



Gravity Prediction Method and the Application of Full-Space Tunnels in the Deep Extension of High-Density Vertical Orebodies

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A full-space tunnel gravity prediction method for evaluating the deep extension of high-density vertical orebodies has been developed based on the significant density differences between surrounding rocks and their characteristics as large-scale orebodies. This method uses the residual gravity anomaly parameters in the tunnel to determine whether the orebody is upright based on the symmetry of the anomaly. Moreover, it identifies the deep extensional characteristics of high-density upright orebodies based on the positive and negative anomalies. This method has characteristics such as strong anti-interference ability, slight influence from the terrain, and high detection accuracy. It has produced promising outcomes in the detection of deep, concealed orebodies in the Huangshaping skarn-type tin-tungsten polymetallic deposit located in southern Hunan. This method not only resolves the problem of rapid identification of the deep extension of single-hole high-density orebodies and provides reference basis for deep exploration but also proposes a new method for mineral resource evaluation.

Keywords: tunnel gravity exploration, high-density vertical orebody, extension prediction of the deep orebody, huangshaping tin-tungsten polymetallic deposit, pingbao mining, southern Hunan

INTRODUCTION

When an orebody is discovered through trenching during the exploration of ore deposits, estimating the depth of extension of orebody is typically concerning. The most direct way to evaluate the deep extension of an orebody is by constructing a trench or drill hole. However, the deep extension of the orebody should be predicted before designing a deep prospecting trench or drill hole. Therefore, physical prospecting is generally conducted before designing the exploration project. Common geophysical exploration methods include gravity, magnetic, electrical, seismic, radioactive, and geothermal exploration. Seismic, radioactive, and geothermal exploration methods are typically ineffective for non-radioactive metal orebodies, and magnetic and electrical exploration methods cannot be employed owing to electromagnetic interference in the pit (Zeng, 2008). However, the gravity measurement in the pit is unaffected by electromagnetic interference. When there is a significant density difference between a certain size of the orebody and surrounding rock, good detection results can be obtained via gravity method.

With improvements in the accuracy of gravity instruments and the wide range of research targets, gravity exploration methods have transitioned from traditional surface gravity exploration to underground-, aviation-, satellite-, and marine gravity-based exploration methods (Zeng, 2008).

Currently, the theory of gravity exploration is relatively well developed as the factors of interference and terrain influence have been extensively studied (Fu, 2010; Hu, 2015; Zhang et al., 2015). Research on the instruments (Cai, 1980; Zeng, 1999; Geng, 2016; Lu et al., 2018), operating methods, data processing, forward and inverse calculations, non-ferrous metal detection, and coal mine safety (Goodacre, 1973; Beyer, 1979; Okabe, 1979; Schmoker, 1980; Casten and Gram, 1989; Li and Chouteau, 1998; Nagy et al., 2000; Yang et al., 2000; Fan, 2007; Luo, 2008; Yi et al., 2008; Cao, 2011; Zhdanov et al., 2011; Howell et al., 2012; Luo and Liu, 2012; Zhang, 2012; Zhang et al., 2012; Teng et al., 2013; Liu and Liu, 2014; Wang et al., 2014; Zhang, 2017; Ming et al., 2021) has been conducted worldwide.

In the study of underground gravity, Smith (1950) described the methodologies to solve various problems and interference factors, which should be considered in underground gravity measurements. Hammer (1950) comprehensively discussed the methodology of underground gravity measurement and proposed potential factors that may trigger errors, basic approaches for data processing, and formulas for calculating the density from underground gravity measurements. Xu and Zhou (1989) translated their studies and illustrated the characteristics of borehole gravity. Nevertheless, these studies reveal a greater inclination toward the field of borehole gravity rather than tunnel gravity.

Studies on the tunnel gravity detection related to the construction method, interference factors, near-area ground improvement, tunnel correction, and abnormal characteristics were conducted according to previous results from (Wu, 2011). Tao (2017) suggested the tunnel gravity terrain correction and summarized the influencing rules followed by the three kinds of models: positive, negative, and combined terrains under the conditions of one, two, and three dimensions, respectively, by applying forward modeling. The patent applied by Han Runsheng and Li Wenyao (Han and Li, 2014; Li and Han, 2014; Han et al., 2021) resolved the problem of orienting the high-density or high-density upright cuboid metal orebodies in tunnel exploration by locating and detecting deeply concealed orebodies using four main parameters; these included the gravity anomaly Δg and gradient anomalies in the X, Y, and Z directions. Furthermore, a full-space tunnel gravity method that could determine the location of high-density sloping concealed orebodies was introduced by Han et al. (2020). They described the methodology of systematic tunnel gravity exploration of concealed orebodies that has been implemented in lead-zinc deposits of Huize and Maoping, Yunnan (Han et al., 2014) with remarkable results. However, it cannot be applied in a wide range owing to the observation accuracy of instruments. Thus, further research using few parameters and simple approaches to illustrate the deep extension characteristics of the orebody is urgently required. This paper proposes a method to tackle the problem of evaluating (Wang, 2019) the deep extension of orebodies depending solely on the residual gravity anomaly (Δg), providing a valuable reference basis for deep ore prospecting.

PREDICTION PRINCIPLE OF THE DEEP EXTENSION OF HIGH-DENSITY VERTICAL OREBODY

Formula for the Abnormal Weight of the Upright Rectangle

Assuming that the orebody is a rectangular parallelepiped with a uniform residual density σ , the half-length, half-width, and half-height of the orebody are a , b , and c , respectively. The length, width, and height of the orebody are parallel to the x , y , and z axes of the spatial rectangular coordinate system, respectively. The center point coordinates of the orebody are (ξ_0, η_0, ζ_0) , as shown in Figure 1.

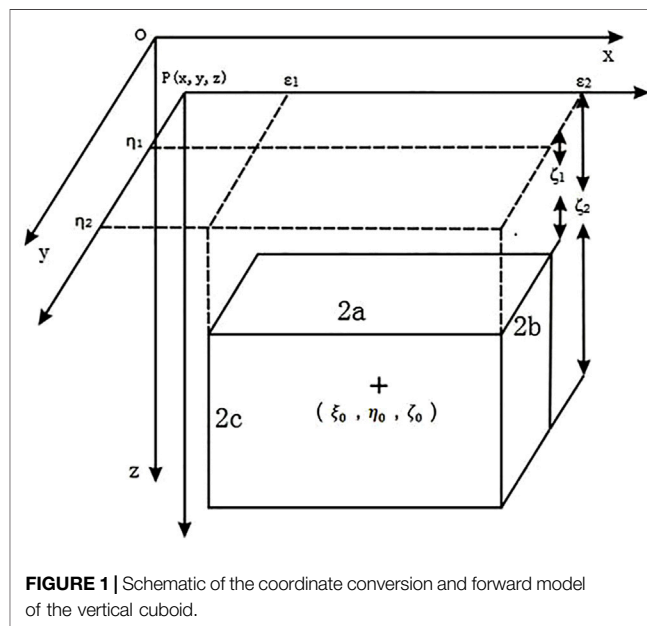


FIGURE 1 | Schematic of the coordinate conversion and forward model of the vertical cuboid.

According to Newton's law of universal gravitation, the gravity anomaly for a cuboid with a uniform residual density σ in the geological body can be expressed as follows (Zeng, 2008):

$$\Delta g(x, y, z) = Vz = G\sigma \int_{\xi_1}^{\xi_2} \int_{\eta_1}^{\eta_2} \int_{\zeta_1}^{\zeta_2} \frac{\zeta - z}{\rho^3} d\xi d\eta d\zeta \quad (1)$$

where G is the gravitational constant, and (x, y, z) are the coordinates of the observation point.

The distance from the center point of the orebody to the measuring point can be expressed as follows:

