



Entropy Analysis of a Valveless Piezoelectric Pump With Hyperbolic Tubes

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In order to solve the problems of large reflux and low output performance of the valveless piezoelectric pump with conical tubes, the valveless piezoelectric pump with hyperbolic tubes was proposed. However, the previous research only paid attention to the difference of forward and reverse flow resistance inside the tubes. For this reason, the prototypes of conical tube valveless piezoelectric pump and hyperbolic tubes valveless piezoelectric pump were fabricated in this paper. The maximum output flow rate of the valveless piezoelectric pump with hyperbolic tubes was 54 ml/min. Subsequently, the size and position of entropy production of the hyperbolic tube and the conical tube were numerically calculated based on entropy production theory in this paper. The results show that the entropy production rate of the hyperbolic tube was significantly lower than that of the conical tube, which is consistent with the experimental results. This research analyzed the energy loss inside the valveless piezoelectric pump by using the entropy generation theory, and provided a new design and research method for improving the output performance of the valveless piezoelectric pump in the future.

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INTRODUCTION

As the core component of the micro-fluid control system, the output performance of the micro pump determines the characteristics of the entire system. According to different driving methods, micro pumps can be divided into electromagnetic pumps (Mi et al., 2020; Rusli et al., 2018) heat pumps (Feng et al., 2021; Skye and Wu, 2021), pneumatic pumps (Hamid et al., 2017; Ma et al., 2019), shape memory pumps (Robertson et al., 2016) and piezoelectric pumps (Huang et al., 2020; Ma et al., 2016; Ma et al., 2015). Among them, piezoelectric pumps have attracted more and more attention for their outstanding advantages such as high energy density, fast response speed, large driving force, and easy integration. According to whether there are valves inside, piezoelectric pumps can be divided into valve piezoelectric pumps and valveless piezoelectric pumps (Singhal et al., 2004; Li et al., 2021). The valve piezoelectric pump has a movable valve, which can avoid backflow and generate a larger output pressure. However, due to the existence of micro-valves, processing and assembly are difficult. In addition, the valve is easy to wear and there is a risk of fatigue damage, which greatly reduces the reliability of the valve piezoelectric pump. Compared with the former, the valveless piezoelectric pump uses a fluid diode to replace the active valve, which simplifies the structure of the valveless piezoelectric pump, makes it easy to miniaturize, and improves the system integration (Zhang et al., 2017).

The valveless piezoelectric micropump is usually composed of a piezoelectric vibrator, a pump body and a fluid diode. The fluid diode has the characteristics of unequal flow resistance between the

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forward and reverse directions, so it is the most important part of the valveless pump. Among them, the cone tube, as the most classic fluid diode, was first used in valveless piezoelectric pumps (Stemme and Stemme, 1993; Olsson et al., 1997). Because there is no movable valve, there are shortcomings of large output return and low output performance for the valveless piezoelectric pump. To this end, many researchers have proposed valveless piezoelectric pumps such as Tesla tube (Morris and Forster, 2003; Gamboa et al., 2005), Y-shaped tube (Zhang et al., 2008), multi-stage bifurcated tube0 (Fadl et al., 2010; Huang et al., 2014), spiral tube (Leng et al., 2013), and vortex tube (Huang et al., 2019; Huang et al., 2021) to reduce backflow and increase output flow.

However, the conical tubes valveless piezoelectric pump is still widely studied due to its simple structure. Among them, Zhang et al. (Bian et al., 2020) proposed the hyperbolic tubes valveless piezoelectric pump based on the conical tube, aiming to improve the output performance of the valveless piezoelectric pump by reducing the vortex during the flow of the piezoelectric pump. However, this study only used software simulation (Tang et al., 2021) and combined with experimental research on the internal flow field distribution of the hyperbolic tubes to illustrate the flow characteristics of the pump (Wang et al., 2020). Therefore, this article will analyze minutely the internal flow and energy loss of the hyperbolic tubes valveless piezoelectric pump based on the entropy generation theory. Subsequently, the size and position of entropy generation in the conical tube and the hyperbolic tube are obtained by numerical calculation. Finally, the output performance experiments of the prototypes are carried out. The calculation and analysis of fluid diodes using entropy generation theory can provide new ideas for improving the output performance of valveless piezoelectric pumps.





