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Calculation method for internal force and deformation of the prestressed I-beam on the elastic foundation

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The elastic foundation beam theory has been widely used in civil engineering, including railway, tunnel, and building foundations. With the development of fabricated structures, more elastic foundation beams need to be prestressed. In order to explore the frame foundation beam with the fabricated anchor-cable in the slope reinforcement project, in this article, prestress is applied to both ends of the beam. Then, according to the calculation method of internal force and deformation of the beam under concentrated force and the equivalent load theory of the prestressed structure, three new methods: the finite difference method (FDM) of the Euler-Bernoulli beam on the Winkler foundation, FDM of the Euler-Bernoulli beam on the Pasternak foundation, and theoretical analytical solution of the Timoshenko beam on the Winkler foundation, are deduced to calculate internal force and deformation under prestress force and concentrated force, respectively. Typical calculation parameters are selected for design and verification via the three new methods and GEO5, respectively. The results show that the calculated values of the three new methods are basically consistent with those calculated using GEO5 software, which verifies the feasibility of the three new methods.

KEYWORDS

internal force and deformation, prestressed I-beam, elastic foundation, calculation method, Winkler foundation, Pasternak foundation, Timoshenko beam

Introduction

The cast-in-suit anchor frame beam was introduced in China in the 1990s by the China Academy of Railway Sciences (CARS). Over the past decades, a relatively complete set of cast-in-suit anchor frame beams has been formed, including design and construction (Wu et al., 2011; Zhang et al., 2022). Moreover, it has been widely employed in slope reinforcement of railway, highway, and water conservancy projects and other infrastructure engineering (Song and Zhou, 2004; Xia and Zhou, 2006; Wu and

Huang, 2008; Li et al., 2019) and has acquired good reinforcing effects (Cao et al., 2016; Zhao et al., 2018; Zhao et al., 2021). However, some problems have also been exposed in engineering applications: 1) longer construction periods; 2) instable construction quality and insufficient durability; 3) high labor costs and personnel security issues. Instead, prestressed precast anchor frame beams, with a higher degree of mechanization, a lower risk of steep slope operation, better molding quality, and reinforcing while excavating, become a new direction of slope reinforcement technology.

Compared with the traditional cast-in-suit anchor frame beam, the pre-tensioning method or post-tensioning method is adopted to apply prestress to the prefabricated beam body before applying anchor loads to the prestressed precast anchor frame beam. The design and construction methods of traditional cast-in-suit anchor frame beams are no longer applicable. As a new slope reinforcement technology, the first task for the design and construction of prestressed precast anchor frame beams is how to take the prestressed loads of the beam and anchor cable into consideration in the design calculation and then calculate the internal force and deformation of the beam properly. There are two main types of elastic foundation beam models regarding the elastic foundation. One is the Winkler foundation model, which is assumed to have displacement immediately in the load zone, while no displacement in the unload area. The other is the Pasternak foundation model. In this model, it is assumed that the spring elements connect with one layer that is incompressible and vertical and only appears as transverse shear deformation, resulting in the shearing effect among each spring element in the Winkler elastic foundation model. Two main theories are applied to the beam in the elastic foundation beam model. One is the Euler-Bernoulli beam theory, which ignores transverse shear deformation, and the cross-section of the beam is perpendicular to the central axis before and after deformation. The other is the Timoshenko beam theory, which takes the transverse shear deformation into account. Numerous research results have been gained on the calculation of internal force and deformation of cast-in-suit frame foundation beams with the prestressed anchor-cable under load. The main calculation methods include the upside beam method (Ma et al., 2004), the finite difference method (FDM) (Li et al., 2009) and the finite element method (FEM) (Liu et al., 2004) of the Euler-Bernoulli beam on the Winkler elastic foundation, and the finite element method of the Timoshenko beam on the Winkler elastic foundation (Xia et al., 2010). A large number of field tests (Shi et al., 2019; Ye et al., 2019; Li et al., 2020) and measurements have verified the validity of the elastic foundation beam theory on internal force and deformation of cast-in-suit frame foundation beams with the prestressed anchor-cable.

As a novel slope reinforcement technology, few studies have been conducted on the internal force and deformation of the reinforced concrete beam on an elastic foundation under the influence of anchorage force and beam prestress (Zhu et al., 2004; Qin et al., 2008). Some research institutions have carried out numerical simulation analysis of foundation beams with prefabricated prestressed anchor cables (Liu, 2011; Xiao et al., 2013; Sheng and Lu, 2021) and applied that to engineering practices. Yan et al. (2012) and Wang (2012) studied the stress behavior of precast anchor piers with ANSYS. Zhang et al. (2015) and Yang et al. (2015) successfully employed the self-developed precast anchor pier for landslide treatment. Cheng and Wang, (1997) and Wu and Sun, (1999) described the application of the prefabricated cross beam in the reinforcement of the Three Gorges and Yunnan slopes. Leng and Liao, (2007) studied the numerical analysis and engineering application of the prefabricated cross-shaped laminated plate. It can be seen that most precious studies focused on the engineering applications and complex numerical calculations of the foundation beam with the prefabricated prestressed anchor cable, and little attention has been devoted to the simple and accurate calculation formula of stress and deformation of I-beams decomposed by prefabricated cross beams, which is difficult to be applied to the engineering design.

