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# Subduction-zone parameters that control slab behavior at the 660-km discontinuity revealed by logistic regression analysis and model selection

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The potential mechanisms that drive the behavior of subducted oceanic plates at the 660-km discontinuity are subject to debate. Here we conduct logistic regression analysis and model selection to determine the key subduction-zone parameters in natural subduction zones that discriminate the plate behavior along the discontinuity. We select the key variables based on three information criteria: leave-one-out cross-validation score (LOO), Akaike Information Criterion (AIC), and Bayesian Information Criterion (BIC). Among the 17 subduction-zone parameters analyzed, only the trench velocity, convergence rate, and trench width are selected in the simplest model that minimizes BIC. The thermal parameter and several other variables are also selected to minimize AIC and LOO. Our results suggest that a stagnant slab occurs along the 660-km discontinuity when there is a narrow oceanic plate and a retreating trench in natural subduction zones, which has also been modeled in previous numerical simulations. Neither the stress nor the deformation rate of the upper-plate margin is selected in the three optimal models, which suggests that back-arc spreading in natural subduction zones does not globally characterize plate behavior at the 660-km discontinuity, although back-arc spreading and a stagnant slab coincide in some numerical simulations. The combination of subduction-zone data analysis and numerical simulations will therefore provide deep insights into the dynamics of Earth's deep interior.

## KEYWORDS

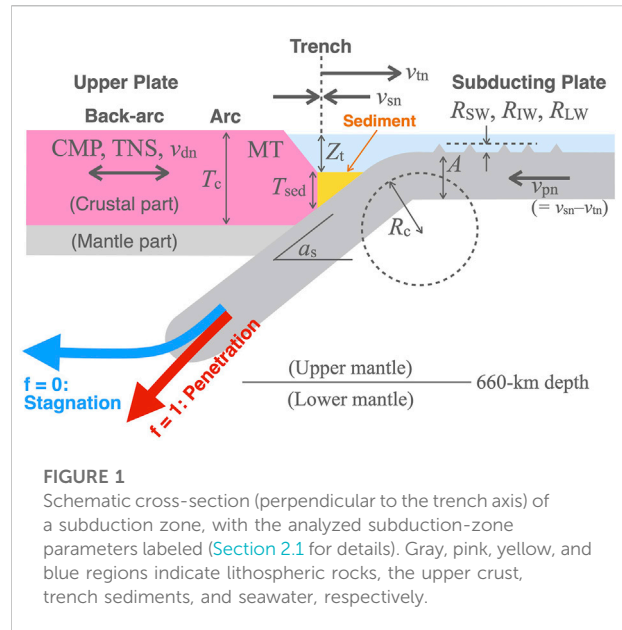
plate tectonics, Subduction Zone, stagnant slab, slab penetration, discriminant analysis, sparse modeling

## 1 Introduction

The 660-km discontinuity is a major seismic discontinuity at approximately 660 km depth that marks the boundary between the upper and lower mantle. This boundary has been inferred to act as a partial barrier to subducted oceanic plates (slabs) and is associated with a phase change in mantle minerals (e.g., Ringwood, 1991). Seismic tomography has imaged the stagnation of some slabs above this boundary (so-called stagnant slabs), such as that observed beneath Eastern Asia, with others, such as that beneath the Mariana Islands, penetrating this boundary (e.g., Fukao et al., 2009). The fate of subducting slabs is a topic of particular interest among geophysicists, as it affects the thermal and chemical evolution of the entire mantle and influences various tectonic processes, such as intra-plate volcanism that is triggered by water derived from the stagnant slab (e.g., Zhao et al., 2007; Richard and Iwamori, 2010).

The mechanisms that generate the variations in slab behavior at the 660-km discontinuity are still debatable, although numerical experiments have been conducted to clarify the relationship between given tectonic conditions and the resultant slab behavior. Christensen (1996) and subsequent studies have shown that seaward trench motion (or trench retreat) coincides with plate stagnation along the 660-km discontinuity. Additional controlling mechanisms have been proposed, such as the thermal structures of the subducting and upper plates, the width of the subducting oceanic plate, Clapeyron slopes and viscosity jumps at the 410- and 660-km discontinuities, the rheology of the back-arc basin and subducting plates, and the water content of the subducting plate (e.g., Torii and Yoshioka, 2007; Stegman et al., 2010; Nakakuki and Mura, 2013; Čížková and Bina, 2013; Garel et al., 2014; Nakao et al., 2016, 2018; Agrusta et al., 2017). This suite of potential mechanisms that can be used to explain the slab behavior at the 660-km discontinuity therefore makes it necessary to determine the predominant mechanisms by analyzing data obtained from natural subduction zones.

