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RECEIVED 02 May 2023 ACCEPTED 06 July 2023 PUBLISHED 28 July 2023

CITATION

Ağca S, Çankaya G and Sonmezoglu S (2023), Impact of tellurium as an anion dopant on the photovoltaic performance of wide-bandgap Cu(In,Ga)Se₂ thin-film solar cells with rubidium fluoride post-deposition treatment. *Front. Energy Res.* 11:1215712. doi: 10.3389/fenrg.2023.1215712

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Impact of tellurium as an anion dopant on the photovoltaic performance of wide-bandgap Cu(In,Ga)Se₂ thin-film solar cells with rubidium fluoride post-deposition treatment

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The development of wide-bandgap Cu(In,Ga)Se₂ thin films is crucial in order to reach the theoretical Shockley-Queisser limit values in single-crystal solar cells. However, the performance of solar cells based on wide-bandgap thin film absorbers has lagged significantly compared to that of their narrow-bandgap counterparts. Herein, we develop a feasible strategy to improve the photovoltaic performance of wide-bandgap Cu(In,Ga)Se₂ chalcopyrite thin-film solar cells by simultaneously doping with both RbF PDT and Te²⁻ anions as dopants in the absorber layer during the three-stage co-evaporation process. Besides inducing significant change in the GGI gradient, the synergistic effect of the Te²⁻ anion dopant is rather beneficial in terms of controlling grain size, defects in grain boundaries, and charge carrier lifetime for encouraging charge separation and extraction, which contributes to simultaneously boosting short-circuit current density and fill factor. Te-poor devices afford an impressive efficiency of 9.58%, compared to 6.43% for control devices. More importantly, the efficiency and Voc values obtained for wide-bandgap-based thin-film solar cells containing Te anions were the highest compared to their counterparts as reported in the literature. These results demonstrate the role of Te²⁻ anions in wide-bandgap absorber thin films on the photovoltaic performance of thin-film solar cells and the potential of this approach for use in reasonable and effective design of highly efficient wide-bandgap thin-film solar cells.

KEYWORDS

 $\label{eq:culling} Cu(In,Ga)(Se,Te)_2 \ solar \ cells, \ Te \ anions, \ RbF \ post-deposition \ treatment, \ three-stage \ co-evaporation, \ wide-bandgap \ chalcopyrite \ absorber \ thin-films, \ GGI \ gradient$

1 Introduction

Given the emergence of climate change as one of the most crucial environmental challenges in today's world, solar energy has become one of the most important candidate alternatives to fossil fuels in combating climate change (Scheer et al., 2004;

Sönmezoğlu Ates et al., 2016; Kamikawa et al., 2022; Kaya et al., 2022; Liu et al., 2022; Mabvuer et al., 2022; Oliveira et al., 2022). Among the various types of solar cells, one of the most important technologies in the field of solar energy, Cu(In,Ga)Se₂ thin-film solar cells (TFSCs) have emerged as currently one of the most promising photovoltaic technologies, owing to their very high power conversion efficiency (PCE) of 23.4%, high absorption coefficient, effective usage of raw materials, and tunable direct bandgap (Wei et al., 1998; Guillemoles, 2002; Lundberg et al., 2005; Theelen and Daume, 2016; Naghavi et al., 2017). In spite of their high efficiency, Cu(In,Ga)Se2 TFSCs still remain less efficient compared to both their theoretically achievable efficiency (33.7%) based on the Shockley-Queisser (S-Q) limit for single-junction solar cells (Shockley and Queisser, 1961) and other types of solar cells, such as single-crystal silicon (26.80%) and perovskite solar cells (26.08%). This is attributable to the preference for narrow-bandgap materials (1.0-1.1 eV) instead of materials with the ideal bandgap value (1.4 eV) as defined by the S-Q limit, causing major loss in open-circuit voltage (Voc) and reducing short-circuit current (Jsc) and fill factor (FF) (Polman et al., 2016; NREL Best Research-Cell Efficiency Chart, 2022). In order to reach theoretical efficiency, it is of major importance to choose suitable wide-bandgap absorber materials, with a bandgap between 1.3 and 1.4 eV.

