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# Pore-scale phenomena in carbon geological storage (Saline aquifers—Mineralization—Depleted oil reservoirs)

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The injection of CO<sub>2</sub> into geological formations triggers inherently coupled thermo-hydro-chemo-mechanical processes. The reservoir pressure and temperature determine the CO<sub>2</sub> density, the CO<sub>2</sub>-water interfacial tension, and the solubility of CO<sub>2</sub> in water (hindered by salts and competing gases). The CO<sub>2</sub>-water interface experiences marked pinning onto mineral surfaces, and contact angles can range from the asymptotic advancing to receding values, in contrast to the single contact angle predicted by Young's equation. CO<sub>2</sub> dissolves in water to form carbonic acid and the acidified water dissolves minerals; mineral dissolution enhances porosity and permeability, triggers settlement, may couple with advection to form "wormholes", produces stress changes and may cause block sliding and shear bands. Convective currents can emerge beneath the CO<sub>2</sub> plume and sustain CO<sub>2</sub> and mineral dissolution processes. On the other hand, mineralization is a self-homogenizing process in advective regimes. The crystallization pressure can exceed the tensile capacity of the host rock and create new surfaces or form grain-displacive lenses. Within the rock matrix, coupled reactive-diffusion-precipitation results in periodic precipitation bands. Adequate seal rocks for CO<sub>2</sub> geological storage must be able to sustain the excess capillary pressure in the buoyant CO<sub>2</sub> plume without experiencing open-mode discontinuities or weakening physico-chemical interactions. CO<sub>2</sub> injection into depleted oil reservoirs benefits from time-proven seals; in addition, CO<sub>2</sub> can mobilize residual oil to simultaneously recover additional oil through oil swelling, ganglia destabilization, the reduction in oil viscosity and even miscible displacement. Rapid CO<sub>2</sub> depressurization near the injection well causes cooling under most anticipated reservoir conditions; cooling can trigger hydrate and ice formation, and reduce permeability. In some cases, effective stress changes associated with the injection pressure and cooling thermoelasticity can reactivate fractures. All forms of carbon geological storage will require large reservoir volumes to hold a meaningful fraction of the CO<sub>2</sub> that will be emitted during the energy transition.

## KEYWORDS

porous media, carbon geological storage, dissolution, mineralization, oil recovery, thermal effect, seals

## Introduction

We emit ~40 billion tons of CO<sub>2</sub> per year, however, emissions should decrease to ≤18 billion tons per year to prevent further global warming (IPCC, 2021). This scenario will require reduced fossil fuel consumption and CO<sub>2</sub> capture and geological storage. Currently, the total capacity of operating CO<sub>2</sub> storage facilities is 150 million tons/year (GCCSI, 2021), i.e., three orders of magnitude lower than needed. Therefore, thousands of new geological storage projects must come online in the near future.

The injection of CO<sub>2</sub> into geological formations triggers inherently coupled thermo-hydro-chemo-mechanical processes. In this manuscript, we present a comprehensive review of the underlying physical processes, summarize the properties of bulk phases and interfaces, and analyze the most relevant pore-scale phenomena that can affect the long-term geological storage of CO<sub>2</sub>.

## CO<sub>2</sub> and aqueous solutions: Properties and interactions

Fluids involved in carbon geological storage exhibit different pressure and temperature-dependent bulk properties, solubilities and capillary response.

### Density

The pressure and temperature-dependent CO<sub>2</sub> density determines storage capacity, gravity-driven displacement, instabilities, and capillary pressure against the seal (Van der Meer, 1993; DiCarlo, 2013). Figure 1 shows the density  $\rho$  [kg/m<sup>3</sup>] of water, oil and CO<sub>2</sub> as a function of pressure and temperature: the density of CO<sub>2</sub> varies significantly but remains a buoyant phase, while water is the heavier fluid in high-temperature reservoirs.

Geological storage projects can involve gas, liquid and supercritical CO<sub>2</sub> (Figure 2). The gas-liquid phase boundary  $P = 0.0008(\frac{T}{C})^2 + 0.1(\frac{T}{C}) + 3.5$  MPa applies until the critical point ( $T = 31^\circ\text{C}$  and  $P = 7.37$  MPa). In the supercritical regime, beyond the critical point, the Widom line divides the lighter gaseous-like regime from the denser liquid-like regime (Simeoni et al., 2010; Houben et al., 2021). Density gradients in  $P$  or  $T$  are pronounced across the Widom line, particularly near the critical point; for example, the CO<sub>2</sub> density changes from 661 kg/m<sup>3</sup> to 285 kg/m<sup>3</sup> when the temperature increases from 35°C to 50°C at a constant pressure of  $P = 8.5$  MPa (see Figure 2).

### Viscosity

Liquid CO<sub>2</sub> and water are Newtonian fluids, yet most crude oils exhibit some shear thinning. Figure 3 shows the reduction in viscosity  $\mu$  [μPa·s] with temperature for different types of crude oils, water and CO<sub>2</sub> (Note: viscosity shown in log-scale). Molecular vibration acts against intermolecular attraction hence the viscosity of liquids decreases with higher temperature while pressure has a second-order effect. On the other hand, the viscosity of gases tends to increase with temperature and pressure because the collision frequency increases. The viscosity of CO<sub>2</sub> is much lower than water and oil, and it experiences sharp changes across the liquid-gas LG phase boundary and the Widom line WL in agreement with density changes.

### Interfacial tension

Interfacial tension  $\gamma$  [mN/m] and contact angle  $\theta$  [°] determine capillary phenomena between multiphase fluids in porous media. The capillary pressure  $\Delta P$  is the pressure difference between the non-wetting  $P_{nw}$  [Pa] and the wetting  $P_w$  [Pa] fluid; for a spherical interface of radius  $r$  [m]

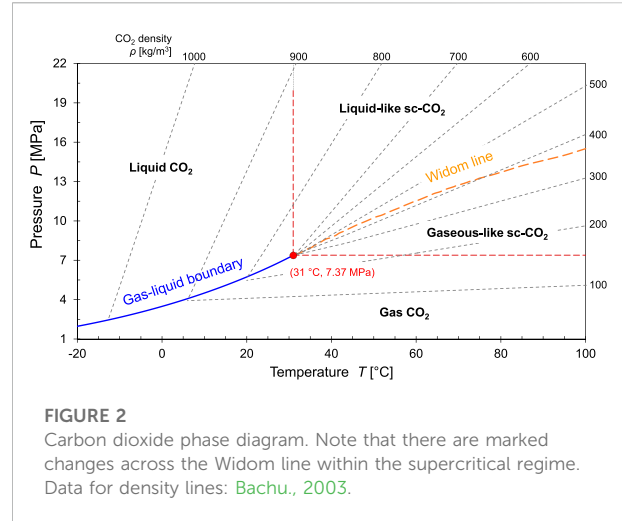
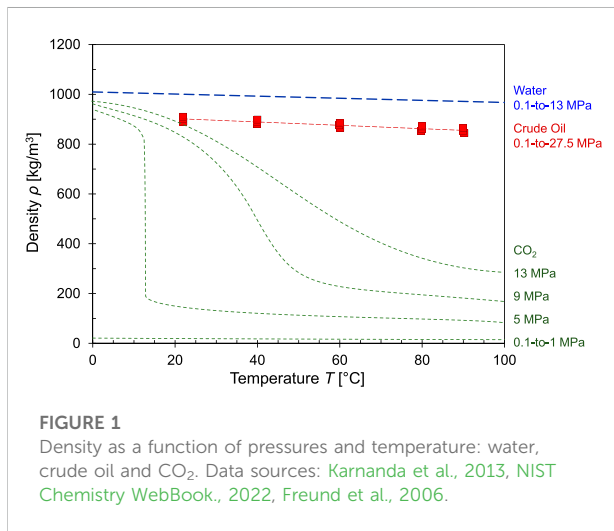
$$P_{nw} - P_w = P_c = \frac{2\gamma \cdot \cos \theta}{r} \quad (1)$$

The interfacial tension between CO<sub>2</sub> and water is reported as a function of pressure and temperature in most cases. However, intermolecular distance  $l_0$  [m] determines molecular interaction at the interface, therefore the CO<sub>2</sub> density  $\rho$  [kg/m<sup>3</sup>] which scales as  $\rho \propto l_0^{-3}$  is a more meaningful parameter to predict the interfacial tension. Figure 4 plots compiled CO<sub>2</sub>-water interfacial tension data gathered from the literature and our own studies as a function of CO<sub>2</sub> density (computed from reported pressure and temperature conditions): there is a clear decrease in the interfacial tension with increasing CO<sub>2</sub> density. The remaining spread in the data indicates that thermal vibration affects interfacial tension beyond its role on density (in addition to measurement variability among devices and test protocols).

### Wettability—contact angles

Wettability captures the interaction between two fluids and the mineral surface. Wettability plays a critical role in capillary phenomena, pore-scale fluid arrangement, and macroscopic multiphase flow in CO<sub>2</sub> geological storage.

Contact angle measurements provide an apparently simple assessment of the otherwise complex physical processes that



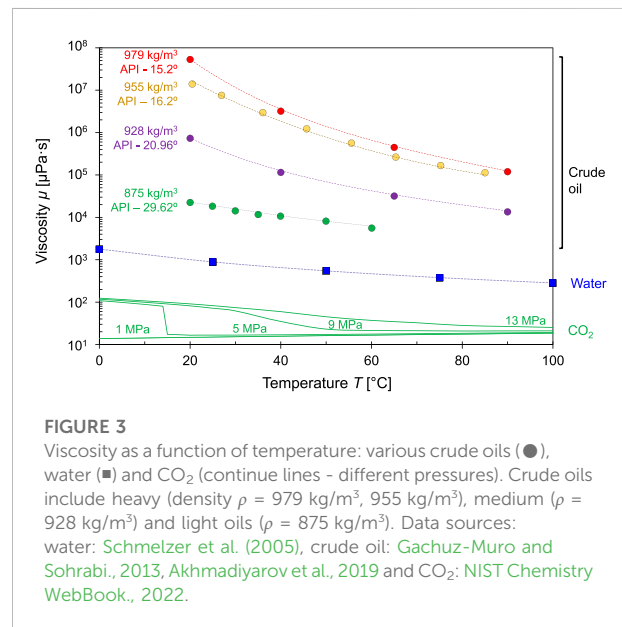
underline wettability (Sargent et al., 2005; Drelich et al., 2020). Force equilibrium parallel to the mineral surface relates the contact angle to the interfacial tensions  $\gamma$  [N/m] between the fluid  $F$ , liquid  $L$  and solid  $S$  (Young, 1805):

$$\sum F_H = (\gamma_{FS} - \gamma_{LS}) - \gamma_{LF} \cos \theta \quad (2)$$

On the other hand, force equilibrium normal to the substrate surface allows us to anticipate the adhesion  $A$  [N/m] between the liquid-fluid interface and the mineral substrate  $A = \gamma_{LF} \sin \theta$ . Equilibrium contact angle measurements show that both calcite and quartz are strongly water-wet (Figure 5, Bikkina, 2011; Arif et al., 2017a).

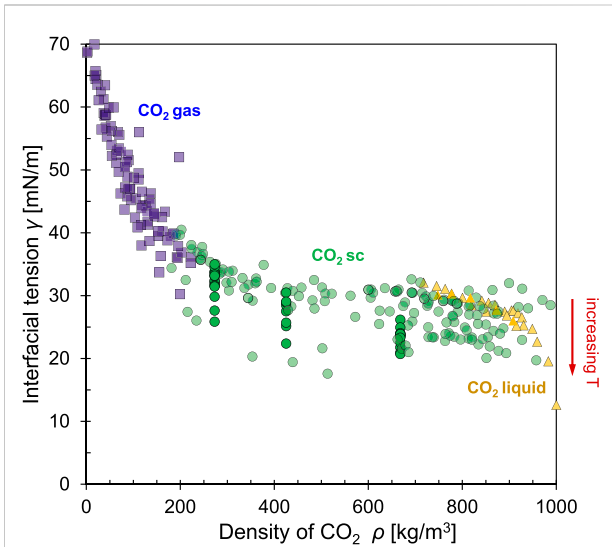
In contrast to the single thermodynamic contact angle predicted by Young's equation, interfaces can exhibit a range of contact angles between two asymptotic values: the advancing contact angle  $\theta_A$  and the receding contact angle  $\theta_R$  when it is measured in a quasi-static condition. For example, a deionized water droplet resting on a horizontal smooth calcite surface surrounded by supercritical CO<sub>2</sub> at  $P = 9$  MPa and  $T = 40^\circ\text{C}$  forms an advancing contact angle  $\theta_A = 138^\circ$  and a receding contact angle  $\theta_R = 20^\circ$  (quasi-static injection/extraction measurements conducted for this study). Even dynamic measurements show marked hysteresis; for example, a water droplet released on a smooth calcite surface in a CO<sub>2</sub> environment at  $P = 20$  MPa and  $T = 35^\circ\text{C}$  exhibits an advancing contact angle  $\theta_A = 122^\circ$  and a receding contact angle  $\theta_R = 108^\circ$  (see inset in Figure 5—Note: dynamic contact angle measurements based on a sliding droplet do not capture the full pinning effect of the contact line measured in quasi-static tests at the verge of sliding, thus, hysteresis in dynamic tests is smaller than in quasi-static tests—Huhtamäki, et al., 2018).

X-ray CT measurements confirm these observations and show that a wide range of contact angles exist *in-situ* possibly



due to (Andrew et al., 2014): 1) the contact line adhesion  $A$ , 2) surface roughness and chemical inhomogeneities on the solid surface (de Gennes, 1985; Eral et al., 2013), and 3) the fact that *in-situ* contact angles rarely represent equilibrium conditions.