$$\rho = \sqrt{(\xi_0 - x)^2 + (\eta_0 - y)^2 + (\zeta_0 - z)^2}$$

Eq. 1 can be transformed into

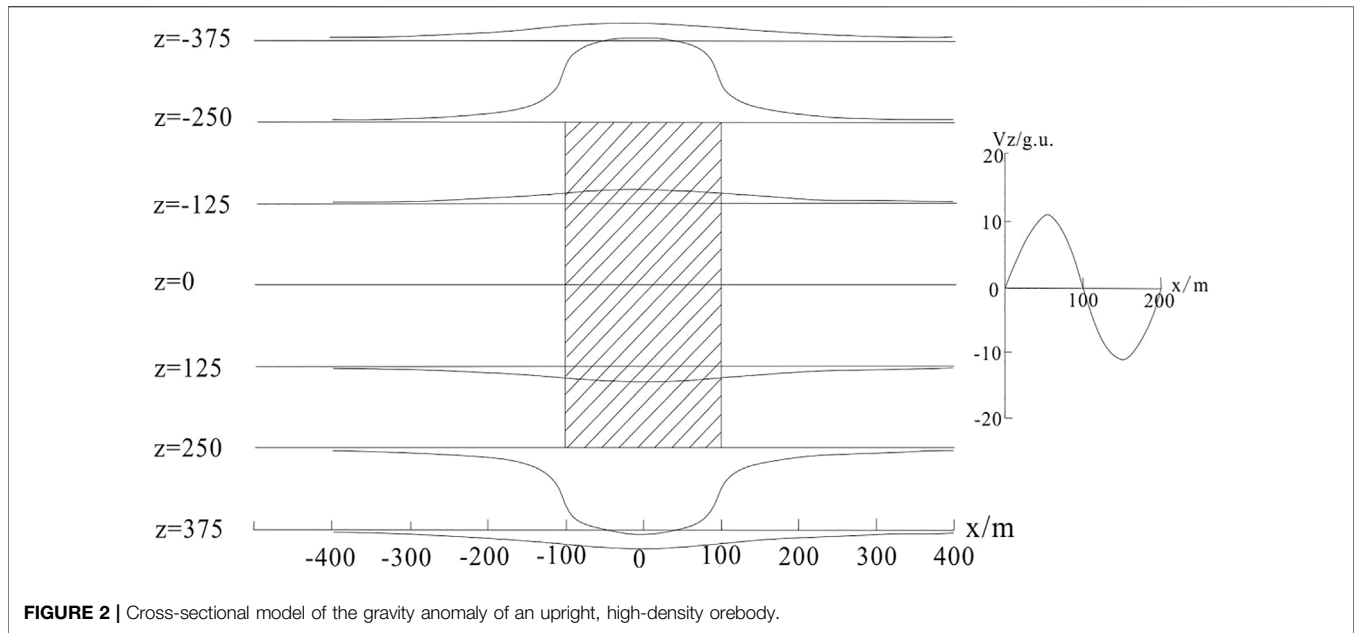


FIGURE 2 | Cross-sectional model of the gravity anomaly of an upright, high-density orebody.

$$\begin{aligned}
 g(x, y, z) = V_z = & -G\sigma \left\{ \left[\left[\xi_2 \ln(\eta + \rho) + \eta \ln(\xi_2 + \rho) + \zeta \arctan \frac{\zeta \rho}{\xi_2 \eta} \right] \right. \right. \\
 & \left. \left. - \left[\xi_1 \ln(\eta + \rho) + \eta \ln(\xi_1 + \rho) + \zeta \arctan \frac{\zeta \rho}{\xi_1 \eta} \right] \right] \right\}_{\eta_1}^{\eta_2} \Big|_{\zeta_1}^{\zeta_2} \\
 = & -G\sigma \left\{ \left[\left[\xi_2 \ln(\eta_2 + \rho) + \eta_2 \ln(\xi_2 + \rho) + \zeta \arctan \frac{\zeta \rho}{\xi_2 \eta_2} \right] \right. \right. \\
 & \left. \left. - \left[\xi_1 \ln(\eta_2 + \rho) + \eta_2 \ln(\xi_1 + \rho) + \zeta \arctan \frac{\zeta \rho}{\xi_1 \eta_2} \right] \right] \right\} \\
 & - \left\{ \left[\left[\xi_2 \ln(\eta_1 + \rho) + \eta_1 \ln(\xi_2 + \rho) + \zeta \arctan \frac{\zeta \rho}{\xi_2 \eta_1} \right] \right. \right. \\
 & \left. \left. - \left[\xi_1 \ln(\eta_1 + \rho) + \eta_1 \ln(\xi_1 + \rho) + \zeta \arctan \frac{\zeta \rho}{\xi_1 \eta_1} \right] \right] \right\}_{\zeta_1}^{\zeta_2} \\
 = & -G\sigma \left\{ \left[\left[\xi_2 \ln(\eta_2 + \rho) + \eta_2 \ln(\xi_2 + \rho) + \zeta_2 \arctan \frac{\zeta_2 \rho}{\xi_2 \eta_2} \right] \right. \right. \\
 & \left. \left. - \left[\xi_1 \ln(\eta_2 + \rho) + \eta_2 \ln(\xi_1 + \rho) + \zeta_2 \arctan \frac{\zeta_2 \rho}{\xi_1 \eta_2} \right] \right] \right\} \\
 & - \left\{ \left[\left[\xi_2 \ln(\eta_1 + \rho) + \eta_1 \ln(\xi_2 + \rho) + \zeta_2 \arctan \frac{\zeta_2 \rho}{\xi_2 \eta_1} \right] \right. \right. \\
 & \left. \left. + \left[\xi_1 \ln(\eta_1 + \rho) + \eta_1 \ln(\xi_1 + \rho) + \zeta_2 \arctan \frac{\zeta_2 \rho}{\xi_1 \eta_1} \right] \right] \right\} \\
 & - \left\{ \left[\left[\xi_2 \ln(\eta_2 + \rho) + \eta_2 \ln(\xi_2 + \rho) + \zeta_1 \arctan \frac{\zeta_1 \rho}{\xi_2 \eta_2} \right] \right. \right. \\
 & \left. \left. + \left[\xi_1 \ln(\eta_2 + \rho) + \eta_2 \ln(\xi_1 + \rho) + \zeta_1 \arctan \frac{\zeta_1 \rho}{\xi_1 \eta_2} \right] \right] \right\} \\
 & + \left\{ \left[\left[\xi_2 \ln(\eta_1 + \rho) + \eta_1 \ln(\xi_2 + \rho) + \zeta_1 \arctan \frac{\zeta_1 \rho}{\xi_2 \eta_1} \right] \right. \right. \\
 & \left. \left. - \left[\xi_1 \ln(\eta_1 + \rho) + \eta_1 \ln(\xi_1 + \rho) + \zeta_1 \arctan \frac{\zeta_1 \rho}{\xi_1 \eta_1} \right] \right] \right\} \quad (2)
 \end{aligned}$$

which is

$$V_z = -G\sigma \sum_{i=1}^2 \sum_{j=1}^2 \sum_{k=1}^2 (-1)^{i+j+k} \left[\xi_i \ln(\eta_j + \rho) + \eta_j \ln(\xi_i + \rho) + \zeta_k \arctan \frac{\zeta_k \rho}{\xi_i \eta_j} \right] \quad (3)$$

where $\rho = \sqrt{\xi_i^2 + \eta_j^2 + \zeta_k^2}$,

The coordinates of the observation point are (x, y, z) :

$$\begin{aligned}
 \xi_i &= \xi_0 + (-1)^i a - x \quad (i = 1, 2); \\
 \eta_j &= \eta_0 + (-1)^j b - y \quad (j = 1, 2); \\
 \zeta_k &= \zeta_0 + (-1)^k c - z \quad (k = 1, 2);
 \end{aligned}$$

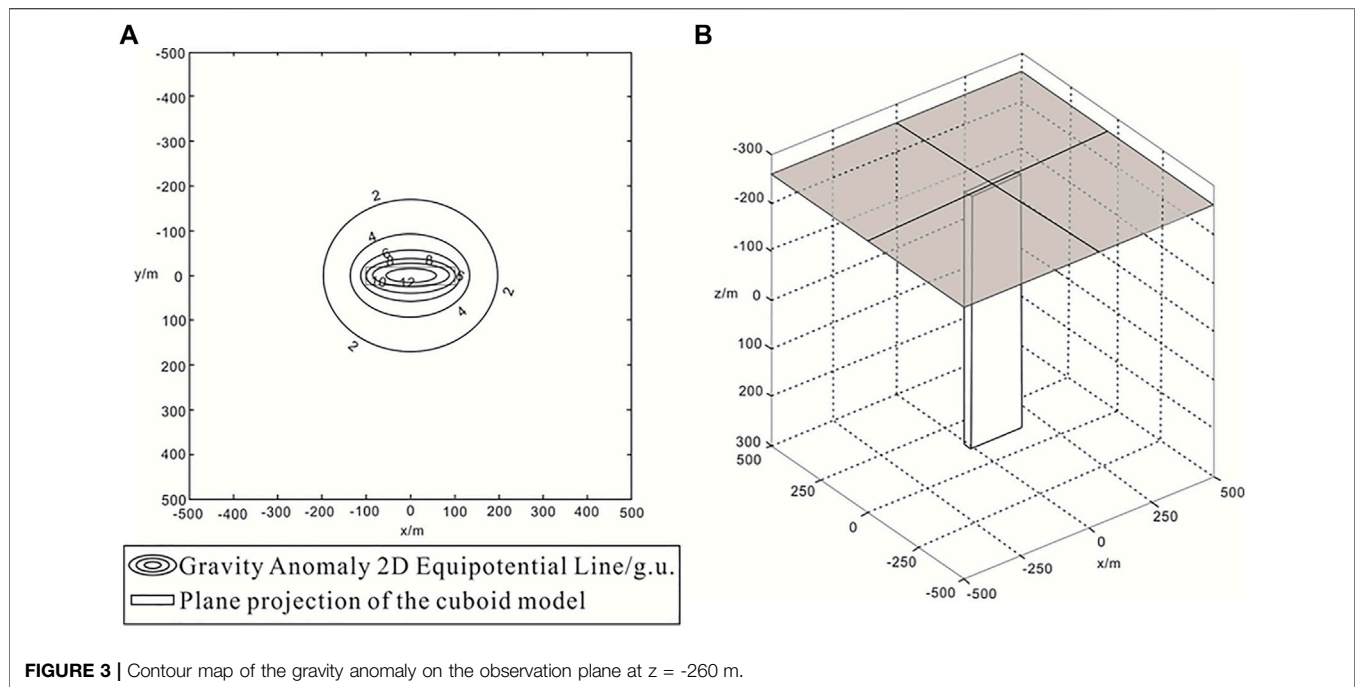
Eqs. 2, 3 are used to calculate the weight of the upright rectangle.