To adjust the bandgap of Cu(In,Ga)Se₂ TFSCs, the most effective approaches are as follows. First, by changing the (Ga)/(In+Ga) (GGI) ratio, the bandgap can be varied from 1.02 eV (for CuInSe₂) to 1.68 eV (for CuGaSe₂) (Gloeckler and Sites, 2005; Witte et al., 2015; Zahedi-Azad et al., 2020). However, research to date has shown that the TFSCs with the highest efficiency are based on bandgaps in the 1.2-1.3 eV range achieved in films where the GGI ratio is between 0.25 and 0.30. As the bandgap is increased further, Jsc falls off precipitously while Voc remains nearly constant. This limits the use of Cu(In,Ga)Se₂ films with high Ga concentration as a wide-bandgap absorber layer in TFSCs. Second, by altering the anion elements (the Se/S ratio), the bandgap of these materials can be controlled in the range of 1.04 eV (CuInSe₂) to 1.53 eV (CuInS₂). Nevertheless, when Ga is incorporated into these matrices, the bandgap increases considerably to values between 1.68 eV (CuGaSe₂) and 2.43 eV (CuGaS₂) (Ishizuka et al., 2018). At the same time, precise control of the S/Se ratio is very difficult because of the thermal annealing process during fabrication, which poses major challenges at this stage, such as the formation of nonstoichiometric compositions (Kim et al., 2012; Kobayashi et al., 2015; Ju et al., 2017). Based on the aforementioned results, we see that gradients of the Ga/In cation and Se/S anion are a major limitation on improvements to the performance of widebandgap absorber-based cells; research into alternative approaches to effective control of bandgaps therefore needs immediate attention.

Another promising approach for control of the bandgap is to use tellurium (Te) as an anion dopant, leading to a wider distribution of anisotropic electrons and thus an enlarged electropositive region via increase in the strength of the chalcogen bond, as a result of the lower electronegativity and higher radius compared to those of Se and S (Rosenfield et al., 1977; Sönmezoğlu, 2014; Akın et al., 2017; Steinke et al., 2023). Another outstanding property is that Te has a lower effective mass (0.45 m_o) than selenium (1.40 m_o),

which implies an increase in mobility of the hole carriers of absorber materials (Caldwell and Fan, 1959; Beyer et al., 1971; Madelung, 2012; Sönmezoğlu et al., 2013; Sönmezoğlu and Akman, 2014). A detailed literature search revealed that there have been many attempts to examine the effects of a Te2- anion on the microstructural, morphological, optical, and electrical properties of chalcopyrite thin films (Mise and Nakada, 2010a; Fiat et al., 2013a; Fiat et al., 2013b; Fiat et al., 2014a; Fiat et al., 2014b; Karatay et al., 2017; Atasoy et al., 2018); surprisingly, however, few studies have focused on the impact of Te²⁻ anions on the photovoltaic parameters of wide-bandgap chalcopyrite absorber-based TFSCs. The first results on the utilization of a Te²⁻ anion in wide-bandgap absorber materials for solar cells were reported by Kim et al. (2013). They reported on the synthesis of colloidal CuInTe2-xSex gradient-alloyed quantum dots (QDs) via a simple hot injection method and their use as a p-type wide-bandgap absorber layer (1.40 eV) in heterojunction solar cells. However, their device architecture was designed as a heterojunction structure with a CuInTe_{2-x}Se_x p-type absorber layer and n-type TiO2 materials instead of TFSCs, and therefore, the solution-processed heterojunction solar cell based on the gradient-alloyed Cu_{0.23}In_{0.36}Te_{0.19}Se_{0.22} QDs remained at a low efficiency of 3.8%. In a second study, Hamid et al. (2019) investigated the use of single-step evaporation-processed CuInTe₂ doped with manganese element as a wide-bandgap absorber layer (1.45 eV) in TFSCs, achieving a PCE of 1.8%. In addition to the poor performance exhibited by their TFSCs, the existence of surface disorder and amorphous structure originating from the fabrication process created negative effects. Thus, in order to design wide-bandgap absorber materials that are suitable for highly efficient TFSC applications, it is necessary to understand the role of Te²⁻ anions in chalcopyrite thin film with respect to device performance.