In this article, the most basic structural unit of the prefabricated frame beam—I-beam is taken as an example. Based on the calculation method of the internal force and deformation of the beam under concentrated force and the equivalent load theory of the prestressed structure, three methods, namely, FDM of the Euler–Bernoulli beam on the Winkler foundation, FDM of the Euler–Bernoulli beam on the Pasternak foundation, and the theoretical analytical solution of the Timoshenko beam on the Winkler foundation, are deduced to calculate the internal force and deformation under prestress and concentrated force, respectively. Finally, the rationality of these three methods is verified by GEO5.

Methodology

Theoretical model of the Timoshenko beam on the Winkler foundation

The Winkler foundation is a simple model widely used for decades to account for the surface response of elastic bodies. It assumes that the normal surface displacement on the surface and restoring force per unit area are related linearly and completely locally (Yin, 2000a; Yin, 2000b). The Timoshenko beam theory based on the rigid cross-section assumption considers the influence of shear deformation on the additional deflection of the beam. Taking the effect of double-layer reinforcement, prestressed reinforcement of prefabricated beams, and anchorage force of the anchor cable into consideration, a theoretical model of the Timoshenko beam on the Winkler foundation can be obtained to calculate the internal force and deformation of the prestressed I-beam.

The governing differential equation of the Timoshenko beam can be expressed as

$$D\frac{d^4\omega}{dx^4} - \frac{k_s D}{C}\frac{d^2\omega}{dx^2} + k_s\omega = q - \frac{D}{C}\frac{d^2q}{dx^2}.$$
 (1)

If the displacement of beam ω is known, the rotation angle φ can be obtained as follows:

$$\varphi = \left(1 - \frac{Dk_s}{C^2}\right)\frac{d\omega}{dx} + \frac{D}{C}\frac{d^3\omega}{dx^3} + \frac{D}{C^2}\frac{dq}{dx},$$
(2)

$$D = E \left(\frac{H^3}{12} + H y_c \right)^2 + E_g \left(y_g - y_c \right)^2,$$
(3)

where *C* is the shear stiffness and $C = \kappa GA$, κ is the Timoshenko shear factor and $\kappa = 10(1 + v)/(12 + 11v)$, and *v* represents the Poisson ratio; *G* is the shear modulus of the beam; *A* is the crosssectional area of the beam, *q* is the arbitrary pressure on the beam; *E_g* is the tensile modulus of reinforcement; *E* is the Young's modulus of the concrete; y_c is the distance between the central axis of the beam and the neutral axis and $y_c = E_g y_g/(E_g + EH)$; y_g is the distance between the reinforcement and the central axis; *H* is the beam height; k_s is the foundation coefficient; and *D* is bending stiffness.

The pressure on the beam just presents in the form of not only a point but also an interval. Thus, the pressure on the beam q should be considered, and it can be described in Eq. 4, which can be expressed in terms of Fourier series as follows:

$$q = f(x) = \begin{cases} 0 \ 0 \le x < r, \\ q_0 \ r \le x \le r + B \\ 0 \ r + B < x \le L, \end{cases} A_n \cos \frac{n\pi}{L} x, \quad (4)$$
$$A_0 = \frac{1}{L} \int_0^L f(x) dx, \\A_n = \frac{2}{L} \int_0^L f(x) \cos \frac{n\pi}{L} x dx, \end{cases}$$
(5)

where *L* represents the length of the beam, *r* is the start point of the pressure q_0 , *B* is the width of the pressure, and q_0 is the value of the pressure when $r \le x \le r + B$.

Substituting Eq. 4 into Eq. 1, the governing differential equation of the reinforced Timoshenko beam can be rewritten as

$$D\frac{d^4\omega}{dx^4} - \frac{k_s D}{C} \frac{d^2\omega}{dx^2} + k_s \omega = A_0 + \sum_{n=1}^{\infty} A_n \left[1 + \left(\frac{n\pi}{L}\right)^2 \frac{D}{C} \right] \cos \frac{n\pi}{L} x.$$
 (6)

The solution of ω can be obtained in Eq. 7 as follows:

$$\omega = e^{\alpha x} \left(c_1 \cos \beta x + c_2 \sin \beta x \right) + e^{-\alpha x} \left(c_3 \cos \beta x + c_4 \sin \beta x \right) + \frac{A_0}{k_s} + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi}{L} x,$$
(7)

When $k_s < 4C^2/D$, and the expressions of α , β , and a_n are shown in Eq. 8:

$$\alpha = \sqrt[4]{\frac{k_s}{D}} \sqrt{1 + \frac{D}{C} \sqrt{\frac{k_s}{D}}},$$

$$\beta = \sqrt[4]{\frac{k_s}{D}} \sqrt{1 - \frac{D}{C} \sqrt{\frac{k_s}{D}}},$$

$$a_n = \frac{A_n \left[1 + \left(\frac{n\pi}{L}\right)^2 \frac{D}{C}\right]}{k_s + \frac{k_s D}{C} \left(\frac{n\pi}{L}\right)^2 + D\left(\frac{n\pi}{L}\right)^4},$$
(8)

where c_1 , c_2 , c_3 , and c_4 are four constants determined by the boundary conditions. In general, the stiffness of the beam is greater than that of the foundation, and therefore $k_s < 4C^2/D$ can satisfy the normal practice of the construction of beams, reinforced pavement, or granular base over the soft ground.

