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# Revealing the shallow soil structure of the Yigong Lake in the Tibetan Plateau using a portable distributed acoustic sensing interrogator

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The harsh and extreme environmental and near surface conditions of the Tibetan Plateau have limited the conventional electrical-based seismic instruments from obtaining high-quality seismic data through long-term and continuous observations, setting challenges for environmental seismology study and natural hazard monitoring in this area. Distributed acoustic sensing (DAS) is an emerging technique based on optical fiber communication and sensing. It provides a possible solution for subsurface imaging in extreme conditions at high spatiotemporal resolution by converting fiber-optic cables into dense seismic strainmeters. We deploy two survey lines with armored optical fiber cables in the Yigong Lake area, Southeastern Tibetan Plateau, to record ambient noise for a week. The DAS interrogator is specifically designed in a portable size with very low power consumption (25 W/h). Hence, we can use a 12V-DC battery for power supply to adjust the power limitation during the field recording. Ambient noise interferometry and multichannel analysis of surface waves are used to get 2D shear wave velocity profiles along the fiber paths. The results highlight the great potential of DAS for dynamic monitoring of the geological evolution of lakes and rivers in areas of extreme environments as in the Tibetan Plateau.

#### KEYWORDS

distributed acoustic sensing, sedimentary thickness, Tibetan Plateau, ambient noise tomography, seismic interferometry

# Introduction

Seismic data acquisition is the foundation of data processing and interpretation in seismological research. High-quality seismic data may significantly improve the accuracy of geological understanding. Recently, seismologists have tried to use denser and denser observational networks to obtain highspatiotemporal-resolution seismic data for a better knowledge of the Earth's dynamic processes and monitoring of natural hazards (Karplus and Schmandt, 2018; Nishikawa et al., 2019; Kohler et al., 2020). Although conventional electrical seismic instruments can be deployed in most areas for seismological research, it is very challenging to meet the aims of long-term monitoring in harsh environments, such as the volcano, glaciers, and plateaus due to the high cost of deployment and maintaining.

Distributed acoustic sensing (DAS) is an emerging seismic measurement technique that benefits from the development of optical fiber communication and sensing. Typically, a DAS system is composed of a standard optical fiber cable and an interrogation unit (IU). The IU lights the fiber using short laser pulses and measures Rayleigh back-scattering (RBS) along the fiber through a form of coherent optical time domain interferometry. A phase shift in the RBS is caused by the compression or elongation along the fiber, which is proportional to axial strain along the cable. Thus, DAS is working as a dense network for dynamic strain measurements (Parker et al., 2014; Hartog, 2017). DAS has initially been used for downhole applications in oil and gas exploration (Mateeva et al., 2014; Daley et al., 2016; Lellouch and Biondi, 2021). Currently, it has been widely adapted to seismological investigations, such as earthquake detection (Lindsey et al., 2017; Jousset et al., 2018; Wang et al., 2018; Yu et al., 2019), near-surface characterization (Dou et al., 2017; Fang et al., 2020; Spica et al., 2020; Shao et al., 2022), observing oceanic and atmospheric phenomena (Lindsey et al., 2019; Sladen et al., 2019; Williams et al., 2019; Zhu and Stensrud, 2019; Shinohara et al., 2022), monitoring volcano and glacier (Booth et al., 2020; Walter et al., 2020; Klaasen et al., 2021; Nishimura et al., 2021), and so forth. For the most up-to-date review of DAS application in seismology, we refer the readers to Lindsey and Martin (2021).

In this study, we use DAS to record continuous ambient noise in the Yigong Lake area in Tibetan Plateau. Yigong Lake is a barrier lake in the Yigong village, Bome County, Nyingchi, Tibet, characterized by a high altitude, a weak geological foundation, and an active modern crustal movement. It is located in the southeast of the Tibetan Plateau, about 2,600 m above sea level and spanned about 20 square kilometers (Figure 1). The strong uplift of the Earth's crust makes it the heaviest and wettest region on the Tibetan Plateau and prone to natural disasters such as landslides, avalanches, and debris flows. Historic records show that two giant landslides occurred in 1,900–2,000 resulting a deposit of total volume about  $5 \times 10^8$  m<sup>3</sup> and  $3 \times 10^8$  m<sup>3</sup>,





