-Agx phase) immersed in a conductive matrix (R-Agx phase). The indentation moduli of both phases are visible in **Figure 3A**. The variation of modulus along the white line in **Figure 3A** is shown in **Figure 3B**. The modulus of the R-Agx phase, i.e., $\sim 9 \text{ GPa}$, corresponds to the domains of the line from 0 to $\sim 0.5 \mu\text{m}$ and from ~ 1.3 to $2 \mu\text{m}$. The modulus for the R-Agx phase between ~ 0.5 and $\sim 1.3 \mu\text{m}$ has a value of $\sim 7.2 \text{ GPa}$.

Three different values of indentation modulus for each of these phase-separated glasses were extracted from the PF-QNM experiments. The first one corresponds to an average value and is evaluated by taking into account the whole region included in the blue parallelepiped in **Figure 4**, the second one corresponds to the indentation modulus for the R-Agx phase and is obtained from regions included in the red parallelepiped in **Figure 4**, the third one corresponds to the P-Agx phase (regions in the green parallelepiped in **Figure 4**). The average indentation modulus, estimated from data recorded over the whole $5 \times 1.5 \mu\text{m}$ image is equivalent to a “macroscopic measurement.” At least thirty different regions were taken into in the PF-QNM image to estimate the value of and the uncertainty on the indentation modulus for each phase.



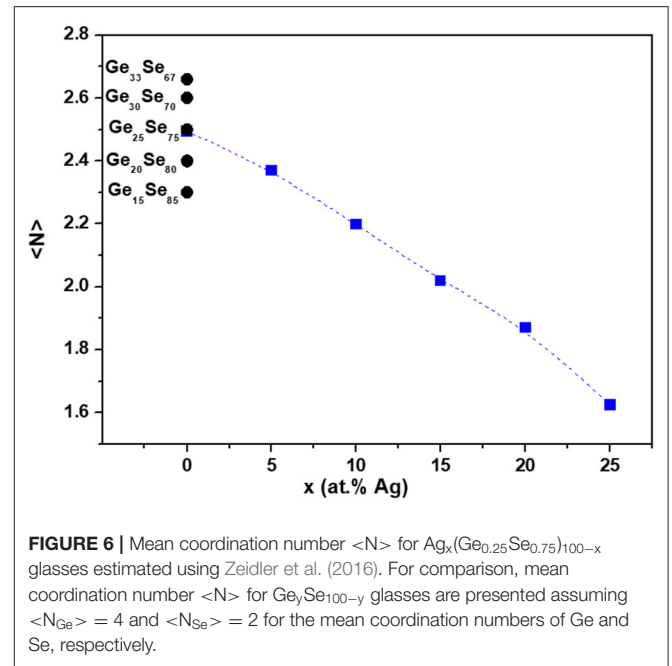
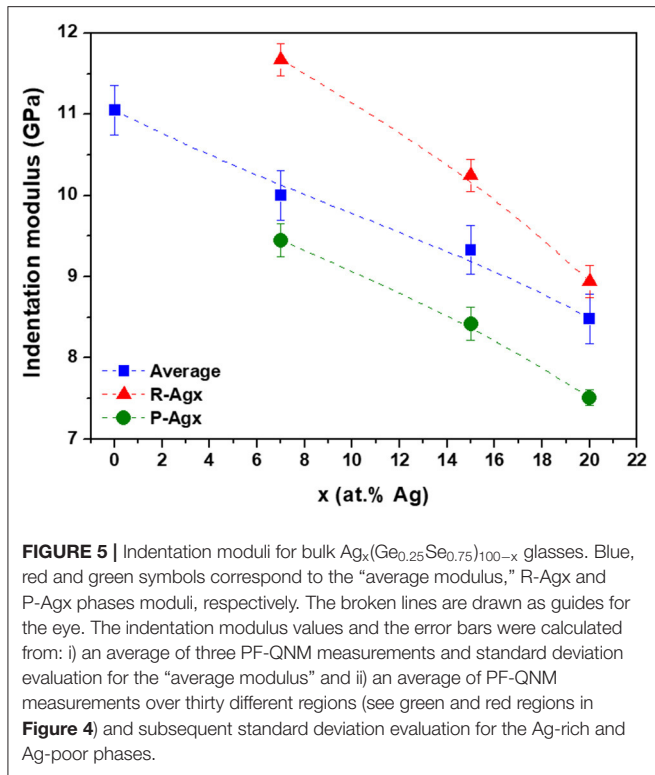
the $\text{Ge}_{25}\text{Se}_{75}$ glass. Introduction of silver affects the tetrahedral network through changes in bond angles, eventual ES to CS tetrahedra transformation associated to network disruption and selenium chains breaking. At high silver content, a high flexibility of the network is expected as a consequence of depolymerization of the glass network. Based on this statement, the local mechanical measurements present an apparent contradiction since the indentation modulus value is higher for the R-Agx phase than for the P-Agx phase (it will be discussed in the next section). A decrease of both local indentation moduli with the increase of the silver content in the glass is also observed.

