



# New Advances in Dial-Lidar-Based Remote Sensing of the Volcanic CO<sub>2</sub> Flux

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We report here on the results of a proof-of-concept study aimed at remotely sensing the volcanic CO<sub>2</sub> flux using a Differential Adsorption lidar (DIAL-lidar). The observations we report on were conducted in June 2014 on Stromboli volcano, where our lidar (Light Detection And Ranging) was used to scan the volcanic plume at ~3 km distance from the summit vents. The obtained results prove that a remotely operating lidar can resolve a volcanic CO<sub>2</sub> signal of a few tens of ppm (in excess to background air) over km-long optical paths. We combine these results with independent estimates of plume transport speed (from processing of UV Camera images) to derive volcanic CO<sub>2</sub> flux time-series of ~16–33 min temporal resolution. Our lidar-based CO<sub>2</sub> fluxes range from 1.8 ± 0.5 to 32.1 ± 8.0 kg/s, and constrain the daily averaged CO<sub>2</sub> emissions from Stromboli at 8.3 ± 2.1 to 18.1 ± 4.5 kg/s (or 718–1565 tons/day). These inferred fluxes fall within the range of earlier observations at Stromboli. They also agree well with contemporaneous CO<sub>2</sub> flux determinations (8.4–20.1 kg/s) obtained using a standard approach that combines Multi-GAS-based in-plume readings of the CO<sub>2</sub>/SO<sub>2</sub> ratio (~8) with UV-camera sensed SO<sub>2</sub> fluxes (1.5–3.4 kg/s). We conclude that DIAL-lidars offer new prospects for safer (remote) instrumental observations of the volcanic CO<sub>2</sub> flux.

**Keywords:** volcanic CO<sub>2</sub>, DIAL-lidar, Stromboli, remote sensing, CO<sub>2</sub> flux

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## INTRODUCTION

A major step forward in ground-based volcano monitoring has recently arisen from the advent of modern instrumental techniques and networks for volcanic gas observations (Galle et al., 2010; Oppenheimer et al., 2014; Saccorotti et al., 2014; Fischer and Chiodini, 2015). Such technical advances provide improved temporal resolution relative to traditional direct sampling techniques (Symonds et al., 1994; Giggenbach, 1996). As longer-term volcanic gas records increase in number and quality, full empirical evidence is finally emerging for increased CO<sub>2</sub> flux emissions prior to eruption of mafic to intermediate volcanoes (Aiuppa, 2015). Precursory plume CO<sub>2</sub> flux increases have been now detected at several volcanoes, including Etna (Aiuppa et al., 2008; Patanè et al., 2013), Kilauea (Poland et al., 2012), Redoubt (Werner et al., 2013), Turrialba (de Moor et al., 2016a), and Poas (de Moor et al., 2016b).

At Stromboli (in Italy), however, CO<sub>2</sub> flux observations have been particularly valuable for interpreting, and eventually predicting, the volcano's behavior (Aiuppa et al., 2010a, 2011). On

Stromboli, the “regular” mild strombolian activity is occasionally interrupted by larger-scale vulcanian-style explosions, locally referred as “*major explosions*” or (in the most extreme events) “*paroxysms*” (Rosi et al., 2006, 2013; Andronico and Pistolesi, 2010; Pistolesi et al., 2011; Pioli et al., 2014). These explosions, although short-lived (tens of seconds to a few minutes), represent a real hazard for local populations, tourists and volcanologists, since they produce fallout of coarse pyroclastic materials over wide dispersal areas (Rosi et al., 2013). In addition, such events are not anticipated by any detectable anomaly in the geophysical or volcanological record, perhaps because they originate deep in the crustal roots of the volcano’s plumbing system (Bertagnini et al., 2003; Métrich et al., 2005, 2010; Allard, 2010). Observational evidence suggests, however, that “*major explosions*” (Aiuppa et al., 2011) and “*paroxysms*” (Aiuppa et al., 2010a) are both systematically preceded by days/weeks of anomalous CO<sub>2</sub>-rich gas leakage from Stromboli’s deep (8–10 km) magma storage zone (Aiuppa et al., 2010b). CO<sub>2</sub> flux emissions from the open-vent crater plume have become, therefore, a unique monitoring tool for volcanic hazard assessment and mitigation on the volcano.

