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Influences of suction pipe structures on hydraulic performance and internal flow of electric coolant pumps

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An electric coolant pump (ECP) serves as a critical component in the thermal management of electric vehicles. To fulfill the requirements of pressurizing and circulating coolant for various components, a complex structure with multiple pipes is integrated into the pump inlet. This study focuses on the design and analysis of three suction pipe structures: a straight pipe (Case A), a bend (Case B), and a combination of a bend with manifolds (Case C). The objective of this study is to explore the impacts of suction pipe structures on the hydraulic performance, flow pattern, temperature distribution, and vorticity of ECP. Taking into account the variability of coolant physical parameters with temperature, ECP is numerically simulated using the unsteady Reynolds-averaged Navier-Stokes (RANS) equation and the shear stress transport $k-\omega$ turbulence model. The experimental and numerical results exhibit good agreement. Case A demonstrates the highest efficiency, Case B follows as the second most efficient, and Case C displays the lowest efficiency. However, the pressure rise remains essentially consistent in all cases. The average efficiencies of Cases B and C are 1.18% and 2.13% lower than that of Case A. The temperature of ECP increases with an increase in the coolant temperature. The temperature of the printed circuit board (PCB) surpasses that of the motor. Case A exhibits the most favorable flow pattern, while Case C demonstrates the least favorable. The bend introduces secondary flow, further intensified by the manifold, leading to an increase in vorticity. The high-vorticity zones expand as the flow rate increases. This study offers valuable insights into the optimization of the ECP suction pipe structure.

KEYWORDS

electric coolant pump, numerical simulation, experiment, suction pipe, hydraulic performance, vorticity

1 Introduction

During the operation of electric vehicles for various modes of transportation, modules such as batteries, motors, and brakes generate substantial heat, necessitating efficient thermal management to mitigate temperature levels and ensure the safety and optimal functioning of electric vehicles (Zhang et al., 2022). The electric coolant pump (ECP) plays a pivotal role in this thermal management process by pressurizing the coolant for circulation. To address the diverse cooling requirements of various modules, multiple pipes are integrated into the ECP inlet, which are categorized into main pipes and manifolds. The convergence of incoming flows from different pipes induces interactions, resulting in flow irregularities and heightened energy dissipation (Kumar Samal and Moharana, 2021; Minocha and

JyeshtharajJoshi, 2020; Liu et al., 2021), consequently impacting the operational efficacy of ECP. The motor and the printed circuit board (PCB) constitute the heat-generating components of ECP, and variations in incoming flow temperatures from different pipes directly influence the heat dissipation of these components. Additionally, temperature fluctuations affect the physical properties of the coolant, thereby exerting an influence on the hydraulic performance of ECP (Jiang et al., 2019; Li et al., 2021). In this research, due consideration is given to the significant variations in coolant physical property parameters at different temperatures, elucidating the relationship between coolant physical properties and temperature in the simulation model (Johann Friedrich Gülich, 2008).

ECP falls within the category of centrifugal pumps characterized by compact overall dimensions, exhibiting a complexity surpassing that of conventional centrifugal pumps. Primary hydraulic components, including the impeller and volute, stand out as pivotal factors influencing performance. Previous studies by Li et al. (2019) and Yuan et al. (2022) involved the optimization of centrifugal pump blade parameters, resulting in improvements in head, efficiency, and resistance to cavitation. Lu et al. (2021) deduced that augmenting the tongue enhances the hydraulic efficiency of centrifugal pumps, particularly under overload conditions. Lin et al. (2022) proposed an impeller design with a sinusoidal trailing edge, effectively mitigating overall energy losses in centrifugal pumps. Wang et al. (2023) observed that a judicious selection of the number of centrifugal pump blades enhances the flow pattern within the pump. Qi et al. (2023) employed a particle swarm algorithm for impeller optimization, leading to increased efficiency in centrifugal pumps. Gu et al. (2024) highlighted that trimming the rear shroud results in decreased performance of centrifugal pumps. Furthermore, cavitation exacerbates flow instability within the pump (Fecser and Lakatos, 2021) and increases energy losses (Li et al., 2023a). Cui et al. (2019) introduced an injection device to ameliorate the cavitation performance of centrifugal pumps. Jia et al. (2023) emphasized that a reduced incident angle of the sealing ring clearance diminishes flow losses in centrifugal pumps and suppresses cavitation occurrences. Ye et al. (2023) concluded that heightened friction losses on the blade surface constitute the primary factor contributing to the hump in the pump head characteristic curve. Deng et al. (2019) introduced a novel turbulence model that markedly enhances the computational precision of thrust bearings. Li et al. (2023b) proposed a transient head prediction method for mixed-flow pumps, demonstrating good agreement with experimental results. Gu et al. (2023) introduced a novel dynamic mesh simulation method to explore the impact of seal gasket fractures on the hydraulic performance of multistage centrifugal pumps. In summary, numerous factors influence the hydraulic performance of pumps, and the effects of different suction pipe structures on the hydraulic performance of ECP are inherently variable.