DISCUSSION

Mean Coordination Number and Composition

The results above show a decrease in the “average” indentation modulus M (symbols in blue, **Figure 5**) when the Ag content increases in the glasses. However, for a given Ag content, the

The indentation moduli for R-Agx and P-Agx phases as well as the “average modulus” as a function of the silver content are presented in **Figure 5**. The “average modulus” decreases with the increase of silver concentration in the glass, which is in agreement with previous Vickers microhardness measurements (Piarristeguy et al., 2018). This decrease is a consequence of the structural changes induced by the introduction of Ag in



R-Agx phase systematically exhibit a higher modulus than the P-Agx phase. The explanation has to be searched in difference in compositions and, consequently in average coordination numbers, between R-Agx and P-Agx.

In $\text{Ge}_y\text{Se}_{100-y}$ glasses, assuming $\langle N_{\text{Ge}} \rangle = 4$ and $\langle N_{\text{Se}} \rangle = 2$ for the mean coordination numbers of Ge and Se, respectively, the mean coordination number can be estimated

$$\langle N \rangle_{\text{Ge-Se}} = \frac{y}{100} \langle N_{\text{Ge}} \rangle + \frac{100 - y}{100} \langle N_{\text{Se}} \rangle \quad (1)$$

In $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses, when ionic bonds are participating to the structure, the determination of the mean coordination is more complex. Zeidler et al. (2018) have estimated the nearest-neighbor coordination numbers $\langle N_{ij} \rangle$ extracted from neutron diffraction experiments on glasses along the $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ tie line. Taking into account that the first coordination sphere implies only $\langle N_{\text{GeSe}} \rangle$, $\langle N_{\text{SeSe}} \rangle$, and $\langle N_{\text{SeGe}} \rangle$, the mean coordination number for this family of glasses can be estimated

$$\langle N \rangle_{\text{Ag-Ge-Se}} = \frac{0.25(100 - x)}{100} \langle N_{\text{GeSe}} \rangle + \frac{0.75(100 - x)}{100} [\langle N_{\text{SeGe}} \rangle + \langle N_{\text{SeSe}} \rangle] \quad (2)$$

Figure 6 shows the mean coordination number $\langle N \rangle$ for $\text{Ge}_y\text{Se}_{100-y}$ and $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses. As expected, we observe that the mean coordination number decreases monotonously with the increase of silver content in the glass,

with values comprised between $\langle N \rangle = 2.50$ for Ag0 to $\langle N \rangle = 1.62$ for Ag25. Silver incorporation is expected to depolymerize the network, searching for selenium to compensate its positive charge, leading to inter-tetrahedral Ge-Se-Ge, and selenium chains breaking. Therefore, the introduction of silver is expected to introduce some flexibility into the network.

An attempt to gain a further insight into the mean coordination numbers of the R-Agx and P-Agx phases in $\text{Ag}_x(\text{Ge}_z\text{Se}_{1-z})_{100-x}$ glasses is described below.

