



Study of Supercritical CO₂ Physical Property Calculation Models

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The supercritical CO_2 (SC-CO₂) thermodynamic cycle has been a promising and revolutionary technology. Related research has put higher demands on CO_2 physical property calculations in terms of computational accuracy and speed. In this study, the deviations between the experimental data and predicted data calculated by several models of CO_2 physical property calculation were obtained. According to the comparison results, the Span–Wagner equation of state, the Vesovic model, and the free-volume viscosity model are selected to construct a set of high-precision CO_2 property calculation models. For the time-consuming problem of commonly used models, new models were developed by using the polynomial fitting and interpolation method, which improved the speed of the physical property calculation by five orders of magnitude while ensuring high accuracy. On this basis, a physical property calculation program for the SC- CO_2 thermodynamic cycle could be developed, which can meet the demands of engineering applications in accuracy and calculation speed.

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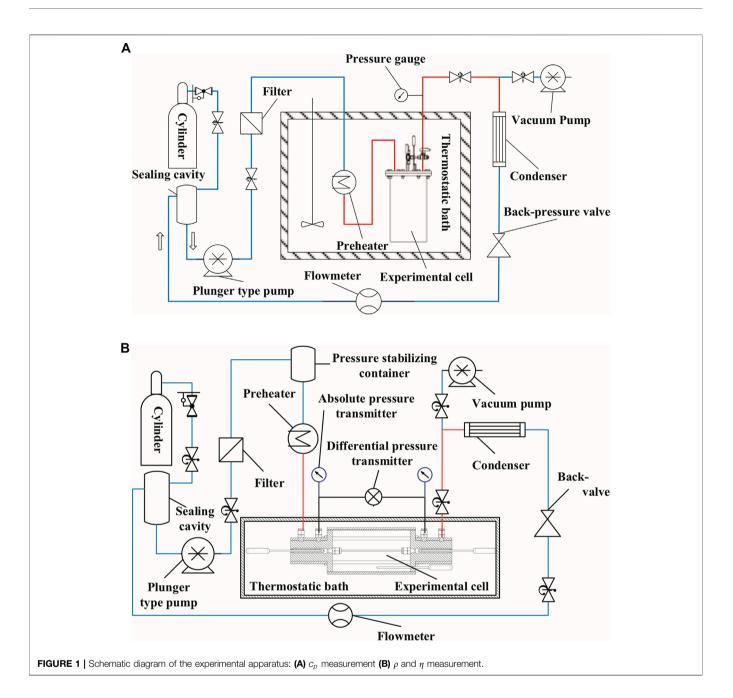
Min L, Liu M, Xi D, He M, Liu X and Huang Y (2022) Study of Supercritical CO₂ Physical Property Calculation Models. Front. Energy Res. 10:835311. doi: 10.3389/fenrg.2022.835311 Keywords: supercritical CO₂, thermodynamic properties, transport properties, property models, polynomial fitting, interpolation method

INTRODUCTION

In recent years, supercritical fluid (SCF) techniques have played a catalytic role in scientific research and engineering applications. In the food, pharmaceutical, and cosmetic industries, SCF extraction technology enables soluble components to be extracted from raw materials with high speed and efficiency, taking advantage of the unique properties above the critical point (Sodeifian et al., 2016a; Sodeifian et al., 2016b). SCFs have also been introduced into the biochemical and pharmaceutical industries to improve the bioavailability of drugs by producing nano- and micro-sized drug particles (Sodeifian et al., 2017; Sodeifian and Sajadian 2018; Sodeifian et al., 2018; Sodeifian et al., 2019) and to make it simpler to control the characteristics of the ultimate product by supercritical solvent impregnation (Ameri, Sodeifian, and Sajadian, 2020). Meanwhile, supercritical dyeing technology has been used to prevent environmental pollution (water) and improve the performance of the pigments during the dyeing process (Ardestani et al., 2020).

In the field of energy and power, ultra-supercritical units have drawn much attention in recent years due to their energy-saving and emission reduction advantages (Wang 2015). In related studies, supercritical carbon dioxide (SC-CO₂) has been introduced into the new energy conversions as a low-pressure alternative to supercritical water. The critical pressure of carbon dioxide is moderate. In the pseudocritical region, CO₂ has the characteristics of large specific heat, low viscosity, and high compressibility, the utilization of which makes the SC-CO₂ thermodynamic cycle system compact

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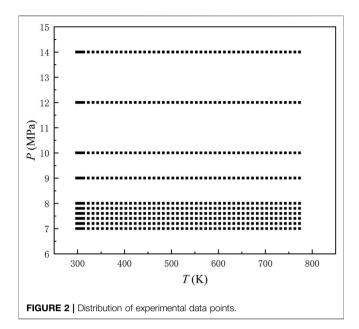


and efficient, showing a promising future in the fields of advanced nuclear power (Huang and Wang 2012) and industrial waste heat recovery (Tao and Tao 2016).