respectively (Shang et al., 2003; Zhou et al., 2016). A recent study utilizing radiocarbon dating of detrital materials provides pieces of evidence that there were at least eight large landslide events in this area over the past 5,500 years (Guo et al., 2020). As a result, these repeated landslides can change the sedimentary structure and further influence the development of river profiles. Revealing the sedimentary thickness could help us to better understand the geological evolution of the Yigong Lake. However, we should note that most of the previous studies in this area were focused on using satellite remote sensing technologies to study landslides, and these methods were hard to detect the underground structures. Due to the high content of water in the soil, it is infeasible to use the traditional electronic seismic instruments for data acquisition (Zhang et al., 2021). In addition, acquiring drilling data is very expensive. As a result, there is no previous study in this area for shallow structure characterization using seismic data. The emergence of DAS enables us to obtain field seismic data for the first time.

In the following parts, we first introduce the field deployment of seismic data acquisition using DAS. Thereafter, seismic interferometry is applied to the preprocessed data to obtain the noise correlation functions (NCFs) following the workflow of Bensen et al. (2007). After extracting surface waves from NCFs, multichannel analysis of surface waves (MASW) (Park et al., 1999) is utilized to retrieve dispersion curves and invert the shear wave velocity. According to the variations of the shear wave velocity, the sedimentary thickness in the study area can be inferred. These results confirm the reliability of DAS in capturing data and performing near-surface characterization in extreme environments.

## DAS data acquisition and preprocessing

The ambient noise data acquisition is located between a highway and the Yigong Lake, as shown in Figure 2. In this area, it is not readily available to get access to the industrial power supply, and there is no facility room for equipment installation. To perform seismic data acquisition under such an extreme condition, the DAS IU is specially manufactured to have a portable size of  $150 \text{ mm} \times 300 \text{ mm} \times 110 \text{ mm}$  (Figure 3A), much smaller than the already commercialized IUs, such as from Silixa (444 mm × 518 mm × 178 mm) or Schlumberger  $(559 \text{ mm} \times 955 \text{ mm} \times 244 \text{ mm})$ . The full-load power consumption is 25 W, lower than an ordinarv laptop. Because of these adjustments, we are able to use a 12V-DC battery for the power supply, as shown in Figure 3B. More details about the IU are referred to Xu T et al. (2021) and Zhang et al. (2021).

DAS data were acquired during the local daytime from 10 a.m. to 7 p.m. on 28 October 2021 to 3 November 2021. The maximum continuous recording length is 8 h because of running out of battery power. We installed two trenched DAS arrays. The surface trenches, which are 10cm–20 cm deep, are dug with hand tools. The optical fiber cables were backfilled with excavated soil to ensure coupling between DAS and the ground, as shown in Figure 4A. Line I is 1.355 km long with 755 channels. A 0.497 km-long session of the line was settled on the ground while the rest of the line was deployed underwater across the Yigong Lake. This line was originally installed in April 2021, when the Yigong Lake is in the dry season so that we could easily get access to the lake bottom to deploy optical fiber cables. During the data acquisition time, the Yigong Lake is in the wet season and optical fiber cables are





buried underwater. At the transition zone from the land to the water, parts of optical fiber cables are suspended in the air, as shown in Figure 4B. Consequently, the DAS data are discontinuously recorded during the test stage. After realizing this, we buried the suspended parts using the sand nearby, as shown in Figure 4C. Line II is 0.66 km long with 410 channels. Most parts are located on land, and only a small portion was across a water pool.

We select a 5-h DAS dataset in Line I and an 8-h DAS dataset in Line II for analysis, respectively. The sampling rate is 1000 Hz, and the gauge length is 10 m. The channel spacing is 1.6 m. We drop the data in the first 10 channels of Line I and the first 50 channels of Line II because of the contamination by the wandering man-walks.

Following Bensen et al. (2007), we process the DAS data with the workflow as summarized in Figure 5. The long-time raw data are split into short-time windows of 30-min long.