According to $\varphi = (1 - \frac{Dk_s}{C^2}) \frac{d\omega}{dx} + \frac{D}{C} \frac{d^3\omega}{dx^3} + \frac{D}{C^2} \frac{dq}{dx}$, the rotation angel φ can be obtained in Eq. 9as follows:

$$\begin{split} \varphi &= e^{\alpha x} \left(c_5 \cos \beta x + c_6 \sin \beta x \right) + e^{-\alpha x} \left(c_7 \cos \beta x + c_8 \sin \beta x \right) \\ &+ \sum_{n=1}^{\infty} \left\{ \left[\frac{D}{C} \left(\frac{n\pi}{L} \right)^3 - \left(1 - \frac{Dk_s}{C^2} \right) \frac{n\pi}{L} \right] a_n \\ &- \frac{n\pi D}{LC^2} A_n \right\} \sin \frac{n\pi}{L} x, \end{split}$$
(9)
$$c_5 &= c_1 \left[\alpha \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\alpha^3 - 3\alpha\beta^2 \right) \right] + c_2 \left[\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(3\alpha^2\beta - \beta^3 \right) \right], \\ c_6 &= c_1 \left[-\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\beta^3 - 3\alpha^2\beta \right) \right] + c_2 \left[\alpha \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\alpha^3 - 3\alpha\beta^2 \right) \right], \\ c_7 &= c_3 \left[-\alpha \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(3\alpha\beta^2 - \alpha^3 \right) \right] + c_4 \left[\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(3\alpha^2\beta - \beta^3 \right) \right], \\ c_8 &= c_3 \left[-\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(-3\alpha\beta^2 + \beta^3 \right) \right] + c_4 \left[-\alpha \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(-\alpha^3 + 3\alpha\beta^2 \right) \right], \end{split}$$
(10)

where c_5 , c_6 , c_7 , and c_8 are four constants, which can be figured out when c_1 , c_2 , c_3 , and c_4 are known.

According to $M = -D\frac{d\varphi}{dx}$, the bending moment *M* can be obtained in Eq. 11as follows:

$$\begin{split} M &= -D \frac{d\varphi}{dx} \\ &= -D \left\{ e^{\alpha x} \left(c_9 \cos \beta x + c_{10} \sin \beta x \right) + e^{-\alpha x} \left(c_{11} \cos \beta x + c_{12} \sin \beta x \right) \right. \\ &+ \sum_{n=1}^{\infty} \left[\left. -a_n \left(\frac{n\pi}{L} \right)^2 - a_n \frac{k_s}{C} + \frac{A_n}{C} \right] \cos \frac{n\pi}{L} x \right\}, \end{split}$$
(11)
$$c_9 &= c_1 \left[\left(\alpha^2 - \beta^2 \right) \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\alpha^4 - 6\alpha^2 \beta^2 + \beta^4 \right) \right] + c_2 \left[2\alpha\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{4\alpha\beta D}{C} \left(\alpha^2 - \beta^2 \right) \right], \\ c_{10} &= c_1 \left[-2\alpha\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{4\alpha\beta D}{C} \left(-\alpha^2 + \beta^2 \right) \right] + c_2 \left[\left(\alpha^2 - \beta^2 \right) \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\alpha^4 - 6\alpha^2 \beta^2 + \beta^4 \right) \right], \\ c_{11} &= c_3 \left[\left(\alpha^2 - \beta^2 \right) \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\alpha^4 - 6\alpha^2 \beta^2 + \beta^4 \right) \right] + c_4 \left[-2\alpha\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{4\alpha\beta D}{C} \left(-\alpha^2 + \beta^2 \right) \right], \\ c_{12} &= c_3 \left[2\alpha\beta \left(1 - \frac{Dk_s}{C^2} \right) + \frac{4\alpha\beta D}{C} \left(\alpha^2 - \beta^2 \right) \right] + c_4 \left[\left(\alpha^2 - \beta^2 \right) \left(1 - \frac{Dk_s}{C^2} \right) + \frac{D}{C} \left(\alpha^4 - 6\alpha^2 \beta^2 + \beta^4 \right) \right], \end{split}$$
(12)

where c_9 , c_{10} , c_{11} , and c_{12} are four constants, which can be figured out when c_1 , c_2 , c_3 , and c_4 are known.

According to $Q = C(\frac{d\omega}{dx} - \varphi)$, the shear force Q can be obtained in Eq. 13as follows:

$$Q = De^{\alpha x} (c_{13} \cos \beta x + c_{14} \sin \beta x) + De^{-\alpha x} (c_{15} \cos \beta x + c_{16} \sin \beta x) + D\sum_{n=1}^{\infty} \left[-a_n \frac{n\pi k_s}{LC} - a_n \left(\frac{n\pi}{L} \right)^3 + \frac{n\pi}{CL} A_n \right] \sin \frac{n\pi}{L} x,$$
(13)
$$c_{13} = c_1 \left(\alpha \frac{k_s}{C} - \alpha^3 + 3\alpha\beta^2 \right) + c_2 \left(\beta \frac{k_s}{C} - 3\alpha^2\beta + \beta^3 \right),$$
(13)
$$c_{14} = c_1 \left(-\beta \frac{k_s}{C} + 3\alpha^2\beta - \beta^3 \right) + c_2 \left(\alpha \frac{k_s}{C} - \alpha^3 + 3\alpha\beta^2 \right),$$
(14)
$$c_{15} = c_3 \left(-\alpha \frac{k_s}{C} + \alpha^3 - 3\alpha\beta^2 \right) + c_4 \left(\beta \frac{k_s}{C} - 3\alpha^2\beta + \beta^3 \right),$$
(14)
$$c_{16} = c_3 \left(-\beta \frac{k_s}{C} + 3\alpha^2\beta - \beta^3 \right) + c_4 \left(-\alpha \frac{k_s}{C} + \alpha^3 - 3\alpha\beta^2 \right),$$

where c_{13} , c_{14} , c_{15} , and c_{16} are four constants, which can be figured out when c_1 , c_2 , c_3 , and c_4 are known.