In this study, a series of RbF post-deposition-treated Cu(In,Ga)Se₂ wide-bandgap absorbers with Te²⁻ anions incorporated as a dopant were grown via three-stage co-evaporation, and the effects of the Te2- anion on the photovoltaic performance of these wide-bandgap absorber layer-based TFSCs were explored for the first time. The incorporation of the Te2- anion into the Cu(In,Ga)Se₂ thin films encouraged more uniform morphology and well-controlled growth of large grains with high crystallinity. In addition, it was observed that the use of Te²⁻ anions as an additive reduced hole-trap density in the bulk materials and led to enhanced carrier lifetime, thus giving rise to pronounced enhancement of Jsc and FF. As a result, Te-poor Cu(In,Ga)(Se,Te)₂-based TFSCs exhibited a high PCE of 9.58% compared to the 6.43% of controlbased TFSCs; to the best of our knowledge, this is the highest PCE compared to Te-incorporated wide-bandgap TFSCs reported in the literature.

2 Experimental methods

Our group purchased all materials commercially (Cu and Te from Alfa Aesar, In and Se from Evochem, and Ga from Haines & Maassen) and used them without further purification, unless stated otherwise. We used 400-nm Mo-coated 3-mm-thick soda-lime glass (SLG; obtained from NICE Solar Energy GmbH) without a diffusion barrier as a substrate. Since there was no diffusion

TABLE 1 Se and Te evaporation temperatures and Te concentrations of samples.

Sample	T _{Subs.} (°C)	Se T _{Evap} (°C)	Te T _{Evap} (°C)	Te (at%)
CuInGaSe ₂	620	260		0.0
CuInGa(Se _{0.995} Te _{0.005}) ₂	620	240	385	0.3
$CuInGa(Se_{0.99}Te_{0.01})_2$	620	235	395	0.6

barrier layer on the SLG, Na could diffuse from the glass to the absorber during the deposition process. The absorber layer was applied using a Balzers BAK600 multi-source evaporation chamber under vacuum with $2 \cdot 10^{-5}$ Pa pressure. Cu(In,Ga)(Se,Te)₂ absorber layers were approximately 2 µm thick and were deposited by the three-stage co-evaporation method at a substrate temperature of 620°C. Cu, In, Ga, Se, and Te sources were filled equally at each deposition stage to eliminate variation in the evaporation rates. The composition and thickness of deposited films were controlled by an LLS system in situ. Cu was evaporated only in the second stage, In was evaporated in both the first and the third stages, and Ga was evaporated in all stages. RbF PDT was applied without breaking the vacuum after the third stage. Se and Te were evaporated during the three-stage process and the PDT process. Different Te concentrations were obtained by varying Se and Te evaporation temperatures. The (Se+Te)/(Cu+In+Ga) ratio was kept constant at approximately 10. Cu(In_{0.3}Ga_{0.7})(Se_(1-x),Te_x)₂ absorbers having three different x values (x = 0, x = 0.005, and x = 0.01) were produced with a 0.9 Cu/(Ga+In) ratio (CGI) and 0.7 GGI, which corresponds to a bandgap of approximately 1.4 eV. Table 1 shows the Se and Te evaporation temperatures and Te concentrations of samples obtained from energy-dispersive spectroscopy (EDS) measurements.