The two ends of the segmented time series became discontinuous. We performed 10% tapering to reduce the numerical oscillation known as the Gibbs effect when filtering and performing Fourier transforms and to minimize the impact of discontinuities at both ends of the signal (Figure 6) (Xu Y et al., 2021). The influence of the DC component and linear trend is removed by de-meaning and de-trending. According to the spectrum of the original data in Figure 6, the energy is mainly concentrated below 25 Hz. We down-sample the data to 50 Hz to reduce the computational time and the memory requirements. The corresponding resampled waveform and its spectrum are shown in Figure 7. When seismic data are collected in the field, it is inevitable to be affected by environmental and human factors, resulting in a high amplitude in the time-domain signal, as shown in the red box in Figure 7.







The introduction of temporal normalization can eliminate the effect of strong signals and improve the quality of the correlation results (Bensen et al., 2007). By selecting a time window of a certain length, the average of the absolute amplitude value is obtained in the time window, which is used as the weight of the center of the time window. The weight of all data is calculated by sliding the window, and then the original data is divided by the corresponding weight. Finally, the normalized sequence in the time domain is obtained. This procedure can be formulated as follows:

$$w(i) = \frac{1}{2N+1} \sum_{j=i-N}^{j=i+N} d(i), \qquad (1)$$



### FIGURE 8

(A) Waveform, and (B) corresponding spectrum of the 30min record of Channel 200 in Line II after temporal normalization. The red box denotes that the signal still has unbalanced high and low-frequency energy after the temporal normalization.



$$\bar{d}(i) = \frac{d(i)}{w(i)},\tag{2}$$

in which d(i) is the seismograph that has been band-pass filtered between 0.3 and 23Hz,  $\bar{d}(i)$  is the temporal normalized seismograph, and 2N + 1 is the number of points in the normalization window. Figure 8A shows that the strong energy effect has been effectively reduced after the runningabsolute-mean normalization. After the temporal normalization, the signal still has unbalanced high and low-frequency energy, and the signal spectrum is not flat, as shown in the red box in Figure 8B. Therefore, it needs to be "flattened", also called spectrum whitening. The Hanning window W(f) with a 0.3–23 Hz pass-band is applied to the smoothed complex spectrum D(f) of the signal time series d(t):

$$Y(f) = W(f) \times D(f) / |\bar{D}(f)|, \qquad (3)$$

in which  $|\overline{D}(f)|$  is the smoothed amplitude spectrum. The Fourier transform Y(f) is inversely transformed to the time domain for subsequent data processing. Figure 9 shows the final processed waveform and its spectrum. The resulting signal is more like white noise with a "flat" spectrum in the frequency domain.

### **Methods**

# Ambient noise interferometry for surface wave reconstruction

Surface waves are reconstructed by using ambient noise interferometry, which has been of great interests in variant research areas (Claerbout, 1968; Weaver and Lobkis, 2001; Shapiro et al., 2005). By cross-correlating the continuous recorded ambient noise at two receiver locations, empirical Green's functions between these two receivers can be obtained. After temporal stacking of the correlations in a long range, coherent waveforms will emerge (Campillo and Paul, 2003; Shapiro and Campillo, 2004). In the frequency domain, the cross-correlation can be formulated as:

$$C_{AB}(\omega) = \langle u_A^*(\omega) u_B(\omega) \rangle, \qquad (4)$$

where  $u_A(\omega)$  and  $u_B(\omega)$  denote the pre-processed wavefields in the frequency domain recorded at receiver A and receiver B, respectively, following the workflow in Figure 5. The symbol  $\omega$  is the angular frequency, \* denotes complex conjugate, and  $\langle \cdot \rangle$ denotes temporal average.

# Multichannel analysis of surface waves for shear wave velocity inversion

After the surface wave reconstruction using ambient noise interferometry, we apply MASW (Park et al., 1999) to extract dispersion curves. First, the data in the time-offset domain are slant-stack-transformed into the frequency-velocity domain as follows:

$$\mathcal{E}(\omega,c) = \int e^{-\frac{i\omega}{c}x} \frac{U(x,\omega)}{|U(x,\omega)|} dx,$$
(5)



in which  $E(\omega, c)$  is the dispersion image, c is the phase velocity at the angular frequency  $\omega$ ,  $U(x, \omega)$  is the amplitude spectrum of the timedomain data u(x, t) through Fourier transform. Next, the dispersion curves are identified by picking the maximum value in the dispersion image  $E(\omega, c)$  at each frequency-velocity pair.