Let's modify Young's force equilibrium analysis parallel to the mineral surface to take into consideration the adhesion force  $A$  acting against both advancing on receding directions; then the anticipated contact angle hysteresis is  $\cos \theta_R - \cos \theta_A = 2A/\gamma_{LF}$ . Conversely, we can recover the adhesion force from contact angle hysteresis measurements  $A = \gamma_{LF} (\cos \theta_R - \cos \theta_A)/2$ . Evidently, the adhesion force scales with the liquid-fluid interfacial tension  $\gamma_{LF}$  and should not be ignored in wettability analyses. For example, consider the water-CO<sub>2</sub>-calcite system reported above: the computed adhesion force  $A = 18.8$  mN/m while the



**FIGURE 4**

The interfacial tension between CO<sub>2</sub> and water as a function of CO<sub>2</sub> density. The dataset combines results for CO<sub>2</sub> in gas (■), super critical (●) and liquid states (▲). Note: density determines molecular proximity for electrical interactions, thus, it emerges as the governing variable. Still, the data spread is due to thermal effects: interfacial tension decreases with increased random vibration. Data sources: Hough et al., 1952, Park et al., 2005, Sutjiadi-Sia et al., 2008, Chun and Wilkinson., 1995, Da Rocha et al., 1999, Hebach et al., 2002, Chiquet et al., 2007, Kvamme et al., 2007, Espinoza and Santamarina., 2010, Georgiadis et al., 2010, and measurements by the authors for this study (2022).

water-CO<sub>2</sub> interfacial tension is  $\gamma_{LF} = 22.39 \text{ mN/m}$  ( $\theta_A = 138^\circ$ ,  $\theta_R = 20^\circ$ ).

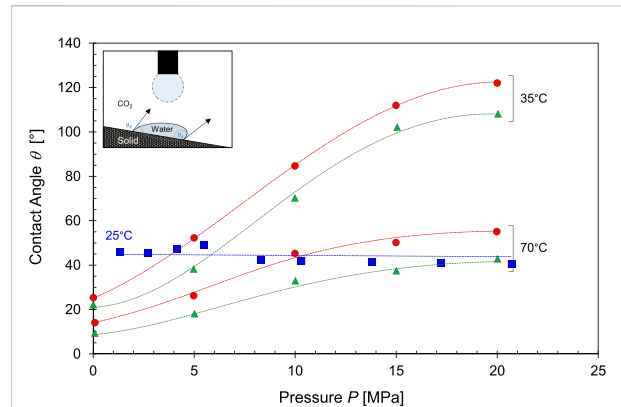
### CO<sub>2</sub> solubility in aqueous solutions

While two phases may exhibit interfacial tension, their mutual solubilities can have important implications. For example, CO<sub>2</sub> dissolves in water and enhances carbon geological storage, on the other hand, minerals dissolve in CO<sub>2</sub>-acidified water and may trigger internal instabilities (details follow).

The solubility of CO<sub>2</sub> in water increases with pressure but decreases with temperature. Henry's law relates solubility  $\chi$  [mol/l] to the partial pressure of CO<sub>2</sub>  $P_{CO_2}$  [kPa] through the temperature dependent Henry's constant  $k_H$  [mol L<sup>-1</sup> kPa<sup>-1</sup>]:

$$\chi_{CO_2} = k_H P_{CO_2} \quad (3)$$

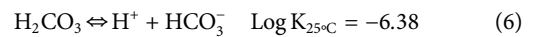
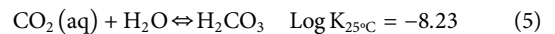
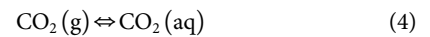
where  $k_H = 3.4 \times 10^{-4}$  [mol L<sup>-1</sup> kPa<sup>-1</sup>] at temperature  $T = 298.15 \text{ K}$ . Figure 6A shows the evolving CO<sub>2</sub> solubility in water as a function of pressure for different temperatures.



**FIGURE 5**

Contact angle formed by a water droplet on calcite in a CO<sub>2</sub> atmosphere as function of CO<sub>2</sub> pressure. Advancing contact angle (●), receding contact angle (▲) and static contact angle (■). Data sources: Bikkina, 2011, Arif et al., 2017a. For clarity, only a subset of the dataset is presented in this figure.

CO<sub>2</sub> dissolves in water to form aqueous carbon dioxide CO<sub>2</sub> (aq). The ionization of carbonic acid produces H<sup>+</sup> and decreases the pH:



Therefore, the equilibrium pH for water with dissolved CO<sub>2</sub> is also a function of pressure and temperature (Figure 6B).

### Effect of salinity

High salinity is a common condition in reservoirs considered for CO<sub>2</sub> geological storage, such as saline aquifers, depleted hydrocarbon reservoirs and coalbeds. In general, CO<sub>2</sub> solubility in water may either increase (salting-in effect, e.g. NaClO<sub>4</sub> and K<sub>3</sub>PO<sub>4</sub>) or decrease (salting-out effect, e.g., salts from strong bases and strong acids such as NaCl, CaCl<sub>2</sub> and MgCl<sub>2</sub> (Yasunishi and Yoshida 1979; García, 2005; García et al., 2005). In most geological settings, CO<sub>2</sub> solubility decreases with salt concentration as water mobility decreases, particularly for small high-valence ions. Consequently, the CO<sub>2</sub> solubility in salt solutions at the same ionic concentration follows the following sequence: highest in KCl solutions → NaCl → CaCl<sub>2</sub> → lowest in MgCl<sub>2</sub> solutions. Figure 7 shows that the solubility of CO<sub>2</sub> decreases by ~50% for a 3 mol NaCl solution compared to deionized water. Therefore, the storage capacity for dissolved CO<sub>2</sub> decreases by half and the needed reservoir size doubles in brine saturated formations.

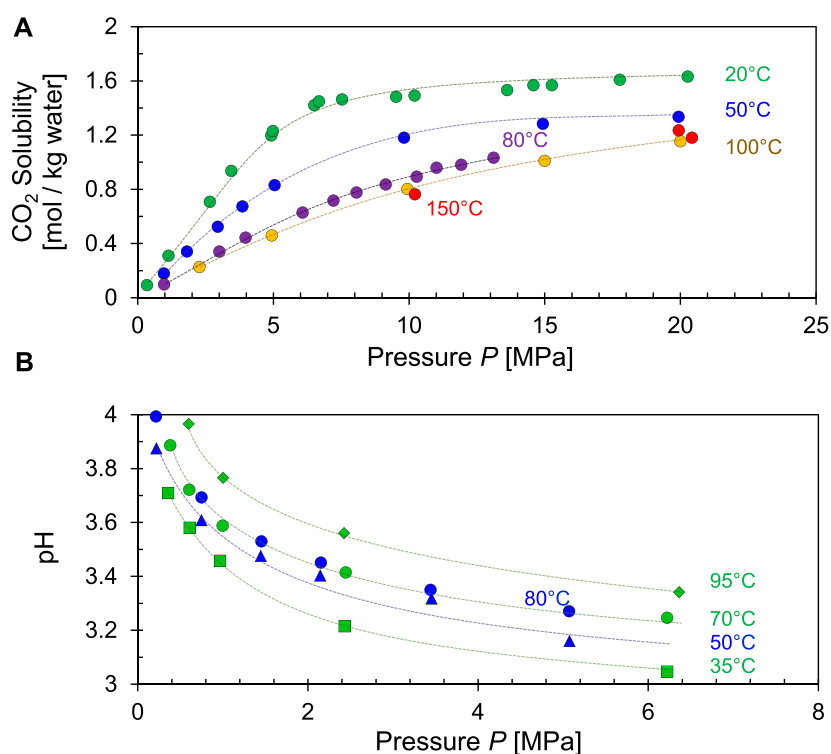


FIGURE 6

Solubility of CO<sub>2</sub> in water and change in pH. (A) CO<sub>2</sub> solubility increases with pressure but decreases with temperatures—Data sources: Duan and Sun., 2003, Sanaei et al., 2019. (B) Equilibrium pH for CO<sub>2</sub> saturated water as a function of pressure and temperatures—Data sources: blue = Haghi et al., 2017; green = Peng et al., 2013.

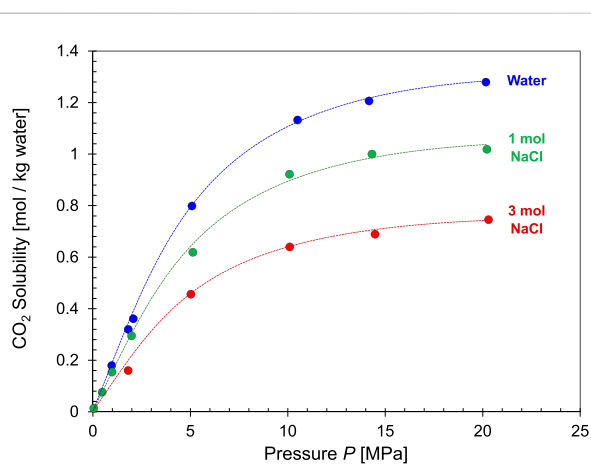


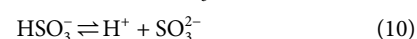
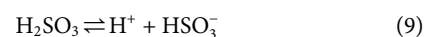
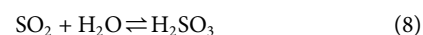
FIGURE 7

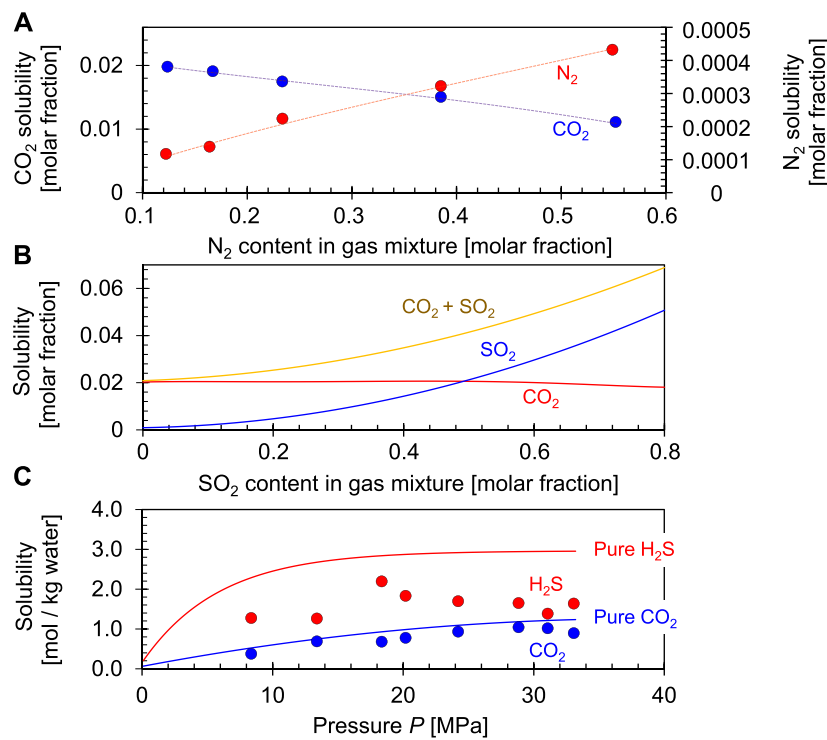
The effect of salinity on CO<sub>2</sub> solubility in NaCl solutions as a function of pressures ( $T = 50^{\circ}\text{C}$ —Data source: Koschel et al., 2006). Data at  $p = 101.3\text{ kPa}$  were gathered at  $T = 25^{\circ}\text{C}$  thus solubility values are higher than the rest of the trend for  $T = 50^{\circ}\text{C}$  (Data sources: Markham and Kobe., 1941; Koschel et al., 2006; Duan et al., 2006; Kiepe et al., 2002).

### Effect of mixed gases

Flue gas released at power plants and industrial sources contains N<sub>2</sub> (~66%), CO<sub>2</sub> (10%–15%), water vapor, other gases such as SO<sub>2</sub>, H<sub>2</sub>S, NO<sub>2</sub>, and light hydrocarbons such as CH<sub>4</sub> (Bell et al., 2021). Separation is costly and the presence of competing gases affects the solubility of CO<sub>2</sub> in water mainly because: (1) the solubility of CO<sub>2</sub> in water depends on its partial pressure which decreases in a mixed gas at constant total pressure, and (2) competing ions. In the case of a N<sub>2</sub>-CO<sub>2</sub>-water system, N<sub>2</sub> does not react with water, therefore the solubility of CO<sub>2</sub> only depends on its partial pressure or mole fraction in the mixed gas (Figure 8A).

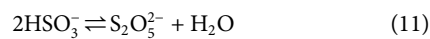
Similar to salt, coexisting gases can either promote or inhibit CO<sub>2</sub> solubility, yet competition is more common in carbon storage. For example, the dissolution of SO<sub>2</sub> in water involves the following reactions (Goldberg and Parker., 1985):



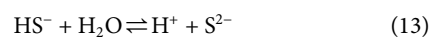


**FIGURE 8**

Effects mixed gas of CO<sub>2</sub> solubility. (A) CO<sub>2</sub> and N<sub>2</sub> solubilities in the N<sub>2</sub>+CO<sub>2</sub>+H<sub>2</sub>O systems at a T = 34.85°C and p = 8.0 MPa. Data source: Nakhaei-Kohani et al., 2022. (B) CO<sub>2</sub> and SO<sub>2</sub> solubilities in the SO<sub>2</sub>+CO<sub>2</sub>+H<sub>2</sub>O systems at a T = 50°C and p = 10 MPa. Data source: Miri et al., 2014. (C) CO<sub>2</sub> and H<sub>2</sub>S solubilities in the H<sub>2</sub>S + CO<sub>2</sub>+H<sub>2</sub>O systems at a T = 120°C. Trends and data from Savary et al., 2012.



The generated H<sup>+</sup> competes with the ionization of carbonic acid and decreases the CO<sub>2</sub> solubility (Figure 8B). Similarly, H<sub>2</sub>S dissolution in water inhibits CO<sub>2</sub> dissolution as well (Figure 8C):



## Fluid-Mineral interaction: Dissolution and precipitation

Water acidifies in the presence of CO<sub>2</sub> and prompts mineral dissolution. The consequences of reactive fluids on the reservoir and seal stability require careful consideration.