Distribution of the Abnormal Sections of Vertical High-Density Orebodies

The gravity anomaly of an upright high-density orebody can be calculated using Eq. 3. Through forward calculation, the gravity of the xOz section of the vertical high-density cuboid orebody is found to be abnormal, as shown in. Model parameters are as follows: length $2a = 200 \text{ m}$; width $2b = 40 \text{ m}$; height $2c = 500 \text{ m}$; and density difference $\sigma = 1.0 \times 10^3 \text{ kg/m}^3$.

Figure 2 shows the gravity anomaly characteristics on the $y = 0$ and xOz sections as follows:

- (1) When $z = -375 \text{ m}$, The section line is above the model. As the model body is below the section line, the gravity anomaly is positive and symmetrical along $x = 0$, and its value is largest at $x = 0$ (approximately 3.0 g.u.). When x approaches $\pm \infty$, the outlier infinitely approaches 0 g.u.
- (2) When $z = -250 \text{ m}$, The section line is at the top of the rectangular parallelepiped model. As the model body is completely below the section line, the gravity anomaly field direction is downward, and the gravity anomaly is positive and symmetrical along $x = 0$; its value is largest at $x = 0$ and is approximately 12.0 g.u. When x approaches $\pm \infty$, the outlier infinitely approaches 0 g.u.



- (3) When $z = -125$ m, The section is in the upper middle part of the rectangular parallelepiped model. As the outliers generated under the model are larger than those above the model, the direction is opposite; specifically, the positive outliers are greater than the negative outliers. Thus, the anomalies are positive and symmetrical along $x = 0$. Furthermore, the outlier value is the largest when $x = 0$, which is approximately $3.0g.u.$ When x approaches $\pm\infty$, the outlier infinitely approaches $0g.u.$
- (4) When $z = 0$ m, the section line is in the middle of the rectangular parallelepiped model, and the gravity anomaly values generated on the center line of the model are equal in magnitude and opposite in direction, and the anomalies cancel each other; thus, the gravity anomaly is $0g.u.$
- (5) When $z = 125$ m, The section line is in the lower middle part of the model. As the outlier generated below the model is smaller than that above it, the direction is opposite; namely, the negative outlier is greater than the positive outlier. Thus, the anomaly is negative and symmetrical along $x = 0$. When $x = 0$, the outlier is the lowest, that is, approximately $-3.0g.u.$ When x is close to $\pm\infty$, the outlier is infinitely close to $0g.u.$
- (6) When $z = 250$ m, The section line occurs at the bottom of the rectangular parallelepiped model. As the model body is completely above the section line, the gravity anomaly field direction is upward. The gravity anomaly is negative and symmetrical along $x = 0$. The anomaly value at $x = 0$ when the time is the smallest is approximately $-12.0g.u.$ When x approaches $\pm\infty$, the outlier infinitely approaches $0g.u.$
- (7) When $z = 375$ m, The section line is located below the model. As the model body is above the profile line, the gravity anomaly is negative and symmetrical along $x = 0$, and its value is the smallest at $x = 0$ (approximately $-3.0g.u.$). When x approaches $\pm\infty$, the outlier infinitely approaches $0g.u.$

The abnormal value on the center plane is 0, and the gravity anomaly values above and below the center plane are positive and negative, respectively. The outlier is the most positive on the top plane of the model and most negative on the bottom plane of the model. The abnormal value of the same profile line is the largest above the center point, and the abnormal curve is symmetrical. When x approaches $\pm\infty$, the abnormal value is infinitely close to 0. Below the center point, the abnormal value is the smallest, and the abnormal curve is symmetrical; when x approaches $\pm\infty$, the outlier is infinitely close to $0g.u.$

Spatial Distribution of Anomalies in Vertical High-Density Orebodies

Assuming that the length $2a$, width $2b$, and height $2c$ are 200, 40, and 500 m, respectively, for an upright cuboid, the residual density $\sigma = 1.0 \times 10^3 \text{ kg/m}^3$. The center point is $(0, 0, 0)$ and passing through different positions parallel to the xOy plane section provides the spatial distribution of the gravity anomaly of the vertical cuboid.

(1) $z = -260$ m indicates that the observation plane is located on the upper part of the model, as shown in **Figure 3B**.

Figure 3A shows the gravity anomaly characteristics of the vertical cuboid on the observation plane at $z = -260$ m:

- 1) The abnormal contour is elliptical; it is dense near the cuboid and sparse in the area away from the cuboid.
- 2) The entire upright cuboid is located below the observation plane; the residual density of the model is greater than 0, and the gravity anomaly is positive.
- 3) The contour map of the gravity anomaly is symmetric about the x and y axes, and the anomaly profile is a single-peak curve.

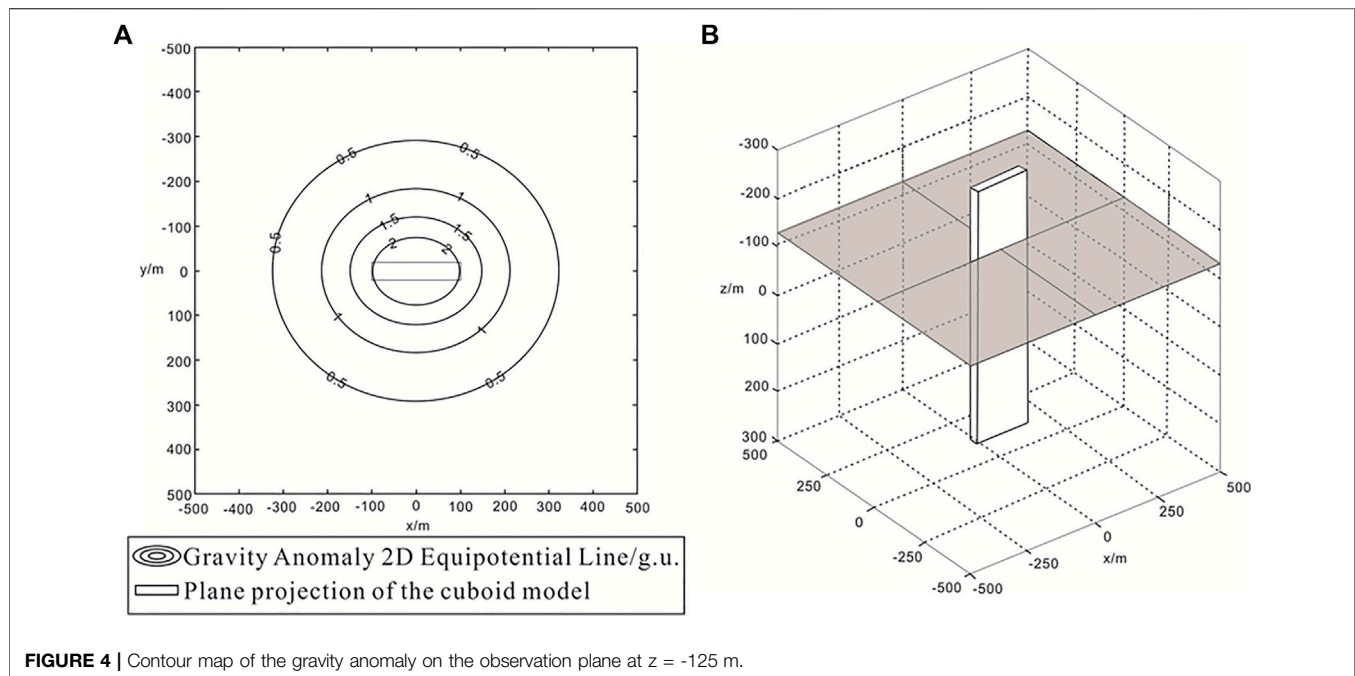


FIGURE 4 | Contour map of the gravity anomaly on the observation plane at $z = -125$ m.