Unsteady flow constitutes the primary catalyst for vibration and flow-induced noise (Zeng et al., 2020; Cheng et al., 2021). Li et al. (2020) conducted an examination of the internal flow field within a mixed-flow pump, identifying a pronounced entrainment effect at the impeller inlet. Ji et al. (Leilei et al., 2021) asserted that the internal flow field, as calculated using the Wray–Agarwal turbulence model, exhibits a closer alignment with experimental results. Feng et al. (2021) observed that centrifugal pump internal flow experiences significant separation following sudden power failure, leading to increased vortices within the impeller channel. AhmedAl-Obaidi (2021) ascertained that pressure and velocity within centrifugal pumps increase proportionally with the impeller speed. Guo et al. (2022) scrutinized the flow field within an axial pump, highlighting the conspicuous influence of impeller rotation and end-wall effects. Rao et al. (Zhi et al., 2023) demonstrated that the implementation of a double-tongue volute enhances hydrostatic pressure and velocity gradients in the central region of the centrifugal pump, concurrently mitigating turbulence at the tongue under partial load conditions. In summary, the internal flow field serves as a reflection of the pump flow characteristics, and a thorough analysis of this internal flow field forms the foundational step for investigating the underlying flow mechanisms within the pump.

The generation and shedding of vortices contribute to pronounced pressure fluctuation (Li et al., 2022), exacerbating the unsteady flow within centrifugal pumps. The vortex analysis of centrifugal pumps emerges as a pivotal method for comprehensively investigating the mechanisms underlying unsteady flow phenomena (Yuan et al., 2021; Lv et al., 2022). Yasuyuki et al. (Nishi and Noji, 2020) posited a correlation between impeller-induced losses due to rotation and the distribution of vortices within the volute. Bai et al. (2022) emphasized the substantial impact of volute geometry on the evolution and formation of vortices. Vorticity, serving as a quantifier of vortex strength and fluid flow conditions, plays a crucial role in this analysis (Zhang et al., 2020). Shi et al. (2021) identified conspicuous variations in the internal flow field resulting from differing impeller tip clearances, with larger vorticity resulting from tip leakage. Introducing obstacles within the impeller channel, as demonstrated by Zhao and Guo (2021), led to a significant attenuation of vorticity. Dong et al. (2023) compared the effects of starting schemes on the internal flow field of a centrifugal pump. Ding et al. (2023) reported a substantial reduction in vorticity within the impeller after modifying the trailing edge of the blade, thereby enhancing the stability of centrifugal pumps. In the context of the suction pipe, bends and manifolds induce vortex motion, and the collision between the backflow orifice jet and inflow generates vortices, thereby intensifying flow instability. Consequently, the development of vortex analysis for ECP represents an indispensable prerequisite for the comprehensive study of flow stability.

This study conducts numerical simulations of ECP featuring diverse suction pipe structures. The accuracy of the numerical simulation results is validated through experimental verification. Subsequently, an in-depth analysis is performed to elucidate the impacts of suction pipe structure variations on the hydraulic performance and internal flow field of ECP.