### Mineral dissolution in CO<sub>2</sub> acidified water

A mineral's crystal structure determines its susceptibility to dissolution in CO<sub>2</sub> acidified water. In particular, crystallographic

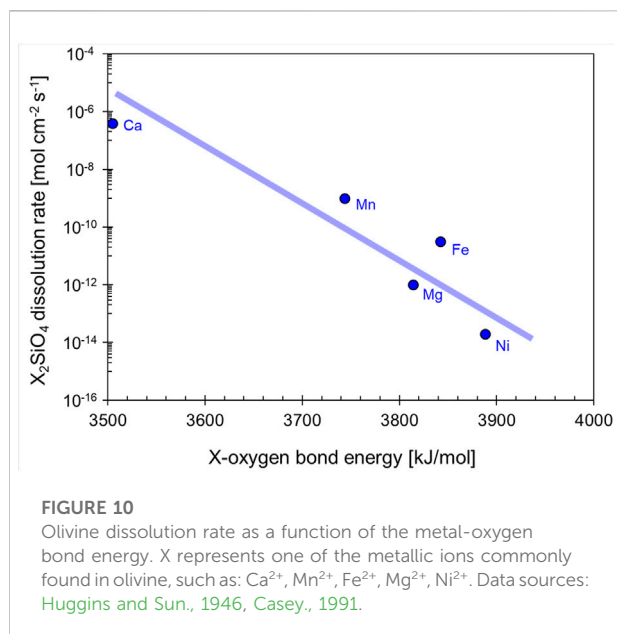
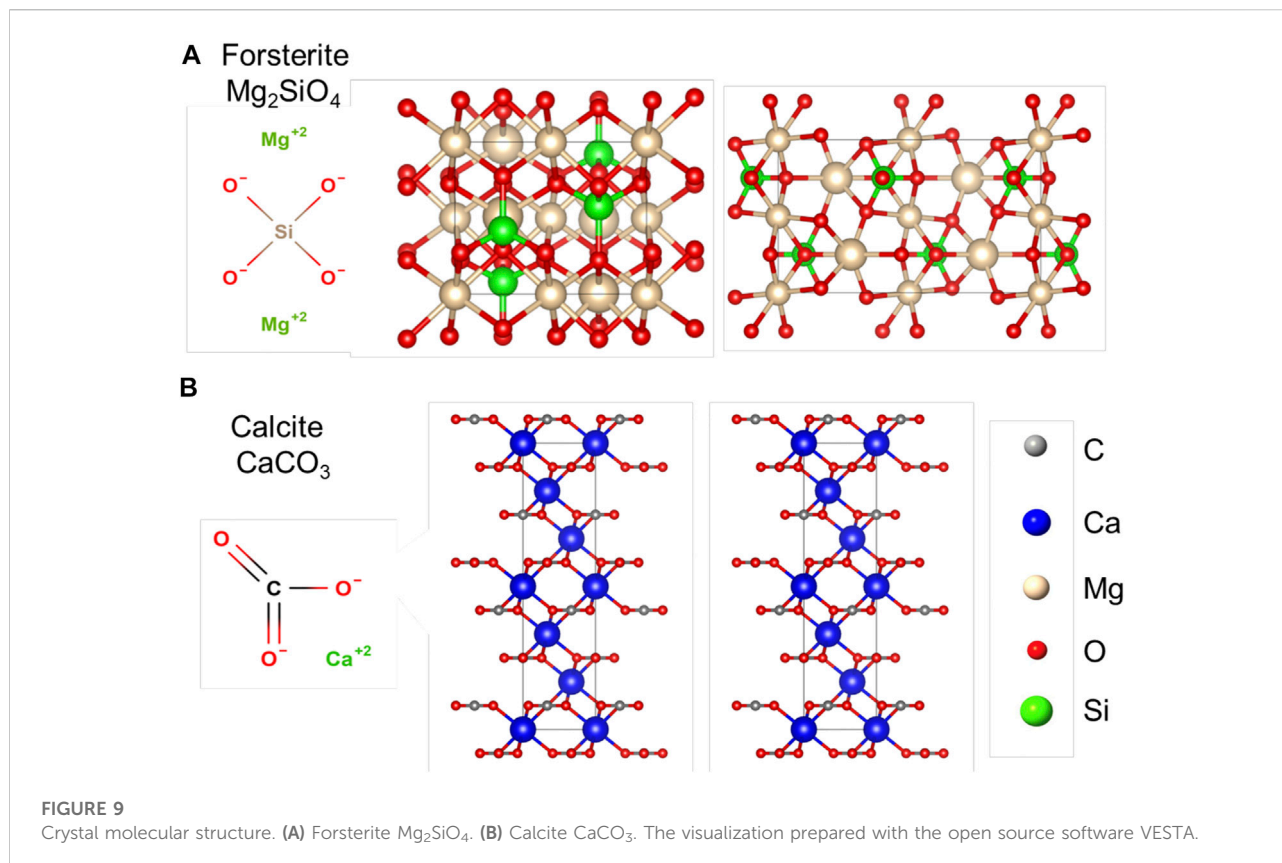
orientation may cause spatially heterogeneous dissolution (Daval et al., 2013; Pollet-Villard et al., 2016; Oelkers et al., 2018).

### Crystal structure

Silicates and carbonates are the most common minerals in CO<sub>2</sub> geological reservoirs. Others, minerals include oxides, sulfides, sulfate and halide.

Silicates are abundant and chemically reactive with carbonated water, thus, they have become the most promising mineral class for carbon mineralization (O'Connor et al., 2002). The fundamental unit in silicate structures is the silicon-oxygen tetrahedron (SiO<sub>4</sub>)<sup>4-</sup> (see Figure 9A). Metal ions are present in all silicates except quartz. Other cations with similar dimensions such as Al<sup>3+</sup>, Mg<sup>2+</sup>, Fe<sup>2+</sup>, Fe<sup>3+</sup>, Mn<sup>2+</sup> and Ti<sup>4+</sup> can replace Si<sup>4+</sup> from tetrahedral groups by isomorphic substitution (Brigatti et al., 2006).

Carbonates are common in most formations currently considered for CO<sub>2</sub> geological storage. The carbonate ion (CO<sub>3</sub><sup>2-</sup>) is the basic structural unit of carbonate minerals (see Figure 9B). This anion forms ionic bonds with cations such as Ca<sup>2+</sup>, Mg<sup>2+</sup>, Fe<sup>2+</sup>, Mn<sup>2+</sup> as well as rare-earth elements e.g., Ba<sup>2+</sup>, Ce<sup>3+</sup>. Calcite CaCO<sub>3</sub>, magnesite MgCO<sub>3</sub>, rhodochrosite MnCO<sub>3</sub>, siderite FeCO<sub>3</sub>, and dolomite CaMg(CO<sub>3</sub>)<sub>2</sub> are some of the most common carbonate minerals in nature (Effenberger et al., 1981).



## Chemical bonds in minerals

Oxygen and metals Ca, Mg, Fe, Mn, Ni become oppositely charged by electron transfer and form ionic bonds, whereas

silicon Si, carbon C, sulfur S and oxygen O bond together by sharing electron pairs in covalent bonds. Covalent bonds have higher bond energy than ionic bonds, for example, the bond energy in the covalent bond formed by silicon-oxygen is  $\sim 13,100$  kJ/mol, and the bond energy in the ionic bond formed by magnesium-oxygen is  $\sim 3,816$  kJ/mol (Huggins and Sun., 1946). Therefore, metal-oxygen bond breakage prevails in mineral dissolution and controls the dissolution rate (Figure 10). In addition, amorphous minerals dissolve more readily than their crystalline counterparts, such as basaltic glass compared with crystalline basalt.

## Mineral dissolution

In the presence of water, a mineral liberates ions (reactants) that move away from the mineral surface to the bulk solution mainly due to electrical interaction within the Helmholtz layer and random vibration. Ions experience the greatest potential change across the Helmholtz layer, and eventually form activated complexes and new reaction products in the bulk solution (Hayes and Katz., 1996; Crundwell., 2014). Common mineral and glass dissolution reactions in  $CO_2$  acidified water are summarized in Table 1 (see also Espinoza et al., 2011; Kim and Santamarina., 2014). Figure 11 shows the dissolution of forsterite in  $CO_2$  acidified water. Dissolution involves the breaking of magnesium-oxygen bonds and the release of  $Mg^{2+}$  and  $SiO_4^{4-}$ ;

TABLE 1 Common mineral dissolution reactions in CO<sub>2</sub> acidified water.

Mineral		Reaction	
Silicates	Forsterite	$Mg_2SiO_4 + 4H^+ \rightarrow 2Mg^{2+} + SiO_2 + 2H_2O$	
	Olivine	$(Mg_{0.80}Fe_{0.20})_2SiO_4 + 4H^+ \rightarrow 1.60Mg^{2+} + 0.40Fe^{2+} + SiO_2 + 2H_2O$ $(Mg_{0.43}Fe_{0.57})_2SiO_4 + 4H^+ \rightarrow 0.86Mg^{2+} + 1.14Fe^{2+} + SiO_2 + 2H_2O$	
	Fayalite	$Fe_2SiO_4 + 4H^+ \rightarrow 2Fe^{2+} + SiO_2 + 2H_2O$	
	Enstatite	$MgSiO_3 + 2H^+ \rightarrow Mg^{2+} + SiO_2 + H_2O$	
	Orthopyroxene	$Mg_{0.38}Fe_{0.62}SiO_3 + 2H^+ \rightarrow 0.38Mg^{2+} + 0.62Fe^{2+} + SiO_2 + H_2O$	
	Ferrosilite	$FeSiO_3 + 2H^+ \rightarrow Fe^{2+} + SiO_2 + H_2O$	
	Anorthite	$CaAl_2Si_2O_8 + 8H^+ \rightarrow Ca^{2+} + 2Al^{3+} + 2SiO_2 + 4H_2O$	
	Ablite	$NaAlSi_3O_8 + 4H^+ \rightarrow Al^{3+} + Na^+ + 3SiO_2 + 2H_2O$	
	Wollastonite	$CaSiO_3 + 2H^+ \rightarrow Ca^{2+} + SiO_2 + H_2O$	
	Clinocllore	$Mg_5Al_2Si_3O_{10}(OH)_8 + 16H^+ \rightarrow 5Mg^{2+} + 2Al^{3+} + 3SiO_2 + 12H_2O$	
	Daphnite	$Fe_5Al_2Si_3O_{10}(OH)_8 + 16H^+ \rightarrow 5Fe^{2+} + 2Al^{3+} + 3SiO_2 + 12H_2O$	
	Epidote	$Ca_2FeAl_2Si_3O_{12}(OH) + 13H^+ \rightarrow 2Ca^{2+} + Fe^{2+} + 2Al^{3+} + 3SiO_2 + 7H_2O$	
	Ferroactinolite	$Ca_2Fe_5Si_8O_{22}(OH)_2 + 14H^+ \rightarrow 2Ca^{2+} + 5Fe^{2+} + 8SiO_2 + 8H_2O$	
	Prehnite	$Ca_2Al_2Si_3O_{10}(OH)_2 + 10H^+ \rightarrow Ca^{2+} + 2Al^{3+} + 3SiO_2 + 6H_2O$	
	Tremolite	$Ca_2Mg_5Si_8O_{22}(OH)_2 + 14H^+ \rightarrow 2Ca^{2+} + 5Mg^{2+} + 8SiO_2 + 8H_2O$	
	Wairakite	$CaAl_2Si_4O_{10}(OH)_4 + 8H^+ \rightarrow Ca^{2+} + 2Al^{3+} + 4SiO_2 + 6H_2O$	
	Forsterite	$Mg_2SiO_4 + 4H^+ \rightarrow 2Mg^{2+} + SiO_2 + 2H_2O$	
	Diopside	$MgCaSi_2O_6 + 4H^+ \rightarrow Mg^{2+} + Ca^{2+} + 2SiO_2 + 2H_2O$	
	Carbonates	Calcite	$CaCO_3 + H^+ \rightarrow Ca^{2+} + HCO_3^-$
		Magnesite	$MgCO_3 + H^+ \rightarrow Mg^{2+} + HCO_3^-$
Siderite		$FeCO_3 + H^+ \rightarrow Fe^{2+} + HCO_3^-$	
Dolomite		$CaMg(CO_3)_2 + 2H^+ \rightarrow Ca^{2+} + Mg^{2+} + 2HCO_3^-$	
Oxides	Hematite	$Fe_2O_3 + 6H^+ \rightarrow 2Fe^{3+} + 3H_2O$	
	Magnetite	$Fe_3O_4 + 8H^+ \rightarrow 2Fe^{3+} + Fe^{2+} + 4H_2O$	
	Corundum	$Al_2O_3 + 6H^+ \rightarrow 2Al^{3+} + 3H_2O$	
Sulfide	Pyrite	$FeS_2 + H_2O \rightarrow Fe^{2+} + 0.25SO_4^{2-} + 0.25H^+ + 1.75HS^-$	
	Pyrrhotite	$FeS + H^+ \rightarrow Fe^{2+} + HS^-$	
Sulfate	Gypsum	$CaSO_4 \cdot 2H_2O \rightarrow Ca^{2+} + SO_4^{2-} + 2H_2O$	
Halide	Halite	$NaCl \rightarrow Na^+ + Cl^-$	
Basalt glass (Gysi and Stefánsson, 2012)		$K_{0.008}Na_{0.08}Ca_{0.27}Mg_{0.26}Mn_{0.003}S_{0.002}Ti_{0.02}Fe_{0.169}^{2+}Fe_{0.012}^{3+}Al_{0.35}Si_{1.00}O_{3.327} + 1.128H^+ + 2.201H_2O = 0.008K^+ + 0.08Na^+ + 0.27Ca^{2+} + 0.26Mg^{2+} + 0.003Mn^{2+} + 0.002HS^- + 0.02Ti(OH)^4 + 0.169Fe^{2+} + 0.012Fe(OH)^4 + 0.35Al(OH)^4 + H^4SiO^4$	

subsequently, Mg<sup>2+</sup> ions form hydrated ions with water, while SiO<sub>4</sub><sup>4-</sup> ions form activated complexes with protons, HSiO<sub>4</sub><sup>3-</sup>.