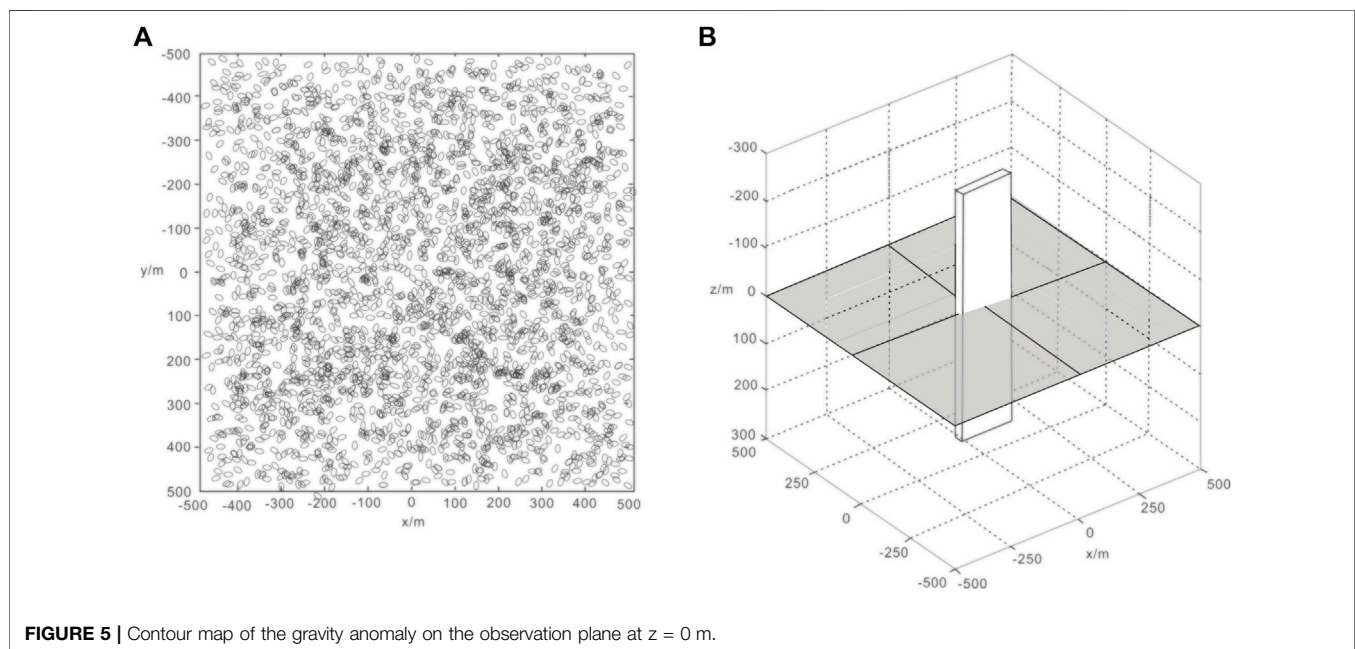


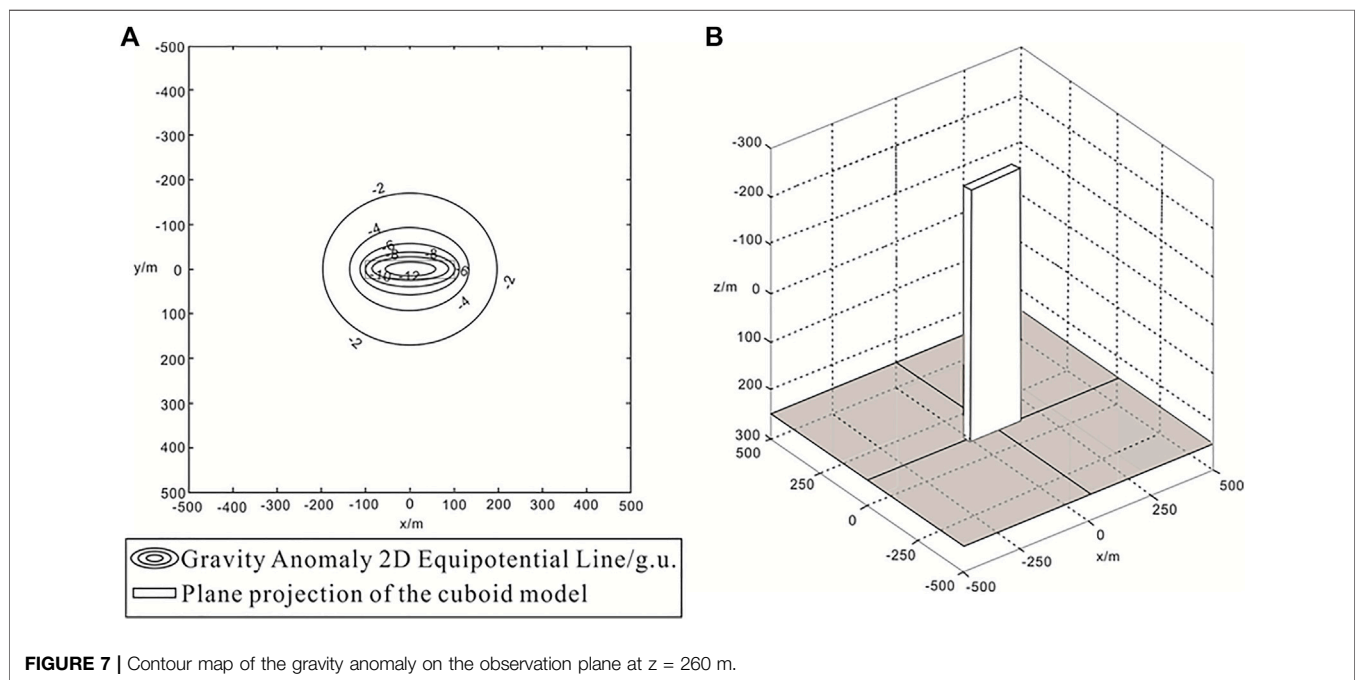
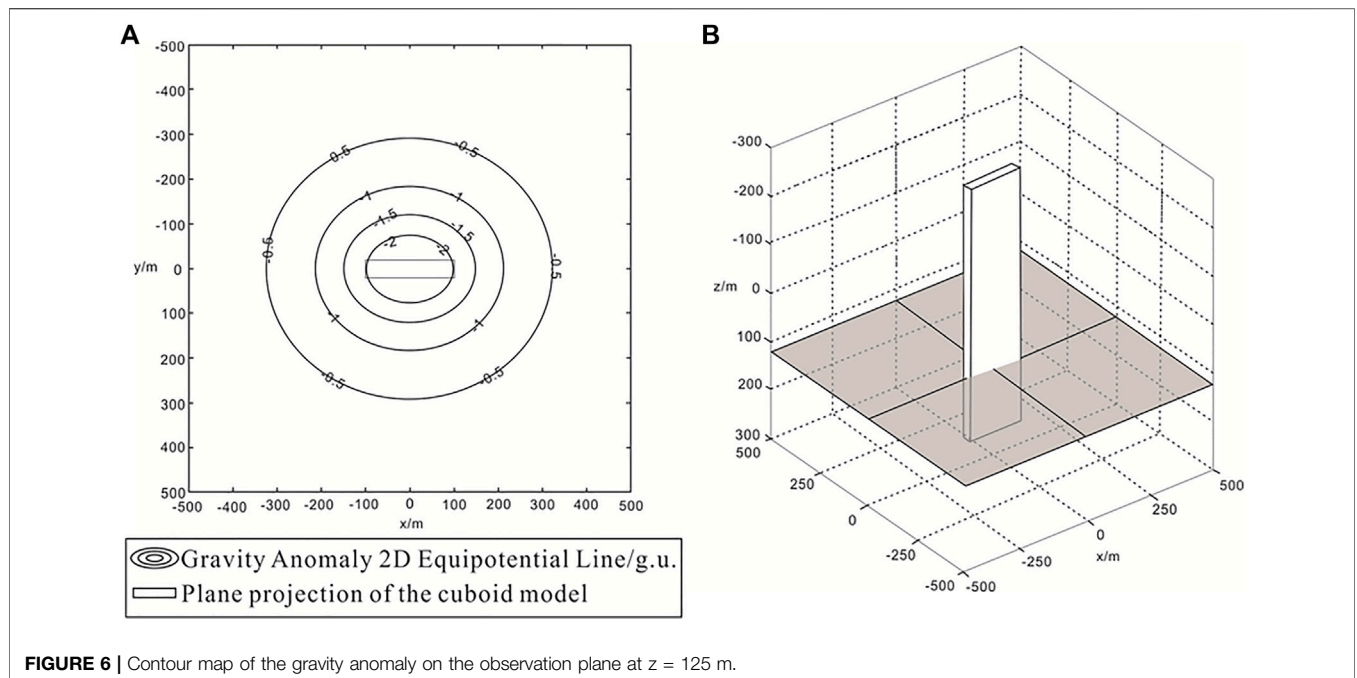
FIGURE 5 | Contour map of the gravity anomaly on the observation plane at $z = 0$ m.

4) The observation plane is 260m away from the center of the cuboid, and the maximum value of the anomaly is above the upright cuboid, approximately 12g.u.

(2) $z = -125\text{m}$ indicates that the observation plane is in the upper middle part of the model, as shown in **Figure 4B**.

Figure 4A shows the gravity anomaly characteristics of the oblique cuboid on the observation plane at $z = -125\text{m}$:

- 1) The abnormal contours near the cuboid are dense, and those far away from the cuboid are sparse.
- 2) The observation plane passes through the upright cuboid; the part of the upright cuboid below the observation plane is larger than that in the upper part; the residual density of the model is greater than 0, and the gravity anomaly is positive.
- 3) The contour map of the gravity anomaly is symmetrical about the x and y axes, and the anomaly curve is a bimodal curve.



4) The observation plane is 125 m away from the center of the cuboid, and the maximum value of the anomaly is at the center of the cross-section between the vertical cuboid and the observation plane, which is approximately $2.5g.u.$

(3) $z = 0m$ indicates that the observation plane is in the middle of the model, as shown in **Figure 5B**.

