



First-Principles Design of Rutile Oxide Heterostructures for Oxygen Evolution Reactions

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The oxygen evolution reaction (OER) plays a key role in the determination of overall water-splitting rate. Lowering the high overpotential of the OER of transition metal oxides (TMOs), which are used as conventional OER electrocatalysts, has been the focus of many studies. The OER activity of TMOs can be tuned via the strategic formation of a heterostructure with another TMO substrate. We screened 11 rutile-type TMOs (i.e., MO_2 ; $\text{M} = \text{V}, \text{Cr}, \text{Mn}, \text{Nb}, \text{Ru}, \text{Rh}, \text{Sn}, \text{Ta}, \text{Os}, \text{Ir}, \text{and Pt}$) on a rutile (110) substrate using density functional theory calculations to determine their OER activities. The conventional volcano approach based on simple binding energies of reaction intermediates was implemented; in addition, the electrochemical-step symmetry index was employed to screen heterostructures for use as electrode materials. The results show that RuO_2 and IrO_2 are the most promising catalysts among all candidates. The scaling results provide insights into the intrinsic properties of the heterostructure as well as materials that can be used to lower the overpotential of the OER.

Keywords: density functional theory, oxygen evolution reaction, rutile-type oxide, heterostructure, scaling relation

INTRODUCTION

Green hydrogen production remains a challenge that must be overcome to achieve a hydrogen economy (Turner, 2004). Water electrolysis is one of the approaches toward hydrogen eco-friendliness. It is based on the hydrogen evolution reaction (HER) and oxygen evolution reaction (OER) as cathodic and anodic reactions, respectively. However, the slow kinetics of the OER limit the commercialization of this approach. Oxide systems based on Ir or Ru (i.e., IrO_x or RuO_x) are known to have the best OER performances (Lee et al., 2012; Frydendal et al., 2014; Suen et al., 2017). However, these materials are expensive. Therefore, many studies have been conducted to identify a cost-effective alternative oxide material with high activity. For example, transition-metal substitution (García-Mota et al., 2011) and the introduction of oxygen vacancies (Xiao et al., 2020) have been studied as means to control the compositions of expensive materials. However, transition-metal substitution and oxygen vacancies only locally affect the active sites.

To activate the metal oxide, support materials are mixed with the oxide to enhance the electrical conductivity and activity of the electrocatalytic reaction (Kumar et al., 2016; Qingxiang et al., 2018; Bu et al., 2019; Lu et al., 2019). In addition, research on heterostructures is conducted with the goal of activating the material on the substrate. Modifications are mainly based on the strain effect and charge transfer due to the formation of an interface between the support material and catalyst. For instance, the heterostructure of $\text{La}_{0.5}\text{Sr}_{0.5}\text{CoO}_{3-\delta}$ and MoSe_2 induces the phase transition of MoSe_2 from the 2H to the 1T phase (Oh et al., 2019). In addition, the charge transfer from the Co ion to Mo ion improves the electrochemical activity. The results of another study showed that the heterostructure of IrO_x and SrIrO_3 outperforms iridium or ruthenium oxide systems (Seitz et al.,

2016). Based on density functional theory (DFT) calculations, IrO_3 or anatase IrO_2 motifs are formed during strontium leaching in the outermost surface layers of SrIrO_3 and contribute to the high activity. Another strength of the heterostructure is that it can possibly contribute to reducing the use of precious metals by replacing them into low-cost metals while maintaining similar intrinsic activity of each active site (Esposito et al., 2010; Zhou et al., 2014; Wang et al., 2015; Jin et al., 2016). Rutile (TiO_2) is a well-known substrate material for metal oxide systems because of its high structural stability and simple structure (Hanaor et al., 2012). It is suitable for the growth of oxide films, facilitating the fabrication of heterostructures with various metal oxides. In addition, rutile affects the catalytic activity of metal oxides and has a high cost effectiveness (Seitsonen and Over, 2010; Stacchiola et al., 2013; Wei et al., 2015; Sun et al., 2016; Li et al., 2017). However, the effect of the rutile support on the activity remains unclear due to difficulties with respect to the characterization of the heterostructure during experiments (Stacchiola et al., 2013).