STRUCTURAL DESIGN AND WORKING PRINCIPLE

Design of Hyperbolic Tube

Figure 1 shows a schematic diagram of the conical tube structure, in which *L* is the length of the conical tube, *D* is the diameter of the throat of the conical tube, and θ is the cone angle. The hyperbolic tube is made by changing the structure of the conical tube, that is, the gradually expanding straight wall of the conical tube is changed to a hyperbolic wall, and the throat of the conical tube is changed to a circular arc transition section. The left side of **Figure 2** shows the wall curve of the hyperbolic tubes, where h_1 is the height of the arc section, h_2 is the height of the hyperbola section, and *r* is the minimum radius of the flow channel of the tube. The hyperbolic function is:

$$\frac{x^2}{a^2} - \frac{y^2}{b^2} = 1$$
 (1)

Among them, *a* is equal to *r*, and *b* is the only variable in **Formula** (1). Constrain the hyperbola with two straight lines y = 0 and y = 9. The arc part at the throat is constructed by **Formula** (2):

$$(x-c)^{2} + (y-d)^{2} = r$$
 (2)





Finally, the obtained hyperbolic arc streamline is rotated 360° around the axis to obtain the inner flow passage of the hyperbolic tube as shown in the right side of **Figure 2**.

Working Principle

The valveless piezoelectric pump with hyperbolic tubes is composed of a piezoelectric vibrator, a pump body containing two hyperbolic tubes, and a pump cover, as shown in **Figure 3**. In this study, the fluid flowing from the arc section into the hyperbola is defined as the forward flow, and the fluid flowing from the hyperbolic section into the arc section is defined as the reverse flow. The piezoelectric vibrator receives the current

TABLE 1 Analysis of grid independer	TABLE 1	Analysis	of grid	independent
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Tube1			Tube2		
N/10 ⁵	Q/kg⋅s ⁻¹	r/ %	N/10 ⁵	Q/kg⋅s ⁻¹	r/ %
1.35	14285.7	_	2.62	13454.4	_
1.84	14416.4	0.91	3.03	13539.9	0.63
2.8	14532.2	0.8	3.76	13594.6	0.4
4.0	14644.9	0.77	3.54	13625.7	0.23

voltage signal to generate up and down vibrations to change the volume of the pump cavity and change the pressure inside the pump cavity, so as to realize the unidirectional transportation of the fluid on the macroscopic level.

The circular piezoelectric vibrator is fixed by adhering its periphery to the pump body, and AC voltage is applied to generate periodic vibration, and its vibration shape is approximately parabolic. By establishing a polar coordinate system, the origin of the coordinate system is the center of the piezoelectric vibrator, then the polar coordinate equation of the parabola:

$$W(r) = \omega_0 \left(1 - \frac{r_2}{R_2} \right) \tag{3}$$

In the formula, Ω_0 is the maximum amplitude, *R* is the radius of the piezoelectric vibrator, and the maximum volume change of the pump cavity is:

$$\Delta V = 2\pi \int_{0}^{R} \omega_0 \left(1 - \frac{r^2}{R^2} \right) = \pi \omega_0 \frac{R^2}{2}$$
(4)

The flow resistance experienced by the fluid in the hyperbolic tube is mainly local resistance and friction resistance, as shown below:

$$\gamma_{+} = \gamma_{\tau+} + \gamma_{D+} \tag{5}$$





TABLE 2 | Entropy production of conical tube under different pressure difference.