As a first approximation, the dependence of the indentation modulus M_{Ge} on Ge content in $\text{Ge}_y\text{Se}_{100-y}$ glasses, shown in **Figure 7** (black symbols), can be taken as linear with

$$\Delta M_{\text{Ge}} = 14.04 \Delta \langle N \rangle \quad (3)$$

In the same way, the dependence of the indentation modulus M_{Ag} on Ag content for $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses, shown in **Figure 7** (blue symbols), can be taken as linear with

$$\Delta M_{\text{Ag}} = -3.94 \Delta \langle N \rangle \quad (4)$$

On the whole, the effect of Ge modulation is almost three times more impacting on the indentation modulus than the evolution of the average Ag content and evolves in the opposite direction.

With the assumption that the difference of indentation moduli $\Delta M_{\text{R-P}} = M(\text{R-Agx}) - M(\text{P-Agx})$ between the Ag-rich phase ($\text{Ag}_{x'}(\text{Ge}_{z'}\text{Se}_{1-z'})_{100-x'}$) and the Ag-poor phase ($\text{Ag}_{x''}(\text{Ge}_{z''}\text{Se}_{1-z''})_{100-x''}$), for a given glass composition, can be related to the decoupled effect of both the variation in Ag content, $\Delta x'x'' = x' - x''$, and the variation in Ge content, $\Delta z'z'' = 100.z' - 100.z''$, one can write

$$\Delta M_{\text{R-P}} = \Delta M_{x'x''} + \Delta M_{z'z''} \quad (5)$$

where $\Delta M_{x'x''}$ is the modulus difference related to a Ag variation $\Delta x'x''$ and $\Delta M_{z'z''}$ a modulus difference related to Ge variation $\Delta z'z''$.

Taking into account Equations (3–4), the equation (5) can be written

$$\Delta M_{R-P} = 14.04 \Delta \langle N \rangle - 3.94 \Delta \langle N \rangle = 10.06 \Delta \langle N \rangle \quad (6)$$

The modulus difference between R-Agx and P-Agx phases, ΔM_{R-P} , for a given glass, can be calculated from the experimental data shown in **Figure 5**. These values are 2.33, 1.85, 1.43, for Ag7, Ag15, and Ag20 glasses respectively, which corresponds (see Equation 6) to $\Delta \langle N \rangle$ values of 0.23, 0.18, and 0.14 with the Ag content increasing in the glass.

Additional information can help to estimate the mean coordination number for R-Agx ($\langle N_R \rangle$) and P-Agx ($\langle N_P \rangle$) phases. The PF-QNM images (**Figures 2b,d,f**) indeed allow to calculate the surface percentage for the R-Agx (% R) and P-Agx (% P) phases. Assuming that the volume percentage is equivalent to the surface percentage, the mean coordination number is therefore a linear combination of the mean coordination number in R-Agx and P-Agx phases:

$$\langle N \rangle = (\%P) \langle N_P \rangle + (\%R) \langle N_R \rangle = (\%P) \langle N_P \rangle + (\%R)(\langle N_P \rangle + \Delta \langle N \rangle) \quad (7)$$

where we have assumed that $\langle N_R \rangle = \langle N_P \rangle + \Delta \langle N \rangle$ because R-Agx phase presents a high indentation modulus in comparison with the P-Agx phase. The mean coordination numbers $\langle N_R \rangle$ and $\langle N_P \rangle$ are presented in **Figure 7**.

It can be seen that, consistently with the average coordination number, both $\langle N_R \rangle$ and $\langle N_P \rangle$ decrease with increasing Ag content. These decreases in coordination number can be understood by a decrease in Ge (z' , z'') and/or an increase of

Ag (x' , x''). The differences $\langle N_R \rangle - \langle N_P \rangle$ for a specific Agx composition decreases when Ag content increases (respectively equal to 0.23, 0.18, 0.14 for Ag7, Ag15, and Ag20). Interestingly, depending upon the preparation procedure, the glass Ag 25, not measured here, is either homogeneous or shows only very small inhomogeneities.