The predominant mechanisms driving slab behavior at the 660-km discontinuity can be constrained by statistically clarifying the relationship between the slab behavior at the boundary (i.e., stagnation or penetration) and the observed physical parameters along all of the subduction zones (hereafter “subduction-zone parameters”). This problem can be reduced to a general discriminant problem to determine the key explanatory variables that effectively discriminate an objective variable into labeled classes (e.g., Kuwatani et al., 2014; Ueki et al., 2018) by defining the slab behavior (stagnation or penetration) as an objective variable that can be characterized using subduction-zone parameters. Sparse modeling, which expresses an objective variable using a limited number of variables, is a suitable solver for this



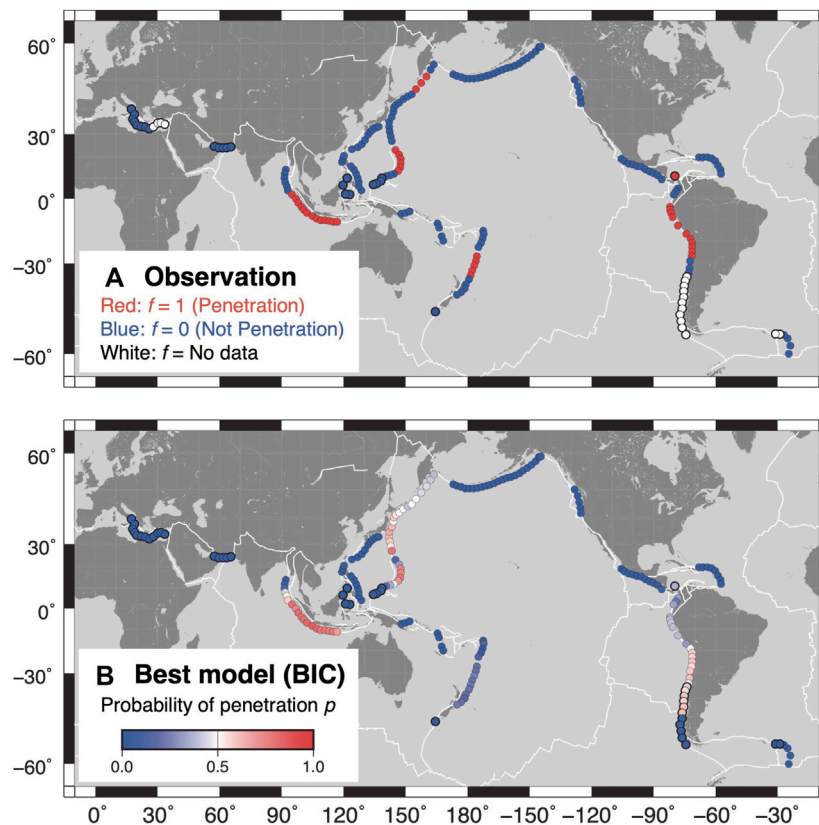
discriminant problem because it provides key variables and can assist in identifying the important mechanisms. Furthermore, sparse modeling does not require a large data volume and is applicable to sparse datasets, such as the observed subduction-zone parameters to date. We therefore apply sparse modeling for this discriminant problem owing to its explainability of the subduction-zone parameter results and the current sparsity of the subduction-zone parameter dataset.

Here we identify the subduction-zone parameters that shape the slab behavior at the 660-km discontinuity *via* discriminant analysis and variable selection. We then discuss the predominant mechanism driving this slab behavior by comparing our analytical results with published numerical solutions for plate subduction.

## 2 Methods

### 2.1 Data

We investigated 17 types of subduction-zone parameters, which are visualized in Figure 1, as explanatory variables of the subducted oceanic plate behavior at the 660-km discontinuity. We used the subduction-zone parameters that were observed at 163 subduction-zone locations (circles with thin black lines in Figure 2) as training data for the logistic regression analysis (Section 2.2), with the remaining 40 or less locations, which lack observations for some of the subduction-zone parameters, serving as testing data for the resultant regression model (circles with thick black lines in Figure 2).