In the SC-CO₂ thermodynamic cycle, the distortions of CO₂ properties in the pseudocritical region such as the isothermal compression coefficient are beneficial to reduce the power consumption of the compressor and improve the system efficiency but also pose a greater challenge for the calculation of the properties. Moreover, cycle analysis, the design of key thermodynamic processes, and components all place high demands on the physical properties of CO₂ (Yang et al., 2020). Therefore, how to calculate the thermodynamic and

transport properties of CO_2 within the cycle conditions quickly and accurately has been a research focus.

However, the existing CO_2 property calculation models and CO_2 property calculation programs are no longer well suited to the needs of the rapid development of SC-CO₂ research. For example, the Span–Wagner equation of state (Span and Wagner 1996), recognized as the equation of state with the highest accuracy, is limited by its complex form, which means that it is time-consuming. The commonly used property database NIST REFPROP, which is based on the Span–Wagner equation, is hardly available to large-scale calculations owing to the low computational speed. In previous studies, some scholars have



noticed such problems and have used polynomial fitting methods (Bahadori et al., 2009; Ouyang 2011) and neural network algorithms (Zhang et al., 2019) to obtain property prediction models with higher accuracy and calculation speed. Nevertheless, these available models do not correspond to the actual needs of the SC-CO₂ thermodynamic cycle in terms of applicable working conditions and solving methods, and little research work has concentrated on the efficiency evaluation of these models to prove their superiority of computational efficiency. Thus, it is necessary to develop new models differing from the existing ones and pay attention to computational speed comparison of the commonly used prediction models with that of the new models.

This study aims to meet the demands for accuracy and efficiency of CO_2 property calculation in the thermodynamic cycle. Meanwhile, a CO_2 physical property measurement experiment was conducted to obtain precise property data. On this basis, we compared the deviations between the experimental and predicted values of various CO_2 physical property calculation models and finally selected the Span–Wagner equation of state, the Vesovic model (Vesovic et al., 1990), and the free-volume viscosity model (Liu et al., 2013) to make up a set of high-precision CO_2 property calculation models. Also, to solve the time-consuming problem of traditional high-precision models, a high-speed calculation model of CO_2 physical properties was constructed by using the polynomial fitting and interpolation method, which also satisfies the requirement for accuracy.

RESEARCH ON HIGH-PRECISION CO₂ PROPERTY MODELS Experiment in CO₂ Physical Property Measurement

Although much research has been done on CO_2 property data, few experimental studies have been conducted in the range of

TABLE 1 | Uncertainties of the experiment.

U(p)/MPa	U(7)/K	U(c _p)/%	U(p)/%	U(η)/%
0.003	0.04	0.9	0.06	1.4

supercritical conditions. Thus, to obtain precise experimental data in the supercritical region, the experiment of CO_2 physical property measurement was carried out so that the accuracy of available models could be checked.

In this work, the density (ρ), isobaric specific heat (c_p), and viscosity (η) of CO₂ temperatures from 298.15 to 773.15 K and pressures from 7 to 14 MPa were measured through the U-tube vibration method (Jiang 2011), flow calorimeter (Liu et al., 2014; He et al., 2015), and capillary method (Liu et al., 2011), respectively. The physical property measurement experiment apparatus is shown in **Figure 1**.

Several key technical problems have been solved in this experiment. First, a double-layer vacuum cylinder was used to ensure thermal insulation and prevent high temperature instability. Second, in order to minimize the heat loss, the fluid was heated by a microheater inside the tube assembly, and the heating tubes were fine. In addition, the horizontal placement of the experimental body proposed by Liu et al. (2014) was also used to avoid the vertical density stratification produced because of gravity.