The extracted dispersion curves are further used to invert the shear wave velocity. The misfit function is defined as minimizing the difference between the picked dispersion curves and the theoretical dispersion curves calculated using the fast delta-matrix algorithm (Buchen and Ben-Hador, 1996). Considering the non-linearity and non-convexity of the inverse problem, a Mente Carlo-based global search method is used to potentially escape from being stuck at the local minimum. Detailed description of the software used in this study can be found in Olafsdottir et al. (2018) and Olafsdottir et al. (2020).

# Results

In Line I, data from Channel 11 to Channel 360 are used for calculating NCFs. From Channel 11 to Channel 185, we select every channel as a virtual source, and the other 175 channels on the right of the virtual source are cross-correlated with the virtual source. From Channel 186 to Channel 360, we select every channel as a virtual source, and the other 175 channels on the left of the virtual source are cross-correlated with the virtual source are cross-correlated with the virtual source are 175 channels on the left of the virtual source are cross-correlated with the virtual source. In Line II, data from Channel 51 to Channel 350 are used for calculating NCFs. From Channel 51 to

Channel 200, we select every channel as a virtual source, and the other 150 channels on the right of the virtual source are cross-correlated with the virtual source. From Channel 201 to Channel 350, we select every channel as a virtual source, and the other 150 channels on the left of the virtual source are cross-correlated with the virtual source. Next, we perform linear stacking of all the time segments for each cross-correlation pair to average the effect of temporal noise and spatial irregularity. Finally, we get 350 virtual gathers for Line I and 300 virtual gathers for Line II.

Figure 10 shows the NCFs in Line I at Channel 11 (Figure 10A) and Channel 175 (Figure 10B), and in Line II at Channel 51 (Figure 10C) and Channel 200 (Figure 10D) as virtual sources, respectively. From the reconstructed virtual gathers, it can be seen that NCFs are bilateral time functions divided into positive and negative half lags, representing positive and negative wave propagation directions, respectively. There is an obvious asymmetry that the dispersive Rayleigh-wave on the positive lag is much stronger than that on the negative lag, suggesting that the azimuth of the ambient noise is not uniformly distributed. According to the schematic layout in Figure 2, we can infer that the higher energy results from the traffic noise on the highway. Similar asymmetry is also observed by Zeng et al. (2017) from field DAS data.

Before the dispersion curve extraction, the positive and negative half branches of NCFs are stacked in reverse order, and the positive half lags are reserved for subsequent processing to improve the signalto-noise ratio (Huang et al., 2021). Such a procedure can help to



mitigate the influence of uneven energy distribution of noise sources on surface wave velocity measurements.

The initial phase velocity picks are automatically chosen as the maximum value of the dispersion image in each frequency bin. Then, they are manually assigned with mode numbers and spurious picks are removed.

For the inversion of each 1D dispersion curve, we set the depth of the bottom of the interface to 90 m, and the whole model is divided into six layers. The thickness of the first layer is 8 m, the thickness of the second layer is 10 m, and the thickness of all other layers is 18 m. Since both P-wave velocity and density can affect the dispersion curve, the empirical formula of Brocher (2005) is used to convert the shear wave velocity into P-wave velocity and density. Finally, we obtain the 2D shear wave velocity profiles using cubic spline interpolation.

Figure 11 shows the dispersion images of NCFs at Channel 11 (Figure 11A) and Channel 175 (Figure 11B) in Line I, and at Channel 51 (Figure 11C) and Channel 200 (Figure 11D) in Line II, respectively. Although fundamental- and higher-mode surface waves are visible, only the fundamental mode dispersion curves, the black lines in Figure 11, are used for inversion.

Figure 12 shows the 1D dispersion curve inversion results of the fundamental mode in Figure 11. The red lines denote the best inverted velocities which have the lowest misfit values. The other lines denote the top 2% best fitting models. To evaluate the accuracy of the inverted velocity models, we calculate the theoretical dispersion curves using the best inverted velocities in Figure 12. The corresponding results are shown in Figure 13 as the red lines, and they match well with the picked dispersion curves marked by the blue circles in Figure 13. Such a comparison suggests that the inverted velocity models could interpret the observed data in a good performance.