**FIGURE 2**  
**(A)** Plate behavior at the 660-km discontinuity,  $f$ , which is based on the maximum plate depth defined by SubMap 4.3 (Heuret and Lallemand, 2005). **(B)** Probability of the subducted oceanic plate penetrating the 660-km discontinuity,  $p_{BIC}$ , which is calculated using the optimal model that yields the minimum BIC (black bars in Figure 3). Circles with thin black lines in both figures indicate the samples used for training, and those with thick black lines indicate test samples. White lines mark the plate boundaries.

We reduced our problem to a two-class discriminant problem by labeling the objective variable, the plate behavior, beneath the  $j$ -th sample location as:

$$f_j = \begin{cases} 1 & (d_j > 670 \text{ km}) \\ 0 & (d_j \leq 670 \text{ km}), \end{cases} \quad (1)$$

where  $d_j$  is the maximum depth of the subducting oceanic plate beneath location  $j$ , which is taken from SubMap 4.3 (Heuret and Lallemand, 2005). Therefore,  $f = 1$  indicates that the oceanic plate penetrates the 660-km discontinuity, whereas  $f = 0$  indicates that the oceanic plate either stagnates at the 660-km discontinuity or does not reach this boundary. The global distribution of  $f$  used in the analysis is shown in Figure 2A; 24.5% of the test data are labeled  $f = 1$ , and the rest are labeled  $f = 0$ .

The 17 explanatory variables used to discriminate  $f$  are almost the same as the subduction-zone parameters in Nakao et al. (2022). These 17 variables are as follows: 1)  $A$ , the age of the subducting oceanic plate at the trench (Müller et al., 1997); 2)  $a_s$ , the mean dip angle of the subducting oceanic plate over the

0–125 km depth range (Heuret and Lallemand, 2005); 3) CMP (compression) and 4) TNS (tension), dummy variables that express the strain state of the upper plate, and are based on earthquake focal mechanisms (“UPS” in Heuret et al., 2011): (CMP, TNS) = (1, 0) for compressible upper plates, (0, 1) for tensile upper plates, and (0, 0) for neutral upper plates; 5)  $L_p$ , the width of the oceanic plate defined by Brizzi et al. (2018); 6) MT (margin type), a dummy variable that expresses either the accretionary or erosional conditions of the upper-plate margin: MT = 1 for accretionary margins and MT = 0 for erosional margins; 7)  $R_c$ , the bending radius of the subducting oceanic plate; 8)  $R_{SW}$ , 9)  $R_{IW}$ , and 10)  $R_{LW}$ , the seafloor roughnesses of the subducting oceanic plate for different bathymetric wavelengths, which have been defined by Lallemand et al. (2018): 12–20 km [short wavelengths (SW)], 20–80 km [intermediate wavelengths (IW)], and 80–100 km [long wavelengths (LW)]; 11)  $T_c$ , the upper-plate crustal thickness at the arc, which is taken from CRUST 1.0 (Laske et al., 2013); 12)  $T_{sed}$ , the sediment thickness at the trench, which is taken from GlobSed 3 (Straume et al., 2019); 13)  $v_{sn}$ , the

trench-normal component of the convergence rate at the trench (Lallemand et al., 2008); 14)  $v_{dn}$ , the trench-normal component of the extension rate of the upper-plate margin (Lallemand et al., 2008), which is positive for back-arc spreading and negative for back-arc shortening; 15)  $v_{tm}$ , the trench-normal component of the trench migration rate, which is positive for a retreating trench (oceanward motion) and negative for an advancing trench (continent-ward motion), with the values used in this study based on the Indian hotspot (Steinberger et al., 2004; Lallemand et al., 2008); 16)  $Z_b$ , the trench depth; and 17)  $\Phi$ , the thermal parameter, which is defined as  $\Phi = Av_{sn} \sin a_s$  by Kirby et al. (1991) and corresponds to the coldness of the subducted oceanic plate.