Differential pressure (Pa)	Viscous entropy generation (10 ⁻⁴)	Turbulent entropy production	
500	12.66	9.97	
1,000	23.66	20.08	
2,000	46.79	40.93	
4,000	85.47	107.55	

$$\gamma_{-} = \gamma_{\tau^{-}} + \gamma_{D^{-}} \tag{6}$$

In the formula, γ_+ is the total forward flow resistance coefficient of the hyperbolic tube, and γ_- is the total reverse flow resistance coefficient of the tube. The output flow of the pump can be expressed as:

$$Q = \nabla V f \frac{(\gamma_+ - \gamma_-)}{(2 + \gamma_+ + \gamma_-)}$$
(7)

Where f is the driving frequency of the piezoelectric vibrator.

EXPERIMENTAL SETUP

In this paper, a stereolithography machine (SLA) was used to fabricate the valveless piezoelectric pump with conical tubes and the valveless piezoelectric pump with hyperbolic tubes. As shown in the left of **Figure 4**, the pump body is made of photosensitive resin. The tube parameters are shown in **Table 1**, the pump cavity diameter is 41 mm, the pump cavity depth is 0.8 mm, and the vibrator radius is 41 mm.

The experimental setup of the piezoelectric pumps with two tubes is shown in the right of **Figure 4**. The input AC voltage of the piezoelectric pump is set to 200 Vpp (peak-to-peak voltage). The reservoir contained a certain amount of deionized water. Take the small radius of the flow tube as the inlet and the flow tube with joints as the outlet, and the liquid level of the inlet and outlet are the same during the experiment. The analytical balance is used to measure the output flow of the piezoelectric pump per unit time under different frequencies. The accuracy of the analytical balance is 0.001 g.

SIMULATION ANALYSIS

Entropy Generation Theory

During the operation of the valveless piezoelectric pump, due to the viscosity of the working fluid and Reynolds stress, the consumed mechanical energy is irreversibly converted into internal energy. According to the second law of thermodynamics, the entropy generation theory is used to calculate the energy loss.

For turbulent flow, based on the Reynolds time-average equation, the entropy generation rate can be divided into two parts, the one is caused by the average velocity, called the direct dissipation term; the other is caused by the pulsating velocity, called the turbulent dissipation term. So, define the local total entropy production rate S_D per unit volume as:

$$S_D = S_{D1} + S_{D2}$$
(8)

In the formula, S_{D1} is the entropy production rate caused by the average speed, $W/(m3 \cdot K)$; S_{D2} is the entropy production rate caused by the pulsating velocity, $W/(m3 \cdot K)$.

TABLE 3 | Entropy production of hyperbolic tubes under different pressure difference.

Differential pressure (Pa)	Viscous entropy generation (10 ⁻⁴)	Turbulent entropy production	
500	5.11	5.43	
1,000	14.42	14.2	
2,000	29.07	31.59	
4,000	72.13	85.11	



The entropy production rate caused by the average speed is:

$$S_{D1} = \frac{\mu}{T} \left[\left(\frac{\partial \bar{u}}{\partial y} + \frac{\partial \bar{v}}{\partial x} \right)^2 + \left(\frac{\partial \bar{u}}{\partial z} + \frac{\partial \bar{\omega}}{\partial x} \right)^2 + \left(\frac{\partial \bar{\mu}}{\partial y} + \frac{\partial \bar{\omega}}{\partial y} \right)^2 \right] + 2 \frac{\mu}{T} \left[\left(\frac{\partial \bar{u}}{\partial x} \right)^2 + \left(\frac{\partial \bar{v}}{\partial y} \right)^2 + \left(\frac{\partial \bar{\omega}}{\partial z} \right)^2 \right]$$
(9)