The CBD method was used for deposition of the CdS buffer layer. Cadmium acetate dihydrate (obtained from Alfa Aesar) with a purity of 99.999% was used as a cadmium source. Analytical grade (>99%) thiourea was obtained from Merck and was used as a sulfur source. In addition, 25% ammonia solution was obtained from Merck and was added as a complexing agent. The absorbers were dipped into the CBD solution and stirred at 15 rpm for 8 min at 60°C. The resulting CdS layer was approximately 50 nm thick. ZnO, ITO, and front contact depositions were performed using the Alliance Concept Eva 450 system with combined RF sputtering and e-beam evaporation. The purity of the ZnO and ITO targets was 99.99%. Non-doped ZnO and ITO layers were deposited by sputtering with thicknesses of approximately 100 nm and 200 nm, respectively. The Ni/Al/Ni front contact was added by e-beam evaporation. The thicknesses of the layers were 30 nm, 2 µm, and 30 nm for Ni, Al, and Ni, respectively. Al was obtained from Umicore with a purity of 99.999%; the purity of the Ni source was 99.95%.

A ZEISS Supra 40 VP scanning electron microscope (SEM) was used to evaluate the microstructure, with acceleration voltage set to 5 kV. Microstructure photographs were taken after absorber deposition. The chemical compositions of samples were determined using a Bruker EDS with acceleration voltage set to 15 kV; this was attached to the SEM. X-ray diffraction (XRD) analysis was carried out using a Rigaku MiniFlex 600 XRD with Cu–Kα radiation at 20 values between 20° and 90°. Each absorber layer was also

produced on SLG without a Mo back contact to eliminate the Mo peaks in XRD measurements and to enable investigation of the optical properties of the films. Absorbance measurements were carried out at wavelengths between 600 and 1,400 nm using a Shimadzu UV-2600 UV-Vis spectrophotometer. Time-resolved photoluminescence (TRPL) characterization was carried out after absorber layer deposition in order to understand the back contact recombination rate and to estimate the lifetime of the minority carriers. Glow discharge optical emission spectroscopy (GDOES) measurements were obtained using a Spectruma GDA750 to investigate the compositional depth profile. J-V measurements were performed under AM1.5G standard test conditions using a four-point measurement setup at room temperature. The lock-in method was employed to measure external quantum efficiency (EQE) using a xenon lamp equipped with a monochromator. Voc and FF values were obtained from the J-V results. Jsc was calculated via integration of the EQE result with the AM1.5G spectrum.

3 Results and discussion

Figure 1 shows top-view and cross-sectional SEM images of the absorber thin films fabricated via a three-stage co-evaporation method. It can be seen in Figures 1A-C that all absorber layers were uniform, dense, and compact. However, the surface was rougher with incorporation of Te, which may have been caused by large grain formation. All absorber layers exhibited large columnar grains, and the thicknesses were approximately 2 µm in each case, as shown in Figures 1D-F. It appears that larger grains (up to 2 µm) were formed with incorporation of Te into the control film, implying that Te assisted in grain growth. The underlying reason is that Te crystallizes slowly due to its high boiling point, and therefore slower nucleation occurs during the three-stage process as a result of the higher diffusion rate of the Te compared to Se (Kwon and Hyeon, 2008; Gulen et al., 2018; Shukla et al., 2021). The larger grain size growth with incorporation of Te prevents the formation of disorder in grain boundaries (GBs) at the absorber/CdS interface and within the absorber layer, leading to higher Voc and FF; in addition, it would be more favorable for charge extraction and transfer to balance the length of charge carrier diffusion, resulting in higher Jsc (Xing et al., 2013; Yin et al., 2015).