## 2.2 Logistic regression analysis

We apply logistic regression analysis, a simple method for discriminant problems, to express the probability of plate penetration at the 660-km discontinuity (hereafter  $p$ ) using the 17 subduction-zone parameters. The relative contribution of each explanatory variable to  $p$  is evaluated by standardizing the  $i$ -th explanatory variable ( $i = 1, \dots, I$ ) observed at location  $j$  ( $j = 1, \dots, J$ ) as:

$$x'_{ij} = \frac{x_{ij} - m_i}{s_i}, \tag{2}$$

where  $x_{ij}$  is an unstandardized explanatory variable (i.e.,  $A, a_s, \dots, \Phi$ ),  $m_i = \frac{1}{J} \sum_{j=1}^J x_{ij}$  is the empirical mean value of variable  $i$ ,  $s_i = \left\{ \frac{1}{J-1} \sum_{j=1}^J (x_{ij} - m_i)^2 \right\}^{\frac{1}{2}}$  is the empirical standard deviation of variable  $i$ ,  $I$  is the number of explanatory variables ( $I = 17$ ), and  $J$  is the number of locations used for training ( $J = 163$ ).

A generalized linear model is used for two-class discriminant analysis. We assume a Bernoulli distribution, such that the probability of plate penetration beneath location  $j$  can be written as:

$$p_j(\mathbf{a}; \mathbf{c}) = \frac{1}{1 + \exp(-z_j(\mathbf{a}; \mathbf{c}))} \tag{3}$$

$$z_j(\mathbf{a}; \mathbf{c}) = a_0 + \sum_{i=1}^I c_i a_i x'_{ij}, \tag{4}$$

where  $\mathbf{a} = (a_0, a_1, \dots, a_I)$  is a coefficient vector and  $\mathbf{c} = (c_1, \dots, c_I)$  is a parameter vector that controls whether  $x'_{ij}$  is used (i.e.,  $c_i = 0$  or 1). The hyperplane  $z = 0$  corresponds to the discriminant plane of  $f$ .

The coefficient is estimated based on the premise that the likelihood  $L(\mathbf{a}; \mathbf{c})$  is the maximum for a given  $\mathbf{c}$ :

$$\hat{\mathbf{a}} = \arg \max_{\mathbf{a}} L(\mathbf{a}; \mathbf{c}). \tag{5}$$

Here, the likelihood is defined as the joint probability:

$$L(\mathbf{a}; \mathbf{c}) = \prod_{j=1}^J P(f_j | p_j(\mathbf{a}; \mathbf{c})), \tag{6}$$

where  $P(f_j | p_j)$  is the probability that the observed and predicted subducted-plate behaviors are the same for location  $j$ . When  $P(f_j | p_j)$  follows a Bernoulli distribution, the log likelihood can be written as:

$$\ln L(\mathbf{a}; \mathbf{c}) = \sum_{j=1}^J \{f_j \ln p_j(\mathbf{a}; \mathbf{c}) + (1 - f_j) \ln(1 - p_j(\mathbf{a}; \mathbf{c}))\}. \tag{7}$$

$\hat{\mathbf{a}}$  is therefore estimated *via* logistic regression.

## 2.3 Model evaluation

We determined the optimal variable combination  $\mathbf{c}$  based on three criteria: log-likelihood-based leave-one-out cross-validation score (LOO; Stone, 1974), Akaike Information Criterion (AIC; Akaike, 1974), and Bayesian Information Criterion (BIC; Schwarz, 1978). LOO is defined here as:

$$\text{LOO}(\mathbf{c}) = -\frac{1}{J} \sum_{j=1}^J \{f_j \ln p_{-j}(\hat{\mathbf{a}}; \mathbf{c}) + (1 - f_j) \ln(1 - p_{-j}(\hat{\mathbf{a}}; \mathbf{c}))\}, \tag{8}$$

where  $p_{-j}(\hat{\mathbf{a}}; \mathbf{c})$  is the probability of the subducting plate penetrating the 660-km continuity at location  $j$ , which is calculated using the variables at the other  $(J - 1)$  locations and a given variable combination  $\mathbf{c}$ . Cross validation quantifies the generalized performance of a prediction model using only a small number of observations and can therefore be used to evaluate the effect of omitting subduction zones.

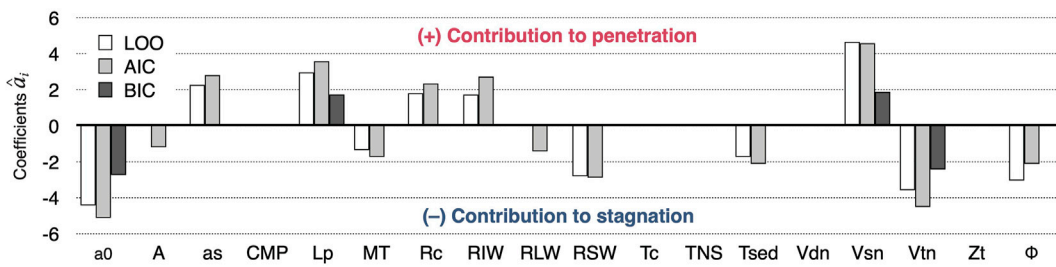
AIC and BIC are defined as:

$$\text{AIC}(\mathbf{c}) = -2 \ln L(\hat{\mathbf{a}}; \mathbf{c}) + 2 \left( \sum_{i=1}^I c_i + 1 \right) \tag{9}$$