If the prestressed reinforcement is arranged in a straight line, the prestressed load can be equivalent to the bending moment M_0 at both ends of the beam. Then, the actual boundary of the internal force calculation model is shown in Eq. 15:

$$\begin{array}{c}
M|_{x=0} = M_{0}, \\
M|_{x=L} = M_{0}, \\
Q|_{x=0} = 0, \\
Q|_{x=L} = 0.
\end{array}$$
(15)

Combine the governing equation and boundary condition, and the results are as follows:

$$c_{9} + c_{11} = -\frac{M_{0}}{D} - \sum_{n=1}^{\infty} \left[-a_{n} \left(\frac{n\pi}{L}\right)^{2} - a_{n} \frac{k_{s}}{C} + \frac{A_{n}}{C} \right], \quad (16)$$

$$c_{13} + c_{15} = 0, (17)$$

 $e^{\alpha L} (c_9 \cos \beta L + c_{10} \sin \beta L) + e^{-\alpha L} (c_{11} \cos \beta L + c_{12} \sin \beta L)$

$$= -\frac{M_0}{D} - \sum_{n=1}^{\infty} \left[-a_n \left(\frac{n\pi}{L}\right)^2 - a_n \frac{k_s}{C} + \frac{A_n}{C} \right] \cos n\pi, \tag{18}$$

$$e^{\alpha L} (c_{13} \cos \beta L + c_{14} \sin \beta L) + e^{-\alpha L} (c_{15} \cos \beta L + c_{16} \sin \beta L) = 0.$$
(19)

The aforementioned expressions can be described as follows:

$$c_{1}\left[\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]+c_{2}\left[2\alpha\beta\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{4\alpha\beta D}{C}\left(\alpha^{2}-\beta^{2}\right)\right]+c_{3}\left[\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)\right]$$
$$+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]+c_{4}\left[-2\alpha\beta\left(1-\frac{Dk_{s}}{C^{2}}\right)\right]$$
$$+\frac{4\alpha\beta D}{C}\left(-\alpha^{2}+\beta^{2}\right)\right]$$
$$=R_{1}-\frac{M_{0}}{D},$$
(20)

$$R_1 = \sum_{n=1}^{\infty} \left[a_n \left(\frac{n\pi}{L} \right)^2 + a_n \frac{k_s}{C} - \frac{A_n}{C} \right], \tag{21}$$

$$c_{1}\left(\alpha \frac{k_{s}}{C} - \alpha^{3} + 3\alpha\beta^{2}\right) + c_{2}\left(\beta \frac{k_{s}}{C} - 3\alpha^{2}\beta + \beta^{3}\right) + c_{3}\left(-\alpha \frac{k_{s}}{C} + \alpha^{3} - 3\alpha\beta^{2}\right) + c_{4}\left(\beta \frac{k_{s}}{C} - 3\alpha^{2}\beta + \beta^{3}\right)$$

= 0, (22)

$$c_{1}\left\{\left[\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{\alpha L}\cos\beta L+\left[\right.\right.\\\left.\left.\left.\left.\left(2\alpha\beta\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{4\alpha\beta D}{C}\left(-\alpha^{2}+\beta^{2}\right)\right]e^{\alpha L}\sin\beta L\right\}+c_{2}\left\{\right.\right.\\\left.\left.\left.\left(2\alpha\beta\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{4\alpha\beta D}{C}\left(\alpha^{2}-\beta^{2}\right)\right]e^{\alpha L}\cos\beta L+\left[\left(\alpha^{2}-\beta^{2}\right)\right.\right.\\\left.\left.\left.\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{\alpha L}\sin\beta L\right\}+c_{3}\left\{\left[\right.\right]\right.\\\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{4\alpha\beta D}{C}\left(-\alpha^{2}+\beta^{2}\right)\right]e^{-\alpha L}\cos\beta L+\left[\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\right.\\\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\\\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\right.\right]\right.\\\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\\\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right]\right.\right.\right]\right.$$