In this study, we performed DFT calculations to theoretically investigate the effect of the TiO_2 substrate on the heterostructure and thus on the OER. We screened the OERs of heterostructures consisting of various rutile-type metal oxides (i.e., VO_2 , CrO_2 , MnO_2 , NbO_2 , RuO_2 , RhO_2 , SnO_2 , TaO_2 , OsO_2 , IrO_2 , and PtO_2) and a TiO_2 substrate. The results show that these rutile heterostructures follow the universal scaling relations of metals and oxides; however, the binding strengths of the O^* intermediates increase due to the TiO_2 substrate. The volcano plot and electrochemical-step symmetry index (ESSI) show that the RuO_2 and IrO_2 are the closest to an ideal catalyst. The results of our computational screening provide insights into the effects of support materials on electrocatalytic reactions.

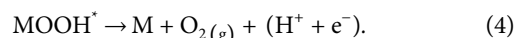
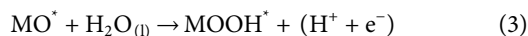
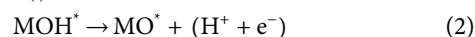
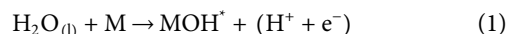
METHODS

Spin-polarized DFT calculations were performed with the projector-augmented wave (PAW) method (Blochl, 1994) and Vienna Ab initio Simulation Package (VASP) (Kresse and Furthmüller, 1996). The electron exchange–correlation energy was treated within the generalized gradient approximation (GGA) and the Perdew–Burke–Ernzerhof functional (Perdew et al., 1997). To determine the trend of the OER activity, the DFT + U method within Liechtenstein’s approach (Anisimovdag et al., 1997) was used and the following correction parameters were employed: $U = 4.95$ eV for Ti, 2 eV for V, 7.15 eV for Cr, 6.63 eV for Mn, 3.32 eV for Nb, 6.73 eV for Ru, 5.97 eV for Rh, 5.91 eV for Ir, and 6.25 eV for Pt (Xu et al., 2015). The energy cutoff for the plane wave basis set was set at 520 eV. The geometry was optimized using the residual minimization method and the direct inversion in the iterative subspace method (RMM–DIIS) algorithm until the net force on each atom was below $0.02 \text{ eV} \cdot \text{Å}^{-1}$, and the total energy was 10^{-6} eV per atom. Dipole slab corrections were also applied to all slab model calculations. The k -point sampling of the Brillouin zone was done with a $4 \times 4 \times 1$ for bulk calculations and $6 \times 6 \times 8$ for slab calculations.

RESULTS AND DISCUSSION

To theoretically investigate the OER activities of rutile-type heterostructures, we considered the heterostructures of 11 rutile-type oxides (i.e., MO_2 , where $M = \text{V, Cr, Mn, Nb, Ru, Rh, Sn, Ta, Os, Ir, and Pt}$) with a TiO_2 substrate. To accurately illustrate the OER activity, the following magnetic structures were used for all calculations according to Xu et al.’s work: nonmagnetic (NM) for TiO_2 , NbO_2 , RuO_2 , RhO_2 , IrO_2 , and PtO_2 , and ferromagnetic (FM) for CrO_2 and MnO_2 (Xu et al., 2015). Furthermore, the ground state magnetic configurations of the other candidate models were identified as FM for VO_2 and NM for SnO_2 , TaO_2 , and OsO_2 , respectively (Supplementary Figure S1). Surface models, that is, six-layer stoichiometric slabs, were built by using the 2×1 supercell of the optimized unit cell. The vacuum of $\sim 15 \text{ Å}$ was applied in the (110) direction, which is the most stable facet of rutile-type oxides (Figure 1A; Kung, 1989). Note that a tri-layer consisting of oxygen–metal–oxygen (O–M–O) atomic layers was considered to be a single layer in our slab models. The top four layers of the slab models were allowed to fully relax. The two layers at the bottom were fixed to represent the bulk state. For heterostructures, the top three layers of the TiO_2 slab model were replaced with MO_2 layers while maintaining the cell dimensions of the TiO_2 substrate. The coordinatively unsaturated sites (CUS) of the metal atoms at the top surface were considered to be the adsorption sites for each intermediate (i.e., OH^* , O^* , and OOH^*) for the OER. To examine the oxygen coverage effect on the OER activity, the pristine and fully O^* covered surfaces were representatively compared (i.e., denoted as 2O_b and $2\text{O}_b 2\text{O}_c$, where the subscripts “b” and “c” represent the bridge sites and CUS, respectively).