### Kinetics of mineral dissolution

The rate of dissolution far-from-equilibrium  $r_d$  [mol/s] can be estimated as (Arthur et al., 2000):

$$r_d = k \cdot s \cdot a_{H^+}^{n_{H^+}} \quad (14)$$

where  $k$  [mol m<sup>-2</sup> s<sup>-1</sup>] is the rate constant for mineral dissolution,  $s$  [m<sup>2</sup>] is the surface area of minerals in a reference volume of solution, and  $n_{H^+}$  [ ] is the reaction order with respect to the activity of protons  $a_{H^+}$  [ ]. The mineral surface area varies during dissolution (Grandstaff, 1978; Gautier et al., 2001; Fischer et al., 2012).



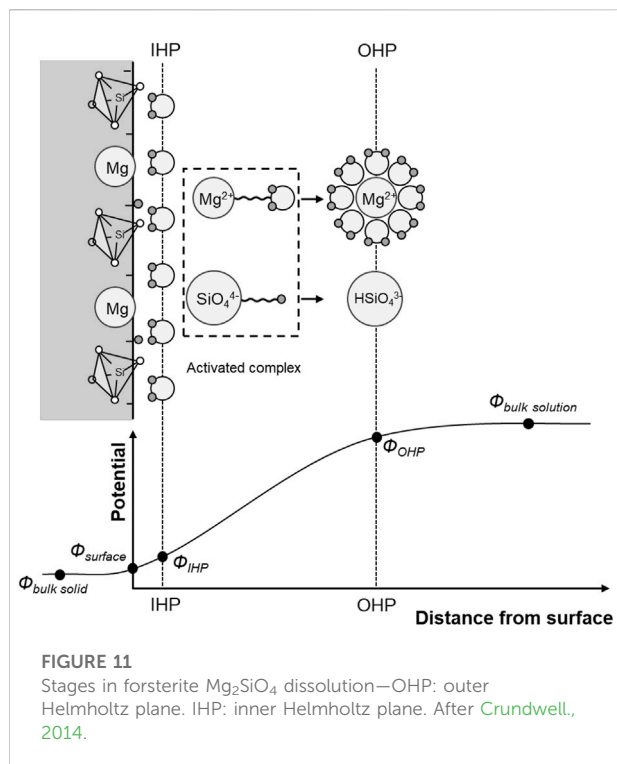


Table 2 lists kinetic dissolution rates calculated using Eq. 14 and kinetic parameters compiled from the literature for common silicate and carbonate minerals in acid solutions. Clearly, carbonates are more reactive than other minerals, in agreement with their bond energies reported above. Anorthite, fayalite and forsterite have higher reactivity than other silicates due to the relative abundance of ionic metal-oxygen bonds and the absence of silicon-oxygen covalent bonds in these minerals ([Angel, 1988](#); [Velbel, 1999](#)). Other parameters affect mineral dissolution rates as well, for example: temperature, pH,  $CO_2$  partial pressure ([Stumm and Morgan, 1996](#)), salinity ([Pokrovsky et al., 2005](#)), organic ligands ([Fredd and Fogler, 1998](#)), and catalysts (e.g., sodium bicarbonate—[O'Connor et al., 2001](#)).

## Mineral dissolution: Hydro-chemo-mechanical coupling

Minerals dissolve during  $CO_2$  injection and produce species that are carried away by advection and diffusion. Consequently, mineral dissolution enhances porosity and permeability, and causes settlement and stress changes.

The evolution of mineral dissolution in a porous medium is inherently related to pore size variability and the presence of preferential flow pathways which deliver the majority of the reactants in advective regimes: 10% of pores may be responsible for more than 50% of the total flow rate in sediments ([Jang et al., 2011](#)),

and flow channeling is even more pronounced in fractured rock masses ([National Academies of Sciences Engineering and Medicine, 2020](#); [Cardona et al., 2021](#)). Therefore, dissolution and transport are coupled, and may involve fabric changes. Hydro-chemo-mechanical coupling effects triggered by dissolution are explored next.

## Dissolution at constant fabric (no mechanical coupling)

Dissolution and advection lead to increased flow focusing and “wormhole” formation ([Kim and Santamarina, 2015](#); [Kim and Santamarina, 2016](#); [Menke et al., 2017](#); [Derr et al., 2020](#)). Two dimensionless ratios combine the three time scales for advection  $t_{adv}$ , diffusion  $t_{diff}$ , and reaction  $t_r$ ; for a channel length  $L_{ch}$

$$\text{Damköhler number } Da = \frac{t_{adv}}{t_r} = \frac{kL_{ch}}{\nu} \quad (15)$$

$$\text{Péclet number } Pe = \frac{t_{diff}}{t_{adv}} = \frac{\nu L_{ch}}{D} \quad (16)$$

where  $\nu$  [m/s] is the average flow velocity in pores,  $D$  [ $m^2/s$ ] is the molecular diffusion coefficient, and the kinetic rate  $\kappa$  [ $1/s$ ] =  $kS_sM_m$  is a function of the rate constant for mineral dissolution  $k$  [ $mol\ m^{-2}\ s^{-1}$ ], the mineral specific surface area  $S_s$  [ $m^2/g$ ], and the mineral molar mass  $M_m$  [ $g/mol$ ]. Various dissolution topologies become apparent in the Péclet-Damköhler  $Pe$ - $Da$  diagram in [Figure 12](#): compact dissolution, conical wormhole, dominant wormhole, ramified wormhole and uniform dissolution regimes. Low advection velocities, i.e., long  $t_{adv}$  and small  $Pe$ , allow reactant consumption near the inlet and instabilities cannot develop; conversely, very high advection velocities, i.e., short  $t_{adv}$ , high  $Pe$  and small  $Da$ , carry reactants all along the porous medium from inlet to outlet and favor a stable homogeneous dissolution. Wormholes develop at intermediate flow rates, i.e., large  $Da$  and  $Pe$ , so that reactants act preferentially along flow channels.

Similarly, various dissolution regimes emerge when reactive fluids flow through rock fractures ([Kim and Santamarina, 2015](#); [Shafabakhsh et al., 2021](#)). A low  $Da$  causes a more uniform aperture enlargement along the length of the fracture. On the other hand, most reactants are consumed near the inlet in large  $Da$  and low  $Pe$  regimes (highly reactive low-advection conditions). Finally, a large fracture aperture slows the diffusive transport of reactants to the mineral fracture surface, and higher reactant concentrations remain along the centerline.

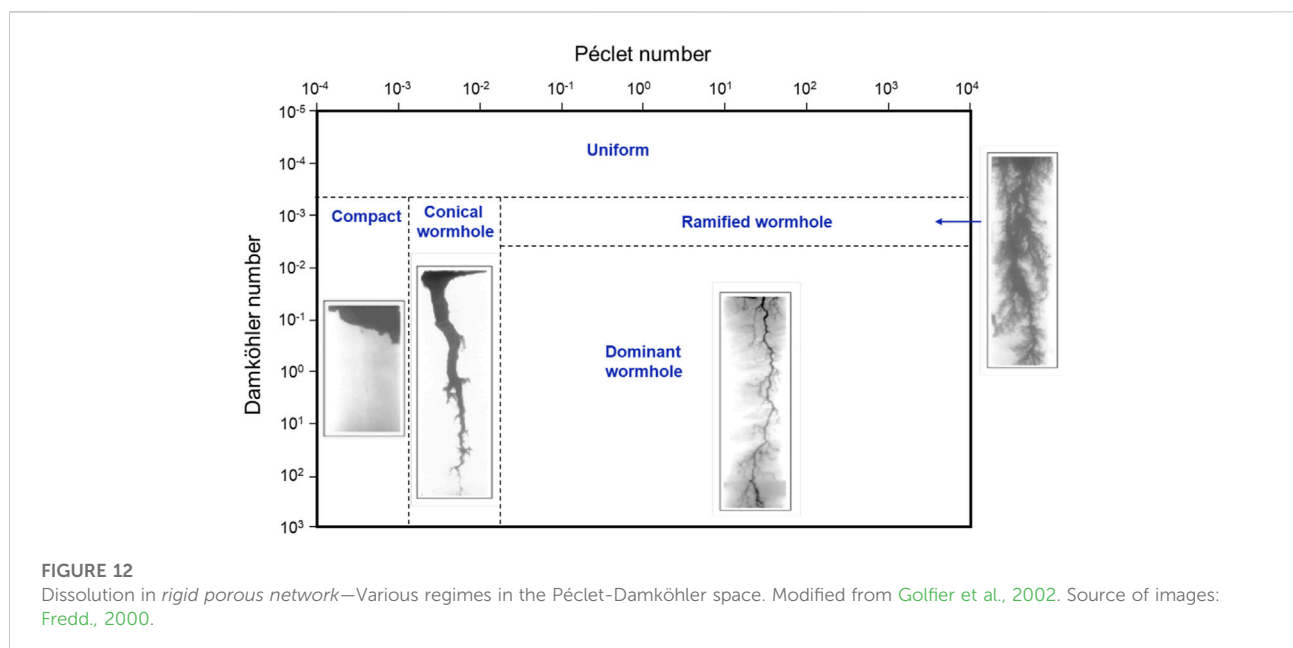
## Dissolution-triggered fabric changes (mechanically coupled effects)

Mineral dissolution during  $CO_2$  injection can cause fabric changes in sediments, block sliding in fractured rocks and stress changes in both cases.

The coefficient of lateral earth pressure at rest  $k_0 = \sigma_h'/\sigma_v'$  [ $\varepsilon_h = 0$ ] is the stress ratio between the effective horizontal stress  $\sigma_h'$  and the effective vertical stress  $\sigma_v'$  under zero lateral strain conditions  $\varepsilon_h = 0$ . Both experimental and numerical results show that even a

TABLE 2 Kinetic rate parameters for mineral dissolution in acid solutions ( $T = 25^{\circ}\text{C}$ ) and the calculated dissolution rates  $r$  (Eq. 14—Calculations assume  $25^{\circ}\text{C}$ , 0.1 MPa, the pH of  $\text{CO}_2$  water is 4.5, the surface area is fixed at  $2\text{ m}^2$ ).

Minerals		Log $k$ ( $\text{mol m}^{-2}\text{ s}^{-1}$ )	$n$	pH	References	$r$ ( $\text{mol s}^{-1}$ )
Silicates	$\text{Mg}_2\text{SiO}_4$ (Forsterite)	-6.85	0.470	0	Palandri and Kharaka., 2004	$1.09 \times 10^{-4}$
	$\text{Fe}_2\text{SiO}_4$ (Fayalite)	-4.8	1.0	acid	Sverdrup., 1990	$1.00 \times 10^1$
	$\text{MgSiO}_3$ (Enstatite)	$-9.3 \pm 0.7$	$-0.25 \pm 0.14$	2–7	Furrer et al., 1993	$4.23 \times 10^{-11}$
	$\text{Mg}_{0.38}\text{Fe}_{0.62}\text{SiO}_3$ (Orthopyroxene)	$-9.8 \pm 0.3$	$-0.49 \pm 0.13$	<5	Schott and Berner., 1983	$6.40 \times 10^{-13}$
	$\text{FeSiO}_3$ (Ferrosilite)	$-9.3 \pm 0.3$	$-0.33 \pm 0.06$	2.1–7.1	Banfield et al., 1995	$1.53 \times 10^{-11}$
	$\text{CaAl}_2\text{Si}_2\text{O}_8$ (Anorthite)	-5.87	1.12	$\leq 5$	Brady and Walther., 1989	$3.90 \times 10^0$
	$\text{NaAlSi}_3\text{O}_8$ (Albite)	-9.5	0.5	acid	Sverdrup., 1990	$7.96 \times 10^{-9}$
	$\text{CaSiO}_3$ (Wollastonite)	$-8.0 \pm 0.2$	$-0.24 \pm 0.04$	$\leq 7.2$	Xie., 1994	$9.57 \times 10^{-10}$
	$\text{Ca}_2\text{FeAl}_2\text{Si}_3\text{O}_{12}(\text{OH})$ (Epidote)	-10.60	0.338	0	Palandri and Kharaka., 2004	$3.63 \times 10^{-9}$
	$\text{Ca}_2\text{Fe}_3\text{Si}_8\text{O}_{22}(\text{OH})_2$ (Ferroactinolite)	$-8.9 \pm 0.2$	$-0.75 \pm 0.06$	3.6–4	Zhang., 1990	$1.89 \times 10^{-13}$
	$\text{Ca}_2\text{Al}_2\text{Si}_3\text{O}_{10}(\text{OH})_2$ (Prehnite)	-10.66	0.256	0	Palandri and Kharaka., 2004	$1.12 \times 10^{-9}$
	$\text{Ca}_2\text{Mg}_5\text{Si}_8\text{O}_{22}(\text{OH})_2$ (Tremolite)	-11.5	0.11	1–6	Schott et al., 1981	$2.55 \times 10^{-11}$
	$\text{Mg}_2\text{SiO}_4$ (Forsterite)	-10.85	-0.47	0	Palandri and Kharaka., 2004	$7.35 \times 10^{-14}$
	$\text{MgCaSi}_2\text{O}_6$ (Diopside)	$-9.4 \pm 0.4$	$-0.22 \pm 0.1$	2–6	Knauss et al., 1993	$4.91 \times 10^{-11}$
	Carbonates	$\text{CaCO}_3$ (Calcite)	-0.30	1.000	0	Palandri and Kharaka., 2004
$\text{MgCO}_3$ (Magnesite)		-6.38	0.589	0	Palandri and Kharaka., 2004	$1.45 \times 10^{-3}$
$\text{FeCO}_3$ (Siderite)		-3.19	0.500	0	Palandri and Kharaka., 2004	$7.26 \times 10^{-1}$
$\text{CaMg}(\text{CO}_3)_2$ (Dolomite)		-3.76	0.500	0	Palandri and Kharaka., 2004	$1.95 \times 10^{-1}$
Oxides	$\text{Fe}_2\text{O}_3$ (Hematite)	-9.39	1.000	0	Palandri and Kharaka., 2004	$2.58 \times 10^{-4}$
	$\text{Fe}_3\text{O}_4$ (Magnetite)	-8.59	0.279	0	Palandri and Kharaka., 2004	$1.76 \times 10^{-7}$
Sulfide	$\text{FeS}_2$ (Pyrite)	-7.52	-0.500	0	Palandri and Kharaka., 2004	$1.07 \times 10^{-10}$
	$\text{FeS}$ (Pyrrhotite)	-8.04	-0.597	0	Palandri and Kharaka., 2004	$9.50 \times 10^{-12}$