2 Electric coolant pump and numerical simulation

2.1 Pump model

The design specifications for ECP include a design flow rate (Q_d) of 130 L/min, a rotational speed (n) of 5,500 rpm, a design pressure

Parameter	Symbol	Value
Impeller inlet diameter	d_1	34 mm
Impeller outlet diameter	d_2	57 mm
Impeller outlet width	b_2	7.5 mm
Number of blades	z_b	7
Volute inlet diameter	<i>d</i> ₃	58 mm
Volute inlet width	b_3	11 mm

TABLE 1 Main structural parameters of ECP.

rise (ΔP) of 110 kPa, and a specific rotational speed ($n_q = nQ^{0.5}/(\Delta P/\rho g)^{0.75}$) of 238 rpm. Here, ρ represents the density of the coolant, considered at 358 K. ECP comprises essential components, namely, a suction pipe, an integrated system consisting of an impeller and motor rotor, a volute, a motor stator, and a PCB. The key parameters of ECP are detailed in Table 1.

This study employs a numerical simulation approach to investigate the impacts of different suction pipe structures on the hydraulic performance and internal flow field of ECP. Three designated suction pipe configurations, denoted as Case A, Case B, and Case C, are devised for examination. Case A features a straight pipe positioned directly opposite the impeller inlet, providing optimal inflow conditions, as illustrated in Figure 1A. In Case B, the suction pipe is inclined at an angle of 30° to the plane of the impeller inlet and connected through a bend, as depicted in Figure 1B. Case C extends the design by incorporating two manifolds in the foundation of Case B, with three pipes intersecting at the bend. This configuration is primarily tailored to address the cooling and heat dissipation requirements of different components simultaneously, as shown in Figure 1C. In this setup, the main pipe has a diameter of 34 mm, while Manifolds A and B have diameters of 14 mm and 12 mm, respectively.

Furthermore, a backflow orifice is strategically positioned at the impeller inlet, establishing a connection between the impeller inlet and the rear chamber through the hollow shaft. This configuration facilitates the dissipation of heat originating from the motor and PCB. The precise location of the backflow orifice is illustrated in Figure 1D.

2.2 Computational mesh

Figure 2 displays the mesh configuration for ECP. Critical hydraulic components, such as the impeller and volute, are discretized using hexahedral meshes with incorporated boundary layers. (Cui et al., 2021; Deng, 2024). Simultaneously, tetrahedral





meshes are applied to the remaining components. This hybrid meshing approach strikes a balance between computational accuracy and meshing cost (Liu et al., 2015), rendering it well suited for the numerical simulation of intricate fluid machinery.

Figure 3 depicts the correlation between the cell number in Case A and the hydraulic performance of ECP. Six cell numbers, namely, 1.97 million, 3.39 million, 4.91 million, 6.64 million, 8.23 million, and 9.57 million, are employed in the assessment. It is observed that, at the cell number of 4.91 million, the efficiency and pressure rise exhibit a consistent trend. Hence, in light of considerations regarding computational accuracy and efficiency, a cell number of 6.64 million is selected for the numerical calculations.

2.3 Computational setup

The numerical simulations of ECP are conducted utilizing the compressible Reynolds-averaged Navier–Stokes (RANS) equations, coupled with the shear stress transport $k-\omega$ turbulence model, and incorporating temperature calculations. These simulations are carried out using Ansys CFX commercial software (Gu et al., 2022). The continuity equation and momentum equation are expressed in Eqs 1, 2:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i} \left(\rho u_i \right) = 0, \tag{1}$$

$$\frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_j}(\rho u_i u_j) = -\frac{\partial p}{\partial x_i} + \frac{\partial}{\partial x_j}\left(\mu_{eff}\left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i}\right)\right), \quad (2)$$

where t is the time, x stands for the coordinate, u represents the velocity, and p denotes the pressure.

In the energy conservation equation, the effects of pressure, dissipation, and heat source are considered as shown in Eq. 3:

$$\frac{\partial}{\partial t}(\rho C_p T) + \frac{\partial}{\partial x_j}(\rho C_p T u_j) = \frac{\partial}{\partial x_j}\left(\lambda_{eff}\frac{\partial T}{\partial x_j}\right) - \frac{\partial}{\partial x_j}(p u_j) + \Phi + S,$$
(3)

where *T* is the temperature, μ_{eff} denotes the effective viscosity, C_p is the constant pressure specific heat capacity, λ_{eff} represents the effective thermal conductivity, Φ stands for dissipation, and *S* is the heat source. Subscripts *i* and *j* indicate the directions.

