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# [Analysis and modeling magnetic](https://www.frontiersin.org/articles/10.3389/fenrg.2023.1127198/full) [energy harvester with field](https://www.frontiersin.org/articles/10.3389/fenrg.2023.1127198/full) [shaping capacitors](https://www.frontiersin.org/articles/10.3389/fenrg.2023.1127198/full)

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A detailed circuit model is discussed in this paper for the operation of magnetic energy harvesters with field shaping capacitors (FSC) feeding constant voltage load. First an equivalent circuit with nonlinear inductance, current source and a diode bridge was given based on the physical analysis of the harvester. Then detailed analysis of the circuit operation and state transitions under various FSC are provided. A mathematical model is established to replace the circuit model so that the optimization of the system parameters, such as number of turns, load voltage, could be carried out directly. Simulation and experimental results are given in the paper to prove the effectiveness of the proposed circuit analysis method.

#### KEYWORDS

magnetic energy harvester, magnetic saturation, field shaping capacitor, advanced sensing, magnetic saturation analysis

## 1 Introduction

The smart grid requires distributed real-time monitoring network to actively sense the states, to define the problems of the power system, to discover knowledge, and to solve the problems [Zhang et al. \(2021\).](#page-9-0) This means that more and more smart sensor grids need to be integrated into the power system and its apparatus. Even though the power consumption of both the sensors and the microelectronic chips has decreased in the last decades, life long energy supply is still one of the biggest challenges for the widely distributed and wirelessly enabled sensor systems [Moon et al. \(2013\).](#page-9-1) A cost-effective and self-sustained smart sensing node system uses energy harvesting to convert ambient energy sources into electric power to supply the electronic systems [\(Sudevalayam and Kulkarni, 2011;](#page-9-2) [Tang et al., 2018\)](#page-9-3). A variety kinds of energy sources [Tang et al. \(2018\),](#page-9-3) for example, vibrations [Ottman et al. \(2002\),](#page-9-4) RF energy [Facen and Boni \(2006\),](#page-9-5) corona current [Shi et al. \(2023\),](#page-9-6) thermal and solar [Tan and](#page-9-7) [Panda \(2011\),](#page-9-7) could be utilized with different techniques, but magnetic energy harvesters based on electromagnetic induction is still the most reliable and cost-effective way to extract energy from the power system itself [\(Moon et al., 2013;](#page-9-1) [White et al., 2018\)](#page-9-8). The power line goes through the center of a magnetic core and the 50 or 60 Hz alternating current in the power line will induce voltage on the secondary side, similar manner with a current transformer. What is different with a current transformer is that the output of the secondary winding is feeding a low impedance load to extract energy. Usually a rectifier is followed to converter the AC voltage to DC so that it can power the sensors and microelectronics.

One of the most challenging problems for the magnetic energy harvester is the magnetic saturation when the transmission line current is high. Reference [Moon and Leeb \(2015a\)](#page-9-9)



<span id="page-1-0"></span>and [Moon and Leeb \(2015b\)](#page-9-10) have developed a circuit model incorporating the nonlinear behaviour of the core to model and predict the harvester system performance. In the analysis of [Moon](#page-9-9) [and Leeb \(2015a\)](#page-9-9) and [Moon and Leeb \(2015c\),](#page-9-11) a constant voltage load other than a pure resistor was used as the load of the energy harvester system, which is of much more practical meaning since the energy harvested is usually used to charge a battery or supercapacitor load. Several groups continued to research on utilizing the saturation region of the core [Park et al. \(2021\)](#page-9-12) and [Gruber et al. \(2021\).](#page-9-13) Reference [Moon and Leeb \(2016\)](#page-9-14) and [Moon](#page-9-15) [and Leeb \(2014\)](#page-9-15) proposed two methods to increase the energy harvested for a constant voltage load without significantly increasing circuit complexity. One method called transfer window alignment (TWA) is to change the energy harvesting timing while another is to increase the time to reach saturation by shaping the core voltage with a field shaping capacitor (FSC) in series connection with the core. A detailed analysis and control strategy of TWA was given in [Moon](#page-9-10) [and Leeb \(2015b\)](#page-9-10) and the energy extracting capability improved considerably. However, no detailed circuit analysis was done on FSC technology.

In this paper, a comprehensive circuit analysis for the FSC method is introduced based on the nonlinear model of the magnetic core. The optimization of the core and system design parameters are presented. Simulation and experimental results are given as verifications.