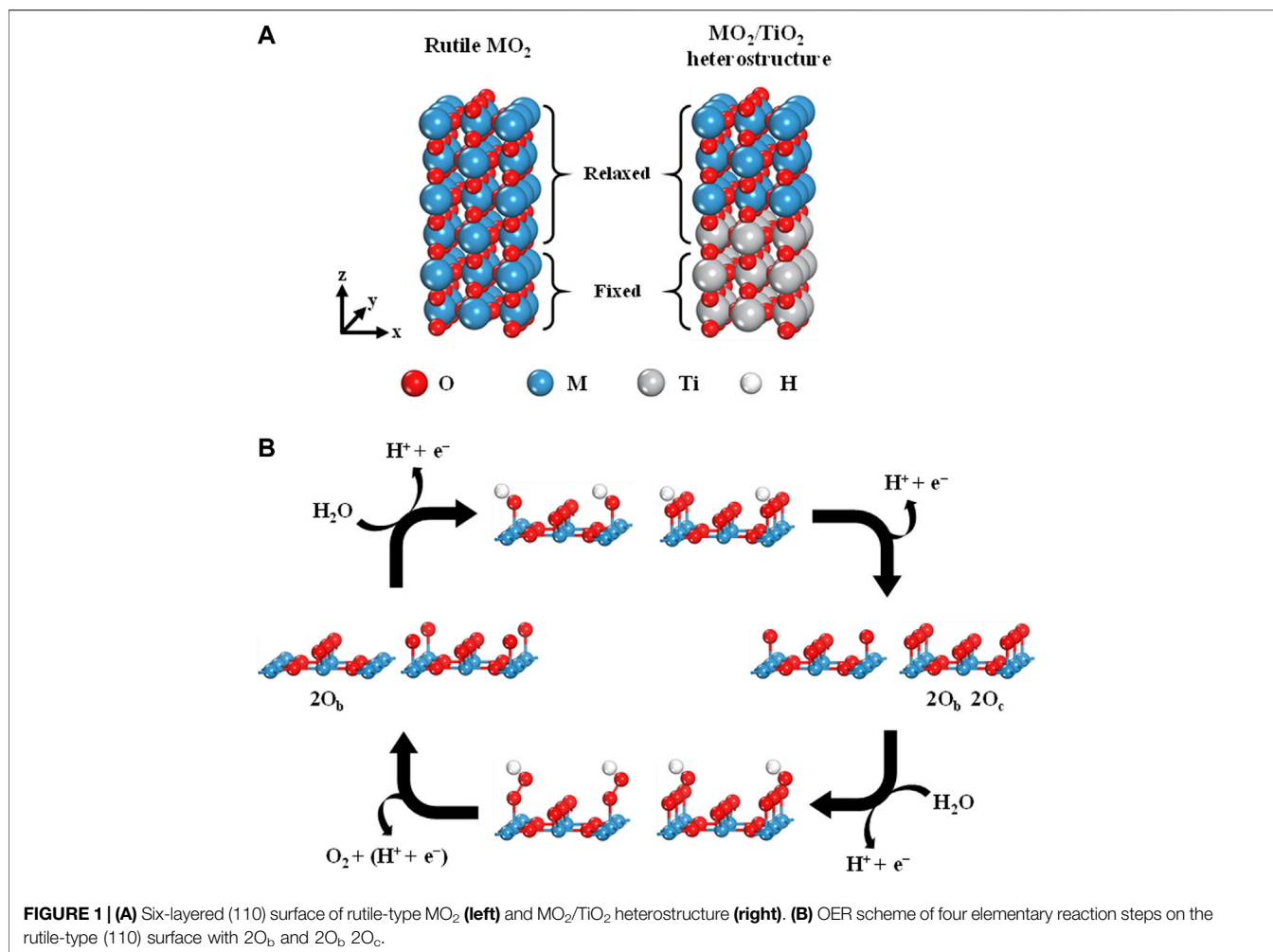
The overall OER process consists of four elementary steps involving proton-coupled electron transfer (PCET; Figure 1B) (Hammes-Schiffer, 2015; Costentin and Savéant, 2017). In this study, we followed the conventional four-electron pathways with respect to the OER of rutile heterostructures, which can be described as follows:



The adsorption free energy was calculated using the following equation:

$$\Delta G = \Delta E + \Delta ZPE - T\Delta S - eU + k_B T \ln a_{\text{H}^+}, \quad (5)$$

where ΔE is the binding energy of each reaction intermediate; ΔZPE is the zero-point vibrational enthalpy; $-T\Delta S$ is the entropic correction at room temperature, $-eU$ is the energy shift by the electrode potential, where U is the electrode applied potential relative to the standard hydrogen electrode (SHE) and e is the elementary charge transferred; and $k_B T \ln a_{\text{H}^+}$ is used as correction for the free energy of H^+ ions, where a_{H^+} is the activity of the proton. In this study, we considered the standard conditions for Gibbs free energy calculations