In the setup of boundary conditions, all cases are configured with a mass flow rate inlet and a total pressure outlet set at 1.5 bar. Uniform total flow rates are maintained for all cases, with Case C distributing the flow rate based on the pipe cross-sectional area. The mass flow rate inlet coolant temperature for Cases A and B is



TABLE 2 Relationship between coolant properties and temperature.

Physical property	Fitting formula	
ρ (kg/m³)	$-3.569 \times 10^{-6}T^4 + 1.56 \times 10^{-3}T^3 - 0.377T^2 + 47.183T - 1256.518$	
μ (kg/m/s)	$1.554 \times 10^{-7} T^2 - 1.266 \times 10^{-4} T + 0.026375$	
C_p (J/kg/k)	$-7.669 \times 10^{-6} T^4 + 3.929 \times 10^{-3} T^3 - 1.14314 T^2 + 187.3563 T - 10743.206$	
λ (W/m/k)	$3.259 \times 10^{-9} T^4 - 1.569 \times 10^{-6} T^3 + 4.242 \times 10^{-4} T^2 - 0.06047T + 3.868$	

established at 358 K, while the inlet coolant temperature for the main pipe of Case C is also set to 358 K; however, for the manifolds, it is set to 398 K. The thermal flux intensities on the motor and PCB heat dissipation surfaces are specified as 16,211.5 W m² and 22,556.4 W m², respectively. The remaining walls are configured as adiabatic and non-slip. The relationship between coolant physical properties and temperature is modeled using the CFX Expression Language embedded in Ansys CFX, with the fitting equations detailed in Table 2.

The interface between the rotating and stationary components is configured as a frozen rotor for steady simulations and a transient frozen rotor for unsteady simulations. The accuracy and stability of calculations are the numerical contingent upon the Courant-Friedrichs-Lewy number, which is meticulously maintained at approximately 5 throughout the calculations (Lu et al., 2023). For unsteady calculations, the convergence residual is set at 10⁻⁴, and the time step corresponds to the duration required for a 3° rotation of the impeller, totaling 6 cycles of impeller rotation. The final result used for analysis is the average of the outcomes from the last revolution of the calculation.

3 Experimental setup and validation

3.1 Experimental setup

The experimental setup for assessing the efficiency and pressure rise of ECP is illustrated in Figure 4. The experimental bench primarily comprises a test pump, an auxiliary pump, a pressure sensor, a flowmeter, a tank, a tachometer, and a wattmeter. Notably, the tank serves the purpose of heating the coolant and maintaining a constant temperature. The measurement errors for the pressure sensor, flowmeter, tachometer, and wattmeter are specified as 0.1%, 0.3%, 0.1%, and 0.1%, respectively. The pump performance characteristics are computed following the procedures outlined by Johann Friedrich Gülich (2008), and the experiment is conducted in triplicate to mitigate experimental errors and ensure result accuracy. The uncertainties associated with pressure rise and efficiency do not exceed 0.17% and 0.35%, respectively.

In the experiment, the pressure rise is computed based on the differential between inlet and outlet pressures, while the efficiency is determined using Eq. 4:



$$\eta = \frac{\Delta PQ}{P_i},\tag{4}$$

where P_i represents the input power of the ECP rotor.

3.2 Comparisons of hydraulic performance

The pressure rise and efficiency of Case A are tested in the experiments, and the comparison of the experimental and simulation results is shown in Figure 5. Experimentally, the optimum efficiency reached 64.95%, corresponding to a pressure rise of 87.34 kPa. In numerical simulations, the highest efficiency occurs at Q_d , registering at 70.3%, with a pressure rise of 93.24 kPa.

The experimentally obtained optimal efficiency is 5.35% lower than that of the numerical simulation, with the pressure rise being 0.063% lower. These disparities are attributed to external factors such as tube resistance and measurement errors influencing the experiments. However, despite the differences, the performance curve derived from the numerical simulation exhibits a congruent trend with the experimental data, affirming the accuracy of the numerical simulation method.