As a matter of fact, compositions of the P-Agx and R-Agx phases should not only differ in Ag composition (x' , x''), but also in Ge composition (z' , z'') as already suggested by the results from CR-AFM data shown in a previous (Piarristeguy et al., 2018). Indeed, one expects a R-Agx phase with a higher Ge content (z') than the P-Agx one (see **Figure 8**).

In the following, we aim at gaining a further insight in the compositions of the R-Agx phase ($Ag_{x'}(Ge_{z'}Se_{1-z'})_{100-x'}$) and of the P-Agx phase ($Ag_{x''}(Ge_{z''}Se_{1-z''})_{100-x''}$). As already mentioned, according PF-QNM images (**Figures 2b,d,f**), the average composition is a linear combination of the compositions in R-Agx and P-Agx phases:

$$Ag_x(Ge_{0.25}Se_{0.75})_{100-x} = (\%R) Ag_{x'}(Ge_{z'}Se_{1-z'})_{100-x'} + (\%P) Ag_{x''}(Ge_{z''}Se_{1-z''})_{100-x''} \quad (8)$$

with

$$x = (\%R) x' + (\%P) x'' \quad (9)$$

$$0.25(100 - x) = (\%R) z'(100 - x') + (\%P) z''(100 - x'') \quad (10)$$

where (% R) and (% P) correspond, respectively to the surface percentage for R-Agx and P-Agx phases.

Combining these equations leads to a great number of potential compositions, however further restrictions on the

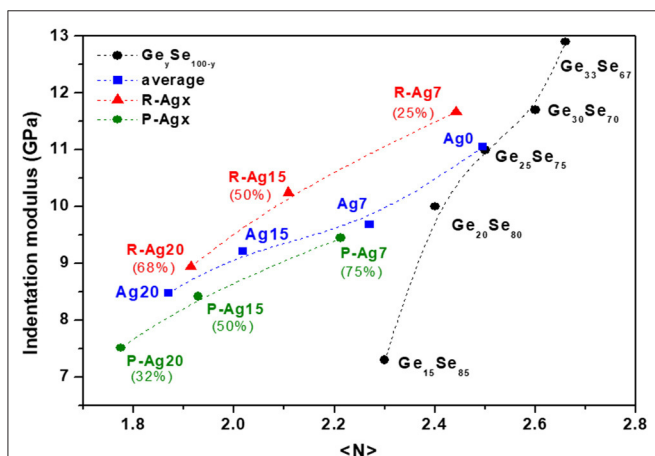


FIGURE 7 | Indentation moduli in function of the mean coordination number for Ge_ySe_{100-y} glasses (black symbols), $Ag_x(Ge_{0.25}Se_{0.75})_{100-x}$ glasses (bleu symbols), and R-Agx (red symbols) and P-Agx (green symbols) phases. Percentages correspond to the proportion of R-Agx and P-Agx phase for each composition Agx. The broken lines are drawn as guides for the eye.

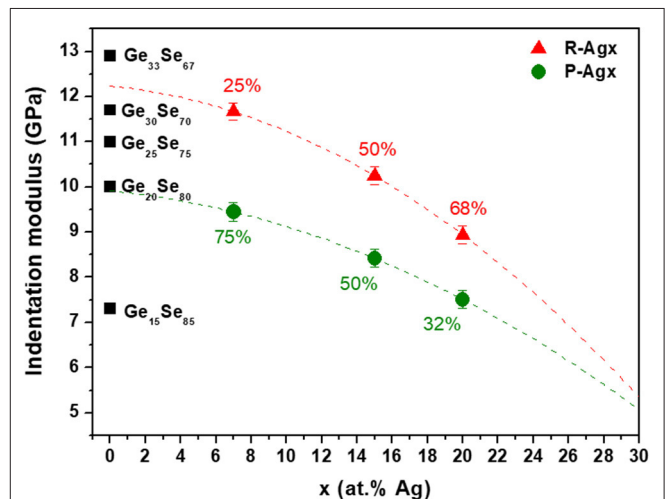


FIGURE 8 | Indentation moduli for R-Agx (red symbols) and P-Agx (green symbols) phases in $Ag_x(Ge_{0.25}Se_{0.75})_{100-x}$ glasses and for Ge_ySe_{100-y} glasses (black symbols). Percentages correspond to the proportion of R-Agx and P-Agx phases for each composition Agx. The broken lines are drawn as guides for the eye.