(i.e., $T = 298$ K and $\text{pH} = 0$). Free energy correction values were taken from Valdes et al.'s work (Valdes et al., 2008). The binding energy for each reaction intermediate (i.e., ΔE_{OH^*} , ΔE_{O^*} , and ΔE_{OOH^*}) was calculated as follows:

$$\Delta E_{\text{OH}^*} = E(\text{OH}^*) - E(*) - [E(\text{H}_2\text{O}) - 0.5E(\text{H}_2)] \quad (6)$$

$$\Delta E_{\text{O}^*} = E(\text{O}^*) - E(*) - [E(\text{H}_2\text{O}) - E(\text{H}_2)] \quad (7)$$

$$\Delta E_{\text{OOH}^*} = E(\text{OOH}^*) - E(*) - [2E(\text{H}_2\text{O}) - 1.5E(\text{H}_2)], \quad (8)$$

where $E(\text{OH}^*)$, $E(\text{O}^*)$, and $E(\text{OOH}^*)$ represent the total energies of the slab models for each adsorbate, $E(*)$ is the total energy of the bare slab, and $E(\text{H}_2\text{O})$ and $E(\text{H}_2)$ represent the total energies of an isolated water molecule and hydrogen gas, respectively. The differences in the Gibbs free energy (ΔG) of each step were calculated as follows:

$$\Delta G_1 = \Delta G_{\text{OH}^*} \quad (9)$$

$$\Delta G_2 = \Delta G_{\text{O}^*} - \Delta G_{\text{OH}^*} \quad (10)$$

$$\Delta G_3 = \Delta G_{\text{OOH}^*} - \Delta G_{\text{O}^*} \quad (11)$$

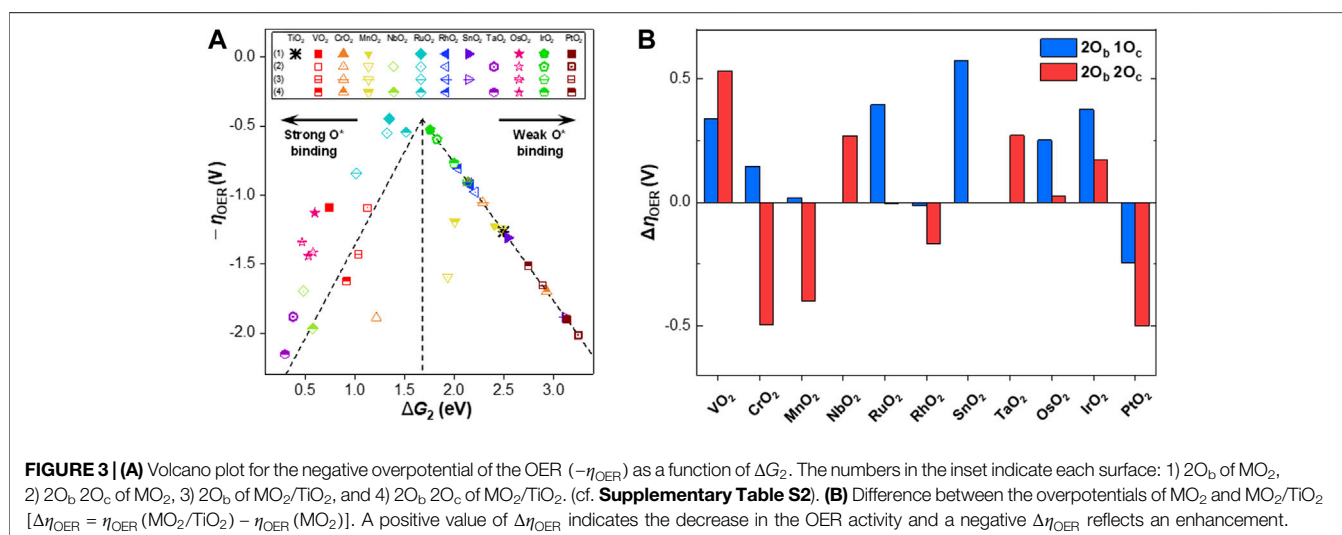
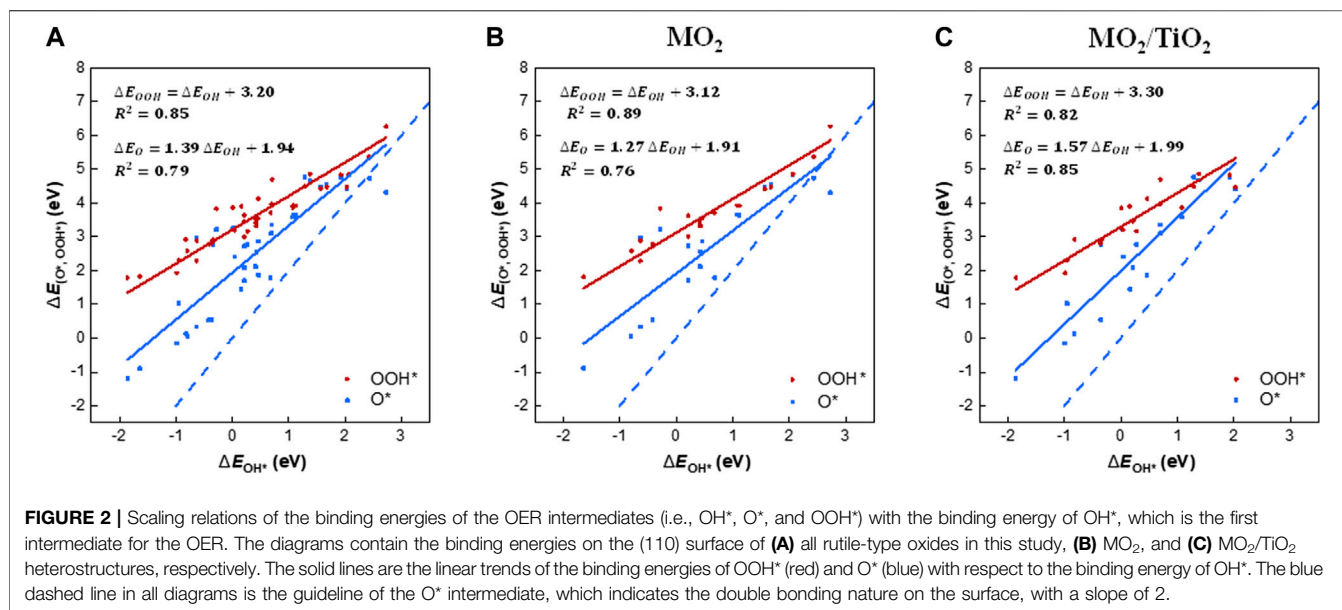
$$\Delta G_4 = 4.92 - \Delta G_{\text{OOH}^*} \quad (12)$$

Finally, the theoretical overpotential of the OER (η_{OER}) can be calculated as follows:

$$\eta_{\text{OER}} = \max([\Delta G_1, \Delta G_2, \Delta G_3, \Delta G_4]/e) - U_{\text{eq}}, \quad (13)$$

where U_{eq} indicates the equilibrium potential of the OER (i.e., 1.23 V vs. reversible hydrogen electrode). **Eqs 1–13** assume that thermodynamics of the reaction steps is a valid descriptor for reaction kinetics based on the Brønsted–Evans–Polanyi (BEP) relations (Vojvodic et al., 2011), which refer to that a free-energy change in the transition states (kinetics) follows a change in the reaction heat (thermodynamics). Note that the connection between thermodynamics and kinetics is not always established (Kuo et al., 2017; Kuo et al., 2018), which requires further kinetic experiments or microkinetic modeling based on the activation barrier calculations for all plausible transition states.

We examined the scaling relations between the adsorption free energies of the reaction intermediates (i.e., OH^* , O^* , and OOH^* ; **Supplementary Table S2**) for all rutile-type catalysts of interest (**Figure 2A**). The binding energies of OOH^* and OH^* are linearly correlated, with an offset of 3.20 eV. Note that our scaling trend is