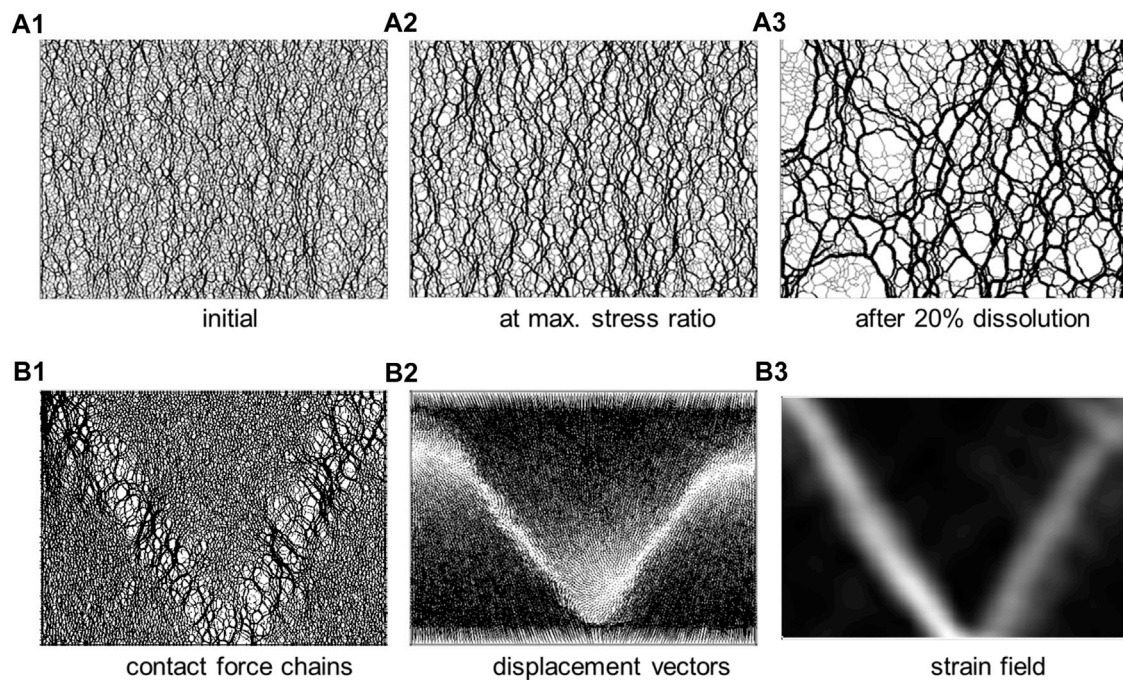


FIGURE 13

Dissolution in *granular materials* under constant vertical stress and zero lateral strain. (A) Force chain evolution during the dissolution of randomly distributed soluble grains: the evolving brittle “honeycomb” structure exhibits pronounced force chain arching (20% of the grains are soluble—Shin et al., 2008). (B) Pressure-solution whereby the particle dissolution rate is proportional to the normal force: the emergence of shear bands, Cha and Santamarina, 2019.

small amount of particle dissolution can produce a pronounced drop in  $k_0$  and the state of stress may even reach the Coulomb failure envelope, i.e., Rankine active coefficient  $k_0 \rightarrow k_a$  (mixture of soluble and insoluble grains—Shin and Santamarina, 2009). At the grain scale, forces initially carried by dissolving particles get transferred to less soluble neighboring grains, the particle coordination number decreases, and force chains form a brittle honeycomb-like fabric (Figure 13A—Shin et al., 2008; Cha and Santamarina, 2014; Kim et al., 2019). After reaching its lowest value, the stress ratio may gradually recover upon further dissolution to reach a “terminal state”.

Stress concentration at grain contacts favors localized dissolution, i.e., “pressure solution” (even in single-mineral systems—Rutter, 1983; Etheridge et al., 1984; Fowler and Yang, 1999). Strong intergranular force chains and force-dependent contact dissolution combine to create a positive feedback mechanism that can lead to shear localization (Cha and Santamarina, 2016—Figure 13B). This sequence of events explains the nontectonic origin of shear discontinuities and polygonal fault systems observed in marine sediments and lacustrine deposits (Cartwright et al., 2003; Shin et al., 2008; Shin et al., 2010).

Let’s assume a carbonate reservoir (mineralogy = 100% calcite, porosity = 0.3) at a depth of 800 m ( $p = 8$  MPa,  $T =$

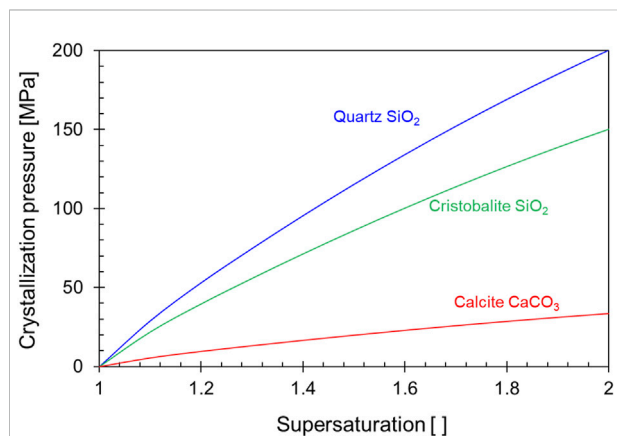
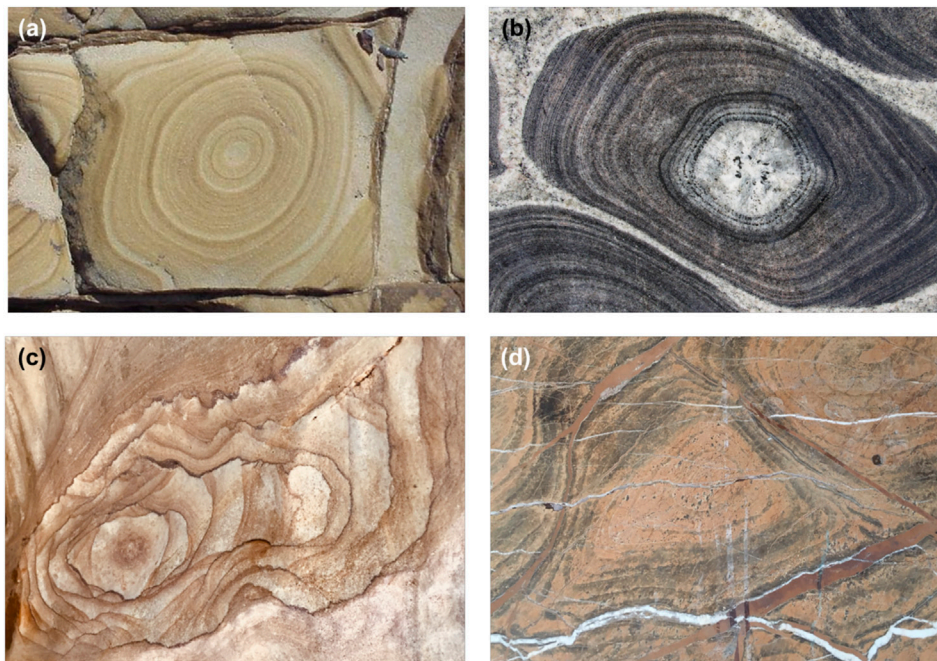


FIGURE 14

Crystallization pressures of quartz, cristobalite, and calcite as a function of supersaturation Trends from Wiltshko and Morse, 2001; formulation in Maliva and Siever, 1988.

60 °C) filled with CO<sub>2</sub> saturated water (pH ≈ 3.21, calcite solubility ≈ 4.41 × 10<sup>-5</sup> mol/L). Under these conditions, 1.3 g ~ 0.5 cm<sup>3</sup> of calcite will dissolve for every pore volume of CO<sub>2</sub> saturated water, and cause a volumetric strain of  $\epsilon = 0.5 \times 10^{-6}$ . The induced isotropic stress change is  $\Delta\sigma = E \times \epsilon \approx 1$  kPa per



**FIGURE 15**

Diffusive-reactive transport: Liesegang bands. (A) in sandstone - from T McCarron. (B) Orbicular granite—Lahti et al., 2005. (C) Sandstone—Wadi Lajab, Saudi Arabia. (D) Jordanian carbonate.

pore volume of fluid exchange (the Young's modulus is  $E \approx 5$  GPa for a carbonate with porosity 0.3—Ng and Santamarina 2022). Thus, the stress change for a single pore volume is small; however, ensuing convective transport and dissolution localization may eventually trigger mechanical instabilities.

## Precipitation: Crystallization pressure and matrix filling

Precipitation is a self-homogenizing process in porous media: precipitates reduce the permeability of preferential flow paths and reactive fluids deviate to other pathways. Still, precipitation has other important consequences related to crystallization pressure and ensuing pore habits. These processes play a central role in CO<sub>2</sub> mineralization.

### Crystallization pressure: new fractures

New minerals nucleate on preexisting surfaces (heterogeneous nucleation), grow into the pore space, and exert crystallization pressure when growth is constrained by pore and fracture walls (Lee and Kurtis, 2017). A thin film between the crystal and the wall and corner flow (along roughness and crevices which are not occupied by the growing mineral) maintain connectivity with the saturated

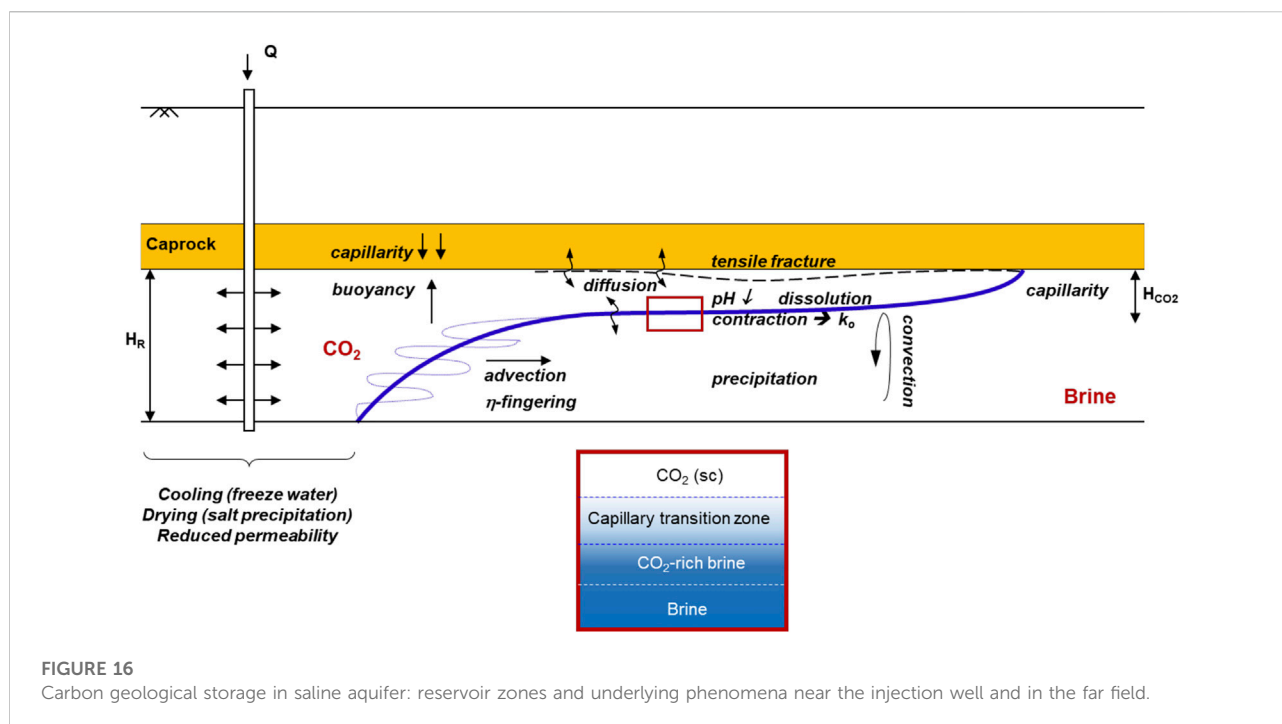
solution and sustain crystal growth (Becker and Day, 1916; Taber, 1916; Correns, 1949).

The crystal experiences an internal excess pressure  $\Delta P = P_{crys} - P_{sob}$  where  $P_{crys}$  is the pressure in the growing crystal and  $P_{sol}$  is the hydrostatic pressure of the pore solution. This excess pressure  $\Delta P$  is a function of the solution supersaturation  $S = c/c_0$ , where  $c$  and  $c_0$  are the solute concentrations in the supersaturated solution and in the saturated solution respectively:

$$\Delta P = P_{crys} - P_{sol} = \frac{RT}{V_m} \ln S \quad (17)$$

Other parameters include: the gas constant  $R$  [J mol<sup>-1</sup> K<sup>-1</sup>], the absolute temperature  $T$  [K], and the molar volume of the precipitated solid phase  $V_m$  [m<sup>3</sup>/mol]. Figure 14 shows the calculated crystallization pressure of quartz SiO<sub>2</sub>, cristobalite SiO<sub>2</sub> and calcite CaCO<sub>3</sub> for different degrees of supersaturation (Wiltchko and Morse, 2001).