Ge content can be considered and reduce the number of suitable compositions.

At first, the Ge content for the R-Agx phase (z') has to be larger than 0.25 since it was shown that the R-Agx phase is richer in Ge than the P-Agx one, and obviously to the average composition (z). A second set of restrictions can be estimated from the difference of indentation moduli $\Delta M_{R-P} = M(R-Agx) - M(P-Agx)$ between R-Agx and P-Agx phases, assuming once again that the modulus dependence on Ag and Ge can be decoupled (Equation 5).

$$\Delta M_{R-P} = -0.12\Delta_{x'x''} + 0.28\Delta_{z'z''} \quad (11)$$

where $\Delta M_{Ge} = 0.28 \Delta y = 0.28 \Delta_{z'z''}$ and $\Delta M_{Ag} = -0.12 \Delta x = -0.12 \Delta_{x'x''}$ were obtained from a linear dependences of the indentation modulus M_{Ge} on Ge content in Ge_ySe_{100-y} glasses (Figure 1) and the indentation modulus M_{Ag} on Ag content for $Ag_x(Ge_{0.25}Se_{0.75})_{100-x}$ glasses, shown in Figure 5 (blue symbols).

With the above information, it is possible to access to the possible values $\Delta_{z'z''}$ for each Agx composition. Figure 9 shows the indentation modulus differences ΔM_{R-P} predicted as a function of $\Delta_{x'x''}$ in the form of several linear curves, each one corresponding to possible values of $\Delta_{z'z''}$. Intercepts between experimental ΔM_{R-P} and the line network give access to the potential $\Delta_{z'z''}$ values for each Agx composition. One can deduce that

- For Ag 7: $\Delta_{x'x''} = [0-7]$, $\Delta M_{R-P} = 2.33$ and $\Delta_{z'z''} = [8.2-11.3]$
- For Ag 15: $\Delta_{x'x''} = [0-15]$, $\Delta M_{R-P} = 1.85$ and $\Delta_{z'z''} = [6.5-13]$
- For Ag 20: $\Delta_{x'x''} = [0-20]$, $\Delta M_{R-P} = 1.43$ and $\Delta_{z'z''} = [5-13.7]$

Taking into account all the previous constraints that include $z' > 0.25$, $x' > x$, $0 < x'' < x$, Equations (8–11) and $\Delta_{z'z''}$ estimated from Figure 9, several compositions for the R-Agx phase (x' and z') and for the P-Agx phase (x'' and z'') can be proposed. They are reported in Figure 10. Among these compositions, some might

not be realistic, in particular if they do not contain a sufficient Se ratio to compensate the positive charges brought by Ge and Ag (compositions not shown in Figure 10).

Therefore, a last restriction is proposed inspired from the limit of $Ag_x(Ge_{0.25}Se_{0.75})_{100-x}$ glass forming ability proposed by Zeidler et al. (2018). This limit corresponds to the consumption of all Se-Se homopolar bonds, can be estimated when the concentration of Ge, Se, and Ag lead to the mean coordination number associated to the number of Se-Se bonds $\bar{n}_{Se}^{Se} = 0$. Then \bar{n}_{Se}^{Se} can be calculated from this relation

$$\bar{n}_{Se}^{Se} = Z_{Se} - Z_{Ge} * \left[\frac{y}{(1-y)} \right] - 2\bar{N}_{Broken} * \left[\frac{2x}{(1-y)(100-x)} \right] \quad (12)$$

where Z_{Se} and Z_{Ge} are the overall coordination numbers of Se and Ge equal to 2 and 4 respectively, and \bar{N}_{Broken} is the mean number of broken Se-Se bonds per silver atoms estimated around 0.33 (Zeidler et al., 2018). The criteria of $\bar{n}_{Se}^{Se} = 0$ has been used to find the limit between the compositions lying within the glass forming ability and those beyond. The composition limit is shown in Figure 10. Most compositions found for R-Agx and P-Agx phases are in the glass forming compositions.