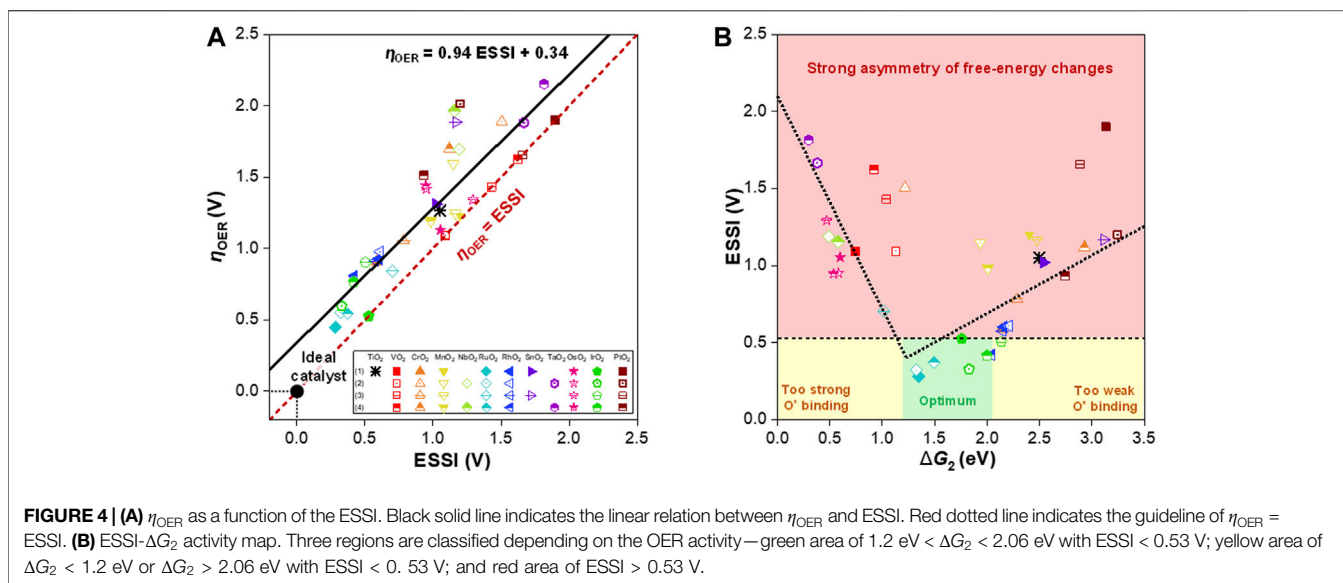


similar to the “universal” scaling relation reported by Man et al. [i.e., $\Delta E_{OOH^*} = \Delta E_{OH^*} + 3.20$ (± 0.20 eV)] (Man et al., 2011), implying that all rutile-type systems, including heterostructures, follow the conventional scaling relations for metals and oxide surfaces. Based on the best fit, 68% of the points are within ± 0.35 eV (1σ) and 95% are within ± 0.70 eV (2σ). The scaling relationship between ΔE_{O^*} and ΔE_{OH^*} exhibits a slope of 1.39 (blue solid line), that is, it is much less steep than the slope of two (blue dashed line), which is the indicator line of the double bonding nature of O* (Rossmeisl et al., 2007).

To analyze the substrate effect of TiO₂ on the heterostructures, we divided the scaling relations into two groups, that is, (110) surfaces of MO₂ and MO₂/TiO₂ heterostructures (Figure 2B,C). Notably, the binding energies of OOH* and O* on the MO₂/TiO₂ (110) surface are more concentrated than those of MO₂. The scaling relation between OOH* and OH* of MO₂/TiO₂ (red solid

line) shows an increased intercept by 0.18 compared with the corresponding scaling relation of MO₂, implying slightly weakened interactions between OOH* and the surface. Note that the scaling relations between the binding strengths of O* and OH* species on MO₂ and MO₂/TiO₂ (blue solid lines) apparently differ. The slopes of MO₂/TiO₂ (110) heterostructures are closer to 2 (i.e., double bonding nature of O*) than those of MO₂ (110) surfaces (Rossmeisl et al., 2007). This is due to the intensified binding strengths of O* intermediates in heterostructures compared with MO₂ surfaces (Divanis et al., 2020). This indicates that the TiO₂ substrate generally stabilizes the O* intermediates, which might lead to a decrease in η_{OER} of rutile oxides involving weakly bound O* intermediates.