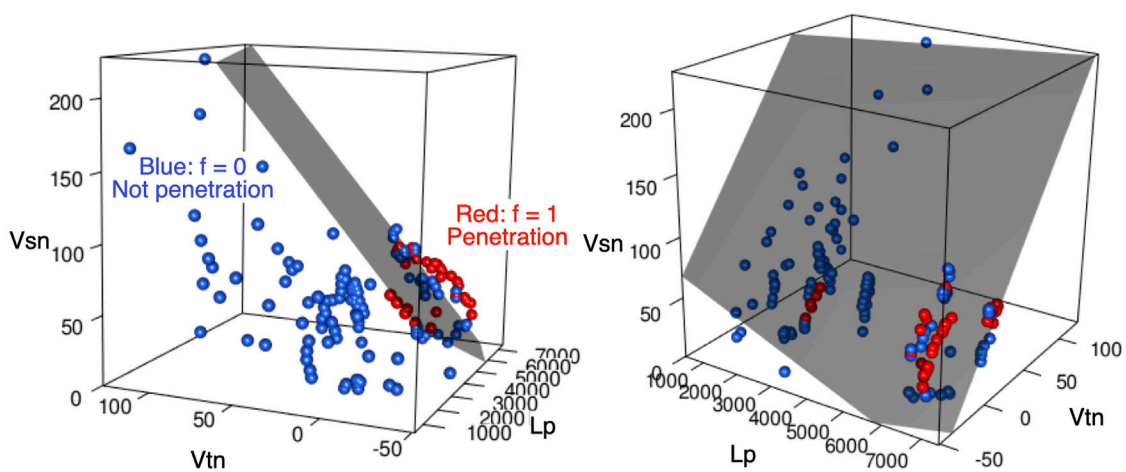
and:

$$\text{BIC}(\mathbf{c}) = -2 \ln L(\hat{\mathbf{a}}; \mathbf{c}) + \left( \sum_{i=1}^I c_i + 1 \right) \ln J, \tag{10}$$

respectively, where  $L(\hat{\mathbf{a}}; \mathbf{c})$  is the maximum likelihood of model  $\mathbf{c}$ . The term  $(\sum_{i=1}^I c_i + 1)$  indicates the number of parameters (i.e., the number of  $a_i$  values used in model  $\mathbf{c}$  and the intercept  $a_0$ ). AIC and BIC are evaluation criteria that balance the misfit with the complexity of a model by imposing a penalty based on the number of variables used in the model. Although both criteria look similar, AIC is proposed in terms of the model predictability, whereas BIC is related to the posterior model probability (e.g., Konishi and Kitagawa, 2008). Therefore, AIC places more weight on the fit between the observation and prediction and prefers a model that contains a larger number of variables than BIC. An exhaustive model evaluation was conducted to determine the optimal models that minimize the three criteria among  $2^I (= 2^{17})$  candidates.



**FIGURE 3** Selected variables and their coefficients  $\hat{a}_i$  in the three optimal models that minimize LOO, AIC, and BIC. Positive and negative values indicate the contributions to penetration and stagnation of the subducted oceanic plates at the 660-km discontinuity, respectively.



**FIGURE 4** Two views of a three-dimensional plot of  $L_p$  (km),  $v_{sn}$  (mm/yr), and  $v_{tn}$  (mm/yr) for all of the training data. Blue and red spheres indicate those labeled as  $f = 0$  (stagnation) and  $f = 1$  (penetration), respectively. The gray plane indicates the estimated discriminant plane that minimizes BIC (i.e.,  $z_{BIC} = 0$  in Eq. 12,  $p_{BIC} > 0.5$  above the plane, and  $p_{BIC} < 0.5$  below the plane).