Figure 2A shows the XRD patterns of the absorber thin films with varying Te dopants grown on SLG via the three-stage process. The three dominant peaks related to the Cu(In,Ga)Se₂ chalcopyrite phase, during which 0.3 GGI (JCPDS 00-035-1102) appeared in the XRD patterns with peaks of (112), (220/204), and (312/116). It can clearly be seen that the intensity of the (112) peak became weaker compared to that of the control absorber thin films; conversely, the intensities of the (220/204) and (312/116) peaks became stronger, which means that crystallization quality and orientation can be tuned through the incorporation of Te⁺² anions during the three-stage co-evaporation process. In particular, the increase along the (220/204) and (312/116) orientations is regarded as the favorable crystallographic direction, because this is beneficial for carrier transportation and achieving a low density of non-radiative recombination centers in the





Cu(In,Ga)Se₂ absorber layer, leading to improvement in the charge transport mechanism in TFSCs (Contreras et al., 2006; Siebentritt et al., 2006). As expected, a reasonably strong shift in the strong (112) peak toward the lower angle was observed with an increase in Te dopants in the absorber thin film, as shown in Figure 2B. This is because of the greater ionic radii of Te^{2-} (221 pm) anions in comparison to Se^{2-} (198 pm) ions, which can lead to induction of stress and increase in the lattice parameters, suggesting that Te²⁻ anions are replaced by Se²⁻ and act as an interstitial impurity inside the thin film (Atasoy et al., 2018). The optical bandgap (Eg) of the absorber thin films with and without Te dopants was calculated by extrapolating the linear range to the energy axis in plots of $(\alpha h \nu)^2$ against energy, shown in Figure 2C. The bandgaps of the Cu(In,Ga)Se₂, Cu(In,Ga)(Se_{0.995}Te_{0.005})₂, and Cu(In,Ga)(Se_{0.99}Te_{0.01})₂ thin films were determined to be 1.40, 1.39, and 1.38 eV, respectively, which would enable more efficient harvesting of solar energy according to the S-Q limit. No significant change was observed in the bandgap value with the incorporation of Te dopant at a low level. A negligible reduction in bandgap can be attributed to the difference between Se and Te chalcogenides in terms of Pauling electronegativity, which leads to an upshift in the band valence (Pauling, 1960). To gain insight into the charge transfer mechanism, we evaluated the hole-extraction ability by examining time-resolved photoluminescence (TRPL) spectra, shown in Figure 2D. As evident from Figure 2D, Cu(In,Ga)(Se_{0.99}Te_{0.01})₂ exhibited the lowest charge lifetime (1.4 ns), as compared to the Cu(In,Ga)(Se_{0.995}Te_{0.005})₂ (2 ns) and control (16 ns) thin films, considering that more PL quenching shows effective hole-extraction across the interface between the absorber and CdS. This can be explained by the reduction of defects in the GBs resulting from the larger grain sizes, consistent with SEM images (Jehl-Li-Kao et al., 2013). Considering these insights, we can surmise that Cu(In,Ga)(Se_{0.99}Te_{0.01})₂ exhibits greater hole-extraction efficiency than that of other options.

Figure 3A shows the effects of the Te dopant on the GGI gradient of absorber thin films. The left side of the graph represents the absorber surface, and the right side represents the Mo back contact. The GGI gradient of control thin films resembled a typical three-stage co-evaporation gradient, with higher Ga content



near the surface and the back contact (Jarzembowski et al., 2015; Schneider et al., 2021). Despite the evaporation of Ga at all stages, the presence of less Ga in the middle region can be explained by insufficient diffusion. When the notch in the middle region is too deep, it affects electron transition and increases the recombination rate (Zahedi-Azad et al., 2020). The addition of Te eliminated the notch in the middle region, as his improved the diffusion of Ga by increasing crystal quality. Figure 3A demonstrates that the Tepoor dopant caused a significant change in the GGI gradients. Although the GGI gradients of Te-doped and non-doped thin films differed markedly, the GGI gradients of both Te-doped thin films were similar. To gain insight into the impact of Te²⁻ anions on the distribution of alkali elements (both sodium (Na) diffused from the substrate and rubidium (Rb) exposed by using postdeposition treatment (PDT)) we examined Na and Rb distributions in absorber thin films with varying Te dopant contents. The Na distributions of control and Cu(In,Ga)(Se_{0.99}Te_{0.01})₂ thin films were similar, and there was a slight increase in the region close to the back contact, as shown in Figure 3B. Since Na was not added from the surface, we speculate that the peak near the surface region in the case of the control thin films was caused by a measurement error. The amount of Na was higher in the region close to the back contact for all absorber thin films because Na diffuses from the substrate. The Na distribution of Cu(In,Ga)(Se_{0.995}Te_{0.005})₂ differed from that of the other thin films in the region close to the back contact. Na was able to diffuse at a high rate up to the middle