## 2 Materials and methods

## 2.1 Circuit model of the magnetic core

The primary side could be viewed as a current source  $i_P = I_P$  sin  $\omega t$  where  $I_P$  is the peak amplitude of the wire current. The secondary side of the core has current  $i_{\mathcal{S}}$ . From Ampere's Law, the



<span id="page-1-1"></span>magnetic field intensity  $H$  at radius  $r$  is:

$$
H(r) = \frac{i_P - Ni_S}{r} = \frac{1}{r} \left( I_P \sin \omega t - Ni_S \right) \tag{1}
$$

where N is the number of turns of the secondary side. The induced AC voltage is rectified to DC voltage by a full bridge rectifier. The full bridge rectifier could be implemented by Schottky diodes or MOSFETs. The rectified DC voltage usually charges an energy storage device, such as a supercapacitor or a battery. Since the charging of an energy storage device is much slower compared with the voltage rectifying dynamics, the load of the harvester could be modelled as a voltage source  $V_L$ .



<span id="page-2-1"></span>

<span id="page-2-2"></span>For a ring shape core, B is almost constant across the range under the assumption that  $H(r)$  changes little between the outer and inner radius. The flux linkage in the core is:

$$
\Gamma = NAB \tag{2}
$$

where A is the cross section area of the core. The induced voltage of the harvester  $v_{core}$  is the time derivative of the flux linkage:

<span id="page-2-0"></span>
$$
v_{core} = \frac{d\Gamma}{dt} \tag{3}
$$

H is calculated from the primary and secondary currents as shown in 1. Using the equations above, the flux linkage is:

$$
\Gamma = NAB = NAf(H) = NAf\left(\frac{1}{r}\left(I_p\sin\omega t - Ni_S\right)\right) \tag{4}
$$

Where  $B = f(H)$  is the magnetizing curve. This equation means that the flux linkage in the core relates with the difference between the primary current and the secondary current. Therefore the current transformer of the magnetic energy harvester could be modelled as an ideal current source in parallel with a non-linear magnetizing inductance  $L_m$ , as shown in [Figure 1.](#page-1-0)



<span id="page-3-0"></span>

<span id="page-3-2"></span>The complicated relationship between B and H determines the harvesting capability of the harvester. In the linear region where  $B =$  $\mu$ H,  $L_m$  is:

<span id="page-3-1"></span>
$$
L_m = \mu \frac{N^2 A}{l} \tag{5}
$$

where  $\mu$  is the permeability of the core material and l is the effective flux length of the core. Since  $N$  is very large,  $L_m$  is large enough so that almost no current flows through it in the linear region. The leakage inductance of the harvester could also be ignored since it is much smaller compared to  $L_m$ .



<span id="page-4-0"></span>

<span id="page-4-1"></span>In the saturation region,  $B \approx B_{sat}$ . From 2,  $\Gamma \approx NAB_{sat}$ . Eq. [3](#page-2-0) shows that the saturation of B or  $\Gamma$  is directly caused by accumulating voltage-seconds which is the time integral of  $v_{core}$ . When feeding a constant voltage load, the time it takes to saturate the core is:

$$
T_{sat} = \frac{2B_{sat}NA}{V_L} \tag{6}
$$

The effective permeability of the core decreases dramatically and most of the induced current  $i<sub>S</sub>$  is bypassed by the magnetizing inductor. Much less energy is harvested at this region. However, crudely dividing the magnetizing curve into linear/non-linear regions is far from accurately predicting the harvested energy. There are several methods to establish mathematical model for a



<span id="page-5-2"></span>non-linear magnetic core. In [Vos \(2020\)](#page-9-16) the model is established by treating the core permeance as a complex parameter that includes loss. [Kuang et al. \(2021\)](#page-9-17) used finite element method in both electromagnetic and electrical system. In [Moon and Leeb \(2015a\),](#page-9-9) the "arctan" function was used to model the saturation where there was one fitting parameter  $α$ . In this research, the Chan model [Chan et al. \(1991\)](#page-9-18) was chosen since it is relatively easy to implement in circuit simulation tools and provides relatively more freedom to emulate the saturation behavior. The Chan model was defined by three parameters:

<span id="page-5-0"></span>
$$
B = \frac{H + H_c}{|H \pm H_c| + H_c \left(\frac{B_{sat}}{B_r} - 1\right)} + \mu_0 H
$$
 (7)

where  $B_{sat}$  is the saturation flux density;  $B_r$  is the remnant flux density and  $H_c$  is coercive force.