An exact value of Ag concentration in R-Agx and P-Agx phases cannot be given from the previous analyses but some interesting information can be brought: (i) the Ge concentration differs in R-Agx and P-Agx phases with Ge concentration being larger in R-Agx phase, (ii) when the average silver content increases, the Ge concentration decreases in both phases, (iii) z' and z'' lie in the range 19–31 at. % Ge.

Conductivity

The PF-QNM images (Figure 2) show a nucleation/growth type phase separation in the $Ag_x(Ge_{0.25}Se_{0.75})_{100-x}$ glasses, in

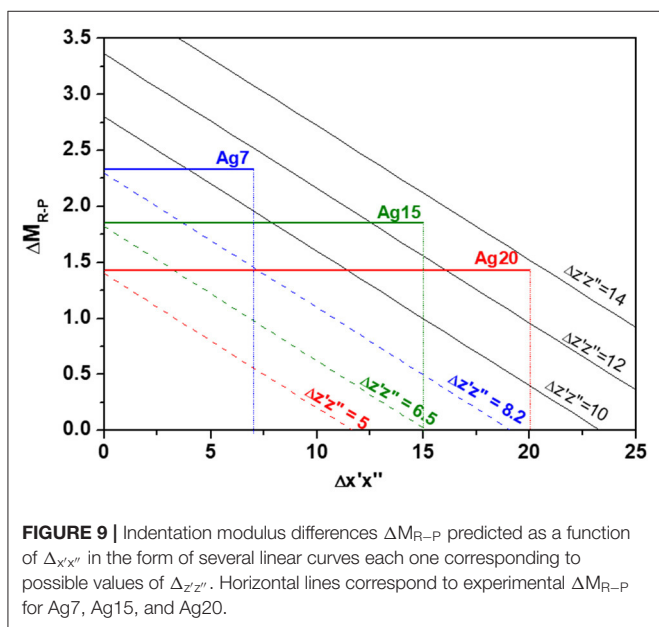


FIGURE 9 | Indentation modulus differences ΔM_{R-P} predicted as a function of $\Delta_{x'x''}$ in the form of several linear curves each one corresponding to possible values of $\Delta_{z'z''}$. Horizontal lines correspond to experimental ΔM_{R-P} for Ag7, Ag15, and Ag20.

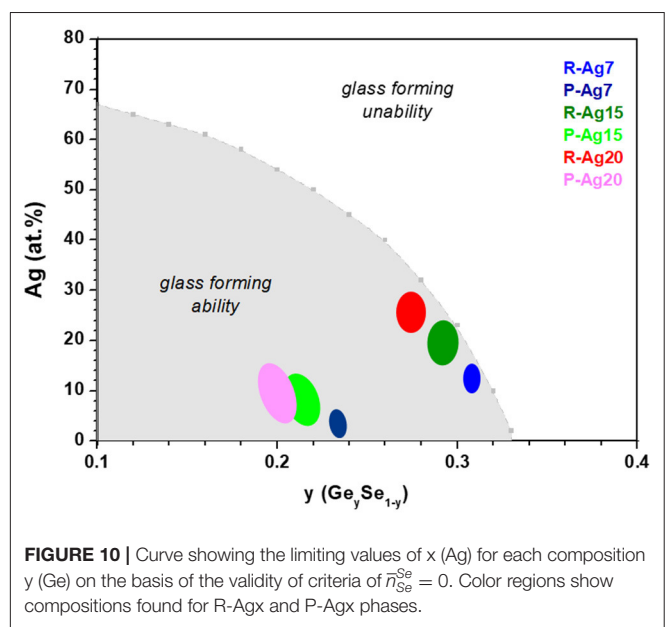


FIGURE 10 | Curve showing the limiting values of x (Ag) for each composition y (Ge) on the basis of the validity of criteria of $\bar{n}_{Se}^{Se} = 0$. Color regions show compositions found for R-Agx and P-Agx phases.