region of the Cu(In,Ga)(Se_{0.995}Te_{0.005})₂ thin films. It is known that the incorporation of a small amount of Na improves the efficiency of Cu(In,Ga)Se2-based solar cells by increasing the hole concentration and conductivity (Urbaniak et al., 2014). However, the addition of excess Na may cause the solar cell parameters to deteriorate. Jackson et al. (2016) showed that PDT of heavier alkalis such as Cs and Rb increases the efficiency of low-GGI Cu(In,Ga)Se₂ solar cells by improving diode quality. In addition, other groups have revealed that RbF PDT increases Voc by reducing recombination (Avancini et al., 2017; Karki et al., 2018; Feurer et al., 2019). In another study, Zahedi-Azad et al. (2019) investigated the effect of these heavier alkalis on the performance of high-GGI wide-bandgap Cu(In,Ga)Se2 solar cells. They found that both alkali groups were beneficial for achieving higher solar cell efficiency. However, the performances of the solar cell with Rb were better than those of a solar cell with Cs. Therefore, we selected RbF for PDT in this work. Figure 3C illustrates the effect of Te dopants on Rb distribution. As can be seen from Figure 3C, all absorber layers exhibited a peak in the surface region (lefthand side) due to non-diffused residual Rb particles. When the Rb distributions of the thin films with and without Te dopants were compared, it was observed that Rb diffused much more successfully in the thin films with Te dopants. As we know that the effects of Na can be seen within the bulk and the surface of the absorber layer, Rb mainly affects the absorber/CdS interface (Nwakanma et al., 2021). The addition of alkalis improves solar



cell efficiency by reducing interface recombination and increasing the incorporation of Cd into the absorber (Chirila et al., 2013; Pianezzi et al., 2014).

Figure 4A shows the photovoltaic performance of our bestperforming cells with Cu(In,Ga)Se2, Cu(In,Ga)(Se0.995 Te0.005)2, and $Cu(In,Ga)(Se_{0.99}Te_{0.01})_2$ thin films. As shown clearly in Table 2, both Te dopant-based devices exhibited lower Voc and higher Jsc values than the one based on the control film. The highest Jsc was obtained in the cell with Cu(In,Ga)(Se_{0.995}Te_{0.005})₂ thin film, and Jsc decreased with further increase of the Te dopant. This decrease may have originated from the high ideality factor, enhanced space charge region recombination, and low carrier collection as a result of the larger amount of Ga in the region close to the back contact (Orgis et al., 2013). Possible mechanisms underlying the enhanced device performance are the faster hole-extraction and transfer with reduction of parasitic optical losses, resulting in higher Jsc and additionally suppressing the defect in GBs through larger grain sizes, leading to higher FF (Dullweber et al., 2001; Gloeckler and Sites, 2005; Liu et al., 2020; Nwakanma et al., 2021). We observed kink effects in the J-V curves of the $Cu(In,Ga)(Se_{0.995}Te_{0.005})_2$ solar cell that may have been the reason for the lower FF. This kink in the current indicates potential barrier(s) to the majority carrier flow (injection barrier), which could be due to window/buffer band alignment or heavy accumulation of excess Na atoms at the interface or in front of the back contact (Weiss et al., 2018;