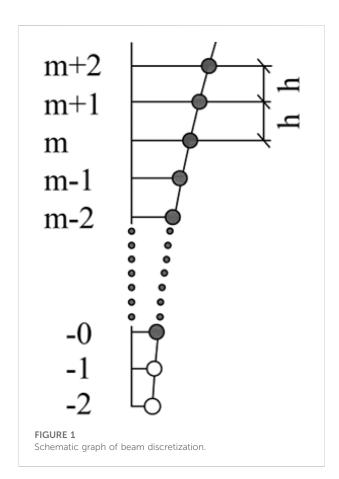
$$\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\right]\right.$$

$$\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\right]\right.$$

$$\left.\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\right]\right.$$

$$\left.\left.\left.\left(\alpha^{2}-\beta^{2}\right)\left(1-\frac{Dk_{s}}{C^{2}}\right)+\frac{D}{C}\left(\alpha^{4}-6\alpha^{2}\beta^{2}+\beta^{4}\right)\right]e^{-\alpha L}\sin\beta L\right\}\right.\right]\right.$$

$$R_3 = \sum_{n=1}^{\infty} \left[a_n \left(\frac{n\pi}{L} \right)^2 + a_n \frac{k_s}{L} - \frac{A_n}{C} \right] cosn\pi,$$
(24)



$$c_{1}\left[\left(\alpha\frac{k_{s}}{C}-\alpha^{3}+3\alpha\beta^{2}\right)e^{\alpha L}\cos\beta L+\left(-\beta\frac{k_{s}}{C}+3\alpha^{2}\beta\right)e^{\alpha L}\sin\beta L\right]+c_{2}\left[\left(\beta\frac{k_{s}}{C}-3\alpha^{2}\beta+\beta^{3}\right)e^{\alpha L}\cos\beta L+\left(\alpha\frac{k_{s}}{C}-\alpha^{3}+3\alpha\beta^{2}\right)e^{\alpha L}\sin\beta L\right]+c_{3}\left[\left(-\alpha\frac{k_{s}}{C}+\alpha^{3}-3\alpha\beta^{2}\right)e^{-\alpha L}\cos\beta L+\left(-\beta\frac{k_{s}}{C}+3\alpha^{2}\beta-\beta^{3}\right)e^{-\alpha L}\sin\beta L\right]+c_{4}\left[\times\left(\beta\frac{k_{s}}{C}-3\alpha^{2}\beta+\beta^{3}\right)e^{-\alpha L}\cos\beta L+\left(-\alpha\frac{k_{s}}{C}+\alpha^{3}-3\alpha\beta^{2}\right)e^{-\alpha L}\sin\beta L\right]\right]$$
$$=0.$$
(25)

Equations 20-25 can be expressed as the following matrix equation:

$$[K]{C} = {R}, (26)$$

$$\{C\} = [c_1, c_2, c_3, c_4]^T, \{R\} = \left[R_1 - \frac{M_0}{D}, 0, R_3 - \frac{M_0}{D}, 0\right].$$
(27)

Stiffness matrix [K] can be solved by Eqs 20-25. For example, K_{11} can be written as $K_{11} = (\alpha^2 - \beta^2)(1 - \frac{Dk_s}{C^2}) + \frac{D}{C}(\alpha^4 - 6\alpha^2\beta^2 + \beta^4)$. If the values of α , β , D, C, k_s , L, and q = f(x) are known, the value of K_{11} can be calculated. Then, Eq. 26 can be solved as:

$$\{C\} = [K]^{-1}\{R\}.$$
 (28)

The parameters $c_5 \sim c_{16}$ can be figured out when the values of c_1 , c_2 , c_3 , and c_4 are known.

Finite difference method of the Euler-Bernoulli beam on the Winkler foundation

Compared with the Timoshenko beam model, the Euler-Bernoulli beam model is also based on a rigid crosssection assumption, while the transverse shear deformation is ignored. Taking the effect of double-layer reinforcement, prestressed reinforcement of prefabricated beams, and anchorage force of the anchor cable into consideration, a finite difference method of the Euler-Bernoulli beam on the Winkler foundation can be obtained to calculate the internal force and deformation of the prestressed I-beam.

The beam is discretized by equal segmentation in finite difference calculation, as shown in Figure 1.

The subgrade reaction p(x) proposed by Winkler is as follows:

$$p(x) = k_s. \tag{29}$$

Therefore, the governing differential equation of the beam on the elastic foundation becomes

$$EI\frac{d^4\omega}{dx^4} + k_s b\omega = 0, (30)$$

where E is the Young's modulus of the beam, I is the inertia moment, and b is the beam width.

Equation 30 can be described in terms of the differential form at the arbitrary point (*m*-point):

$$\frac{EI}{h^4} \left(\omega_{m+2} - 4\omega_{m+1} + 6\omega_m - 4\omega_{m-1} + \omega_{m-2} \right) + k_s b \omega_m = 0, \quad (31)$$

where h represents the fragment length of the beam discretization.

The difference equation of the bending moment M at *m*-point is

$$M = -\frac{EI}{h^2} (\omega_{m+1} - 2\omega_m + \omega_{m-1}).$$
(32)

The difference equation of shear force Q at m-point is

$$Q = -\frac{EI}{2h^3} \left(\omega_{m+2} - 2\omega_{m+1} + 2\omega_{m-1} - \omega_{m-2} \right).$$
(33)

Bending moment at the start point (x=0) is M_0 , and Eq. 32 can be simplified as follows:

$$\omega_{-1} = 2\omega_0 - \omega_1 - \frac{M_0 h^2}{EI}.$$
 (34)

When shear force at the start point (x = 0) is 0, Eq. 33 can be simplified as follows:

$$\omega_{-2} = 2\omega_{-1} - 2\omega_1 + \omega_2. \tag{35}$$

When m = 0, Eq. 31 can be obtained as follows:

$$\frac{EI}{h^4} (\omega_2 - 4\omega_1 + 6\omega_0 - 4\omega_{-1} + \omega_{-2}) + k_s b\omega_0 = 0.$$
(36)