A high permeability amorphous non-crystalline core VITROPERM 500F [VacuumSchmelze \(2022\)](#page-9-19) was used as the magnetic core. The fitted parameters for this core material are:  $B_{\text{sat}} = 1.19$  T,  $B_r = 0.1$  T,  $H_c = 0.3$  A/m. The simulated B-H curve using the Chan model is shown in [Figure 2,](#page-1-1) which is very close to the curve given in the datasheet. Note that the H of the core starts from 0 and is swept forward and backward. The hysteresis of the material is relatively small so the iron loss of the core is ignored in this paper.

## 2.2 Circuit model of the magnetic energy harvester

In the circuit of [Figure 1,](#page-1-0) the voltage of the core  $v_{core}$  equals to the load voltage  $V_L$  during the linear region when neglecting the voltage

drop across diodes. For the rectifier circuit with an FSC as shown in [Figure 3,](#page-2-1)  $v_{core}$  is not constant any more due to the capacitor voltage  $v_C$ . Since $v_C$  lags  $i_S$ ,  $v_C$  might be negative at the start of one period.  $v_{core}$  could be less than  $V_L$  at the beginning of the energy transfer. Therefore the total energy transfer time  $T_{sat}$  may extend after adding an FCS.

The voltage and current waveforms of the harvester with and without FSC are shown in [Figure 4.](#page-2-2) The peak value of the primary current is 14 A, the secondary side has  $N = 200$  turns and the load voltage  $V_I$  = 7 V. The value of the FSC is 20  $\mu$ F. The rectifying diodes are Schottky diodes 1N5819 from DIODES whose forward voltage  $V_F$  is 0.2 V. As expected, with FSC  $v_{core}$  is smaller than  $V_L$  at the beginning of the energy transfer.  $T_{sat}$  is 0.5 ms wider, which means more energy is harvested with FSC. The power harvested at different FSC value (5, 15, and 25  $\mu$ F) is 80, 125, and 87 mW respectively. The power extracted from the core under a certain operating point does not increase monotonically with the value of FSC. There should be an optimal value for FSC. The waveforms of  $v_{core}$  and  $v_c$  are shown in [Figure 5.](#page-3-0) Even though the peak value of  $v_C$  and  $v_{core}$  increases as FSC decreases,  $T_{sat}$  first increases and then decreases. The shape of  $v_C$  also becomes different at 5  $\mu$ F. In order to find out the optimized FSC and load voltage  $V_L$  under a specified design, a detailed circuit analysis should be carried out.

## <span id="page-5-1"></span>2.3 Circuit analysis of the magnetic energy harvester with FSC

Since the circuit behaves symmetrically at  $\omega \in [0, \pi]$  and  $[\pi, 2\pi]$ , the analysis will be focused on the first half of the period. At the beginning of a new sinusoidal period,  $L_m$  is in its linear region and it



<span id="page-6-7"></span>is so large that  $i_{\mathcal{S}}$  entirely flows through FSC. Diodes  $\mathrm{D}_2$  and  $\mathrm{D}_3$  turn on. From KVL:

<span id="page-6-0"></span>
$$
v_{core} = v_C - v_L + i_S R_w - 2V_F \tag{8}
$$

where  $v_C$  and  $v_R$  are the voltages of FSC and wire resistance respectively, as shown in [Figure 3.](#page-2-1) By integrating  $i_S$ ,  $v_C$  is:

<span id="page-6-6"></span>
$$
v_C(t) = V_{C0} + \frac{I_p}{N\omega C} (\cos \omega t - 1)
$$
 (9)

where  $V_{C0}$  is initial voltage of the FSC at  $t = 0$ . Apparently the capacitor voltage at  $t = T_{sat}$  depends on the value of FSC. Therefore, the circuit operation could be divided into two scenarios depending on the value of FSC.