The final 2D shear wave velocity profiles are shown in Figure 14. The shear wave velocity variation above 40 m depth along the two survey lines is small. There is no obvious velocity discontinuity interface, and the shear wave velocity is between 100 and 320 m/s. When the depth exceeds 40 m, the shear wave velocity gradually decreases from west to east. The lateral velocity change in Line II is greater than that in Line I. In Line I, the stratification is obvious, and the shear wave gradually decreases in the horizontal direction at the depth range of 50–70 m. In Line



The inverted 1D shear wave velocities using the fundamental-mode dispersion curves in Figures 11A–D at Channel 11 (A) and Channel 1/5 (B) in Line I, and at Channel 51 (C) and Channel 200 (D) in Line II, respectively. The red lines denote the best fitted models that have the lowest misfit, and the other lines denote the top 2% best fitting models.

II, starting from the depth of 40 m, shear wave velocity presents lateral variation, and the velocity is between 250 and 350 m/s. The lateral variation of shear wave velocity increases with the gradual increase of depth, and the shear wave velocity decreases in the horizontal direction at a depth of 58–60 m.

The maximum velocity of the two profiles above the depth of 80 m is only 700 m/s, which is smaller than the velocity of the shear wave in the bedrock, so it can be inferred that the thickness of the sedimentary layer in this area should be greater than 80 m. In addition, the different features between these two profiles suggest a strong lateral variation of the sedimentary layer along the Yigong Lake.

Considering the location of the lake water (blue line in Figure 2), the different shallow structures and increased thickness of the sediments in the Yigong Lake can be attributed to the accumulation of debris flow alluvium from the surrounding mountains (Zhang et al., 2021).

# Discussion

We have demonstrated the feasibility of using DAS to record ambient noise data in the Tibetan Plateau. By applying the well-established seismic interferometry to the DAS data,



we can reveal the shallow sedimentary structure of the Yigong Lake, which is very helpful to understand the dynamic evolution of this area. Considering the extreme conditions of the Tibetan Plateau, it is difficult, if not possible, to deploy traditional geophones or portable nodes for environmental seismology study, our method can be a promising alternative for the interested readers. We should note that since the surface wave inversion is highly non-unique, further studies are needed to verify the reliability of the inversion results by incorporating other geological information, such as the drilling data. Due to the limited recording time, we could not get enough low frequency information below 1 Hz. Only when these low frequency data are available, we can get access to the



deeper parts to reveal the actual depth of the bedrock. In addition, we only use the land data and a small portion of the underwater data of Line I in this study. The noise characteristics in the rest parts are more complex than those used in this study, which need further processing and evaluation. Future work with integration of the complete DAS dataset might give chances for a better interpretation of the geologic evolution history of the Yigong Lake.

## Conclusion

This study introduces a shallow sedimentary structure imaging method in extreme conditions of the Yigong Lake in the Tibetan Plateau using ambient noise data recorded by DAS. By applying proper data pre-processing and seismic interferometry to the ambient noise data, clear NCFs are obtained and surface waves can be easily identified. Dispersion curves are then extracted using MASW and are inverted to get the shear wave velocity model of the study area. The results show that the lateral variation is not evident in the shallow subsurface layer, and the lateral inhomogeneity of shear wave velocity increases with depth. Since the shear wave velocity is smaller than conventional bedrock at depths greater than 80 m, it can be assumed that the sedimentary thickness is greater than 80 m in the Yigong Lake area. The increased thickness is likely due to debris-flow deposits in the surrounding mountains. The low-cost, high-density acquisition and processing of DAS data demonstrate the effectiveness and practicality of DAS in detecting shallow sedimentary structures in extreme environments in the Tibetan Plateau. We are expecting that DAS will be a

powerful tool for studies of remote and harsh environments in the future.

## Data availability statement

The datasets presented in this article are not readily available because it is still being used for scientific research. Requests to access the datasets should be directed to the corresponding author.

## Author contributions

JY, JZ, and JG studied the method, processed the data and wrote the draft of the manuscript. JY, HZ, and DD acquired the DAS data in the field. HZ planned the project, and TX designed the DAS system. All the authors contributed to the final manuscript.

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# Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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