Substitute Eq. 34 and Eq. 35 into Eq. 36, we get

$$\left(2 + \frac{k_s b h^4}{EI}\right)\omega_0 - 4\omega_1 + 2\omega_2 = -\frac{2M_0 h^2}{EI}.$$
 (37)

 c_0 , a_0 , b_0 , and d_0 are defined as follows:

$$c_{0} = 2 + \frac{k_{s}bh^{4}}{EI},$$

$$a_{0} = 4/c_{0},$$

$$b_{0} = 2/c_{0},$$

$$d_{0} = -\frac{M_{0}h^{2}}{EI}/c_{0}.$$
(38)

Then, Eq. 37 can be simplified as follows:

$$\omega_0 = a_0 \omega_1 - b_0 \omega_2 + d_0. \tag{39}$$

Similarly, assume that the recursion formula at the arbitrary node is

$$\begin{array}{c} \omega_{m} = a_{m}\omega_{m+1} - b_{m}\omega_{m+2} + d_{m}, \\ \omega_{m-1} = a_{m-1}\omega_{m} - b_{m-1}\omega_{m+1} + d_{m-1}, \\ \omega_{m-2} = a_{m-2}\omega_{m-1} - b_{m-2}\omega_{m} + d_{m-2}. \end{array}$$

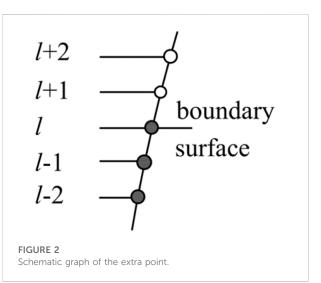
$$(40)$$

Substituting Eq. 40 into Eq. 31, Eq. (41) can be obtained as follows:

$$\omega_{m-2} - [4 + (a_{m-2} - 4)b_{m-1}]\omega_{m+1} + \left[6 - b_{m-2} + (a_{m-2} - 4)a_{m-1} + \frac{k_s b h^4}{EI}\right]\omega_m$$

= $(-a_{m-2} + 4)d_{m-1} - d_{m-2}.$ (41)

Define



$$c'_{m} = 6 - b_{m-2} + (a_{m-2} - 4)a_{m-1} + \frac{k_{s}bh^{4}}{EI},$$

$$a'_{m} = \frac{4 + (-4 + a_{m-2})b_{m-1}}{c_{m}},$$

$$b'_{m} = 1/c_{m},$$

$$d'_{m} = \frac{(4 - a_{m-2})d_{m-1} - d_{m-2}}{c_{m}}.$$
(42)

Then, Eq. 41 can be simplified as

$$\omega_m = a'_m \omega_{m+1} - b'_m \omega_{m+2} + d'_m.$$
(43)

Substituting Eq. 34 into Eq. 43, the values of $c'_{\neg 1}a'_{\neg 1}b'_{\neg 1}$ and d'_{-1} can be obtained in Eq. 44as follows:

$$\begin{array}{c}
c_{-1}' = 1, \\
a_{-1}' = 2, \\
b_{-1}' = 1, \\
d_{-1}' = -\frac{M_0 h^2}{EI}.
\end{array}$$
(44)

The extra point at the end point (x = l) is shown in Figure 2. Presume both end points of the beam are free, meaning that the shear force is zero and the bending moment is the equivalent bending moment M_0 ; therefore, the differential equations of shear and bending moment at *l*-point are obtained as shown in Eq. 45 and Eq. (46):

$$-\frac{EI}{2h^3}\left(\omega_{l+2} - 2\omega_{l+1} + 2\omega_{l-1} - \omega_{l-2}\right) = 0,$$
(45)

$$-\frac{EI}{h^2}(\omega_{l+1} - 2\omega_l + \omega_{l-1}) = M_0.$$
(46)

According to the recursion formula, Eqs 47-49 can be obtained as follows:

$$\omega_l - a'_l \omega_{l+1} + b'_l \omega_{l+2} = d'_l, \qquad (47)$$

$$(a_{l-1}^{'}-2)\omega_{l}-(b_{l}^{'}-1)\omega_{l+1}=-\frac{M_{0}h^{2}}{EI}-d_{l-1}^{'},$$
(48)

$$\begin{pmatrix} b'_{l-2} - a'_{l-1} a'_{l-2} + 2a'_{l-1} \end{pmatrix} \omega_l - (2 - a'_{l-2} b'_{l-1} + 2b'_{l-1}) \omega_{l+1} + \omega_{l+2} = d'_{l-2} + (a'_{l-2} - 2) d'_{l-1}.$$

$$(49)$$

Let $R_2 = a'_{l-2} - 2$, $R_3 = b'_{l-2} - a'_{l-1}a'_{l-2} + 2a'_{l-1}$, $J_2 = b'_l - 1$, $J_3 = 2 - a'_{l-2}b'_{l-1} + 2b'_{l-1}$, $R_{J2} = -\frac{M_0h^2}{EI} - d'_{l-1}$,