Figure 2 presents the distribution of data points, where the horizontal and vertical coordinates indicate the temperature (in K) and pressure (in MPa), respectively. **Table 1** summarizes the total measurement uncertainties of pressure and temperature, and the extended uncertainties at confidence factor k = 2 of c_p , ρ , and η . The uncertainty of c_p comes from the measurement uncertainties of the differential temperature, the mass flow, the heating power, and the pressure. The uncertainty of ρ is mainly from the uncertainty of the density measurement itself, and the uncertainty of η is composed of the uncertainties of the inner diameter of the capillary tube, the mass flow, the differential pressure, the length of the capillary tube, the kinetic energy loss coefficient, the density, and the differential temperature. The equations for calculating synthetic uncertainty of c_p and η are as follows:

$$u_{c}(c_{p}) = \sqrt{\left(\frac{\partial c_{p}}{\partial \Delta T}\right)^{2} u_{\Delta T}^{2} + \left(\frac{\partial c_{p}}{\partial q_{m}}\right)^{2} u_{q_{m}}^{2} + \left(\frac{\partial c_{p}}{\partial Q}\right)^{2} u_{Q}^{2} + \left(\frac{\partial c_{p}}{\partial p}\right)^{2} u_{p}^{2}}, \quad (1)$$

$$u_{c}(\eta) = \sqrt{\left(\frac{\partial\eta}{\partial R}\right)^{2} u_{R}^{2} + \left(\frac{\partial\eta}{\partial q_{m}}\right)^{2} u_{q_{m}}^{2} + \left(\frac{\partial\eta}{\partial \Delta p}\right)^{2} u_{\Delta p}^{2} + \left(\frac{\partial\eta}{\partial L}\right)^{2} u_{L}^{2} + \left(\frac{\partial\eta}{\partial m}\right)^{2} u_{m}^{2} + \left(\frac{\partial\eta}{\partial \Delta T}\right)^{2} u_{\Delta T}^{2}, \quad (2)$$

where ΔT is the differential temperature, q_m is the mass flow, Q is the heating power, p is the pressure, R is the inner diameter of the capillary tube, Δp is the differential pressure, L is the length of the capillary tube, m is the kinetic energy loss coefficient, ρ is the density, and u_{factor} is the the uncertainty of some factors.

Selection and Evaluation Method

In this section, several commonly used models of CO_2 thermodynamic and transport properties were tested. c_p , ρ , and η calculated by different models were compared with

TABLE 2	Comparison between th	predicted and experimental	data out of pseudocritical region.
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Physical quantity	Abbreviation of models	AARE (%)	R ²
Density	FEQ (Span and Wagner 1996)	0.144	1.000
	FEK (Kunz et al., 2007)	0.145	1.000
	FES (Span and Wagner 2003)	26.771	0.782
Isobaric specific heat	FEQ (Span and Wagner 1996)	1.215	0.998
	FEK (Kunz et al., 2007)	1.244	1.000
	FES (Span and Wagner 2003)	1.244	1.000
Viscosity	VS1 (Vesovic et al., 1990)	1.819	0.998
	ECS (Klein, Mclinden, and Laesecke 1997)	3.411	0.991
	VS4 (Quinones-Cisneros and Deiters 2006)	1.713	0.998

TABLE 3 | Comparison between the predicted and experimental data in pseudocritical region.

т/к	p/MPa	η _{exp} /μPa⋅s	AE(%) VS1	VS4
8	307.2	38.17	3.665	4.716
8	309.2	25.72	1.117	2.276
9	311.2	41.47	2.409	3.388
9	313.2	33.88	2.714	3.867
10	323.2	27.95	1.495	2.110
12	323.2	43.53	0.561	1.421
12	333.2	31.57	1.562	2.049
14	333.2	41.99	0.097	0.820
14	343.2	33.76	0.930	1.276

partial experimental data. According to the comparison results, the best performing models were selected. Furthermore, the selected models were evaluated with all the experimental data to reconfirm the accuracy.

Referring to Heidaryan and Jarrahian's (2013) article, this article introduced four statistical parameters to present the comparison results, including the average relative error (ARE), average absolute relative error (AARE), absolute relative error (AE), and correlation coefficient (R^2) (see Eq. 3–6):

$$ARE\% = \frac{100}{N_d} \times \sum_{i=1}^{N_d} \left(\frac{\rho_i^{exp.} - \rho_i^{calc.}}{\rho_i^{exp.}} \right)$$
(3)

AARE% =
$$\frac{100}{N_d} \times \sum_{i=1}^{N_d} \left| \frac{\rho_i^{\text{exp.}} - \rho_i^{\text{calc.}}}{\rho_i^{\text{exp.}}} \right|,$$
 (4)

AE% = 100 ×
$$\left| \frac{\rho_i^{\text{exp.}} - \rho_i^{\text{calc.}}}{\rho_i^{\text{exp.}}} \right|$$
, and (5)

$$R^{2} = 1 - \frac{\sum_{i=1}^{N_{d}} \left(\rho_{i}^{\text{exp.}} - \rho_{i}^{\text{calc.}}\right)^{2}}{\sum_{i=1}^{N_{d}} \left(\rho^{\text{mean}} - \rho_{i}^{\text{calc.}}\right)^{2}},$$
(6)

where ARE indicates the bias between the evaluated data and reference data. The smaller the ARE, the smaller the systematic bias of the data. When ARE is 0, it indicates that the values of evaluated data are randomly distributed around the reference data values. The AARE characterizes the accuracy of the evaluated data compared with the reference data. The maximum value of AE is the maximum deviation of the evaluated data from the reference data. The R^2 represents the precision of the evaluated data. The closer the R^2 is to 1, the stronger the correlation between the evaluated data and the reference data.