Figure 5A shows the gravity anomaly characteristics of the oblique cuboid on the observation plane at $z = 0m$:

1) The observation plane passes through the center of the cuboid; the upright cuboid is equal in the upper and lower parts of the observation plane, and the gravity anomaly is $0g.u.$

- 2) Owing to the presence of random errors, the contour map of the gravity anomaly appears disorderly and irregular.
- 3) The observation plane is located at the center plane of the cuboid, 0 m away from the center of the cuboid, and the anomaly is $0g.u.$

(4) $z = 125\text{ m}$ indicates that the observation plane is in lower middle part of the model, as shown in **Figure 6B**.

Figure 6A shows the gravity anomaly characteristics of the vertical cuboid on the observation plane at $z = 125\text{ m}$:

- 1) The abnormal contours near the cuboid are dense, and those far away from the cuboid are sparse.
- 2) The observation plane intersects the upright cuboid; the lower part of the upright cuboid is smaller than the upper part of the observation plane, and the gravity anomaly is negative.
- 3) The contour map of the gravity anomaly is symmetric about the x and y axes, and the anomaly curve is a single-peak curve.
- 4) The observation plane is 125 m away from the center of the cuboid, and the maximum negative outlier is approximately $-2.5g.u.$

(5) The observation plane at $z = 260\text{ m}$ is in the lower part of the model, as shown in **Figure 7B**.

Figure 7A shows the gravity anomaly characteristics of the upright cuboid on the observation plane at $z = 260\text{ m}$:

- 1) The abnormal contours near the cuboid are dense, and those far away from it are sparse.
- 2) The whole rectangular body is located above the observation plane, and the gravity anomaly is negative.
- 3) The contour map of the gravity anomaly is symmetric about the x and y axes, and the anomaly profile is a single-peak curve.
- 4) The observation plane is 260 m away from the center of the cuboid, and the maximum negative outlier is approximately $-12.0g.u.$

Overall, the anomaly contour is generally elliptical; it is dense near the cuboid, and sparser away from the cuboid. Its equipotential line is symmetric about the x and y axes. The anomaly values are 0, positive, and negative on, above, and below the central plane, respectively. The outlier values are most positive and most negative on the top and bottom planes of the model, respectively. The abnormal value is largest above the center point of the same plane; it gradually decreases from the center to the surroundings to a minimum value of $0g.u.$ The abnormal value is smallest below the center point; it gradually increases from the center to the surroundings to a maximum value of $0g.u.$

PREDICTION METHOD FOR THE DEEP EXTENSION OF A HIGH-DENSITY VERTICAL OREBODY

Method Steps

- (1) Various rocks and ores in the survey area were collected, identified, and their types were determined.
- (2) The densities of rocks and ores in the survey area were determined.
- (3) Gravity observations of the tunnel were made using point coordinate measurements, selecting gravity base points, and collecting point gravity anomaly observations.
- (4) Data processing methods (including the temperature, tilt, solid tide, zero drift, terrain, Bouguer, and latitude corrections, as well as the separation of regional and local fields) were applied for tunnel gravity observations.
- (5) Qualitative predictions of deep vertical high-density orebody extensions were made.
- (6) Quantitatively evaluated the extension of the deep vertical high-density orebodies.

Qualitative Prediction

According to the qualitative analysis from **Eq. 3** and the forward calculations of **Figure 3** to **Figure 7**, we obtained the deep extension characteristics of the high-density vertical orebody, as presented in **Supplementary Table S1**.

- (1) When the observation tunnel is in the upper middle part of the orebody, the residual gravity anomaly Δg is symmetrical and positive.

In particular, when the residual gravity anomaly of the observation tunnel is a symmetrical positive anomaly, the orebody extends to the deeper part, and the depth of the unknown orebody under the tunnel is greater than its height above the tunnel.

- (2) When the observation tunnel is in the middle of the orebody, the residual gravity anomaly Δg is zero.

In particular, when the remaining gravity of the observation tunnel is not abnormal, the depth of the unknown orebody below the tunnel is equal to its height above the tunnel.

- (3) When the observation tunnel is in the lower middle part of the orebody, the residual gravity anomaly Δg is symmetrical and negative.

In particular, when the residual gravity anomaly of the observation tunnel is symmetrical and negative, the extension of the orebody to the deep part is less; specifically, the depth of the unknown orebody below the tunnel is less than the height of the orebody controlled above the tunnel.

Quantitative Evaluation

Eq. 3 can be used to calculate the extended length of the high-density orebody and the number of prospective resources; namely, the half-length a and half-width b of the known high-density upright orebody can be substituted into **Eq. 3** to calculate and plot the residual gravity anomaly. In the residual gravity anomaly map, the value of the half-height c is continuously adjusted, and the shape and intensity of the residual gravity anomaly curve obtained after the calculation and measurement of the residual gravity anomaly curve are observed. When the two

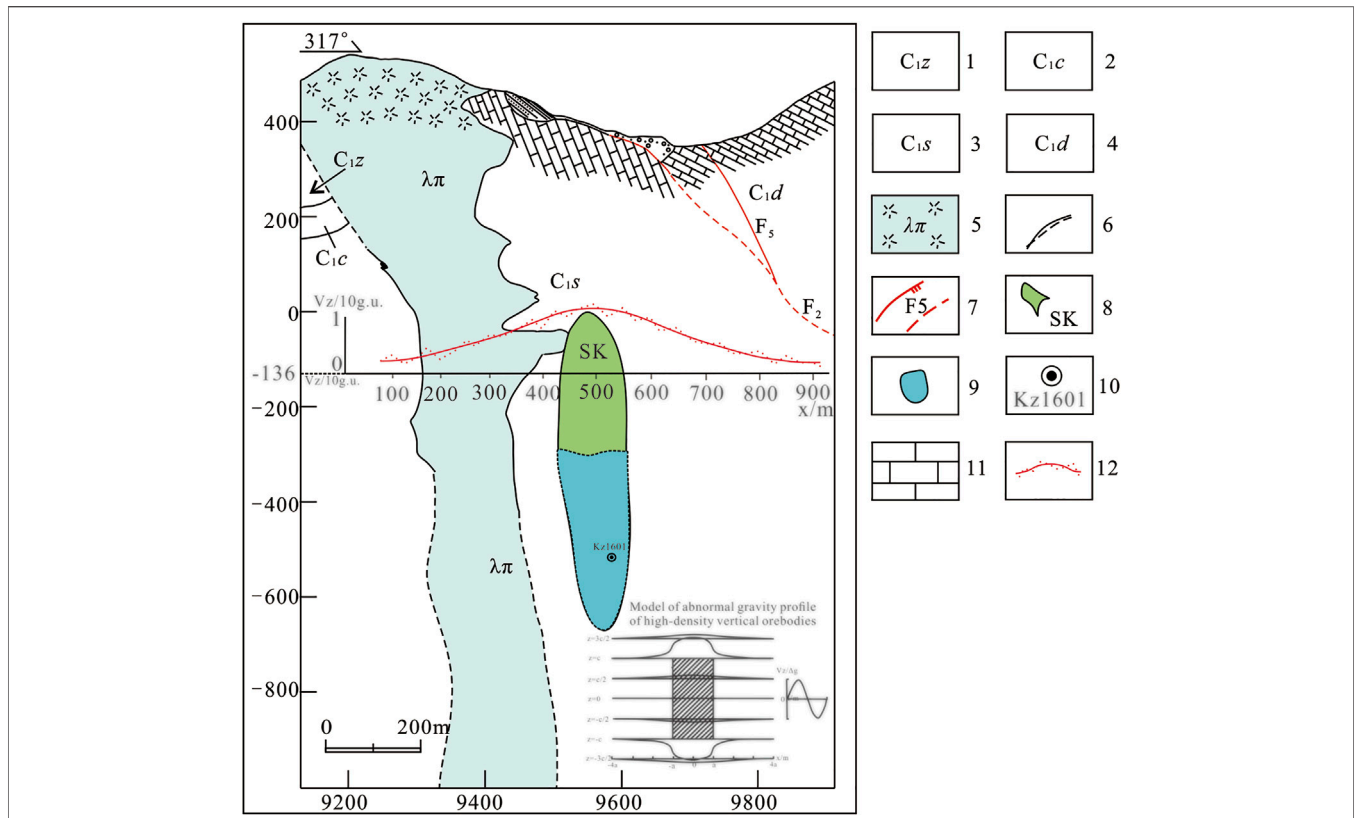


FIGURE 8 | Gravity profile of the geological tunnel at the 109 line in the middle section of -136 m in the Huangshaping mining area.

curves satisfy Eq. 4, the c value in this state represents the semi-extended length of the deep part of the orebody.

$$|V_z - V'_{z,m}|_{max} \leq Z_1 \tag{4}$$

where Z_1 is the threshold.

An estimated range for c can be established using Eq. 3 to calculate multiple estimated gravity anomalies for c within a given range $[c_1, c_2, \dots, c_n]$. The gravity anomalies are marked as $V'_{z,m} (m \in [1, n])$ to determine the difference between the measured gravity anomalies V_z and $V'_{z,m}$ and satisfy Eq. 5; the c value corresponding to the minimum square sum represents the orebody at z , namely, the half extension in this direction.

$$\min\left(\sum (V_z - V'_{z,m})^2\right) \leq Z_2 \tag{5}$$

where Z_2 is the threshold.