4 Results and discussion

4.1 Analysis of performance

The pump performance and internal flow field characteristics derived from unsteady simulation exhibit closer alignment with experimental results than those obtained through steady simulation (Kim et al., 2019). The hydraulic performance of ECP under different conditions is determined through unsteady simulation, as illustrated in Figure 6.

For equivalent flow rates, transitioning from Case A to Case C led to a significant decrease in efficiency, while the pressure rise exhibits only a marginal reduction. At $0.769Q_d$, Case A demonstrates a pressure rise of 108.1 kPa with an efficiency of 67.37%, Case B records a pressure rise of 107.94 kPa, achieving 67.1% efficiency, and Case C yields a pressure rise of 107.85 kPa with 66.53% efficiency. At Q_d , Case A exhibits a pressure rise of 93.24 kPa with an efficiency of 70.3%, Case B shows a pressure rise of 93.16 kPa with 69.42% efficiency, and Case C results in a pressure rise of 93.04 kPa with 68.68% efficiency. At 1.231 Q_d , Case A displays a pressure rise of 73.60 kPa with 67.16% efficiency, Case B presents a pressure rise of 73.45 kPa with 65.68% efficiency, and Case C showcases a pressure rise of 73.30 kPa with 63.56% efficiency.







On the other hand, when transitioning from Case A to Case C, efficiency experiences a more pronounced decline at elevated flow rates. Specifically, at $0.769Q_d$, Case A surpasses Cases B and C in efficiency by 0.27% and 0.84%, respectively. At Q_d , Case A exhibits higher efficiency than Cases B and C by 0.88% and 1.62%, respectively. Moreover, at $1.231Q_d$, Case A outperforms Cases B and C with superiority in efficiency of 1.48% and 3.6%, respectively.

In general, the efficiency of Cases B and C exhibits an average reduction of 1.18% and 2.13%, respectively, compared to Case A, while the pressure rise demonstrates marginal decreases of 0.0016% and 0.0032% within the studied flow rate ranges. The suction pipe structure notably impacts ECP efficiency but has almost negligible effects on pressure rise. The underlying reasons for this phenomenon will be analyzed subsequently.

4.2 Analysis of streamline

The ECP suction pipe structure plays a crucial role in influencing the flow pattern at the impeller inlet, with bends and manifolds impacting the inflow pattern and contributing to increased flow losses (Zhang and Li, 2018). The streamlines for all cases at the design flow rate are depicted in Figure 7.

In general, Case A exhibits the most favorable flow pattern, whereas Case C displays the least favorable. The cooling backflow in ECP is characterized by a pressure-driven flow, generating a jet as it exits the backflow orifice and colliding with the inflow from the suction pipe, thereby disrupting the streamlines. In Case A, the influence of the jet from the backflow orifice intertwines the streamlines in the middle of the pipe, causing a disturbance in



the flow pattern. However, other zones maintain smooth and straight streamlines with a favorable flow pattern. The jet has no impact on the impeller inlet flow pattern, resulting in the highest efficiency. When the suction pipe is configured as a bend, the velocity on the outside of the bend surpasses that on the inside, creating a pressure gradient that induces secondary flow generation (Zhang et al., 2022). Simultaneously, the jet collides with the main flow, disrupting the flow pattern and increasing flow losses, leading to a reduction in the efficiency of Case B. In Case C, a circular cylinder flow manifests at the outlet of Manifold A, and the evident occurrence of flow separation is observed. The outflow from Manifold B is influenced by the inflow from the main pipe, inducing heightened disorder in the secondary flow streamlines and markedly impacting the impeller inlet flow pattern. This exerts the most substantial influence on efficiency, culminating in the lowest attainable efficiency.

The corresponding streamlines for Case C at different flow rates are shown in Figure 7C and Figure 8. The pipe streamlines are basically the same for each flow rate, but the disordered flow caused by the jet impact is not the same. The driving pressure of backflow is the pressure difference between the impeller inlet and outlet, and it is largest when the flow rate is $0.769Q_d$, corresponding to the highest jet intensity, resulting in the most disordered flow pattern. In addition, as the flow rate increases, the secondary flow on the inside of the bend becomes much more serious. This is because the velocity difference between the inside and outside of the bend increases with the flow rate, and the pressure gradient becomes larger, enhancing the secondary flow.