The stress field induced by the crystallization pressure in neighboring pores may trigger an opening mode fracture (Anderson, 2017). Similarly, precipitation may form grain-displacive nodules and lenses in sediments when the crystallization pressure exceeds the *in situ* effective stress (Shin and Santamarina, 2011; Liu and Santamarina, 2022). Note that precipitated carbonates usually occupy a larger volume than the dissolved source minerals (Snæbjörnsdóttir et al., 2018; Snæbjörnsdóttir et al., 2020), therefore volume



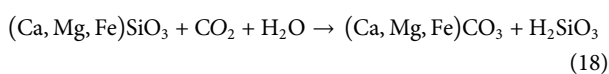
expansion during local dissolution-precipitation will also cause crystallization pressure.

### Recurrent precipitation in the matrix: Liesegang band formation

Crystal growth decreases the local reactant concentration and may locally deplete the reactants that the host medium may contribute. Meanwhile, the reaction front keeps advancing; eventually transported and host rock reactants reach the threshold saturation concentration again and a new round of mineral precipitation starts. Such reactive-diffusion-precipitation process produces a periodic precipitation pattern, known as Liesegang bands (Liesegang, 1896; Rácz, 1999; Panjarian and Sultan, 2001; Nabika et al., 2020). Banded precipitation is found in a wide range of geological settings; Figure 15 shows periodic precipitation patterns in sandstones, carbonate and granite. These field observations are natural analogues for CO<sub>2</sub> mineralization within the rock matrix.

### CO<sub>2</sub> mineralization

The general reaction for CO<sub>2</sub> mineralization starts with calcium, magnesium or iron silicates to produce the corresponding carbonates:



Mafic and ultramafic rocks are silicates; therefore, the ultimate storage ability per cubic meter of rock is

determined by the molar mass ratio between CO<sub>2</sub> and silicate  $M_{\text{CO}_2}/M_{\text{silic}}$ , and the rock density and porosity  $(M_{\text{CO}_2}/M_{\text{silic}}) \cdot (1-\phi) \rho_{\text{rock}}$ . The molar mass ratios  $M_{\text{CO}_2}/M_{\text{silic}}$  are: 0.38 for wollastonite CaSiO<sub>3</sub>, 0.44 for enstatite MgSiO<sub>3</sub> and 0.33 FeSiO<sub>3</sub> for ferrosilite. Then, the theoretical ultimate storage density values can exceed 500 kg/m<sup>3</sup>. However, the reaction is time-dependent and more modest values between 30-and-200 kg CO<sub>2</sub> per cubic meter of basalt could be expected within the first 10-to-100 years after injection (e.g., Xiong et al., 2018).

In fact, the reaction rate depends on the mineral reactivity, the available exposed surface area and fracture spacing, and the evolving efficiency of transport processes. Eq. 14 and kinetic parameters in Table 2 show that the characteristic time for carbonation varies in orders of magnitude for different minerals. For example, consider a rock block with a constant exposed surface of 2 m<sup>2</sup> in contact with CO<sub>2</sub> saturated water ( $p = 0.1$  MPa,  $T = 25^\circ\text{C}$ ,  $\text{pH} = 4.5$ ): it takes around 7 years to dissolve 1 m<sup>3</sup> forsterite Mg<sub>2</sub>SiO<sub>4</sub> ( $\rho = 3,270$  kg/m<sup>3</sup>) but 40,000 years for albite NaAlSi<sub>3</sub>O<sub>8</sub> ( $\rho = 2,620$  kg/m<sup>3</sup>). These times are for reaction-limited processes; advection or diffusion-limited transport may eventually determine the rate of mineralization under reservoir conditions.

### Other phenomena—reservoirs

The geological storage of CO<sub>2</sub> involves other reservoir-dependent processes, such as zone differentiation and convective currents in saline aquifers, excess pressure due to

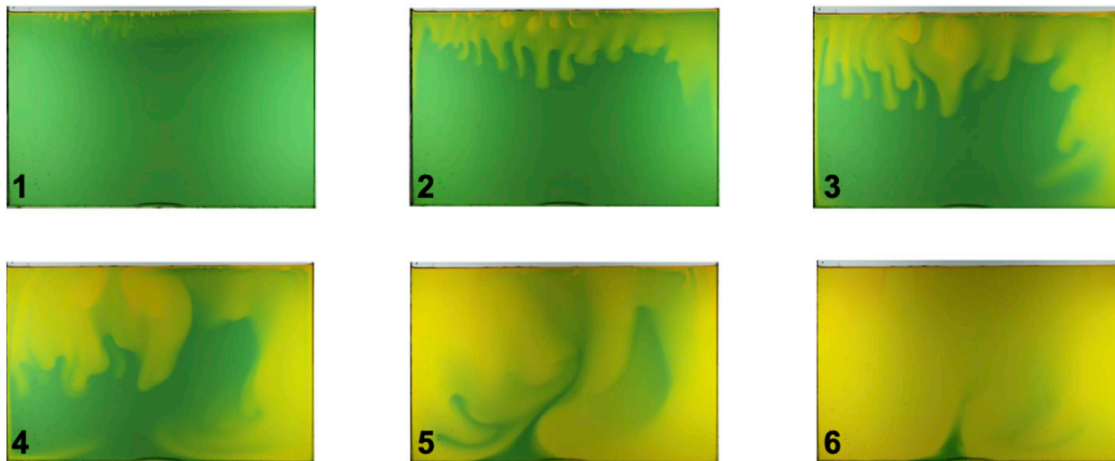


FIGURE 17

CO<sub>2</sub> solute driven convection. Time lapse images of a water bath beneath a CO<sub>2</sub> atmosphere (water with a universal pH indicator—Bang et al., 2012). The time interval between images is 40 min.

the buoyant CO<sub>2</sub> plume and altered seal performance, oil mobilization in depleted reservoirs, and near-well drying and cooling. These phenomena are investigated next.

## CO<sub>2</sub> injection into saline aquifers

Transport during CO<sub>2</sub> injection into saline aquifers involves advection, dispersion, dissolution, diffusion, and convection (Figure 16). CO<sub>2</sub> is lighter and less viscous than water, therefore, CO<sub>2</sub> injection responds to buoyant, capillary and viscous forces: 1) CO<sub>2</sub> invasion next to the well is most sensitive to stratification and invades by viscous fingering through the most pervious layers, 2) the plume becomes gravity-dominant away from the well (Hesse et al., 2008), and eventually 3) it rests against the top seal by buoyancy and capillary trapping. Then, the reservoir develops various zones according to the saturating fluids: the supercritical CO<sub>2</sub> plume, the capillary transition zone, CO<sub>2</sub>-rich brine, and the far-field saturated with the initial brine (Figure 16).

During early storage times, the degree of water saturation in the reservoir is a function of the capillary pressure  $P_c = (\gamma_w - \gamma_{CO_2})H_{CO_2}$ , where  $H_{CO_2}$  is the height of the CO<sub>2</sub> column, as predicted by the reservoir capillarity-saturation response  $S_w = f(\Delta u)$  often captured with Brooks-Corey or van Genuchten models. The degree of water saturation decreases with time as water dissolves in CO<sub>2</sub> and is transported away by diffusion.

Next to the plume lower boundary, stored CO<sub>2</sub> dissolves into the groundwater and convection emerges due to the increased density of water with dissolved CO<sub>2</sub> - between

$\Delta\rho/\rho_0 = 0.1$  and 1% depending on pressure and temperature (Emami-Meybodi et al., 2015. Note: water acidified with carbonic acid dissolves minerals and experiences an even larger increases in density—see previous sections). The Rayleigh number  $Ra$  compares the rate of fluid convection with the rate of diffusive transport:

$$Ra = \frac{k_{hy} g H_R \Delta\rho_0}{\mu_w D} \quad (19)$$

where the governing parameters are the formation permeability  $k_{hy}$  [m<sup>2</sup>],  $g = 9.8$  m/s<sup>2</sup>, the height of the fluid column  $H_R$  [m], the change in fluid density  $\Delta\rho$  [kg/m<sup>3</sup>], the viscosity of water  $\mu_w$  [Pa·s], and the diffusion coefficient of CO<sub>2</sub> in water  $D$  [m<sup>2</sup>/s]. Convection typically occurs when  $Ra > 4\pi^2$  (Kneafsey and Pruess, 2010). Figure 17 shows time-lapse images of emerging convective currents when CO<sub>2</sub> gas rests above an aquifer at atmospheric pressure; the pH indicator changes in color from green to yellow at pH = 5. Small fingers initiate at the interface and begin to grow. Then there is an increase in the finger wavelength and a decrease in the number of fingers.

The convection period is  $t_{cov} = \frac{\mu H_R}{k_{hyd} g \Delta\rho}$  (Jiang et al., 2019) then, for an aquifer of thickness  $H_R = 10$  m and permeability  $k_{hyd} = 20$  md, the convection period is  $t_{cov} = 150$  years assuming  $\Delta\rho = 10$  kg/m<sup>3</sup>.

Finally, the injection of dry CO<sub>2</sub> into a saline aquifer causes salt precipitation as water evaporates into the supercritical CO<sub>2</sub> near the wellbore (Miri et al., 2015). Salt precipitation can reduce CO<sub>2</sub> injectivity into the reservoir; its impact depends on the injection rate, mode (monotonic vs. cyclic), and brine salinity (Talman et al., 2020; He et al., 2022).

## Seal rock: structure trapping

Small pores in high specific surface clay-rich caprocks give rise to the high capillary entry pressure and low permeability needed to hinder the upwards migration of the buoyant CO<sub>2</sub> plume. On the other hand, CO<sub>2</sub> invasion either dissolved in water or as a supercritical phase alters the physiochemical equilibrium in the seal rock.

### Capillary pressure against the seal

The plume height  $h$  [m] determines the excess pressure in the CO<sub>2</sub> plume at the seal-reservoir interface,  $\Delta P = P_{CO_2} - P_w = (\rho_w - \rho_{CO_2})gh$ . On the other hand, the capillary breakthrough pressure  $P_c$  the seal may resist is a function of the interfacial tension  $\gamma$  [N/m] between water and CO<sub>2</sub> at reservoir pressure and temperature conditions, the contact angle  $\theta$  formed by the water-CO<sub>2</sub> interface on the mineral surface, and the characteristic pore diameter  $d^*$  [m] for the percolating path across the seal layer.

$$P_c = \frac{\gamma \cos \theta}{d^*} \quad (20)$$

The characteristic pore diameter  $d^*$  is a function of the specific surface area  $S_s$ , the void ratio  $e$ , the mineral density  $\rho_m$ , a fabric factor  $\psi$  and a  $\beta$  factor related to the pore size distribution along percolating paths:  $d^* = \beta\psi e / (S_s \rho_m)$  (Espinoza and Santamarina, 2010; Espinoza and Santamarina, 2017). Then, the sealing number  $\Pi_{seal}$  compares the capillary breakthrough pressure  $P_c$  to the buoyant pressure  $\Delta P$ :

$$\Pi_{seal} = \frac{P_c}{\Delta P} = \frac{\gamma \cos \theta S_s \rho_m}{\beta \psi e (\rho_w - \rho_{CO_2}) gh} \quad (21)$$

Sealing numbers  $\Pi_{seal} \gg 1$  are expected for safe storage. Clearly, a high sealing number is found for small plume thickness  $h$ , and high specific surface  $S_s$  clayey seals.

On the other hand, the higher CO<sub>2</sub> pressure changes the effective stress can trigger fault reactivation and grain-displacive breakthrough. The stability number  $\Pi_{stability}$  compares the initial horizontal effective stress  $\sigma'_{h0}$  at the reservoir depth  $z$  to the capillary pressure  $\Delta P$  created by a CO<sub>2</sub> plume of thickness  $h$ .

$$\Pi_{stability} = \frac{\sigma'_{h0}}{\Delta P} = \frac{(\sigma_{z0} - P_0)k_0}{hg(\rho_w - \rho_{CO_2})} \quad (22)$$

where the variables are: the initial total vertical stress  $\sigma_{z0}$  and fluid pressure  $P_0$  at the reservoir-seal interface before CO<sub>2</sub> injection, and effective stress ratio  $k_0 = \sigma'_h / \sigma'_z$  at zero lateral strain. Safe CO<sub>2</sub> injection requires reservoirs with a high stability number  $\Pi_{stability}$ . The organic matter in a shale is CO<sub>2</sub>-wet (Arif et al., 2017b), therefore, low organic content shales are preferred for CO<sub>2</sub> seal.

## CO<sub>2</sub> interaction with clay-rich caprock

As anticipated by the sealing number  $\Pi_{seal}$ , clayey shales are promising seal layers for CO<sub>2</sub> geological storage. Hence, we must address the hydro-chemo-mechanical coupling between CO<sub>2</sub> and clay minerals. We can anticipate two salient effects:

- 1) *First contact (before SC-CO<sub>2</sub> invasion)*: CO<sub>2</sub> dissolves in water and acidifies the brine. Low pH promotes protonation, and leads to positively charged surfaces, which will affect the clay fabric.
- 2) *Long term (after SC-CO<sub>2</sub> invasion)*: CO<sub>2</sub> replaces water and becomes the saturating pore fluid. There will be a reduction in osmotic repulsion, dissolution of residual water in CO<sub>2</sub> and salt precipitation, volumetric contraction driven by capillarity, and a three-fold increase in the Hamaker constant for clay-CO<sub>2</sub>-clay as compared to clay-water-clay (Shin and Santamarina, 2010; Espinoza and Santamarina, 2012).

Changes in electrical and capillary forces will induce volumetric changes, thus stress changes in the caprock. Figure 18 shows that a montmorillonite-water paste contracts and experiences desiccation cracks when water evaporates into supercritical CO<sub>2</sub> atmosphere. Sedimentation experiments show that clay particles show strong particle aggregation with heptane and liquid CO<sub>2</sub> which shows that the electrical interaction between clays and water is different than with CO<sub>2</sub> in its different states (Espinoza and Santamarina, 2012).