Based on the scaling relations between the binding energies of the reaction intermediates, a volcano plot was constructed, as



shown in **Figure 3A**. We chose $\Delta G_{\text{O}^*} - \Delta G_{\text{OH}^*}$ (denoted as ΔG_2) as a descriptor, which is commonly used to predict the OER activity of the $4 e^-$ reaction (Man et al., 2011; Krishnamurthy et al., 2018). The plot shows that the O^* bonds of PtO_2 , OsO_2 , TaO_2 , and RhO_2 on the TiO_2 substrate are stronger than those without the TiO_2 substrate, and the activities change along the volcano curve. In the cases of both oxygen coverages of $\text{PtO}_2/\text{TiO}_2$ and $\text{RhO}_2/\text{TiO}_2$, ΔG_2 , which is the potential-determining step for each system, was reduced, exhibiting the improved OER activities (i.e., smaller $|\eta_{\text{OER}}|$) compared to those of PtO_2 and RhO_2 , respectively. In the cases of OsO_2 , TaO_2 , and NbO_2 , the bonds of the O^* intermediates are too strong in the presence of the TiO_2 substrate, resulting in a decreased OER activity.

Similar to the previous studies (Rossmeisler et al., 2007; Man et al., 2011), RuO_2 and IrO_2 were both identified as the active OER catalysts among the MO_2 candidates (i.e., $\eta_{\text{OER}} < \sim 0.5 \text{ V}$, **Figure 3A**). In the case of 2O_b , 2O_c of RuO_2 , the OER activity maintains similarity in the presence of TiO_2 substrate, which is located near the top of the volcano. In general, the TiO_2 substrate stabilizes the adsorbates (i.e., OOH^* , O^* , and OH^*) on 2O_b , 2O_c of RuO_2 , which appears as stronger adsorption free energies by $0.54\text{--}0.70 \text{ eV}$ (**Supplementary Table S1**). Nevertheless, η_{OER} is almost unchanged (i.e., $\Delta\eta_{\text{OER}} = -0.01 \text{ V}$) because the free energy change in the potential determining step (i.e., $\Delta G_{\text{OOH}^*} - \Delta G_{\text{O}^*}$) remains similar (**Supplementary Table S2**). In addition, it is also noteworthy that η_{OER} for 2O_b , 2O_c of IrO_2 is relatively similar, still exhibiting higher activity than other candidates except RuO_2 .

The changes in η_{OER} induced by the TiO_2 substrate are summarized in **Figure 3B**. As the value becomes more negative, the activity of MO_2/TiO_2 improves compared with that of MO_2 . Notably, 2O_b , 2O_c of CrO_2 , MnO_2 , RuO_2 , RhO_2 , and PtO_2 shows an enhanced activity. In addition, 2O_b of RhO_2 and PtO_2 also shows an improved activity on the TiO_2 substrate. However, regardless of the coverage, the TiO_2 substrate decreases the OER activity of VO_2 , OsO_2 , and IrO_2 . In particular, 2O_b , 2O_c

of VO_2 and 2O_b of SnO_2 exhibit significantly increased overpotentials ($\Delta\eta_{\text{OER}}$ of 0.53 and 0.57 V , respectively).

To determine the promising free-energy regime for the OER, we further analyzed the free energies of the steps in terms of the ESSi, as suggested by Calle-Vallejo and coworkers (Govindarajan et al., 2018). The ESSi is an energetic descriptor that indicates the degree of similarity with an ideal catalyst, where all OER steps are perfectly symmetric at 1.23 eV . The ESSi is defined by the following equation and is only applied to steps with $\Delta G_i (i = 1, 2, 3, 4) \geq 1.23 \text{ eV}$ (denoted as ΔG_i^*):

$$\text{ESSi} = \frac{1}{n} \sum_i^n (\Delta G_i^* - 1.23). \quad (14)$$

Figure 4A shows η_{OER} of each model as a function of the ESSi, representing a good linear correlation. 2O_b of RuO_2 is the closest model to the ideal catalyst, followed by 2O_b , 2O_c of RuO_2 and $\text{RuO}_2/\text{TiO}_2$. The 2O_b of PtO_2 appears on the $\eta_{\text{OER}} = \text{ESSi}$ (red dashed line) and has a null bar with respect to the ESSi, which means that only a single step is greater than 1.23 V (i.e., $\Delta G_2 = 3.13 \text{ eV}$). In the presence of the TiO_2 substrate, ΔG_2 for 2O_b of PtO_2 is reduced, leading to a decrease of the ESSi and η_{OER} . Meanwhile, 2O_b , 2O_c of PtO_2 lies relatively far from the line of $\eta_{\text{OER}} = \text{ESSi}$, which has a wide bar with respect to the ESSi, and thus corresponds to a good candidate for optimization (Govindarajan et al., 2018). In practice, the TiO_2 substrate on 2O_b , 2O_c of PtO_2 works for enhancing the OER activity with a decrease in both η_{OER} and ESSi.