The specific selection criteria are as follows. Comparing data out of the pseudocritical region, a smaller AARE value and R^2 value closer to 1 indicate higher accuracy and precision of the model, respectively; if the values of these two statistical parameters are close between different models, then the AE (max) values in the pseudocritical region are compared, and the model with the smaller ones is counted as the better one.

RESULTS AND DISCUSSION

Table 2 shows the comparisons between the predicted data obtained by different commonly used models and the experimental data out of the pseudocritical region. For ρ and c_p calculations, the AARE of the FEQ model (the Span–Wagner equation) is only 0.144 and 1.215%, respectively, and the R^2 are both close to 1. The values of these two parameters in contrast with those of other models indicate a high degree of accuracy and precision of the Span–Wagner equation in ρ and c_p calculations. For the calculation of η , the VS1 model and VS4 model have similar resultant values. The AAREs are 1.819 and 1.713%, and R^2 are both close to 1. Thus, to choose the better viscosity model, their predicted η values are further compared with the experimental data in the pseudocritical region. The VS1 model (also the Vesovic model) is obtained to have a smaller AE value in the pseudocritical region overall (see **Table 3**).

In addition, for η calculation, in order to make up for the shortcomings of the Vesovic model like complexity, the free-volume viscosity model with a simple form is introduced to this study. It is worth mentioning that several parameters of the model are obtained by fitting our experimental data, which means that it concentrates more on the supercritical region. Within that region (298.15 K \leq *T* <773.15 K and seven MP a \leq *p* <14 MPa), its accuracy could be good. The equations are as follows (Liu et al., 2013):

$$\eta = \eta_0 + \Delta \eta, \tag{7}$$

where

$$\eta_0 = 4.0785 \times \frac{\sqrt{MT}}{V_c^2 \Omega^* (T^*)} F_c \text{ and}$$
 (7a)

TABLE 4 Comparison of the calculated values of each model with the experimental data.

Model	Working condition	ARE (%)	AARE (%)	AE _{max} (%)	R ²
Span–Wagner equation of state for ρ	General	-0.0089	0.3126	1.1744	0.9999
	Pseudocritical	0.0972	0.7606	1.2964	0.9988
Span–Wagner equation of state for c_p	General	0.1561	0.4149	4.3631	0.9991
	Pseudocritical	5.478	5.478	8.9908	0.9748
Vesovic model for η	General	0.9963	1.1117	3.8858	0.9974
·	Pseudocritical	-1.5482	1.7743	4.3619	0.9928
Free-volume viscosity model for η	General	-0.0007	0.4528	1.8986	0.9996
· ·	Pseudocritical	-0.1853	1.0962	2.4863	0.9976

No.	<i>T/</i> K	<i>р/</i> МРа	Absolute error/kJ·(kg·K) ⁻¹	No.	<i>T/</i> K	<i>р</i> /МРа	Absolute error/kJ·(kg·K) ⁻¹
1	304.15	7.4	5.998	8	323.15	10	0.072
2	305.15	7.6	5.040	9	323.15	12	0.064
3	305.15	7.8	1.183	10	333.15	12	0.046
4	307.15	7.8	2.088	11	333.15	14	0.040
5	307.15	8	0.918	12	343.15	14	0.013
6	311.15	9	0.277	13	353.15	14	0.026
7	313.15	9	1.229				

TABLE 6 | Specific physical properties represented by x, y, z and the n value of each model.