### 3 Results

We determined the three models that most effectively discriminated the subducted plate behavior  $f$  (Figure 3). LOO is minimized by selecting 10 variables:  $a_s$ ,  $L_p$ ,  $MT$ ,  $R_c$ ,  $R_{IW}$ ,  $R_{SW}$ ,  $T_{sed}$ ,  $v_{sn}$ ,  $v_{tn}$ , and  $\Phi$ . AIC is minimized when  $A$  and  $R_{LW}$  are selected in addition to the 10 LOO variables. The optimal LOO and AIC models are similar since cross validation is asymptotically equivalent to AIC (Stone, 1977). However, BIC is minimized by selecting only three variables:  $L_p$ ,  $v_{sn}$ , and  $v_{tn}$ .  $L_p$ ,  $v_{sn}$ , and  $v_{tn}$  are common variables among the three optimal models. Furthermore, the relative contribution of these three variables to the discrimination is large in the optimal LOO and AIC models compared with the other variables (Figure 3), which indicates that  $f$  is better discriminated in the three-dimensional

(3-D) space of these three variables, as shown in Figure 4), than in the other variable spaces.  $\Phi$  has the third largest contribution to the discrimination in the optimal LOO model; its effects on slab behavior will be discussed along with  $L_p$ ,  $v_{sn}$ , and  $v_{tn}$ .

The probability of slab penetration for the variable combination that minimizes BIC,  $p_{BIC}$ , can be expressed in non-dimensionalized form as:

$$p_{BIC,j} = \frac{1}{1 + \exp(-z_{BIC,j})} \tag{11}$$

$$z_{BIC,j} = -8.81 + 7.68 \times 10^{-4} L_{p,j} + 6.00 \times 10^{-2} v_{sn,j} - 8.50 \times 10^{-2} v_{tn,j}, \tag{12}$$

where  $L_{p,j}$  (km) is the width of the subducting plate to which location  $j$  belongs,  $v_{sn,j}$  (mm/yr) is the trench-normal

convergence rate at location  $j$ , and  $v_{tn,j}$  (mm/yr) is the trench-normal trench retreat rate at location  $j$ . Similarly, the  $p$  that minimizes LOO and AIC can be expressed using the estimated coefficients in [Supplementary Table S1](#). The hyperplane  $z_{BIC} = 0$  corresponds to the estimated discriminant plane of the optimal BIC model to classify  $f$  ([Figure 4](#)).

[Figure 2B](#) shows the global  $p_{BIC}$  distribution for both the training and test data. Even though the optimal BIC model yields a lower likelihood value than the optimal LOO and AIC models, the very small  $p_{BIC}$  values are correctly derived for most of the test locations where there is no 660-km discontinuity penetration, such as those beneath the Mediterranean Sea and Southeast Asian islands. However, a significant discrepancy between  $f$  and  $p_{BIC}$  is observed beneath the southern Tonga region, whereas a better fit between  $f$  and  $p$  is obtained in the optimal LOO and AIC models ([Supplementary Figure S1](#)). The subducted Pacific Plate beneath the southern Tonga region is believed to penetrate the 660-km discontinuity and form a horizontal stagnant slab at 1,000 km depth based on seismic velocity models ([Fukao and Obayashi, 2013](#)), although the optimal BIC model indicates plate stagnation along the 660-km discontinuity.

## 4 Discussion

We propose that  $L_p$ ,  $v_{sn}$ , and  $v_{tn}$ , which are the three variables selected in each of the three optimal models, are the primary subduction-zone features that characterize  $f$ . In addition, LOO and AIC place more weight on model predictability than model simplicity compared with BIC, such that the additional variables selected by LOO and AIC may be considered secondary variables that are important for improving the fit between  $f$  and  $p$ . Here we discuss the likely mechanisms of driving slab behavior in natural subduction zones based on these primary and secondary variables.

$v_{tn}$  is selected as a variable that decreases  $p$ , which is consistent with previous numerical simulations that show relations between slab behavior and the trench motion (e.g., [Christensen, 1996](#)). The relationship between a stagnant slab and trench retreat is governed by the anchoring of the tip of a subducting plate to the 660-km discontinuity, which causes the plate to bend in a convex downward direction, with the negative buoyancy force pulling the upper plate oceanward ([Nakakuki and Mura, 2013](#); [Nakao et al., 2016](#)). Continent-ward counter-flows dominate in the mantle transition zone and extend the subducted plate horizontally during rollback (or trench retreat) ([Nakakuki and Mura, 2013](#); [Nakao et al., 2016](#)). This mechanism may explain the relationship between  $f$  and  $v_{tn}$  in natural subduction zones.