Furthermore, the prospective resources of high-density orebodies can be calculated.

Method Characteristics

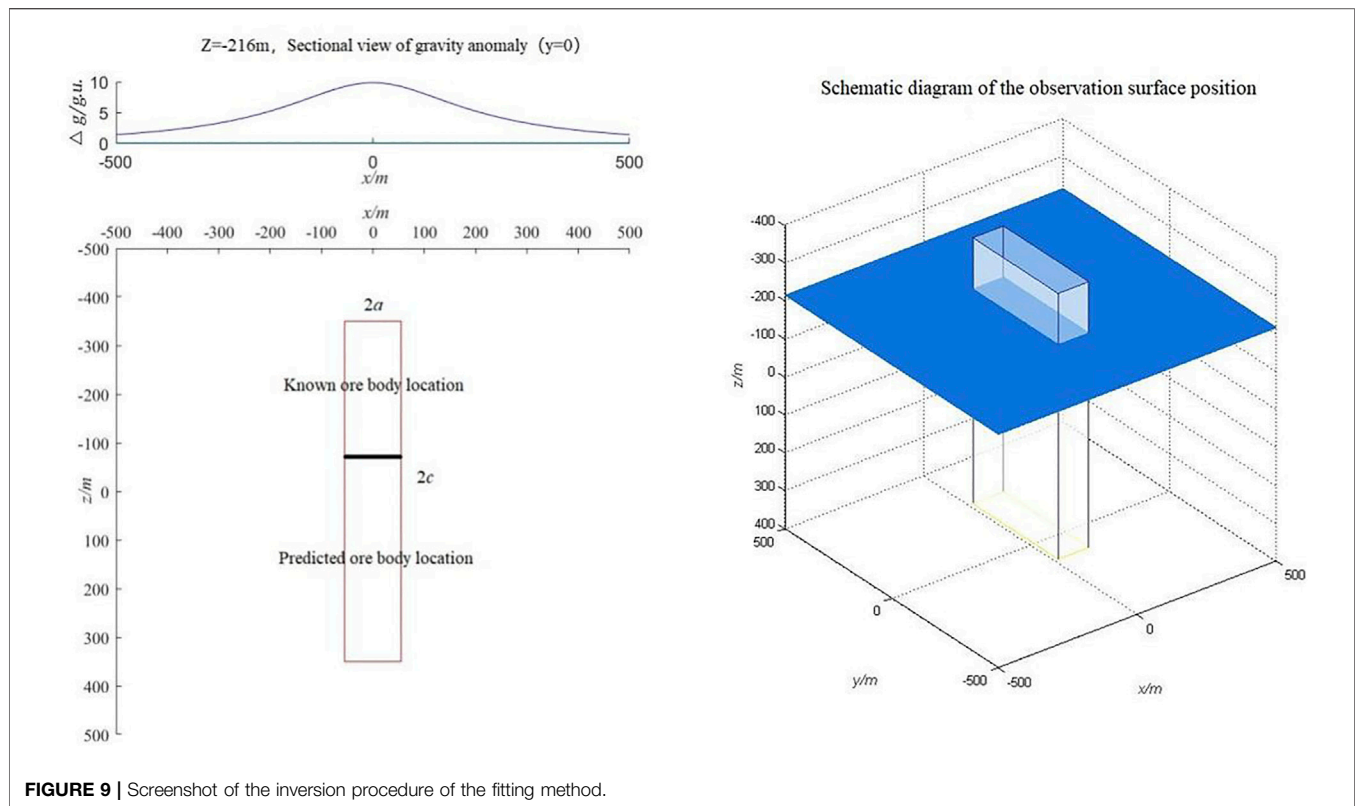
Based on the above principles, the method has the following characteristics:

- (1) Although few parameters are used, only one parameter is used for the remaining gravity anomaly measurements of the pit.

- (2) Good detection results are obtained for orebodies with large density differences between the ore and surrounding rock. Thus, it is applicable to both slab- or vein-type high-density orebodies and columnar high-density orebodies.
- (3) This method can quickly determine the deep extension of high-density upright ore bodies, which can significantly shorten the prospecting cycle and reduce the exploration cost.
- (4) It is unaffected by electromagnetic interference.
- (5) It is only slightly influenced by the surface topography.
- (6) The observation trench is basically flat and is anomalous without curvilinear leveling.
- (7) The gravity method is exceptionally simple and thus can be easily applied.

APPLICATION EXAMPLES

This study conducted gravity detection in the deep pit of the Huangshaping tungsten–tin polymetallic deposit in the Pingbao mining field in south Hunan and achieved good results. This deposit is a representative copper–tin polymetallic deposits in south Hunan, and its formation is closely related to the SK rocks in the contact zone between the early and middle Yanshan quartz porphyry and enclosing rocks in the area. The deposit is primarily controlled by the contact zone, and the main orebodies are in the



form of large veins and irregularities, primarily on both sides of the contact zone between the quartz porphyry and the crystalline tuff in the area. The deposit has been developed and utilized for decades, and the significant resource crisis necessitates the exploration of the deep orebody, which provides the prerequisite for gravity exploration in the pit.

The gravity measurement used a CG-5 high-precision gravimeter manufactured by the Canadian Scintrex Instrument Company. The resolution of the instrument is $0.01g.u.$, standard error is less than $0.05g.u.$, and direct reading range is $80,000g.u.$

The gravimeter was placed in the tunnel gravity measurement after passing the static, dynamic, stabilization time tests after adjusting the measuring range, and grid value calibration test.

The gravity base point of the tunnel was near the tunnel entrance, and the ground was stable and free of interference, which facilitated the observation of the base point in the morning and evening. The morning and evening base points were observed using the observation method of “base point-auxiliary reference point.” Base-measuring points from 1 to n were adopted by the observation method. The gravity measurement points of the tunnel were located using a plan view of the middle section of the tunnel combined with a measuring tape and a laser rangefinder, and the point distance was 25 m.

We performed temperature, tilt, solid tide, zero drift, terrain, Bouguer, and latitude corrections, as well as the regional field and local field separations using the tunnel gravity observation results. Subsequently, field separation of the Bouguer gravity anomaly was conducted. Finally, the results were analyzed to explain the obtained residual gravity anomaly and predict the deep extension of the orebody.

Among which, the correction of temperature, solid tide and tilt is automatically completed by the CG-5 gravimeter. The correction of near area terrain which ranges from 0 to 20 m in the tunnel is carried out by the cylinder formula (Wang, 2003). The height of the cylinder for terrain correction of near area terrain equals to the difference of elevation between the projection point on the ground and the survey point in the tunnel, which original elevation values comes from 1:2000 topographic map.

The cylinder formula of near area terrain correction for tunnel gravity measurement (Wang, 2003) shows below:

$$\Delta g = 2\pi G\sigma \left\{ H - \sqrt{h + H^2 + R^2} + \sqrt{h^2 + R^2} \right\} \times (10^6) \quad (6)$$

Where Δg is the near area terrain correction value of the tunnel (g.u.);

G is the universal Gravitational Constant, $6.67 \times 10^{-11}m^3/(kg \cdot s^2)$;

σ is the residual density, generally $2.67 \times 10^3 kg/m^3$;

h is the upper top depth, the unit is m;

H is the height of the cylinder (difference of elevation), the unit is m;

R is the radius of the cylinder, approximately 20 m.

The method of sector division correction is applied on middle area terrain correction. The elevation value of middle area terrain correction origins from the 1:2000 topographic map. The correction conducted by RGIS software which ranges from 20 to 2000m.

Formulas for sector division correction (Hammer, 1950) and supplementary angle correction show below:

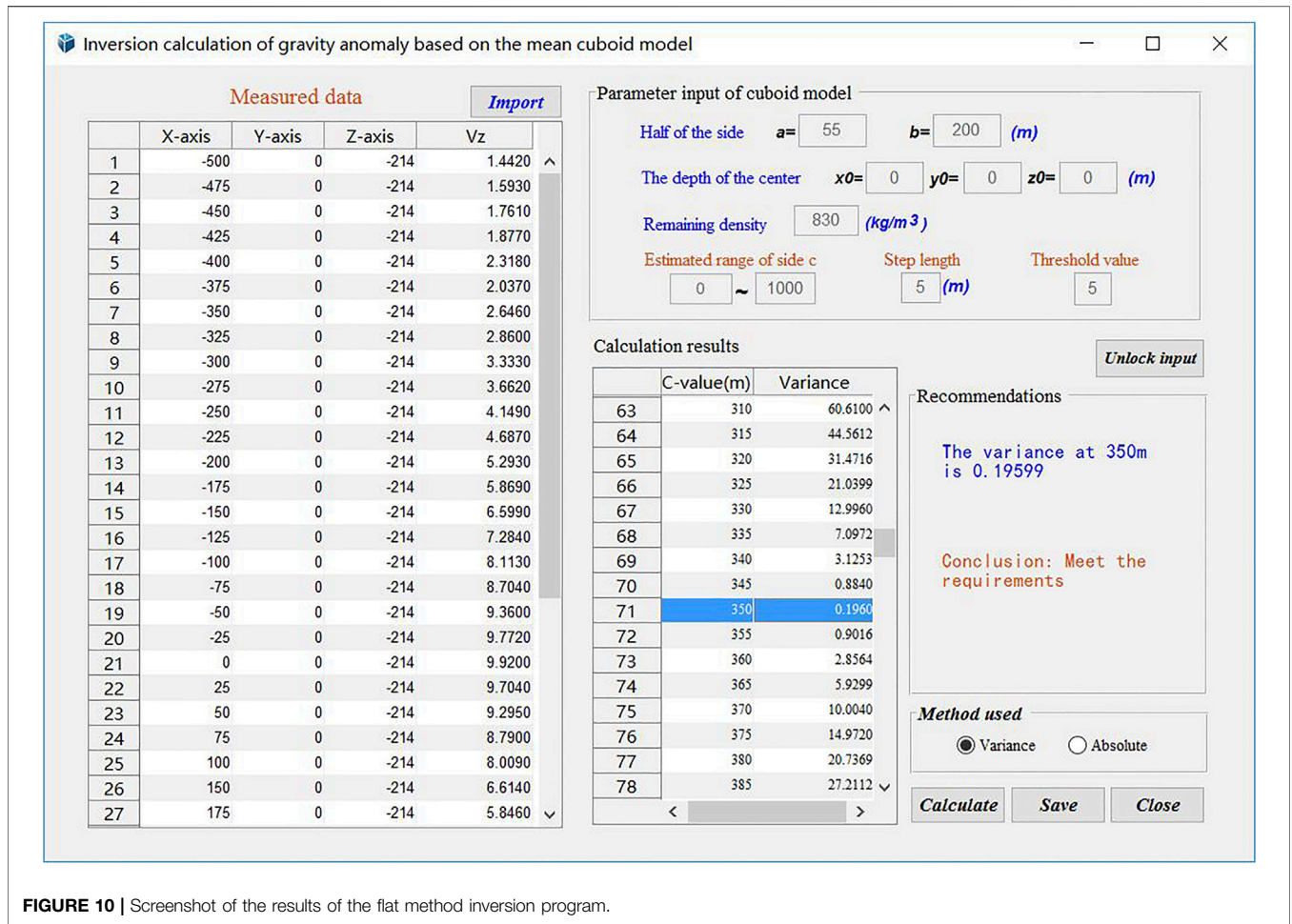


FIGURE 10 | Screenshot of the results of the flat method inversion program.

$$\delta g_s = \frac{2\pi G\sigma}{n} \left(R_{m+1} - R_m + \sqrt{R_m^2 + \Delta H^2} - \sqrt{R_{m+1}^2 + \Delta H^2} \right) (\times 10^6) \quad (7)$$

Where R_m is the inner radius of sector, which the unit is m;
 R_{m+1} is the outer radius of sector, which the unit is m;
 ΔH is the difference of elevation between the average elevation of the sector and the measuring point, which the unit is m;
 n is the number of orientations, that is the number of small sectors divided from near area terrain when the observation point is supposed to be vertex

$$\Delta g_2 = \frac{G\sigma S}{R} \left[1 - \frac{R}{\sqrt{R^2 + \Delta H^2}} \right] \quad (8)$$

Where S is the area of the supplementary angle, $S = \frac{1}{4} \cdot (4 - \pi)r^2$;
 R is the radius of inscribed circle, $1.105r = R$;
 ΔH is the difference of elevation between the average elevation of the supplementary angle and the measuring point, which the unit is m.

The elevation of remote area terrain correction origins from the RGIS software and is able to be adjusted by the RGIS software, which from 2 to 166.7 km.

The calculation formula of square field algorithm (Hammer, 1950) for remote area terrain correction shows below:

$$\Delta g_f = G\sigma \int_{\varepsilon_i - \frac{1}{2}\Delta\varepsilon}^{\varepsilon_i + \frac{1}{2}\Delta\varepsilon} d\varepsilon \int_{\eta_j - \frac{1}{2}\Delta\eta}^{\eta_j + \frac{1}{2}\Delta\eta} d\eta \int_0^{h_{ij}} \frac{\zeta}{(\varepsilon^2 + \eta^2 + \zeta^2)^{\frac{3}{2}}} d\zeta \quad (9)$$

Where (ε_i, η_j) are the coordinates of the center point in the square grid located in the i th row and the j th column;
 $\Delta\varepsilon$ is the increment in the x direction;
 $\Delta\eta$ is the increment in the y direction;
 h_{ij} is the elevation of the grid center point.

Gravity measuring point is generally located in the central of tunnel. Therefore the impact triggered by gravitational effect origins from gravity measuring point in the central tunnel is vanishingly small. That means further tunnel correction is no needed in terms of the remote area terrain.

The regional field and the local field of Bouguer gravity anomaly data in the central area of 136-mining area were separated by RGIS trend analysis module. The results express in **Supplementary Table S2**.

Twenty-eight rock and ore specimens were collected in the mining area, including eight skarn-type tin–tungsten polymetallic ores and twenty rocks (seventeen quartz rocks and three crystalline limestones). The average density of the skarn orebodies in the survey area was $3.45 \times 10^3 \text{ kg/m}^3$, whereas that of quartzite and crystalline limestone (wall rock) was $2.62 \times 10^3 \text{ kg/m}^3$; the average density of the ore was $0.83 \times 10^3 \text{ kg/m}^3$ higher than that of the surrounding rock. The orebodies in the measurement area were all high-density vein-like orebodies, with sufficient density difference for the application of this method.

Comprehensive analysis of the gravity profile of the tunnel at the 109-section line in the middle section of -136 m in the Huangshaping mining area (shown in **Figure 8**) revealed that the orebodies above the middle section of -136 m were controlled orebodies. The gravity anomaly curve measured in the middle section of -136 was symmetrical and evidently positive, and its anomalous position corresponded to a skarn orebody in the quartz porphyry contact zone, which is a high-density body. When the measurement section was located above the middle of the skarn-type orebody, a positive anomaly was observed. This indicated that the deep extension of the skarn-type orebody is greater than the extension above the middle section of -136. The occurrence of the orebody was primarily characterized as upright.

According to the known characteristics of the orebody above the tunnel, the strike length, thickness, density difference, and upper control height of the skarn-type orebody were approximately $2a = 110\text{m}$, $2b = 400\text{m}$, $0.83 \times 10^3 \text{ kg/m}^3$, and 136m , respectively. By comparing the two curves, we concluded that when $2c = 700\text{m}$, the outlier was approximately $10g.u.$, which roughly corresponded with the positive outlier shown in the middle section of -136 m (**Figure 9**). Therefore, we inferred that the outlier value at an orebody height of 700 m is approximately $10g.u.$, and the depth of the -136 middle section exceeds 564m .

Through the quantitative flat method inversion of **Eq. 5**, we concluded that when the half-height $c = 350\text{m}$, the difference was the smallest (0.19599) and met the measurement evaluation requirements. At this time, the total height of the orebody was approximately 700m (minus the upper control), with a depth of 136m . Thus, we inferred that the orebody extended to a depth of 564m (**Figure 10**).

When the result of the absolute value method was 350m , the difference was the smallest (0.3187) and met the measurement evaluation requirements. Therefore, the results of the two methods are consistent.

Based on the drilling of KZ1601, the ore was deep; this demonstrated the presence of concealed orebodies and the effectiveness of this method in their detection.

The tunnel-controlled orebody length was $2a = 110\text{m}$, orebody strike extension length was $2b = 400\text{m}$, orebody height calculated by inversion was $2c = 700\text{m}$, and the average density of the skarn-type orebody obtained by the actual density

measurement was $3.45 \times 10^3 \text{ kg/m}^3$. Thus, the ore volume of the orebody can be predicted as $110\text{m} \times 400\text{m} \times 700\text{m} \times 3.45 \times 10^3 \text{ kg/m}^3 \times 10^{-3} = 96712000\text{t}$.

CONCLUSION

The tunnel gravity prediction method for evaluating the deep extension of an orebody is suitable when the ore and surrounding rock have a significant density difference. It can effectively predict the extension depth and number of ores in the deep orebody. This method involves few parameters, is simple and unaffected by electromagnetic interference, and has high detection accuracy. It can overcome the ambiguity of other geophysical methods for the spatial positioning and detection of orebodies. This method was verified by deep exploration in the Huangshaping mining area. Therefore, this method has important application potential and promotion value for predicting prospective deep resources of the same type of high-density vertical orebodies.

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/**Supplementary Material**, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

All authors listed have made a substantial, direct, and intellectual contribution to the work and approved it for publication.

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SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/feart.2022.897909/full#supplementary-material>

REFERENCES

Beyer, L. A. (1979). Terrain Corrections for Borehole Gravity Measurements. *Geophysics* 44 (9), 1584–1587. doi:10.1190/1.1441028

Cai, B. (1980). Introduction to Foreign Well-Well Gravimeters and Well-Well Gravity Measurements[J]. *Geophys. Geochem. Explor.* 4 (1), 61–63.