## CO<sub>2</sub> injection into depleted oil reservoirs

CO<sub>2</sub> injection into depleted oil reservoirs benefits from time-proven seals. In addition, CO<sub>2</sub> can mobilize residual oil; in this case, we can simultaneously sequester CO<sub>2</sub> while recovering additional oil. Numerous pore-scale mechanisms affect CO<sub>2</sub>-oil displacement processes, including oil swelling, ganglia destabilization, coalescence of disconnected oil blobs into a continuous oil phase, light-component extraction, miscible displacement and asphaltene deposition. Details follow.

### Oil density and viscosity

The oil density decreases slightly with temperature, and pressure has a second order effect (Figure 1). On the other hand, the viscosity of crude oils can decrease by an order of magnitude when the temperature increases from 40°C to 80°C (Figure 3).

### Oil-water interfacial tension

The interfacial tension between two non-polar fluids such as CO<sub>2</sub> and oil tends to be smaller than when dipolar water molecules are involved. Figure 19 shows the interfacial tension

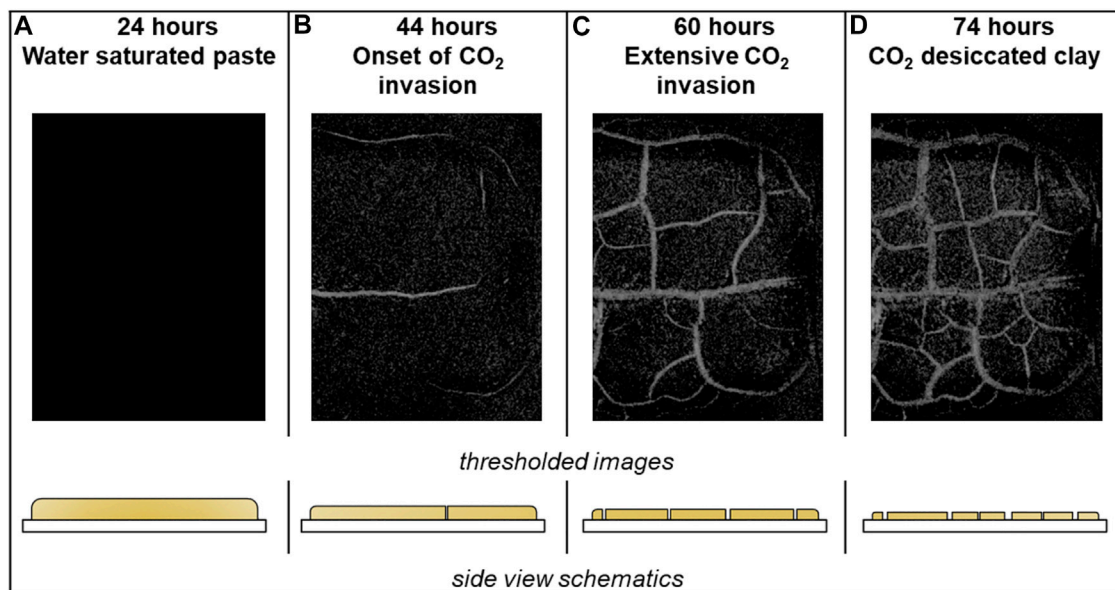


FIGURE 18

Montmorillonite-water paste subjected to a supercritical  $\text{CO}_2$  atmosphere ( $p = 15 \text{ MPa}$  and  $T = 38^\circ\text{C}$ ). The photographs and associated sketches show the evolution of desiccation and the formation of capillary-driven fractures: as water evaporates into the  $\text{CO}_2$  reservoir, suction increases and the water- $\text{CO}_2$  interface “compresses” the sediment; eventually, supercritical  $\text{CO}_2$  invades the sediment and triggers desiccation cracks, [Espinoza and Santamarina., 2012](#).

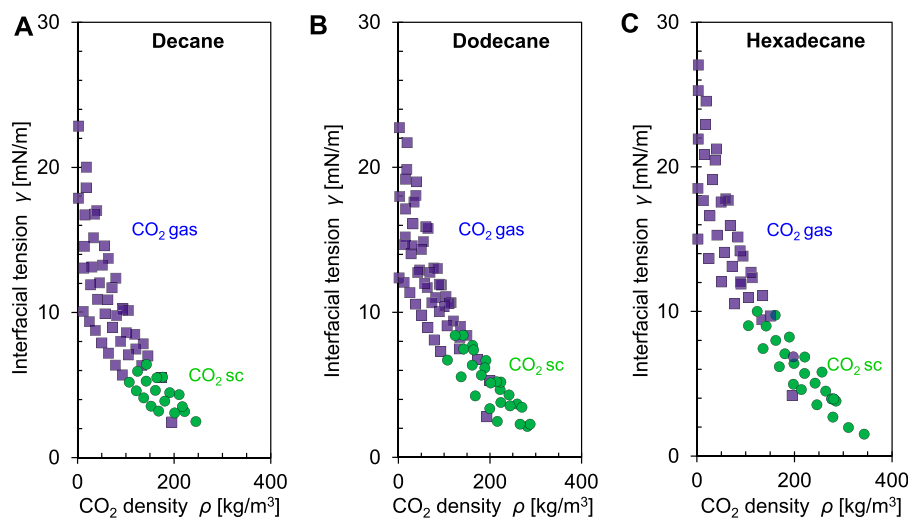


FIGURE 19

Interfacial tension between selected alkanes and  $\text{CO}_2$  as a function of the  $\text{CO}_2$  density at different pressures and temperatures: (A) Decane, (B) Dodecane and (C) Hexadecane. Symbols:  $\text{CO}_2$  gas ( $\blacksquare$ ) and  $\text{CO}_2$  supercritical ( $\bullet$ ). Data source: [Georgiadis et al., 2010](#).

between three different alkanes (decane, dodecane and hexadecane) and  $\text{CO}_2$  as a function of  $\text{CO}_2$  density. The interfacial tension increases with the number of carbon atoms in alkanes and decreases with increasing  $\text{CO}_2$  density.

### $\text{CO}_2$ solubility in crude oil - Viscosity

The solubility of  $\text{CO}_2$  in crude oil is much higher than in water because of their shared non-polar nature ([Schmid, 2001](#)). Solubility is directly dependent on the  $\text{CO}_2$  partial pressure



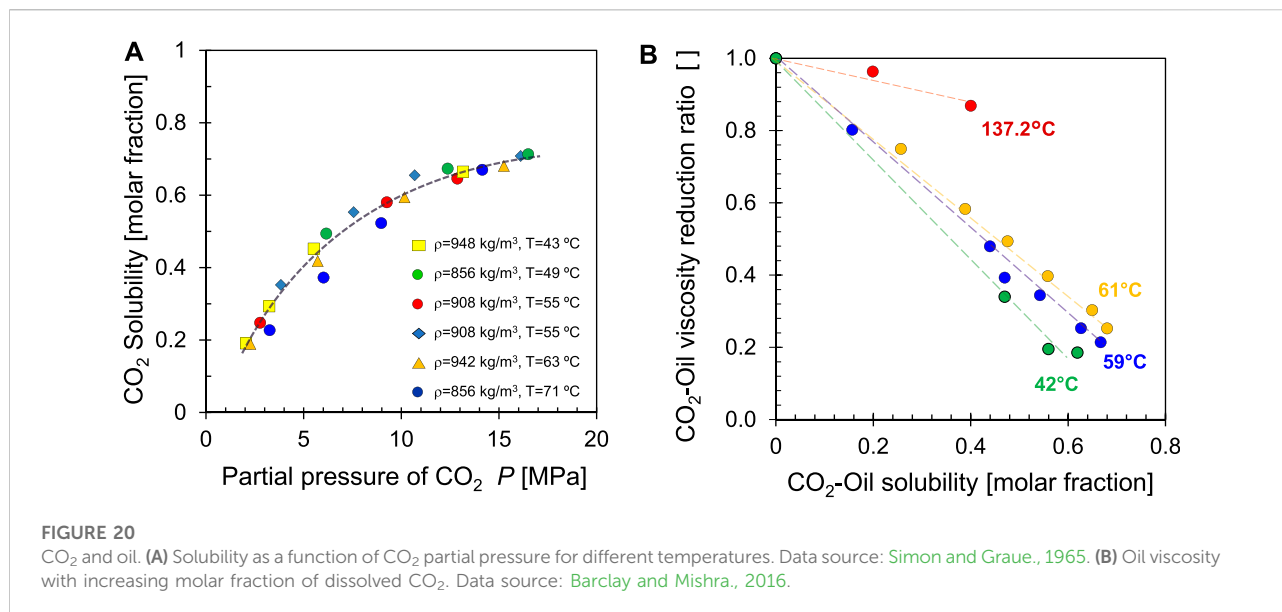


TABLE 3 Minimum miscibility pressure for different crude oils and organic compounds.

Oil	API	Minimum miscibility pressure [MPa]	Temperature [°C]	Method	References
Light density crude oil	41	7.3	27	Vanishing interfacial tension (VIT)	Wang et al., (2010)
Light density crude oil	38	10.6	53	VIT	Gu et al., 2013
Medium density crude oil	29.5	11.7	59	Rising bubble apparatus (RBA)	Dong et al., 2013
Medium density crude oil	25.5	7.25	27	VIT	Nobakht et al., 2008
Heavy density crude oil	10	11.17	44	VIT	Li et al., 2012
Heavy density crude oil	20.2	13.09	60	VIT	Hemmati-Sarapardeh et al., 2014
Heptane	76.7	10	80	VIT	Zolghadr et al., 2013
Hexadecane	52.3	14.2	80	VIT	Zolghadr et al., 2013
Diesel	38.5	13.8	80	VIT	Zolghadr et al., 2013

(Figure 20A), and is also affected by temperature and oil properties. Oil experiences a pronounced decrease in viscosity as CO<sub>2</sub> dissolves in it (Figure 20B).

Crude oil and CO<sub>2</sub> become fully miscible and the interface vanishes when the “minimum miscibility pressure” is reached; Table 3 lists minimum miscibility pressures for various crude oils. Above the minimum miscibility pressure, a transition zone forms in the reservoir. Miscibility is not instantaneous but develops gradually, i.e., “dynamic miscibility” (Verma, 2015).

### Swelling and shrinkage

Miscibility results from (1) the dissolution of CO<sub>2</sub> into the reservoir oil causing *swelling*, and (2) the vaporization of intermediate-molecular-weight hydrocarbons into the injected

CO<sub>2</sub> resulting in *oil shrinkage*. The competitive effects of swelling and shrinkage are evident in the data presented in Figure 21: the oil swelling factor increases with CO<sub>2</sub> pressure until it reaches a maximum volume (at  $p = 8$  MPa when  $T = 50^\circ\text{C}$ , or at  $p = 17$  MPa when  $T = 95^\circ\text{C}$ ); thereafter, vaporization-extraction becomes dominant at high-pressure and oil shrinks.

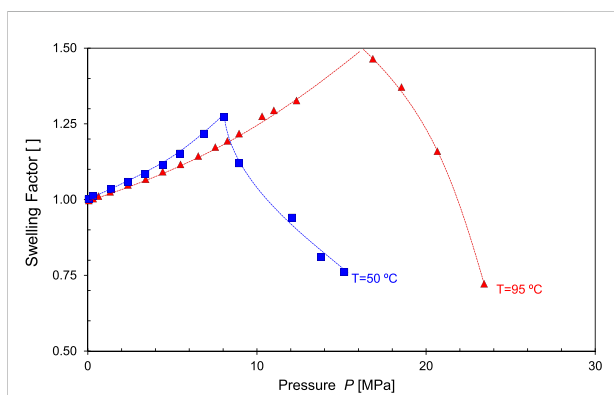
### Asphaltene precipitation

Crude oil contains saturated hydrocarbons (single bonds between carbon atoms and saturated with hydrogen), aromatics (planar rings), resins, and asphaltenes. Asphaltenes are polar, polyaromatic and of high molecular weight substances that are insoluble in n-hexane or n-pentane. The size of asphaltenes ranges from less than 1 nm single molecules in light oils to more than 5 nm

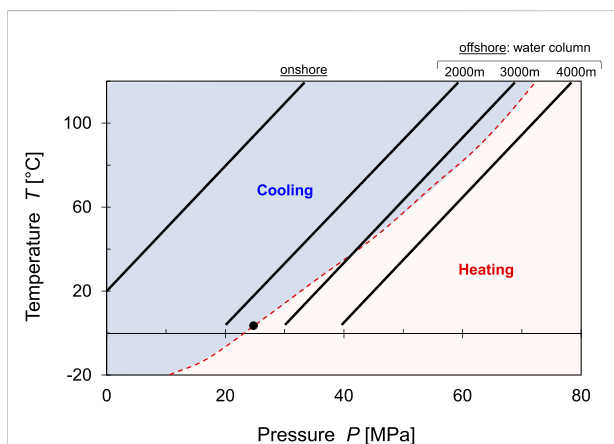
aggregates in black oils and heavy oils (Mullins, 2010). CO<sub>2</sub> dissolution in oil causes asphaltene precipitation; as asphaltene deposits on mineral surfaces, it changes the wettability of the reservoir causing mostly irreversible formation damage. Experimental studies show that ionic liquids prevent asphaltene precipitation (Hu and Guo, 2005; Zheng et al., 2019).

### Thermal effects

The injection of CO<sub>2</sub> triggers thermal changes in surface facilities, the injection well and inside the reservoir (Hoteit et al., 2019).



**FIGURE 21** Miscibility and oil volume change. There are two competing processes: CO<sub>2</sub> dissolution in oil prevails at low pressure and causes oil swelling; on the other hand, light hydrocarbons vaporize into CO<sub>2</sub> and oil shrinks at high pressure. Data source: Rezak and Foroozesh (2019).