Our optimal models also show that increasing  $L_p$  result in decreasing  $p$ , which is consistent with 3-D numerical simulations

(e.g., [Schellart et al., 2007](#); [Stegman et al., 2010](#)) which have demonstrated that a narrow oceanic plate contributes to rapid trench retreat and plate stagnation along the 660-km discontinuity. These previous experiments have indicated that toroidal flows [flows in the two-dimensional (2-D) horizontal section] from sub-oceanic to sub-continental mantle regions become dominant compared with poloidal flows (flows in the 2-D vertical section) around narrow plates, and are essential to the formation of a stagnant slab with a retreating trench. Our results indicate that this along-arc mechanism is a significant driver of slab stagnation/penetration in natural subduction zones.

$v_{sn}$  is selected in each optimal model as a variable that increases  $p$ . This result may appear to contradict some numerical experiments (e.g., [Nakakuki and Mura, 2013](#); [Garel et al., 2014](#); [Agrusta et al., 2017](#); [Nakao et al., 2018](#)), which show a high  $v_{sn}$  during the initial stage of subduction that then decreases as the slab penetrates the 660-km boundary. It also appears to be inconsistent with the subduction mode classification of [Schellart \(2011\)](#), whereby a stagnant slab coincides with large  $v_{sn}$ . However,  $v_{sn}$  is an important variable when selected with  $v_{tn}$  and  $L_p$  because the discriminant plane is located along the data trend ([Figure 4](#)). The slightly large  $v_{sn}$  of the penetrating plates relative to the stagnating plates under negative  $v_{tn}$  and large  $L_p$  may be associated with acceleration during the collapse of the stagnant slab or an external force affecting the movement of the oceanic plates towards the trench (e.g., [Fuji et al., 2021](#)).

$\Phi$  is selected in the optimal LOO and AIC models, suggesting that colder plates tend to stagnate along the 660-km discontinuity. The same result has been obtained in previous simulations (e.g., [Garel et al., 2014](#); [Agrusta et al., 2017](#); [Fuji et al., 2021](#)), as colder (older) plates tend to stagnate more readily because they are generally heavier and drag the upper plate oceanward. A subducted cold oceanic plate can stagnate along the 660-km discontinuity because of the buoyancy related to a delay in the phase transition from pyroxene to majoritic garnet ([King et al., 2015](#)). Metastable minerals, such as olivine and pyroxene, may also contribute to the buoyancy of a cold slab in the mantle transition zone, thereby leading to stagnation ([Tetzlaff and Schmeling, 2000](#); [Agrusta et al., 2014](#)). Therefore, a large  $\Phi$  may induce a negative buoyancy effect in the shallow mantle and a positive buoyancy effect in the mantle transition zone, both of which would lead to stagnation of the subducting slab. Conversely, the rheological effects associated with  $\Phi$  may have a complex influence on slab morphology (e.g., [Lee and King, 2011](#); [Dasgupta et al., 2021](#); [Cerpa et al., 2022](#)).

The optimal LOO and AIC models indicate that a small  $T_{sed}$  value and the absence of an accretionary prism lead to increased  $p$ . This can be explained by the numerical simulations of [Brizzi et al. \(2020\)](#), which show the emplacement of a thicker sediment layer overlying the subducting oceanic plate causes the trench to move oceanward and the sediment to be accreted to the upper plate, forming an accretionary prism. This process further

reduces the shallow dip angle of the subducting oceanic plate. A similar trend has been observed in simulations of the South America (Hu et al., 2021). Therefore, this mechanism, which is driven by a large  $T_{\text{sed}}$  value, yields positive MT, small  $a_s$ , and large  $v_{\text{tn}}$  values, and is consistent with our optimal LOO and AIC models.

CMP, TNS,  $v_{\text{dn}}$ , and  $T_c$  are not selected in any of the optimal models, although previous numerical simulations have indicated that a stagnant slab occurs with stretching of the upper plate (e.g., Nakakuki and Mura, 2013; Garel et al., 2014; Nakao et al., 2016, 2018). Our finding may reflect exceptions to the results of the previous numerical simulations, as observed in the Mariana subduction zone, where the subducting Pacific Plate penetrates the lower mantle despite back-arc spreading, because of the tensile stress related to subduction of the overlying Philippine Sea Plate (Čížková and Bina, 2015). Even though the formation of a back-arc basin may coincide with a stagnant slab in some tectonic environments, our result suggests that it does not explain the global distribution of  $f$ .