4.3 Analysis of temperature

In this section, the influences of suction pipe structures on the temperature distribution of ECP is investigated. The temperature distribution corresponding to different suction pipe structures at Q_d is depicted in Figure 9. Due to the integration of electronic components, the PCB surface exhibits unevenness, leading to non-uniform heat dissipation. Consequently, noticeable local high temperatures are observed on PCB in each case, and the overall temperature surpasses that of the motor.

The temperature distribution on the motor heat dissipation surface is generally consistent between Cases A and B. However, in Case B, the local temperature of PCB is marginally higher than that of Case A. For Case C, the temperatures of the manifolds and main pipe differ, and they become more uniform after mixing by the impeller. Nevertheless, it still remains approximately 2.5 K higher than in the other cases, resulting in a weakening of the heat transfer effect. Consequently, the temperatures of the motor and PCB are higher.

Figure 9C and Figure 10 illustrate the temperature distribution of Case C at different flow rates. With an increase in the flow rate, the motor temperature experiences a slight rise, while the local temperature of PCB increases significantly. As the flow rate increases, the impeller inlet temperature mixing diminishes, causing the expansion of the high-temperature zones. Simultaneously, the coolant temperature in the impeller exhibits reduced mixing, leading to noticeable stratification and an overall temperature increase. Consequently, the motor temperature increases as the coolant enters the rear chamber. The reduction in heat transfer efficiency is attributed to the higher thermal flux intensity on the heat dissipating surface of PCB and the prior absorption of heat from the motor by coolant-flowing PCB, yielding a heightened temperature. Subsequently, the temperature rise is notably more pronounced in PCB in contrast to the motor. Hence, PCB is relatively more vulnerable to thermal damage. In the process of optimizing the heat dissipation of ECP, particular attention is given to the thermal management of PCB.

4.4 Analysis of vorticity

Upon analysis of Section 4.2, it is observed that a pressure gradient manifests as the fluid traverses the bend. Notably, two bundles of secondary flow exhibiting relative motion occur on the inner side of the bend, converging in the midpoint of the inner side. This convergence leads to a significant alteration in flow direction, causing a pronounced curling of streamlines and the generation of a vortex (Zuojun et al., 2016). Vorticity is employed to characterize vortex strength and can be visualized in CFD-Post (Zhang et al., 2019). In this section, the vorticity distribution is analyzed for





different suction pipe structures. The vorticity distribution of the cross section of ECP at Q_d is illustrated in Figure 11.

The suction pipe structure significantly influences the vorticity distribution in the pump. At S_1 , the backflow orifice jet collides with the main flow, creating a high-intensity vortex (Zong and Kotsonis, 2020). The high-vorticity zone is the smallest in Case A and the largest in Case C. As depicted in Figure 7, the high-vorticity zones emerge at the bend of Case B due to the secondary flow generation. In Case C, the outflow from Manifold B interacts with the inflow from the main pipe, intensifying the secondary flow and increasing the vorticity at the Manifold B outlet.

At S_2 , Case A continues to exhibit the lowest vorticity. In Cases B and C, the vorticity distribution caused by the bend and jet is essentially identical. However, the outflow from Manifold B in Case C generates high-vorticity zones. Consequently, when compared, the vorticity in Case C is highest at S_2 .

 S_3 is situated at the impeller inlet, where the rotation of the impeller induces significant shear stresses and high-vorticity zones close to the wall. In Case A, the high-vorticity zones are uniformly distributed, slightly reduced in Case B, and minimized in Case C. In Case B, the disordered streamlines influence the impeller inlet flow pattern, resulting in an uneven vorticity



distribution in S_3 . When compared to Case A, the coolant temperature is higher in Case C, and the decreased viscosity leads to a weakened wall shear effect, resulting in the smallest high-vorticity zones.