The entropy production rate caused by the pulsating velocity is

$$S_{D2} = \frac{\mu_{eff}}{T} \left[\left(\frac{\partial \mu'}{\partial y} + \frac{\partial \nu'}{\partial x} \right)^2 + \left(\frac{\partial u'}{\partial z} + \frac{\partial \nu'}{\partial x} \right)^2 + \left(\frac{\partial u'}{\partial y} + \frac{\partial \omega'}{\partial y} \right)^2 \right] + 2 \frac{\mu}{T} \left[\left(\frac{\partial u'}{\partial x} \right)^2 + \left(\frac{\partial \nu'}{\partial y} \right)^2 + \left(\frac{\partial \omega'}{\partial z} \right)^2 \right]$$
(10)

Where \bar{u} , \bar{v} , \bar{w} , are the components of the average velocity in the x, y, and z directions, m/s; u', v', w' are the components of the pulsating velocity in the x, y, and z directions, m/s; *T* is the temperature, K; μ is the hydrodynamic viscosity, Pa·s; μ_{eff} is the effective dynamic viscosity of the fluid, Pa·s, which can be calculated by **Formula** (11), where μ_{t} is the turbulent flow viscosity, Pa·s.

$$\mu_{\rm eff} = \mu + \mu_{\rm t} \tag{11}$$

However, the pulsating velocity component cannot be calculated in the Reynolds average equation. This paper adopts the k-epsilon turbulence (Shi et al., 2020; Han et al., 2021) model. Therefore, the entropy generation rate produced by the pulsating velocity can be approximated by the following formula.

$$S_{D2'} = \beta \frac{\rho f k}{T} \tag{12}$$

In the formula, the empirical coefficient $\beta = 0.09$, f is the characteristic frequency of turbulent pulsation, s⁻¹, k is turbulent kinetic energy, m²/s²; ρ is the density, kg/cm³.

Integrating the above-mentioned parts of the entropy production rate in the integration area, the direct dissipation entropy production $S'_{D1}(W/k)$ and the turbulent dissipation entropy production $S'_{D2}(W/k)$ can be obtained respectively.

$$S'_{D1'} = \int S_{D1} dV$$
 (13)

$$S'_{D2'} = \int S_{D2'} dV$$
 (14)

Therefore, the entropy production of the mainstream area in the entire computational domain is:

$$S_D = S'_{D1'} + S'_{D2'} \tag{15}$$

According to the entropy generation calculation results of the conical tube and the hyperbolic tube, the resistance coefficient ε is used to describe the flow performance and defined in Eq. 16. the impedance ratio is defined in Eq. 17, which is adopted to describe the unidirectional transport performance of the tube.





$$\varepsilon = \frac{2\nabla P}{\rho v^2} \tag{16}$$

$$\delta = \frac{\varepsilon_R}{\varepsilon_F} \tag{17}$$

Hyperbolic Tube Grid Establishment

The finite element models of the tubes are shown in **Figure 5** (the conical tube on the top and the hyperbolic tube on the bottom). The *k*-epsilon turbulence model is adopted, because there is small turbulence in tubes. The wall is a non-slip boundary, the pressure applied to the inlet is 500, 1,000, 2,000, and 4,000 Pa, the outlet of the tube is defined as 0 Pa, and the numerical convergence residual is set to 10^{-4} .

The grid independence of the model is verified. **Table 1** shows the calculation results of grid independence, where N is the number of grids, Q is the flow rate (kg/s), and r is the relative error. When the

mesh numbers of the conical tube and the hyperbolic tube are 375,700 and 279,000 respectively, the error is less than 1%, so the finite element model is used to calculate the entropy generation.

RESULTS AND DISCUSSION

The experimental result of the pump flow rate is shown in **Figure 6**. The output flow rate of the piezoelectric pump with hyperbolic tubes is finally measured at a frequency of 40 Hz, which is 54 ml/min. The maximum output flow rate of the piezoelectric pump with conical tube is 47 ml/min. Therefore, the output performance of the piezoelectric pump with hyperbolic tubes is greater than that of the piezoelectric pump with conical tubes.