We are unable to explain why variables related to seafloor roughness ( $R_{\text{SW}}$ ,  $R_{\text{IW}}$ , and  $R_{\text{LW}}$ ) were selected as explanatory variables of  $f$  in the optimal AIC and LOO models, given the various geophysical and geochemical processes that shape the seafloor topography, including hot-spot and petit-spot volcanism, trench outer-rise faulting, and sedimentation. The relationship between seafloor roughness and slab behavior at the 660-km discontinuity should therefore be investigated in future numerical simulations.

We note that the optimal variable combination could potentially change if more favorable variables are additionally given. We used relatively shallow data in the present study because deeper data were unavailable. However, we will develop a new regression model that incorporates any new data that may be acquired *via* advances in observational techniques, especially in the deep mantle. For example, the sudden increases in viscosity and density at the 660-km discontinuity might influence slab dynamics (Torii and Yoshioka, 2007; Nakakuki and Mura, 2013; Čížková and Bina, 2013), and mantle winds, which are driven by the pressure gradient in the deep mantle, might give rise to a stagnant slab (Peng et al., 2021). The consideration of future observations of the deep mantle will further advance our understanding of the dynamics of subducting slabs.

While we only performed a simple two-class discriminant analysis based on the maximum depth of a subducted slab, which was the objective variable, the slab morphology and its behavior can be expressed by additional factors, such as its buckling wavelength and amplitude (Lee and King, 2011), and the dip angle of deeper sections of the slab (Hu and Gurnis, 2020). However, these morphological factors may vary with physical parameters that are either not selected or not considered in the present analysis, such as the ages of the subducting and upper plates (Dasgupta et al., 2021; Cerpa et al., 2022). If we can

quantify the slab morphology, then multi-class discriminant analysis will provide additional insights into the dynamics of subducting slabs.

## 5 Conclusion

Logistic regression analysis and exhaustive model evaluation were conducted to determine the subduction-zone parameters that increased the probability of the subducting slab penetrating the 660-km discontinuity,  $p$ . The results highlight that the slab behavior observed at this mantle boundary can be effectively classified by considering the combination of the trench retreat rates  $v_{\text{tn}}$ , convergence rates at the trench  $v_{\text{sn}}$ , and the width of the oceanic plate  $L_p$ ; these three variables were selected in each of the three optimal models, which were based on three employed criteria: LOO, AIC, and BIC. Negative  $v_{\text{tn}}$  (trench advance) and large  $L_p$  (wide oceanic plate) both increase  $p$ , consistent with previous numerical simulations. The predominant influences of  $v_{\text{tn}}$  and  $L_p$  on  $p$  are likely due to both variables being related to mantle flow from the sub-oceanic to sub-continental mantle material surrounding the subducting slab, which is necessary to induce slab rollback. The thermal parameter  $\Phi$  and trench sediment thickness  $T_{\text{sed}}$  were selected to obtain a better fit between  $f$  and  $p$ , which is also consistent with previous numerical simulations. Neither the stress nor deformation rate of the upper-plate margin were selected in the three optimal models, thereby suggesting that back-arc spreading does not globally characterize  $f$ . However, the dataset used in this study lacks some information, such as the data of the deep mantle and chemical composition; therefore, further analyses with improved subduction-zone data are needed to advance our knowledge of the dynamics of Earth's deep interior.

## Data availability statement

Publicly available datasets were analyzed in this study. This data can be found here: All of the data used in our analysis were downloaded from SubMap 4.3 (Heuret and Lallemand, 2005) (<https://doi.org/10.1016/j.pepi.2004.08.022>), Supplementary Material S1 of Brizzi et al. (2018) (<https://doi.org/10.1016/j.tecto.2018.01.027>), CRUST 1.0 (<https://igppweb.ucsd.edu/~gabi/rem.html>), and GlobSed 3 (<https://ngdc.noaa.gov/mgg/sedthick/>).

## Author contributions

AN designed the research, collected the data, conducted the analysis, and wrote the original draft under the supervision of TK. TK and KU developed the analytical

method. TK, KU, KY, and TY reviewed and edited the original draft. HH and SA verified the analytical method and reviewed and edited the statistical discussions in the manuscript.

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## Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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## Supplementary material

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/feart.2022.1008058/full#supplementary-material>



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