agreement with previous papers (Balan et al., 2006; Piarristeguy et al., 2007b, 2010a,b, 2018). The heterogeneous structure has been used to explain the sudden jump in conductivity by 8 orders of magnitude at a silver content lying in the range 7–10 at. % Ag. Indeed, while the Ag-rich phase is embedded in a Ag poor phase for the low conductivity region ($x < 7$ –10% at Ag), the reverse is true for the high conductivity region ($x > 7$ –10% at Ag), with the silver rich phase then controlling the silver diffusion throughout the material. Additional nanoscale electrical characterization of the R-Agx phases by conductive atomic force microscopy (C-AFM) showed that the conductivity increased in R-Agx phases with the increase in total Ag content of the glass (x) (Piarristeguy et al., 2010a). Such an increase can be explained by the increase in the Ag content of the R-Agx phase (x') (see blue, dark green, and red regions in **Figure 10**). However, the present data also show that with increasing Ag content, a softening of the network in R-Agx phases occurs due to a decrease in Ge (z'). Such a softening could also contribute to an increase in conductivity due to an increase in charge carrier mobility. This proposal agrees with Micoulaut et al. (2009) that states “The ease of conduction in an ion conducting glass could indeed be related to the mechanical nature of the host network.”

CONCLUSION

Peak Force Quantitative Nano-Mechanical (PF-QNM) mapping was used to characterize a series of binary $\text{Ge}_y\text{Se}_{100-y}$ glasses and phase-separated $\text{Ag}_x(\text{Ge}_{0.25}\text{Se}_{0.75})_{100-x}$ glasses. This technique proved to be efficient to obtain mechanical information at the nanoscale in a system where information at the nanoscale are difficult to obtain due to the fast silver photo-diffusion under energetic (light or electron beam) stimulus, leading to modification of Ag-centered local structures. As expected, a decrease in the indentation modulus when the selenium content increased in the binary glasses and when the silver content increased in phase-separated ternary glasses, as a consequence of the increasing flexibility of the network, was observed. At nanoscale level, the sensitivity of PF-QNM measurements allowed to observe a difference in the indentation moduli of the two phases, i.e., high and low for Ag-rich (R-Agx) and Ag-poor (P-Agx) phases, respectively.

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A model that contributes to a better understanding of the topology of this phase-separated system has been proposed. It is based on the assumption that the change in nanomechanical properties of the Ag-rich and Ag-poor phases depends on the change in Ag and Ge in an independent way. Within this assumption, the mean coordination numbers and compositions for Ag-rich (R-Agx) and Ag-poor (P-Agx) phases can be qualitatively estimated from PF-QNM data. When the silver content in the glass increases, the mean coordination numbers of the R-Agx phase $\langle N_R \rangle$ and P-Agx phase $\langle N_P \rangle$ decrease continuously. This is accompanied with a decrease of the Ge concentration in both phases but with the average Ge concentration remaining equal to 25.

Finally, these data helped in providing new information about the increase of conductivity in the Ag-rich (R-Agx) phase when silver content increases in the glasses, as observed by C-AFM measurements. As a matter of fact, it would not only be explained by an increase in silver content in the Ag-rich phase but also by a softening of the network and consequently by an increase of charge carrier mobility.

These promising results provide prospects likely to contribute to the understanding of the topology of phase-separated glasses as well as other bi-phase systems such as glass ceramics or composite materials.

DATA AVAILABILITY STATEMENT

All datasets generated for this study are available on request to the corresponding author.

AUTHOR CONTRIBUTIONS

All authors listed have made a substantial, direct and intellectual contribution to the work, and approved it for publication.

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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