### 2.3.1 Large FSC

At  $t = T<sub>sat</sub>, L<sub>m</sub>$  enters into saturation region. Its value decreases dramatically and  $i_S$  flows entirely into  $L_m$ . Based on the reference direction definition in [Figure 3,](#page-2-1) the voltage-seconds accumulating in the core drive it from positive saturation to negative saturation:

<span id="page-6-1"></span>
$$
\int_{0}^{T_{sat}} v_{core} dt = -2B_{sat} NA
$$
 (10)

Diodes  $D_2$  and  $D_3$  turn off at  $T_{sat}$ .  $v_C$  keeps unchanged during the rest of the half period since there is no current flowing through it.  $v_C(T_{sat})$  will be the capacitor voltage initial value for  $[\pi, 2\pi]$ . From the principle of symmetry:

<span id="page-6-2"></span>
$$
v_C(T_{sat}) = -V_{C0} \tag{11}
$$

Solving Eqs  $8$ ,  $10$ ,  $11$  numerically, the value of  $T_{sat}$  and  $V_{C0}$  could be obtained. The average power harvested is:

<span id="page-6-5"></span>
$$
P_h = \frac{2}{T} \int_0^{T_{sat}} V_L \frac{I_p}{N} \sin \omega t \mathrm{d}t \tag{12}
$$

## 2.3.2 Small FSC

From [Figure 5,](#page-3-0) the capacitor voltage  $v_C$  is not continuous when  $C = 5 \mu F$ . The capacitor discharges rapidly once the core enters saturation. Therefore the relationship in Eq. **[11](#page-6-2)** is no longer valid. To solve  $T_{sat}$  and  $V_{C0}$ , the discharging process of the FSC should be examined. At  $t = T_{sat}$ , the current source is bypassed by the magnetizing inductance  $L_m$  and diodes  $D_2$  and  $D_3$  turn off. In view of the fact that the value of  $i_S$  and  $L_m$  is small, the voltage drop on  $L_m$ is very small compared to  $v_C(T_{sat})$ . If the absolute value of  $v_C(T_{sat})$ is larger than  $V_L$ , diodes  $D_1$  and  $D_4$  will turn on to discharge the capacitor through the voltage source  $V_L$ , magnetizing inductance  $L_m$ and wire resistance  $R_w$ . This is a classical RLC transient process. The resonance between C and  $L_m$  will end when  $i_C$  goes back to zero and diode  $D_1$  and  $D_4$  will turn off then. Since the resonant frequency of C and  $L<sub>m</sub>$  is large compared with the line frequency so the discharging time of the capacitor is very short. After discharging process,  $v_C$  stays unchanged since all the diodes are off.

After solving the 2nd order ODE of the RLC circuit, the capacitor voltage during discharging is:

<span id="page-6-3"></span>
$$
v_C(t) = e^{-\alpha t} \left( \frac{\alpha}{\omega_d} \left( V_{C0} + V_L + 2V_F \right) \sin \omega_d t + \left( V_{C0} + V_L + 2V_F \right) \cos \omega_d t \right)
$$
\n(13)

where  $\alpha$  is the damping coefficient and  $\omega_d$  is the oscillation frequency. They could be calculated as:

<span id="page-6-4"></span>
$$
\omega_0 = \frac{1}{\sqrt{LC}}
$$
  
\n
$$
\alpha = \frac{1}{2} \frac{R_w}{L}
$$
  
\n
$$
\omega_d = \sqrt{\alpha^2 - \omega_0^2}
$$
\n(14)



<span id="page-7-1"></span>The discharging time for the capacitor is half of the RLC oscillation period  $\pi/\omega_d$ .  $v_C$  drops to  $-V_{C0}$  according to the principle of symmetry. Together solving Eqs **[10](#page-6-1)**, **[13](#page-6-3)**, **[14](#page-6-4)** numerically, the value of  $T_{sat}$  and  $V_{C0}$  could be obtained. The average power harvested during  $D_2$  and  $D_3$  on is still based on Eq. [12](#page-6-5) while the average power

harvested during  $D_1$  and  $D_4$  on is:

<span id="page-7-0"></span>
$$
P_{h,d} = \frac{2}{T} C (V_{C0} + V_L + 2V_F) V_L \left( 1 + e^{-\alpha \frac{\pi}{\omega_d}} \right)
$$
 (15)

It is worth noting that the value of L used in Eqs **[13](#page-6-3)**–**[15](#page-7-0)** comes from calculation by Eqs **[5](#page-3-1)**, **[7](#page-5-0)**.



<span id="page-8-0"></span>A comparison of the harvested power calculated by the mathematical model presented above and circuit simulation is shown in [Figure 6.](#page-3-2) The simulation condition is the same as [Section 2.3:](#page-5-1)  $I_p = 14 \text{ A}$ ,  $N = 200$ ,  $C = 15 \mu \text{F}$ . The two methods output very similar results and the biggest difference is less than 5%. Therefore it is more convenient to use the numerical model to optimize the system design.