In summary, Case C exhibits the highest vorticity, while Case A has the lowest. The generation of vortices is frequently coupled with turbulence, leading to turbulence losses (Deng et al., 2023; Kan et al., 2023). Furthermore, the vortex introduces additional resistance and

increases flow friction, contributing to reduced efficiency. Conversely, pressure rise is primarily associated with the pump geometrical parameters, explaining the significant variations in efficiencies depicted in Figure 6, while the pressure rise remains essentially constant.

Figure 11C and Figure 12 illustrate the vorticity distribution for Case C at different flow rates. With the increase in flow rates, the vorticity intensifies. Despite the diminishing intensity of the backflow



jet with increased flow rates, the augmented inflow from the pipe extends the reach of the jet influence. Concurrently, the augmentation in flow intensifies the secondary flow, resulting in an elevation of S_1 and S_2 vorticities. Analogously, the domain influenced by the shear stress induced by the rotation of the impeller inlet wall expands following the increase in flow rate, thereby causing an escalation in S_3 wall vorticity.

5 Conclusion

In this study, the effects of three suction pipe structures on the hydraulic performance, temperature distribution, flow pattern, and vorticity of ECP are examined through numerical simulations. The experimental and numerical simulation results are in good agreement. The conclusions are as follows:

- (1) The suction pipe structure exerts a notable impact on the efficiency of ECP, with minimal influence on the pressure rise. Compared to Case A, the average efficiency declined by 1.18% for Case B and 2.13% for Case C. Similarly, the average pressure rise experiences a reduction of 0.002% for Case B and 0.003% for Case C.
- (2) The suction pipe structure has a small effect on the heat dissipation of ECP. The uneven surface of PCB produces high localized temperatures. Coolant temperature has a significant effect on ECP heat dissipation. The motor and PCB temperatures in Case C are significantly higher than those in other cases. The motor and PCB temperatures increase slightly when the flow rate is increased.

- (3) Streamlines illustrate that Case A, with a straight suction pipe, exhibits the optimal flow pattern, while Case C, incorporating manifolds, displays the least favorable flow pattern. Secondary flow is induced on the interior of the bend in Case B and is further intensified by Manifold B in Case C. The flow pattern experiences increased disorder with increasing flow rates. As the flow rate increases, the secondary flow within the bend intensifies, resulting in a more disordered flow pattern.
- (4) The vorticity analysis reveals that Case A exhibits the lowest vorticity, while Case C demonstrates the highest vorticity. Higher vorticity is observed on the inside of the bend, and Manifold B further enhances the vorticity in this zone. Vorticity intensifies with an increase in the flow rate.

Data availability statement

The original contributions presented in the study are included in the article/supplementary material; further inquiries can be directed to the corresponding author.

Author contributions

AY: writing-original draft, writing-review and editing, formal analysis, investigation, and software. YG: writing-original draft, writing-review and editing, conceptualization, funding acquisition, methodology, project administration, and supervision. LC: writing-review and editing. WZ: writing-review and editing.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Nomenclature

Symbols

b_2	Impeller outlet width, mm
<i>b</i> ₃	Volute inlet width, mm
C_p	Pressure specific heat capacity, J/kg/K
d_1	Impeller inlet diameter, mm
<i>d</i> ₂	Impeller outlet diameter, mm
<i>d</i> ₃	Volute inlet diameter, mm
n	Rotating speed of the impeller, r/min
n_q	Specific speed for pumps (p, pressure, Pa)
P _i	Input power for the rotor, W
Q	Flow rate, L/min
\$	Heat source, W
Т	Temperature, K
t	Time, s
и	Velocity, m/s
x	Coordinate, m
$z_{\rm b}$	Number of blades
Greek letters	
ΔP	Design pressure rise, Pa
η	Efficiency, -
λ	Thermal conductivity, W/m/K
λ_{eff}	Effective thermal conductivity, W/m/K
μ	Viscosity, Pa-s
μ_{eff}	Effective viscosity, Pa·s
ρ	Density, kg/m ³
Φ	Dissipation, W/m ³
Superscripts	
-	Time-averaged component
3	Fluctuating component
Subscripts	
d	At design condition
i	Free index
j, k	Dummy index
Abbreviations	
ECP	Electric coolant pump
Exp	Experimental
РСВ	Printed circuit board
CFD	Computational fluid dynamics