Model	x	У	Z	n
T (h,p)	h/(990 kJ/kg)	p/ (20 MPa)	T/(780 K)	4
$\rho(h,p)$	h/(990 kJ/kg)	p/(20 MPa)	ρ/ (480 kg/m3̂)	5
c_p (h,p)	h/ (990 kJ/kg)	p/(20 MPa)	$c_p/[2.4 \text{ kJ}/(\text{kg} \cdot \text{K})]$	5
h(p,T)	T/(780K)	p/(20 MPa)	h/ (990 kJ/kg)	4
$\lambda(\rho, T)$	$\rho/(480 \text{kg/m}^3)$	T/(780 K)	λ /[60 mW/(m · K)]	4
$\eta(\rho,T)$	ρ /(480 kg/m ³)	T/(780 K)	η/(40μPa·s)	3

TABLE 7 | Coefficients D_{ij} of T(h, p).

i	D _{i0}	D _{i1}	D _{i2}	D _{i3}	D _{i4}
0	-0.87811	3.58799	-2.79791	1.14141	-0.06412
1	4.19127	-15.32662	19.82254	-9.87656	1.20086
2	1.11324	-14.17288	41.97630	-47.01950	18.13742
3	-6.78545	41.41008	-91.03666	86.16367	-29.79073
4	3.22732	-18.53626	38.97091	-35.64303	11.99363

$$\Delta \eta = \frac{\rho l \left(\alpha \rho + \frac{pM}{\rho} \right)}{\sqrt{3RTM}} \exp \left[B \left(\frac{\alpha \rho + \frac{pM}{\rho}}{RT} \right)^{\frac{3}{2}} \right], \tag{7b}$$

where η_0 is the viscosity of dilute gas, $\Delta \eta$ is the correction term for the viscosity of dense fluid, Ω^* is the collision integral, F_c is the parameter related to dipole moment and eccentricity factor, B is the superimposed part of the free volume, M is the molar mass, $\alpha \rho$ is the energy barrier, $l = L^2/b_f$ is the L^2 is the molecular mean characteristic square length and b_f is the dissipation of potential energy. In the aforementioned equations, α , B, and l are obtained by fitting the experimental data, and the values are as follows: $\alpha = 13.34816839$, $B = 8.04589798 \times 10^{-11}$, and l = 0.03939321.

Finally, the Span–Wagner equation of state, the Vesovic model, and the free-volume viscosity model are chosen to construct a set of high-precision CO_2 property calculation models. The three selected high-precision models are further evaluated to analyze their accuracy. The experimental data (298.15 K \leq T <773.15 K and seven MP a \leq p <14 MPa) are divided into data in the general region and pseudocritical region, and the calculated data and experimental data are compared on this basis, as shown in **Table 4**.

In the general region, the ARE between experimental data and the predicted ρ calculated by the Span–Wagner equation is only -0.0089%, which is lower than the equation uncertainty ($\pm 0.03\%, \pm 0.05\%$). In the pseudocritical region, although the average relative error is slightly larger than the equation uncertainty, it is still acceptable.

For c_p calculation by the Span–Wagner equation, the ARE is only about 0.15% in the general region, while the maximum AE is 4.3631%, and this data point is near the pseudocritical region. The error around this region is generally large, possibly due to the rough partition of the pseudocritical region, which means that this part of general region is still influenced by property distortions. If data in the controversial region are ignored, the AARE then drops to about 1%. However, in the pseudocritical region, the ARE exceeds 5%, and the maximum AE is close to 9%. The absolute error between calculated and experimental values of c_p in the pseudocritical region is presented in **Table 5**. The absolute error is large in the pseudocritical region from 7 to 8 MPa, up to 5.998 kJ/(kgK),

TABLE 8 | Coefficients D_{ij} of $\rho(h, p)$.

	, .					
i	D _{i0}	D _{i1}	D _{i2}	D _{i3}	D _{i4}	D _{i5}
0	5.51141	-30.68877	66.27682	-69.14811	34.62849	-6.57949
1	-31.90696	177.99507	-367.93354	351.83442	-149.98417	20.26986
2	116.86343	-597.69806	1138.11163	-963.95423	317.08842	-10.35475
3	-84.53118	275.28642	-24.97488	-807.41087	1018.07080	-376.60235
4	-26.86934	388.49330	-1538.70210	2628.70703	-2067.58415	616.10422
5	29.46078	-254.14657	812.59139	-1232.66205	899.00831	-254.29788

TABLE 9 | Coefficients D_{ii} of c_p (h, p).

i	D _{i0}	D _{i1}	D _{i2}	D _{i3}	D _{i4}	D _{i5}
0	78.26160	-442.48715	966.46759	-1004.94280	488.94838	-85.68895
1	-505.96403	2776.18208	-5761.21712	5524.98006	-2319.84029	285.24633
2	1283.18400	-6556.89741	12144.03687	-9327.54913	1960.78820	498.64679
3	-1047.60030	3839.11215	-2038.16407	-7565.71111	11258.16620	-4449.55008
4	-8.71962	2422.70087	-12470.44310	23962.92490	-20272.26838	6368.83762
5	209.14598	-2075.59975	7226.02384	-11644.52629	8902.51944	-2618.48425