 $R_{J3} = d'_{l-2} + (a'_{l-2} - 2)d'_{l-1}$, and $H' = J_2 - R_2a'_l - J_2R_3b'_l + J_3R_2b'_l$. The displacement of the extra point and end point can be solved as follows:

$$\omega_{l} = \left(J_{2}d_{1}^{'} - R_{J2}a_{l} - J_{2}R_{J3}b_{1}^{'} + J_{3}R_{J2}b_{1}^{'}\right)/H', \omega_{l+1} = -\left(R_{J2} - R_{2}d_{1}^{'} + R_{2}R_{J3}b_{1}^{'} - R_{3}R_{J2}b_{1}^{'}\right)/H', \omega_{l+2} = \left(J_{2}R_{J3} - J_{3}R_{J2} - J_{2}R_{3}d_{1}^{'} - R_{2}R_{J3}a_{1}^{'} + R_{3}R_{J3}a_{1}^{'}\right)/H'.$$

$$(50)$$

According to the aforementioned derivation process, the expressions of displacement ω_m , shear force Q_m , and bending moment M_m at the arbitrary node can be obtained as follows:

$$\omega_m = a'_m \omega_{m+1} - b'_m \omega_{m+2} + d'_m, \tag{51}$$

$$Q_m = -\frac{EI}{2h^3} \left(\omega_{m+2} - 2\omega_{m+1} + 2\omega_{m-1} - \omega_{m-2} \right), \tag{52}$$

$$M_m = -\frac{EI}{h^2} (\omega_{m+1} - 2\omega_m + \omega_{m-1}).$$
 (53)

Finite difference method of the Euler–Bernoulli beam on the Pasternak foundation

The Pasternak foundation was developed on the basis of the Winkler foundation, and it is assumed that there is shear action between spring elements. Taking the effect of double-layer reinforcement, prestressed reinforcement of prefabricated beams, and anchorage force of the anchor cable into consideration, a finite difference method of the Euler-Bernoulli beam on the Pasternak foundation can be obtained to calculate the internal force and deformation of the prestressed I-beam.

The governing differential equation of the beam on the elastic foundation is shown in Eq. 54:

$$EI\frac{d^4\omega}{dx^4} + k_s b\omega - G_p b\frac{d^2\omega}{dx^2} = 0.$$
 (54)

TABLE 1 Parameters of the beam.

Parameter	Value	Parameter	Value	
<i>L</i> (m)	3.00	υ	0.20	
<i>r</i> (m)	1.45	y_g	0.225	
<i>B</i> (m)	0.10	Ус	0.2172	
<i>H</i> (m)	0.55	D	5.2971×10 ⁵	
<i>b</i> (m)	0.40	$G (kN \cdot m^{-1})$	3.75×10 ⁵	
$k_s \ (\text{kN} \cdot \text{m}^{-3})$	1.50×10 ⁶	<i>h</i> (m)	0.001	
$E (\text{kN} \cdot \text{m}^{-2})$	3.25×10 ⁷	F (kN)	500	
E_g (kN ·m ⁻²)	2.00×10 ⁸	M_0 (kN·m)	-95	

Equation 52 can be described in terms of the differential form, and thus the differential equation at the arbitrary point (m point) is

$$\frac{EI}{h^4} (\omega_{m+2} - 4\omega_{m+1} + 6\omega_m - 4\omega_{m-1} + \omega_{m-2}) + k_m b\omega_m - \frac{G}{h^2} (\omega_{m+1} - 2\omega_m + \omega_{m-1}) = 0.$$
(55)

The derivation process in the Pasternak model is similar to that in the Winkler model, and the displacement ω_m , soil resistance p_m , shear force Q_m , and bending moment M_m are calculated as follows:

$$\omega_{m} = a'_{m}\omega_{m+1} - b'_{m}\omega_{m+2} + d'_{m}, \qquad (56)$$

$$p_m = -k_m b\omega_m + G_p b \frac{d^2 \omega_m}{dx^2},$$
(57)

$$Q_m = -\frac{EI}{2h^3} \left(\omega_{m+2} - 2\omega_{m+1} + 2\omega_{m-1} - \omega_{m-2} \right) + G_p b \frac{\omega_{m+1} - \omega_{m-1}}{2h},$$
(58)

$$M_m = -\frac{EI}{h^2} (\omega_{m+1} - 2\omega_m + \omega_{m-1}).$$
 (59)

Verification of calculation methods

In order to compare and verify the reasonableness of the aforementioned three methods (theoretical model of the Timoshenko beam on the Winkler foundation, FDM of the Euler–Bernoulli beam on the Winkler foundation, and FDM of the Euler–Bernoulli beam on the Pasternak foundation), these three methods were employed to calculate the internal force and deformation of an example of the I-beam applied to an application. In the soft rock slope reinforcement in the Mengping highway project, I-beams and cross-beams were employed. The I-beam was chosen as an engineering example here, whose width was 300 mm, height was 550 mm, and length was 3000 mm under the concentrated force F = 500kN and