Cao, L. (2011). Application and Prospect of Geophysical Methods in Deep Prospecting for Metal Deposits[J]. *Prog. Geophys.* 26 (2), 701–708. doi:10.3969/j.issn.1004-2903.2011.02.040

- Casten, U., and Gram, C. (1989). Recent Developments in Underground Gravity Surveys. *Geophys prospect* 37 (1), 73–90. doi:10.1111/j.1365-2478.1989.tb01822.x
- Fan, X. (2007). Fast Calculation of Terrain Correction in the Vicinity of Regional Gravity Survey [J]. *J. Eng. Geophys.* 4 (6), 560–562. doi:10.3969/j.issn.1672-7940.2007.06.009
- Fu, J.-W. (2010). Progress in Gravity Logging Technology[J]. *Prog. Geophys.* 25 (2), 596–601. doi:10.3969/j.issn.1004-2903.2010.02.032
- Geng, Q. (2016). Representative Foreign Products of Gravity Instruments and the Latest Progress in Domestic Research and Development [J]. *Geol. Equip.* 17 (01), 27–30. doi:10.3969/j.issn.1009-282X.2016.01.006
- Goodacre, A. K. (1973). Some Comments on the Calculation of the Gravitational and Magnetic Attraction of a Homogeneous Rectangular Prism. *Geophys Prospect* 21 (1), 66–69. doi:10.1111/j.1365-2478.1973.tb00014.x
- Hammer, S. (1950). Density Determinations by Underground Gravity Measurements. *Geophysics* 15, 637–652. doi:10.1190/1.1437625
- Han, R., and Li, W. (2014). A Method of Locating and Detecting High-Density Concealed Ore Bodies in the Whole Space Domain of Tunnel Gravity[P].11.
- Han, R., Li, W., Cheng, R., Wang, F., and Zhang, Y. (2020). 3D High-Precision Tunnel Gravity Exploration Theory and its Application for Concealed Inclined High-Density Ore Deposits. *J. Appl. Geophys.* 180, 104119. doi:10.1016/j.jappgeo.2020.104119
- Han, R., Li, W., Cheng, R., and Wang, J. (2021). 3D High-Precision Tunnel Gravity Exploration Method for Locating and Detecting Deeply Concealed Inclined High-Density Orebodies, [P].
- Han, R., Wang, F., Hu, Y., Wang, X., Ren, T., Qiu, W., et al. (2014). Study on the Dynamics of the Metallogenic Structure of the Huize-type (HZT) Germanium-Rich Silver-Lead-Zinc Deposit and its Chronological Constraints[J]. *Tect. Metallogeny* 38 (04), 758–771. doi:10.3969/j.issn.1001-1552.2014.04.003
- Howell, L. G., Heintz, K. O., and Barry, A. (2012). The Development and Use of a High-Precision Downhole Gravity Meter. *Geophysics* 31 (4), 764–772. doi:10.1190/1.1439808
- Hu, M. (2015). “Research and Application of Terrain Correction Methods in Gravity Exploration [D],” (Chengdu, Sichuan, China: Chengdu University of Technology). Dissertation.
- Li, W., and Han, R. (2014). A Method of Locating and Detecting Low-Density Concealed Ore Bodies in the Whole Space Domain of Tunnel Gravity[P].11.
- Li, X., and Chouteau, M. (1998). Three-dimensional Gravity Modeling in All Space [J]. *Surv. Geophys.* 19 (4), 339–368. doi:10.1023/a:1006554408567
- Liu, Z., and Liu, S. (2014). *Mine Geophysical Prospecting[M]*. Wudaokou, Haidian, China: China University of Mining and Technology Press.
- Lu, L., Zhou, M., Bruce, L., et al. (2018). Common Problems and Solutions in Field Work of CG-5 Gravimeter. [J]. *Geol. Equip.* doi:10.3969/j.issn.1009-282X.2018.04.012
- Luo, D., and Liu, Z. (2012). Research on the Forward Modeling Method of Gravity in Borehole Based on the Internal Theory of Field Sources[J]. *Comput. Tech. Geophys. Geochem. Explor.* 34 (6), 656–665. doi:10.3969/j.issn.1001-1749.2012.06.06
- Luo, Y. (2008). Two New Rectangular Gravitational Force Field Forward Expressions and Theoretical Derivation [J]. *J. Eng. Geophys.* 5 (2), 210–214. doi:10.3969/j.issn.1672-7940.2008.02.015
- Ming, Y., Niu, X., Xie, X., Han, Z., Li, Q., and Sun, S. (2021). Application of Gravity Exploration in Urban Active Fault Detection[J]. *IOP Conf. Ser. Earth Environ. Sci.* 660 (1), 012057. doi:10.1088/1755-1315/660/1/012057
- Nagy, D., Papp, G., and Benedek, J. (2000). The Gravitational Potential and its Derivatives for the Prism[J]. *J. Geod.* 74 (7-8), 552–560. doi:10.1007/s001900000116
- Okabe, M. (1979). Analytical Expressions for Gravity Anomalies Due to Homogeneous Polyhedral Bodies and Translations into Magnetic Anomalies. *Geophysics* 44 (4), 730–741. doi:10.1190/1.1440973
- Schmoker, J. W. (1980). Terrain Effects of Cultural Features upon Shallow Borehole Gravity Data. *Geophysics* 45 (12), 1869–1871. doi:10.1190/1.1441071
- Smith, N. J. (1950). The Case for Gravity Data from Boreholes. *Geophysics* 15, 605–636. doi:10.1190/1.1437623
- Tao, Wei. (2017). “Research on High-Precision Gravity Terrain Correction of Tunnels[D],” (Kunming, China: Kunming University of Science and Technology). Dissertation.
- Teng, L., Ni, S., Zhang, B., Di, B., Zhu, H., Yin, Q., et al. (2013). Application of CORS System in Gravity Survey[J]. *Geophys. Geochem. Explor.* 37 (06), 1018–1022. doi:10.11720/j.issn.1000-8918.2013.6.11
- Wang, J. (2019). Influencing Factors of High-Precision Tunnel Gravity Exploration in the Huangshaping W-Sn-Cu-Pb-Zn Polymetallic Mining Area, Southern Hunan. [C]. *Acta Geol. Sin. Engl. Ed.* 93 (Suppl. 2), 418.
- Wang, Q. (2003). *Gravity [M]*. Beijing: Geological Press.
- Wang, W., Ren, F., Wang, Y., Ji, X., Li, J., and Li, Q. (2014). Application of Gravity Exploration in the Investigation of Sedimentary Bauxite Deposits[J]. *Geophys. Geochem. Explor.* 38 (03), 409–416. doi:10.11720/wtyht.2014.3.01
- Wu, Z. (2011). “Tunnel Gravity Research and its Application in the Prospecting of Lead-Zinc Deposits [D],” (Kunming, China: Kunming University of Science and Technology). Dissertation.
- Xu, G., and Zhou, G. (1989). *Gravity Measurement in Wells and tunnels[M]*. Beijing: Beijing Geological Publishing House.
- Yang, H., Ding, H., Wang, Y., and Cao, Q. (2000). Discussion on Several Issues in Gravity Correction in Mountainous Areas [J]. *Pet. Geophys. Prospect.* (4), 479–487. doi:10.3321/j.issn:1000-7210.2000.04.010
- Yi, B., Zeng, Z. F., and Xue, J. (2008). Application of Geophysical Method in City Active Fault Detection. *Prog. Geophys.* 23 (02), 599–604. doi:10.3969/j.issn.1672-562X.2019.04.079
- Zeng, H. (1999). Current Status and Development of Gravimeter[J]. *Geophys. Geochem. Explor.* 23 (2), 84–89. doi:10.3969/j.issn.1000-8918.1999.02.002
- Zeng, H. (2008). *Gravity Field and Gravity Exploration [M]*. Beijing: Geological Publishing House, 85–86.
- Zhang, D. (2017). “Application of High-Precision Gravity in the Prospecting of Xiaozhuqing Survey Area in Huize Lead-Zinc Mining Area,” (Kunming, China: Kunming University of Science and Technology). Dissertation.
- Zhang, G., Zhao, G., Kuang, H., Teng, F., Su, Y., Liang, J., et al. (2015). Some Advances in Ground Gravity Survey Methods and Techniques in Recent Years [J]. *Prog. Geophys.* 30 (1), 386–390. doi:10.6038/pg20150156
- Zhang, Y., Xu, R. G., and Yu, Y. (2012). The Application of High Precision Gravity Prospection in the Detection of Concealed Faults. *Recent Dev. World Seismol.* (05), 9–10. doi:10.3969/j.issn.1674-7801.2012.03.021
- Zhang, Z. (2012). Application of Gravity and CSAMT Method for Lead-Zinc Mining in Caixia Mountain, Xinjiang. *Mineral. Explor.* 3 (3), 389–396. doi:10.3969/j.issn.0235-4975.2012.05.002
- Zhdanov, M. S., Liu, X., Wan, L., Čuma, M., and Wilson, G. A. (2011). “3D Potential Field Migration for Rapid Imaging of Gravity Gradiometry Data — A Case Study from Broken Hill, Australia, with Comparison to 3D Regularized Inversion,” in *SEG Technical Program Expanded Abstracts*. doi:10.1190/1.3628202

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