Based on the scaling relation between the ESSis and η_{OER} 's, we determined a promising group of OER catalysts using the ESSi- ΔG_2 activity map introduced by Exner (Exner, 2019; **Figure 4B**). The activity map is used to identify OER candidates by adjusting the ESSi threshold. The ESSi threshold ($< 0.53 \text{ V}$) was determined from **Supplementary Figure S2** by applying η_{OER} of IrO_2 , which is the conventional OER catalyst. The free-energy regime was set to $1.20 \text{ eV} < \Delta G_2 < 2.06 \text{ eV}$ by

applying the standard deviation of ± 0.43 eV of the scaling relationship between ΔG_2 and ΔG_3 (**Supplementary Figure S3**) on 1.63 eV. The median value of 1.63 eV in the free-energy regime is determined by assuming the threshold electrode potential as the point where the experimental Tafel slope exceeds 59 mV/dec (i.e., $\eta_{\text{OER}} > 0.4$ V), accounting for kinetics (Exner and Over, 2019). The green-colored area includes 2O_b of RuO_2 , 2O_b , 2O_c of RuO_2 , 2O_b , 2O_c of $\text{RuO}_2/\text{TiO}_2$, 2O_b of IrO_2 , 2O_b , 2O_c of IrO_2 , 2O_b , 2O_c of $\text{IrO}_2/\text{TiO}_2$, and 2O_b , 2O_c of RhO_2 , which are the most promising candidates. Particularly, a portion of promising candidates for the OER (i.e., RuO_2 , $\text{RuO}_2/\text{TiO}_2$, IrO_2 , $\text{IrO}_2/\text{TiO}_2$, and $\text{RhO}_2/\text{TiO}_2$) are more evidently classified on the ESSI- ΔG_2 map (**Figure 4B**), while they are somewhat deviated from the apex in the volcano plot (**Figure 3A**). This implies that conventional volcano analysis does not guarantee to predict the most active OER catalyst, whereas the ESSI- ΔG_2 activity map, which is based on the kinetic scaling relations, is a more robust descriptor for the OER (Exner, 2019; Exner and Over, 2019). The 2O_b , 2O_c surface of IrO_2 exhibits high symmetry of the reaction steps comparable to RuO_2 . On the 2O_b and 2O_b , 2O_c surfaces of $\text{RhO}_2/\text{TiO}_2$, the values of ΔG_2 decrease compared to those of RhO_2 . Accordingly, the free-energy differences of the other steps are adjusted to compensate for the decrease in ΔG_2 , while maintaining ΔG_2 as a potential-determining step with reduced η_{OER} (cf. **Supplementary Table S2**). Next, 2O_b of $\text{IrO}_2/\text{TiO}_2$ belongs to the yellow-colored region, which needs to be reconsidered for activity optimization (Govindarajan et al., 2019). The candidates in the region of ESSI > 0.53 V (i.e., red-highlighted region in **Figure 4B**) are classified as an inferior group, showing poor OER activities over all ranges of ΔG_2 due to the highly asymmetric free-energy changes of the OER intermediates (Exner, 2019).

CONCLUSION

We screened a variety of rutile oxide heterostructures based on a TiO_2 substrate using scaling relations and relevant descriptors to identify a promising OER catalyst. The scaling relations between the reaction intermediates demonstrate that the rutile-type MO_2 heterostructures follow the universal scaling relationship of metal oxides. In addition, the TiO_2 substrate stabilizes the O^* bond on the (110) metal oxide surface. Based on the conventional volcano plot, RuO_2 and IrO_2 are found to be highly active OER catalysts as previously reported. Based on the ESSI descriptor, the superior activity of the RuO_2 can be

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attributed to the high symmetry of the reaction steps. Furthermore, based on the ESSI- ΔG_2 activity map, the candidates can be classified into an optimum group, a second promising group of OER catalysts with potential for optimization, and an inferior group that does not require particular attention.

The results of our computational screening using the scaling relations of rutile-type heterostructures provide valuable insights into the effect of the support material on the overpotential and thus guidelines for the design of a promising OER catalyst.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/**Supplementary Material**, further inquiries can be directed to the corresponding authors.

AUTHOR CONTRIBUTIONS

All calculations and data analyses were performed by HYL. All authors contributed to the discussion of the results and wrote and revised the manuscript. SKK and GYJ supervised the project.

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SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fenrg.2021.606313/full#supplementary-material>.

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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