	D _{i0}	D _{i1}	D _{i2}	D _{i3}	D _{i4}
0	-1.42101	9.67303	-18.02611	15.97775	-5.19521
1	8.72986	-48.26339	97.33641	-85.65223	27.84470
2	-19.46200	104.01501	-206.00900	179.22737	-57.8060 ⁻
3	11.08749	-59.74097	119.21196	-104.40354	33.86865
4	-0.97826	5.57283	-11.64132	10.58920	-3.54565

TABLE 11 Coefficients D_{ij} of $\lambda(\rho, I)$.						
D _{i0}	D _{i1}	D _{i2}	D _{i3}	D _{i4}		
1.73319	-36.88377	202.20690	-354.37541	156.08372		
-12.30436	241.20760	-1256.10471	2134.20547	-845.03747		
33.86575	-565.82226	2821.67532	-4583.30612	1513.56125		
-35.65537	562.24250	-2684.92673	4089.19850	-928.21944		
13.23967	-199.52577	908.88515	-1255.87906	66.71366		
	D _{io} 1.73319 -12.30436 33.86575 -35.65537	Dio Di1 1.73319 -36.88377 -12.30436 241.20760 33.86575 -565.82226 -35.65537 562.24250	Dio Di1 Di2 1.73319 -36.88377 202.20690 -12.30436 241.20760 -1256.10471 33.86575 -565.82226 2821.67532 -35.65537 562.24250 -2684.92673	D _{i0} D _{i1} D _{i2} D _{i3} 1.73319 -36.88377 202.20690 -354.37541 -12.30436 241.20760 -1256.10471 2134.20547 33.86575 -565.82226 2821.67532 -4583.30612 -35.65537 562.24250 -2684.92673 4089.19850		

TABLE 12 Coefficients D_{ij} of $\eta(\rho, T)$.					
i	D _{i0}	D _{i1}	D _{i2}	D _{i3}	
0	-0.07622	0.23975	0.04682	0.28700	
1	1.34854	-0.6812	1.31995	-1.11705	
2	-0.54067	0.81125	-1.62115	1.53517	
3	0.12929	-0.32631	0.69507	-0.74812	

indicating the poor accuracy of the Span-Wagner equation in the pseudocritical region.

For the comparison of η , the ARE between the Vesovic model and the experimental result is below the model's minimum uncertainty of 1% in the general region. In the pseudocritical region, it is still lower than the uncertainty of 2% given by the model for this region. However, there are still several data points near the pseudocritical region with large AE up to 4%, suggesting that the η calculation of the Vesovic model in the pseudocritical region is defective.

The free-volume viscosity model has very high accuracy in η calculation throughout the region. The ARE is within 0.001% in the general region and only about 0.18% in the pseudocritical region.

These evaluation results offer some ideas for the improvement of the existing calculation programs. For example, in terms of CO₂ physical property calculation, the widely used program REFPROP is based partly on the Span–Wagner equation and the Vesovic model. The accuracy is relatively high as the evaluation results show. However, for the η calculation in the supercritical region, it could be more accurate than the Vesovic model used by REFPROP when using the free-volume viscosity model.

Research on Fast Calculation CO₂ Property Models

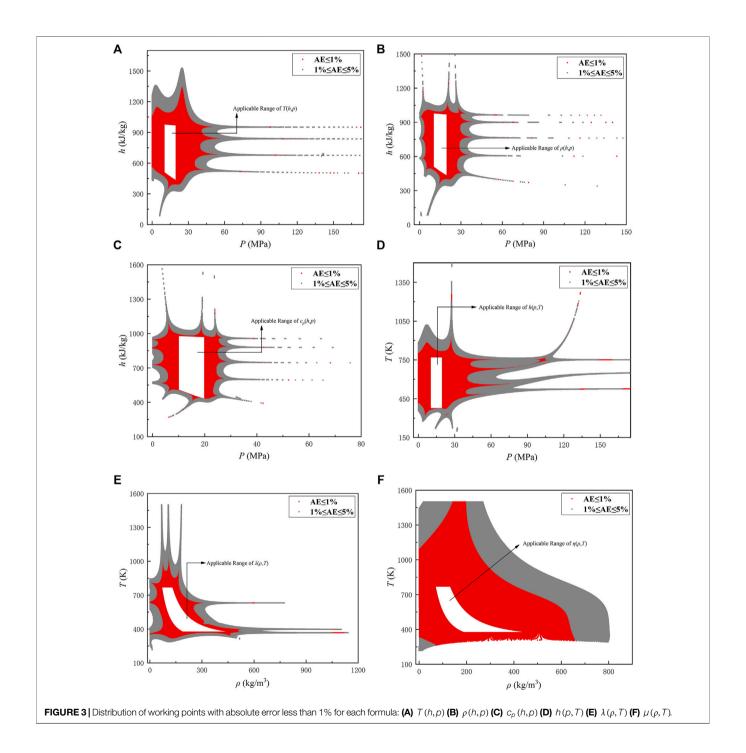
In engineering applications of SC-CO₂, such as sub-channel studies and system analysis, their computational procedures require high precision and efficiency for physical property calculations. However, the aforementioned high-precision computational models are in a relatively complicated form and have low computation speed, which greatly limits their use in engineering. Therefore, models with faster computation speeds need to be studied. In this part, two fast calculation methods are introduced: polynomial fitting and interpolation method.