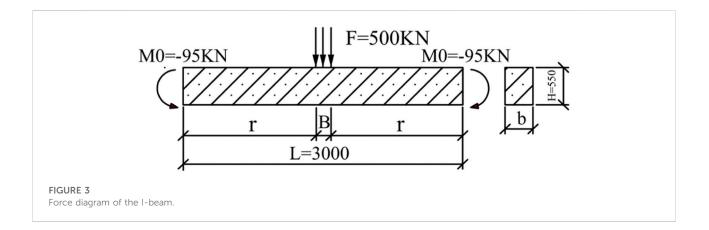
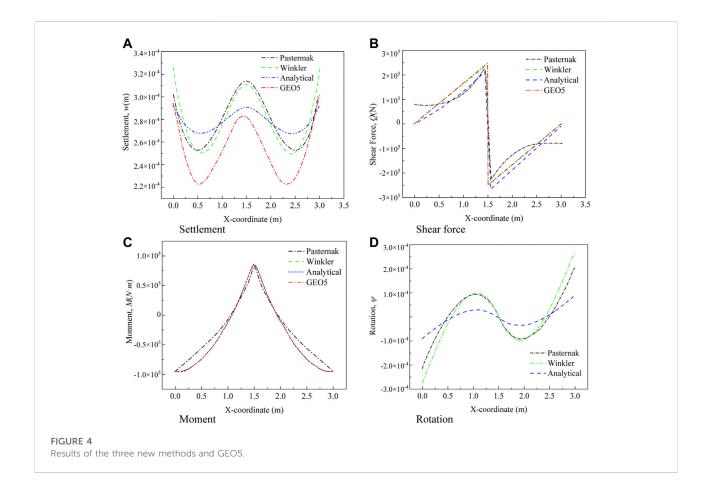


TABLE 2 Calculated results using different models.

Parameter	Analytical solution	Winkler solution	Pasternak solution	GEO5 Pasternak	GEO5 Winkler
ω_{max} (mm)	0.295	0.3244	0.3142	0.3	0.1
Q _{max} (kN)	271.2	249.8	249.6	250	250
M_l (kN·m)	84.93	89.53	82.73	89.16	78.90
φ	8.65×10 ⁻⁵	2.7×10^{-4}	2.08×10^{-4}	_	_



equivalent bending moment $M_0 = -95kN \cdot m$; then, the results of these three models were verified by the numerical solutions calculated by GEO5. The specific calculation parameters are presented in Table 1, and the simplified calculation diagram is shown in Figure 3. The calculation results are demonstrated in Table 2 and Figure 4, respectively.

It is shown that the calculated internal forces and deformations of the prestressed beam under concentrated force using the three models are quite close.

- (1) The sedimentation value calculated by FDM of the Winkler model differs the most, compared with that obtained by GEO5, and the percentage difference is 8.133%. This is partly because the shear deformation was ignored in FDM of the Euler–Bernoulli beam on the Winkler foundation, and the influence of shear deformation on settlement should not be ignored. Minimum variation can be achieved by comparing the sedimentation value calculated by the theoretical analytical solution of the Timoshenko beam on the Winkler foundation with that by GEO5, and the percentage difference is 1.67%.
- (2) The differences between the shear force value calculated by the theoretical analytical solution of the Timoshenko beam on the Winkler foundation and *via* GEO5 are the largest, and the percentage difference is 8.48%. The shear value calculated by FDM of the Winkler model is the closest to that obtained by GEO5, and the percentage difference is 0.08%. This is partly because the shear deformation was ignored in the Winkler foundation, and the value of polynomial terms in the theoretical formula affects the calculated value of the shear force.
- (3) The gap in the bending moment value calculated by FDM of the Pasternak model is the largest in comparison with that obtained using GEO5, and the percentage difference is 7.21%; the shear value calculated by FDM of the Winkler model is the closest to that obtained by GEO5, and the percentage difference is 0.42%.

Conclusion

In this article, in order to calculate the internal force and deformation of the prestressed I-beam with prestressed anchorage loads, three theoretical or FDM methods were deduced and verified using GEO5. Main conclusions are as follows:

(1) Three methods, namely, the theoretical model of the Timoshenko beam on the Winkler foundation, the finite difference method of the Euler–Bernoulli beam on the Winkler foundation, and the finite difference method of the Euler–Bernoulli beam on the Pasternak foundation, are proposed, which can be employed to analyze the internal force and deformation of the prestressed I-beam with prestressed anchorage loads.

- (2) Among the deformation calculation results of GEO5, the results of the theoretical model of the Timoshenko beam on the Winkler foundation are the closest with a difference of 1.67%, while the results of FDM of the Euler–Bernoulli beam on the Winkler foundation have the largest difference of 8.13%.
- (3) Among the shear force calculation results of GEO5, the results of FDM of the Euler-Bernoulli beam on the Winkler foundation are the closest with a difference of 0.08%, while the results of the theoretical model of the Timoshenko beam on the Winkler foundation have the largest difference of 8.48%.
- (4) Among the bending moment calculation results of GEO5, the results of FDM of the Euler–Bernoulli beam on the Winkler foundation are the closest with a difference of 0.42%, while the results of the finite difference method of the Euler–Bernoulli beam on the Pasternak foundation have the largest difference of 7.21%.
- (5) It can be seen that the results obtained by the corresponding solutions of the three models and the GEO5 calculation results fluctuate within 10%. Whether shear deformation is considered or not, it has little influence on the bending moment. It is feasible to use the aforementioned three models to solve the internal force and deformation of the prestressed I-beam under prestressed anchorage loads.

Data availability statement

The original contributions presented in the study are included in the article/Supplementary Material; further inquiries can be directed to the corresponding author.

Author contributions

YY: conceptualization. SW: methodology and writing—original draft preparation. DC: investigation. JY: data curation. PW: writing—original draft preparation. CY and SnW: writing—review and editing.

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Conflict of interest

The authors YY, SW, DC and PW were employed by China Academy of Railway Sciences Corporation Limited, Beijing, China. The authors SW and PW were employed by Beijing Tieke Special Engineering Technological Development Corporation Limited, Beijing, China.

The remaining authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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