# 3 Results

In this section, the optimized field shaping capacitor value C and the harvester core design will be investigated. This equation is valid as long as the primary current is not high enough so the voltage drop on the wire resistance could be ignored. The optimized FSC value with different  $V_I$  under the condition ( $I_p = 14$  A,  $N = 200$ ) is shown in [Figure 7.](#page-4-0) The optimal FSC decreases as the load voltage  $V_L$  increases. Note that the optimal power harvested with FSC under different  $V_L$  is 0.124, 0.125, 0.118, and 0.111 W respectively. Without FSC, the power harvested is 0.109, 0.099, 0.090, and 0.082 W respectively from 6. A proper FSC can improve the energy extraction capability significantly, especially when the load voltage is high. Since higher  $V_L$  results in shorter  $T_{\mathit{sat}},$  the harvester will extract less power as the load voltage increases if no special voltage shaping techniques are used. Higher  $V_{L}$  also means a stronger voltage shaping, which means a smaller FSC is needed.

The optimized FSC value with different N under the condition  $(I_p = 14 \text{ A}, V_L = 6 \text{ V})$  is shown in [Figure 8.](#page-4-1) [Figure 8](#page-4-1) reveals the fact that N has a very complicated impact on the energy harvesting efficiency. When the number of turns is low, the core is easier to get saturated based on 6 while the secondary current would be too small to effectively charge the load if N is large.

The optimized FSC value with different primary current peak value  $I_p$  under the condition ( $V_L$  = 6 V, N = 200) is shown in [Figure 9.](#page-5-2) The maximum amount of energy harvest is proportional to  $I_p$ , as predicted by 12. The optimal FSC changes rapidly as  $I_p$  increases, from 5 to 20  $\mu$ F as  $I_p$  increases from 3 to 20 A. This is owing to the fact that small FSC could be easily charged up by large  $I_p$  based on Eq. [9](#page-6-6), which results a stronger field shaping effect. [Figure 10](#page-6-7) is the power vs. current plot. The power monotonically increases with primary current, as indicated in Eq. **[12](#page-6-5)**.

An experimental prototype was built to verify the analysis. [Figure 11](#page-7-1) shows the current/voltage waveforms of the harvester with different FSC. The current of the constant voltage load was measured through a current sensing resistor and differential amplifier. [Figures 11A, B](#page-7-1) show that the power transfer time increases significantly with a 14.7  $\mu$ F FSC. For small FSC, 4.7  $\mu$ F for example, as in [Figure 11C,](#page-7-1) the current peak is the result of RLC discharging explained in Eq. **[13](#page-6-3)**.

The harvested power from experiments and numerical calculation with different FSC is shown in [Figure 12.](#page-8-0)The experiment condition is:  $I_p = 14$  A,  $N = 200$ ,  $V_L = 5$  V. Simulation and experimental results show the same trend. The difference between the experiments and numerical calculation mainly comes from the discrepancies between the piecewise assumption of the magnetic core's B-H curve in the numerical model and its real saturation curve.

# 4 Discussion

This paper discusses the modeling and optimization process for the magnetic energy harvester with field shaping capacitors in order to enhance its energy extracting capability. The Chan model is used to emulate the saturation of the magnetic core. Then a circuit model is given to study the behaviour of the harvester system feeding a constant voltage load. Based on the circuit model, a detailed mathematical model is developed based on the value of FSC so that the fast calculation and optimization of the system performance is possible through numerical software. Optimization of the harvester parameters under various operating conditions are discussed. The optimal FSC is sensitive to the operation conditions of the energy harvester: load voltage, number of turns and primary side current. Therefore FSC is especially suitable for applications with almost constant current, for example, harvesting energy on the cable of an induction machine. With the model introduced in this paper, designers could maximize the output power through an optimum design of the harvester: number of turns, size of the core and field shaping capacitors.

## Data availability statement

The datasets presented in this article are not readily available because the dataset are not available to the public due to authors' organization policies. Requests to access the datasets should be directed to [minsun@uestc.edu.cn.](mailto:minsun@uestc.edu.cn)

## Author contributions

YM, MS, MZ, and JL contributed to conception and design of the study. GC contributes to the theoretical part. YM designed the experiments. ML and MZ performed simulation and experiments. YM wrote the first draft of the manuscript. MS wrote sections of the

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The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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