Polynomial Fitted Model

In previous studies, many researchers have obtained property calculation models with high prediction accuracy in a certain

TABLE 13 | Average of relative errors between standard data and the fitted models with different maximum number of decimal digits.

Maximum number of decimal digits of coefficien	ts	0.1	0.01	0.001	0.0001	0.00001
Average of relative errors from standard data (%)	h(p,T)	64.62139	6.924697	0.706943	0.091127	0.061878
	T(h,p)	57.79155	5.247569	0.663925	0.19348	0.150844
	$\rho(h,p)$	219.8697	21.96921	1.818924	0.262511	0.140797
	$c_{p}(h,p)$	112.5528	10.94899	1.254259	0.170878	0.130468
	$\lambda(\rho,T)$	20.49005	5.110055	2.711198	2.453105	2.447142
	$\eta(\rho,T)$	17.9272	4.332717	2.351954	2.07974	2.057038



Physical Property Models of S-CO₂

TABLE 14 | Maximum absolute error between the two quick calculation methods and the standard data.

Physical quantity	AE _{max} of the fitted model (%)	AE _{max} of the interpolation method (%)		
T (h, p)	0.1390	0.1945		
ρ(h, p)	0.3000	1.1233		
$c_p(h, p)$	1.1340	2.0117		
h(T, p)	0.7530	0.7591		
$\lambda(\rho, T)$	0.6790	1.1628		
$\eta(\rho, T)$	0.0243	0.1873		

temperature and pressure range, using polynomial fitting methods (Ouyang 2011) and neural network algorithms (Zhang et al., 2019). However, these existing models do not correspond to our actual needs in terms of applicable working conditions and solving methods. Thus, a new model appropriate to related research needs to be fitted.

After trying various equation forms, we found that the fitted model obtained by the polynomial form of **Eq. 8** is the best, and the prediction accuracy is extremely high in the required temperature and pressure range.

where

$$z(x, y) = \sum_{i=0}^{n} C_i(x) y^i,$$
 (8)

$$C_{i}(x) = \sum_{j=0}^{n} D_{ij} x^{i}.$$
 (8a)

According to the actual demand of CO₂ physical properties for sub-channel studies and others, the model is designed to be applicable in the range 373.15 K \leq T <773.15 K, 10 MPa \leq p <20 MPa. Since the Span–Wagner equation and the Vesovic model are highly accurate in this range, the numerical fitting will be performed based on their calculation results.

Six fitted models are finally obtained in the form of **Eq. 8**, and the specific physical properties represented by x, y, z and the value n of each model are shown in **Table 6**. The table of the fitted coefficients D_{ij} can be seen in **Tables 7–12**. According to the comparison results of relative errors between standard data and the fitted model with different maximum numbers of decimal digits (see **Table 13**), it was found that the relative errors generally decrease quickly when the number of decimal places is taken to four digits, and the reduction in relative errors is relatively small when going from 4 to 5 decimal places.

Meanwhile, the accuracy of the result obtained by taking 5 decimal places is sufficient for our needs. Thus, the coefficients are retained to 5 decimal places in this article as shown in the tables.

The deviation between the data by the fitted equations and the combination of the Span–Wagner equation and the Vesovic model outside the given applicable range (373.15 K $\leq T$ <773.15 K, 10 MPa $\leq p <$ 20 MPa) are calculated. **Figure 3** shows that the distribution of the data points outside the given applicable range of the six equations, where the AE is less than 1 and 5%. The red and gray areas represent the range where the deviation is below 1 and 5%, respectively, and the surrounded blank area refers to the models' applicable range. Under the accuracy requirement of 1%, the six equations are still acceptable in a wide range outside the applicable region, indicating that the calculation range of fitted models is expandable.