Calculate the viscous entropy production and turbulent entropy production under different pressure differences



when the two tubes are flowing in the forward direction, as shown in **Table 2** and **Table 3**. It can be seen that under the same working conditions, the turbulent entropy production and viscous entropy production of the hyperbolic tube are smaller than that of the conical tube. Therefore, the output performance of the piezoelectric pump with hyperbolic tubes is better than that of the piezoelectric pump with conical tubes.

The difference between the forward and reverse flow resistance of the hyperbolic tube directly affects the cut-off performance, which in turn affects the output performance of the valveless piezoelectric pump with hyperbolic tubes.

Figure 7 shows that the forward flow resistance of the hyperbolic tube is smaller than the reverse flow resistance, indicating that unidirectional transport of fluid can be achieved. Moreover, the impedance ratio of the hyperbolic tube is larger than that of the conical tube, indicating that the unidirectional transport performance of the hyperbolic tube is better than that of the conical tube.

Figure 8A shows the turbulent entropy production of the conical tube. The turbulent entropy production is mainly caused by turbulent dissipation. It can be seen that the turbulent entropy production is symmetrically distributed on the axis of the conical tube, and it is greater in the reverse flow than that in the forward flow. **Figure 8B** shows the viscous entropy production of the conical tube. Compared with the turbulent entropy production, the viscous entropy production is smaller. The viscous entropy production is mainly caused by the dissipation of the irreversible viscous entropy production. The viscous entropy production is at the smallest radius of the tube. The maximum, decreases as the tube spreads.

Figure 9A shows the viscous entropy production of the hyperbolic tube. Compared with the conical tube, the viscous entropy production area of the hyperbolic tube in the forward and reverse flow is smaller than that of the conical tube. Figure 9B shows the turbulent entropy production of the hyperbolic tube. Compared with the conical tube, the turbulent entropy production of the hyperbolic tube is much smaller than

that of the conical tube, and the viscous entropy production is less dissipated. In summary, the entropy production at the smallest channel diameter of the tube is larger, and both viscous entropy production and turbulent entropy production are positively correlated with the pressure difference, but the turbulent entropy production and viscous entropy production of the hyperbolic tube are both smaller than those of the conical tube.

Figure 10A and **Figure 11A**respectively show the velocity cloud diagrams of the conical tube and the hyperbolic tube. It can be seen that the flow velocity of the fluid in the hyperbolic tube is greater than that in the conical tube. In addition, there is a larger vortex in the conical tube, which is larger than the hyperbolic tube. **Figure 10B** and **Figure 11B** respectively show the pressure cloud diagrams of the conical tube and the hyperbolic tube. It can be seen that the reverse cut-off performance of the hyperbolic tube is better than that of the conical tube.

Therefore, under the same working conditions, the energy loss of the hyperbolic tube is smaller than that of the conical tube, and the output performance of the hyperbolic tube piezoelectric pump is better than that of the conical tube piezoelectric pump.

CONCLUSION

In order to improve the flow characteristics of the valveless piezoelectric pump, two valveless piezoelectric pumps with tubes with different curves as the inlet and outlet pipes were developed. The prototypes of the piezoelectric pumps were fabricated, and the experimental results showed that the hyperbolic tubes had the largest flow rate, and the corresponding valveless piezoelectric pump had better output performance. The entropy generation theory was used to analyze the energy loss of the tube, and the entropy production values of the tubes were calculated. The results show that the total entropy produced by the hyperbolic curve tube is smaller than the entropy produced by the conical tube. The entropy generation theory can effectively predict and analyze the energy loss of the tubes, and provide a reference for the structural design and performance optimization of the piezoelectric pump.

DATA AVAILABILITY STATEMENT

The datasets presented in this study can be found in online repositories. The names of the repository/repositories and

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AUTHOR CONTRIBUTIONS

QY and JC contributed to conception and design of the study. QY and JC performed the simulation analysis. QY wrote the first draft of the manuscript. JC wrote sections of the manuscript. All authors contributed to manuscript revision, read, and approved the submitted version.

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