Werner et al., 2018). In terms of overall evaluation, Te-poor TFSCs afforded an impressive efficiency of 9.58%, proving that a Te anion is a promising dopant for highly efficient wide-bandgap Cu(In,Ga)Se₂ absorber-based solar cells. Figure 4B presents the EQE spectra of the best TFSCs between 350 and 1,050 nm. Both Teincorporated thin films showed very high EQE curves compared to the control thin film, indicating negligible collection losses as a result of the very wide space charge region that is established for close-stoichiometric absorber compositions (Keller et al., 2021). In addition, the decrease in EQE values toward short wavelengths for the control thin -films was caused by reflection losses, providing evidence that a Te anion reduces reflection losses. Beyond this, we evaluated Voc and efficiency as a function of bandgap for Te-incorporated narrow- and wide-bandgap thinfilm-based solar cells that have been reported in the literature in relation to this work. Figure 4C shows Voc as a function of bandgap (Eg) for previously reported solar cells based on narrow- and wide-bandgap chalcopyrite thin films. It is evident that Voc deficiency rapidly decreases with increasing bandgap; as expected, the highest Voc values were reached at bandgap values in the range 1.30-1.45 eV. The best Voc value was achieved at 1.38 eV bandgap in our recent work on Te-poor Cu(In,Ga)(Se,Te)₂based TFSCs. Figure 4D shows the corresponding values to provide a straightforward comparison of the efficiency achieved using different bandgap thin film absorbers. Figure 4D clearly



tellurium-based TFSCs reported in the literature (Mise and Nakada, 2010a, Mise and Nakada, 2010b, Mise and Nakada, 2011; Jehl-Li-Kao et al., 2013; Kim et al., 2013; Hamid et al., 2019; Jia et al., 2019) and in the present article (the latter represented by red stars).

	Jsc (mA/cm ²)	Voc (mV)	FF (%)	PCE (%)
Cu(In,Ga)Se ₂	13.96	777.2	59.3	6.43
Cu(In,Ga)(Se _{0.995} Te _{0.005}) ₂	20.64	728.9	54.7	8.23
Cu(In,Ga)(Se _{0.99} Te _{0.01}) ₂	20.07	688.8	69.3	9.58

TABLE 2	Photovoltaic	parameters	of TESCs.
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shows that our cells based on Te-poor wide-bandgap absorbers exhibited better solar cell performance, achieving 9.58% efficiency, considerably higher than has previously been reported in the literature.

4 Conclusion

We successfully prepared Te-poor CuInGa(SeTe)₂ widebandgap absorbers using RbF PDT via a three-stage evaporation process and examined the impact of Te²⁻ anions on photovoltaic performance in chalcopyrite solar cells. Absorber thin films incorporating Te exhibited effective hole-extraction, a considerable reduction in defects at the surface and grain boundary, lower carrier lifetime, and excellent crystal quality. Compared to control-based TFSCs (6.43%), Cu(In,Ga)(Se_{0.99}Te_{0.01})₂-based cells

exhibited the highest efficiency, at 9.58%. Considering the previous reports in the literature, the Voc and efficiency values of 688 mV and 9.58%, respectively, in the 1.38 eV bandgap represent the highest achieved so far for wide-bandgap TFSCs incorporating Te²⁻ anions. We foresee that future research directions could include the development of Te-based absorber materials with different cations and/or anions and with a device architecture designed with varied buffer layers and/or alkali-fluorine PDT to promote considerable efficiency and Voc for future wide-bandgap TFSCs.

Data availability statement

The original contributions presented in the study are included in the article/Supplementary Material; further inquiries can be directed to the corresponding author.

Author contributions

SA: Investigation, Methodology, Visualization, Formal analysis, Writing - original draft; GC: Supervision, Conceptualization, Funding acquisition, and SS: Project administration, Visualization, Formal analysis, Writing – review and editing. All authors reviewed the article and approved the final version.

Funding

This work was supported by the Scientific and Technological Research Council of Turkey (TUBITAK 2211-C and 2214-A) and the Federal Ministry of Education and Research of Germany (Grant Number: 0324297B).

Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Supplementary material

The Supplementary Material for this article can be found online at: https://www.frontiersin.org/articles/10.3389/ fenrg.2023.1215712/full#supplementary-material

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