Fast Calculation Method: Interpolation

In addition to obtaining new models by polynomial fitting, interpolation is also a commonly used method with high computational speed. The basic idea is to estimate the physical properties through linear interpolation, referring to data tables of physical properties (interpolation nodes). It is fast but strictly limited to the range given by the data table, which means that it is non-expandable.

The applicable range, the physical properties to be calculated, and the corresponding independent variables are all consistent with those in the fitted models.

The values calculated by the fitted model and the interpolation method are compared with the predicted data calculated by the Span–Wagner equation and the Vesovic model. **Table 14** presents the results: the maximum AE of them in the applicable range. The maximum AEs are all very small, within 2.02% (1% if disregarding c_p calculation and the λ interpolation calculation). The results indicate that both fitted models and the interpolation method have high accuracy. It can also be found that the fitted models generally have smaller maximum AE, which means that the accuracy of the fitted model is slightly higher than that of the interpolation method. However, considering the simplicity of the linear interpolation method compared to the fitted model calculation, the accuracy of the interpolation method is already satisfactory.

Physical quantity	Calculation time (ms)			CTR		
	Fitted model	Interpolation method	Span–Wagner + Vesovic model	Fitted model	Interpolation method	
T (h, p)	28.3	3.4	806068	28483	237079	
ρ(h, p)	32.3	2.8	786252	24342	280804	
$c_p(h, p)$	31.7	2.3	808804	25514	351654	
h(T, p)	25.4	2.9	1125590	44315	388134	
$\lambda(\rho, T)$	25.4	2.8	841422	33127	300508	
$\eta(\rho, T)$	13.1	2.5	776509	59275	310604	

Computing Time Comparison With High-Precision Models

The fitted model, interpolation method, and the combination of the Span–Wagner equation and the Vesovic model are separately used to calculate the physical properties under 10⁵ working conditions in the applicable range, and the calculation time is compared as shown in **Table 15**, where CTR is the computing time ratio:

 $CTR = \frac{\text{Calculation time of the Span - Wagner equation}}{\text{Calculation time of the fitted model (or the interpolation method)}}.$ (9)

The computation time of the Span–Wagner equation + Vesovic model is four orders of magnitude more than that of the fitted models. and five orders of magnitude more than that of the interpolation method, showing the superiority of the two fast calculation models (methods) in terms of efficiency. In a crosssectional comparison of the two fast calculation models (methods), it is found that the computation time of the interpolation method is only 1/10 of that of the fitted models. However, given that the overall accuracy of the interpolation method is slightly inferior to that of the fitted models, and the applicability is strictly limited to the given data tables, the choice of which calculation method to use should be combined with the reality in engineering applications.

CONCLUSION

Based on the experiment of CO_2 physical property measurement in the range of temperatures from 298.15 to 773.15 K and pressures from 7 to 14 MPa, some commonly used CO_2 property models are evaluated, and a set of relatively highprecision models is selected, including the Span–Wagner equation of state, the Vesovic model, and the free-volume

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viscosity model. Comparing ρ (the Span–Wagner equation), c_p (the Span–Wagner equation), and η (the Vesovic model, the free-volume viscosity model) calculation results with experimental data, the AARE is 0.3126, 0.4149, 0.9963, and -0.0007, respectively, out of the pseudocritical region.

Meanwhile, to achieve the goal of fast calculation of CO₂ properties, polynomial fitting and interpolation methods are introduced in the present work. The results show that both methods have high accuracy with a relative error of less than 2.02% (1% if disregarding c_p calculation and the λ interpolation calculation) and can improve the calculation speed by 4-5 two magnitudes compared with the Span–Wagner equation + Vesovic model. Additionally, the fitted model has a large expandable region under the 1% accuracy requirement, and the computation time of the interpolation method is only 1/10 of that of the fitted models.

It can be concluded that the fitted models and interpolation method can successfully and quickly predict CO_2 properties within the range of temperatures from 373.15 to 773.15 K and pressures from 10 to 20 MPa. The results are significant for research of the SC-CO₂ Brayton cycle.

DATA AVAILABILITY STATEMENT

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

AUTHOR CONTRIBUTIONS

LM: Software, Formal Analysis, Writing - Original Draft; ML: Methodology, Writing - Review and Editing; DX: Software, Formal Analysis; MH and XL: Investigation, Validation; YH: Conceptualization, Funding Acquisition, Supervision.

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