**FIGURE 22** Joule-Thompson effect—Near well thermal effects. Cooling is anticipated for most onshore and offshore applications. The “inversion curve” predicts heating for seafloor injections beneath > 2,400 m high water columns.

### Joule-Thomson expansion

The CO<sub>2</sub> injection pressure exceeds the formation pressure and rapid quasi-adiabatic expansion of CO<sub>2</sub> takes place near the injection well. When a gas expands from high pressure to low pressure, the temperature change  $\Delta T$  is related to the pressure change  $\Delta P$  by the Joule-Thomson coefficient  $\mu_{JT}$  [K/MPa], which can be expressed in terms of the gas’ molar volume  $V_m$  [m<sup>3</sup>/mol], its molar heat capacity at constant pressure  $C_p$  [J K<sup>-1</sup> mol<sup>-1</sup>], and the coefficient of thermal expansion  $\alpha$ [K<sup>-1</sup>] as (Demirel., 2014):

$$\left(\frac{\Delta T}{\Delta P}\right) \approx \left(\frac{\partial T}{\partial P}\right)_H = \mu_{JT} = \frac{V_m}{C_p} (\alpha T - 1) \quad (23)$$

Gas expansion causes cooling when the Joule-Thomson coefficient  $\mu_{JT} > 0$  and heating when  $\mu_{JT} < 0$ . In general, the Joule-Thomson coefficient decreases with increasing pressure and temperature (Han et al., 2010), and it is zero at the inversion point when  $\alpha T - 1 = 0$  (Note: the coefficient of thermal expansion  $\alpha$  is temperature dependent as well).

The Joule-Thomson coefficient is  $\mu_{JT} \approx 10^\circ\text{C}/\text{MPa}$  for CO<sub>2</sub> at  $T = 40^\circ\text{C}$  and  $p = 5 \text{ MPa}$ , i.e., five times higher than for air (Phuoc and Massoudi, 2021). Figure 22 shows the Joule-Thomson inversion curve for CO<sub>2</sub>; the auxiliary lines correspond to pressure-temperature profiles for a geothermal gradient of 30°C/km. Clearly, CO<sub>2</sub> expansion will cause cooling at most reservoir pressure and temperature conditions, except for injection in the seafloor under deep water columns and beneath permafrost layers.

### Other thermal effects

In addition to the Joule-Thomson effect, endothermic water vaporization into supercritical CO<sub>2</sub> takes place mainly near the injection well (Vilarrasa and Rutqvist, 2017), while exothermic dissolution of supercritical CO<sub>2</sub> into the host brine prevails further away from the well ( $\Delta H = -171.38 \text{ kJ/kg}$  of CO<sub>2</sub>, at  $p = 10.5 \text{ MPa}$  and  $T = 50^\circ\text{C}$ —Koschel et al., 2006). These two competing thermal effects continue until both brine and scCO<sub>2</sub> are saturated with the other phase.

### Field implications

Under most field conditions, CO<sub>2</sub> injection will cause a significant temperature drop. The extent of cooling scales with the pressure drop and is partially compensated by heat transport in the reservoir (Oldenburg, 2007; Gauteplass et al., 2020). Cooling induces thermo-elastic stresses and triggers hydrate or ice formation when PT-conditions cross the corresponding phase boundary; hydrate and ice hinder fluid flow into the formation (Sloan, 2010). Effective stress changes caused by the injection pressure and thermos-elastic effects can reactivate fractures and improve their transmissivity due to shear dilation (Vilarrasa et al., 2017).

Finally, thermal changes affect density (Figures 1, 2), viscosity (Figure 3), interfacial properties (Figures 4, 5), and

solubility and pH (Figure 6). In the supercritical regime, thermally induced density gradients are most pronounced across the Widom line and near the critical point. The induced thermal gradients give rise to thermally driven convection (Islam et al., 2014).

## Storage density—reservoir size

The storage capacity per unit volume of the reservoir or 'storage density'  $SD$  [ $\text{kg CO}_2/\text{m}^3$  of rock] depends on the  $\text{CO}_2$  phase and the reservoir porosity  $\phi$ :

Dissolved:	$SC = M_{\text{CO}_2}\chi_{\text{CO}_2}$	$SD = 4 - \text{to} - 13\text{kg}/\text{m}^3$
Fluid $\text{CO}_2$ :	$SC = \rho_{\text{CO}_2}$	$SD = 50 - \text{to} - 240\text{kg}/\text{m}^3$
Hydrate:	$SC = 0.234\rho_{\text{hyd}}$	$SD = 25 - \text{to} - 77\text{kg}/\text{m}^3$
Mineralized:	$SC = f(\text{mineral, time, flux})$	$SD = 30 - \text{to} - 200\text{kg}/\text{m}^3$

The listed  $SD$  values correspond to: porosity  $\phi = 0.1$ -to- $0.3$ ,  $\text{CO}_2$  solubility in fresh water at common reservoir PT conditions  $\chi_{\text{CO}_2} = 1$  mol/l (reduce it to about half for brine—Figure 7),  $\text{CO}_2$  molar mass  $M_{\text{CO}_2} = 44$  g/mol,  $\text{CO}_2$  mass density  $\rho_{\text{CO}_2} = 500$ -to- $800$   $\text{kg}/\text{m}^3$  (from supercritical to liquid—Figure 1), mass fraction of  $\text{CO}_2$  in hydrate  $\text{CO}_2 \cdot 8\text{H}_2\text{O}$  equal to  $0.234$ , and hydrate density  $\rho_{\text{hyd}} = 1,100$   $\text{kg}/\text{m}^3$ , and attainable mineralization within the first 10-to-100 years after injection.

The annual worldwide  $\text{CO}_2$  emission is  $\sim 40$  Gt/yr. Let's consider storing 1 gigaton of  $\text{CO}_2$ , i.e., a "small fraction" of the total annual emissions: the required volume for  $\text{CO}_2$  dissolved in fresh water is equivalent to a 30 m thick porous reservoir  $50$  km  $\times$   $50$  km wide, assuming high porosity  $\phi = 0.3$  and a storage density  $SD = 13$   $\text{kg}/\text{m}^3$ . Clearly, higher density  $\text{CO}_2$  forms are preferred for efficient storage; still, the storage density  $SD$  values listed above imply that all  $\text{CO}_2$  geological storage forms will require large reservoir volumes.

## Conclusions

The injection of  $\text{CO}_2$  into geological formations triggers numerous coupled thermo-hydro-chemo-mechanical processes.

The gas-liquid phase boundary and the Widom line within the supercritical regime divide the lighter gaseous-like regime from the denser liquid-like regime. The density of  $\text{CO}_2$  has a first-order effect on the  $\text{CO}_2$ -water interfacial tension.

The  $\text{CO}_2$ -water interface pins strongly to mineral surfaces and results in a wide range of contact angles that vary between the advancing and the receding asymptotic values. Soon after injection, the residual water saturation reflects the reservoir capillary pressure vs. saturation response; thereafter, the degree of water saturation decreases with time as water evaporates into  $\text{CO}_2$ .

The solubility of  $\text{CO}_2$  in water increases with pressure but decreases with temperature. In most geological settings,  $\text{CO}_2$

solubility decreases with salt concentration particularly when small high-valence ions are involved. Similarly, the presence of competing gases affects the solubility of  $\text{CO}_2$  in water because of competing ions and a reduction in the partial pressure of  $\text{CO}_2$  (at constant total pressure).

Dissolved  $\text{CO}_2$  acidifies water and prompts mineral dissolution (silicates and carbonates). Dissolution couples with preferential advective flow channels, diffusive transport and mechanical effects to cause various dissolution patterns, internal shear failure (even under zero lateral strain conditions), and strains may localize along shear bands. Next to the plume lower boundary, convective currents may emerge due to the increased density of water with dissolved  $\text{CO}_2$  and minerals.

$\text{CO}_2$  mineralization homogenizes flow pathways. Crystal growth against pore and fracture walls exert crystallization pressure that can create new fracture surfaces or cause grain-displacive mineral lenses (when the crystallization pressure exceeds the rock tensile strength or the far-field confining stress). Precipitation may also take place within the rock matrix by successive dissolution-diffusion-precipitation sequences that result in periodic precipitation patterns, i.e., Liesegang bands.

Buoyant  $\text{CO}_2$  plumes apply excess pressure against the seal caprock and interact with clay minerals through physicochemical processes. The sealing number  $\Pi_{\text{seal}}$  compares the seal capillary breakthrough pressure to the excess buoyant pressure in the  $\text{CO}_2$  plume. On the other hand, the stability number  $\Pi_{\text{stability}}$  compares the initial effective stress to the excess buoyant pressure. Safe  $\text{CO}_2$  storage requires large sealing and stability numbers and careful consideration of physicochemical effects.

$\text{CO}_2$  injection into depleted oil reservoirs benefits from proven seal performance. At the same time, it causes light-component extraction, oil swelling, ganglia destabilization, a reduction in oil viscosity, and miscible displacement. Then, we can simultaneously sequester  $\text{CO}_2$  while recovering additional oil through  $\text{CO}_2$  gravity-driven enhanced oil recovery.

Rapid  $\text{CO}_2$  depressurization near the injection well causes cooling under most anticipated reservoir conditions. Cooling may trigger hydrate and ice formation, and reduce permeability. In some cases, effective stress changes associated with the injection pressure and cooling thermoelasticity can reactivate fractures.

The  $\text{CO}_2$  storage density in terms of  $\text{kg}$  of  $\text{CO}_2$  per cubic meter of the reservoir depends on the  $\text{CO}_2$  phase and the reservoir porosity  $\phi$ . First order estimates are 4-to-13  $\text{kg}/\text{m}^3$  for  $\text{CO}_2$  dissolved in water, 25-to-77  $\text{kg}/\text{m}^3$  for  $\text{CO}_2$  trapped in hydrate, 30-to-200  $\text{kg}/\text{m}^3$  for mineralized  $\text{CO}_2$  (attainable within the first 10-to-100 years of storage), and 50-to-240  $\text{kg}/\text{m}^3$  for supercritical-to-liquid  $\text{CO}_2$  (most efficient but prone to leakage). In all cases, large reservoir volumes will be required to

store a meaningful fraction of the CO<sub>2</sub> that will be emitted during the energy transition (Anderson, 2017; Global status CCS, 2021).

## Author contributions

All authors listed have made a substantial, direct, and intellectual contribution to the work and approved it for publication.

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## Glossary

- $a_{H^+}$  [] Activity of protons
- $c$  [mol/l] Solute concentration (0: at saturation)
- $C_p$  [J/K] Heat capacity
- $D$  [m<sup>2</sup>/s] Molecular diffusion coefficient
- $d^*$  [m] Characteristic pore diameter
- $Da$  [] Damköhler number
- $e$  [] Void ratio
- $E$  [GPa] Young's modulus
- $g$  [m/s] Gravitational acceleration
- $H$  [m] Height of the fluid column (CO<sub>2</sub>: carbon dioxide column; R: reservoir)
- $k$  [mol m<sup>-2</sup> s<sup>-1</sup>] Rate constant of mineral dissolution
- $k_0$  [] Lateral coefficient earth pressure (0: at rest; a: Rankine's active case)
- $k_H$  [mol L<sup>-1</sup> atm<sup>-1</sup>] Henry's constant
- $k_{hyd}$  [m<sup>2</sup>] Permeability
- $L_{ch}$  [m] Channel length
- $l_0$  [m] Intermolecular distance
- $M$  [g/mol] Molar mass (CO<sub>2</sub>; silic: silicate minerals)
- $n_{H^+}$  [] Reaction order with respect to the activity of protons
- $P$  [Pa] Pressure (0: initial pressure; CO<sub>2</sub>: partial pressure CO<sub>2</sub>; nw: non-wetting fluid; w: wetting fluid; crys: on a growing crystal; sol: solution)
- $Pe$  [] Péclet number
- $R$  [J mol<sup>-1</sup> K<sup>-1</sup>] Gas constant
- $Ra$  [] Rayleigh number
- $r$  [m] Radius of the interface
- $r_d$  [mol/s] Mineral dissolution rate
- $S$  [] Supersaturation
- $s$  [m<sup>2</sup>] Surface area
- $S_s$  [m<sup>2</sup>/g] Specific surface area
- $S_w$  [] Degree of water saturation
- $SD$  [kg/m<sup>3</sup>] Storage density in terms of kg of CO<sub>2</sub> per m<sup>3</sup> of the reservoir
- $T$  [K] Absolute temperature
- $T_p$  [m] Plume thickness
- $t$  [s] Time (adv: advection; diff: diffusion; r: reactive)
- $V_m$  [m<sup>3</sup>/mol] Mineral molar volume
- $v$  [m/s] Average flow velocity
- $z$  [m] Height of the sediment column
- $z_w$  [m] Height of the water column above the seafloor
- $\alpha$  [K<sup>-1</sup>] The coefficient of thermal expansion
- $\beta$  [] Factor related to the pore size distribution
- $\gamma$  [N/m] Interfacial tension (FS: fluid-solid; LS: liquid-solid; LF: fluid-liquid)
- $\Delta H$  [kJ/kg] Enthalpy change in phase transformation
- $\epsilon$  [] Strain
- $\theta$  [°] Contact angle (A: advancing; R: receding)
- $\kappa$  [1/s] Kinetic rate of mineral dissolution
- $\mu$  [μPa·s] Fluid Viscosity (w: water)
- $\mu_{JT}$  [] Joule-Thomson coefficient
- $\Pi_{seal}$  [] Dimensionless ratio (seal: sealing number; stability: stability number)
- $\rho$  [kg/m<sup>3</sup>] Density (b: bulk; CO<sub>2</sub>; hyd: hydrate; m: mineral; W: water; 0: initial fluid density)
- $\sigma'$  [Pa] Effective stress (h: horizontal; v: vertical; z<sub>0</sub>: under zero lateral strain conditions)
- $\chi_{CO_2}$  [mol/L] Solubility of CO<sub>2</sub>
- $\phi$  [] Porosity
- $\psi$  [] Geometric fabric factor