On Stromboli, as at other volcanoes, the volcanic gas CO<sub>2</sub> flux is calculated from a combination of co-measured SO<sub>2</sub> fluxes and plume CO<sub>2</sub>/SO<sub>2</sub> ratios (Burton et al., 2013; Aiuppa, 2015). While the SO<sub>2</sub> flux can remotely be sensed by UV spectroscopy (Oppenheimer, 2010; Oppenheimer et al., 2011), measuring the CO<sub>2</sub>/SO<sub>2</sub> ratio requires *in-situ* direct sampling and/or measurements via Multi-GAS (Aiuppa et al., 2010a) or Fourier Transform Infra-Red Spectrometry (La Spina et al., 2013)

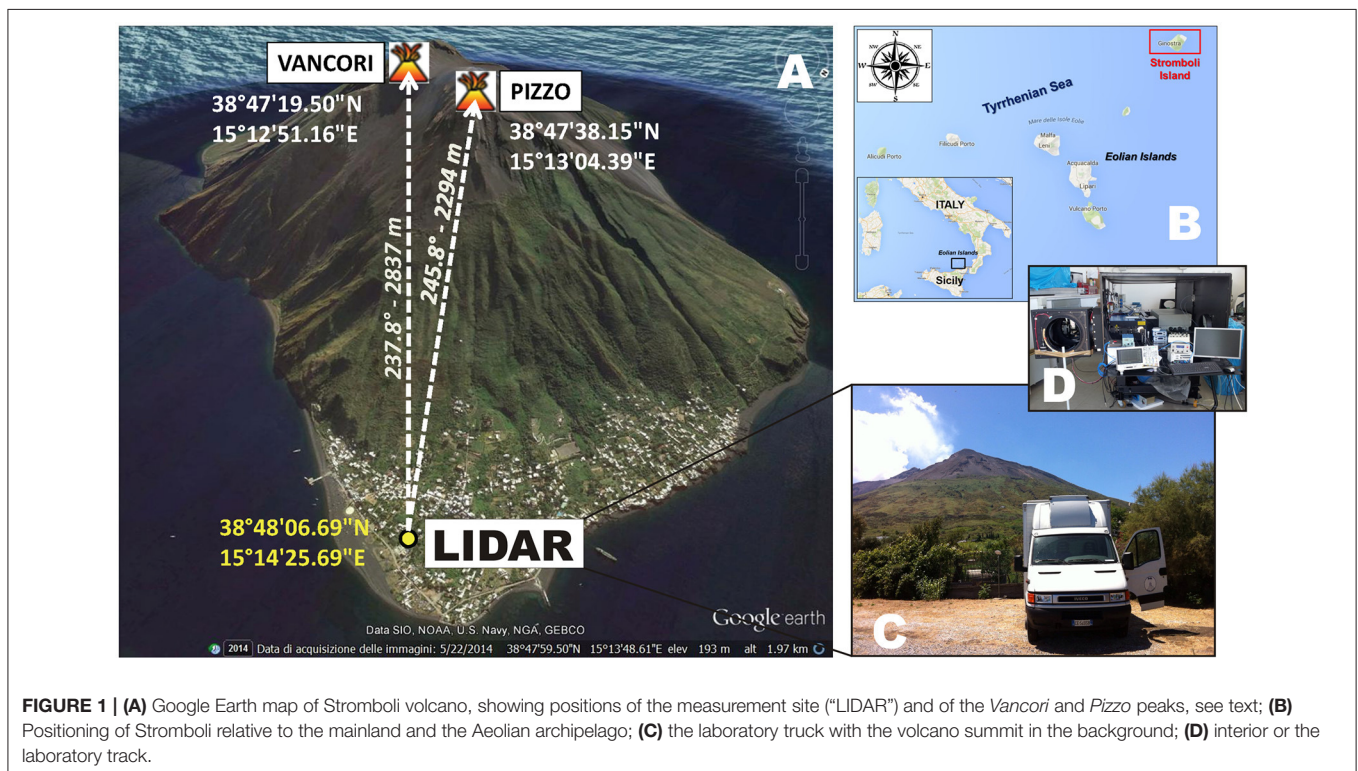
in the vicinity of hazardous active vents. As such, implementation of novel techniques for the remote observation of the volcanic CO<sub>2</sub> flux, from more distal (and safer) locations, remains highly desirable.

New prospects for ground-based *remote* detection of the volcanic CO<sub>2</sub> flux have recently become available from the advent of a new lidar (Light Detection and Ranging) using the DIAL (Differential Absorption lidar) technique (Aiuppa et al., 2015; Fiorani et al., 2015, 2016). DIAL-lidars (Weitkamp, 2005; Fiorani, 2007) use backscattering of artificial light (laser) from atmospheric back-scatterers and/or from the volcanic plume itself, and are therefore potentially ideal for remote volcanic CO<sub>2</sub> detection (Fiorani et al., 2013; Queißer et al., 2015, 2016). In previous work, we demonstrated the ability of our lidar to remotely resolve the volcanic CO<sub>2</sub> flux from a relatively proximal measuring site (<200 m from the source vents) (Aiuppa et al., 2015). Here, we extend this work by reporting on a successful CO<sub>2</sub> flux detection at Stromboli over a far longer optical path (~3 km distance from the vents). Results of this proof-of-concept experiment confirm lidars as promising tools for remote monitoring of the volcanic CO<sub>2</sub> flux (Aiuppa et al., 2015; Queißer et al., 2016).

## MATERIALS AND METHODS

### The Bridge Lidar

Our measurements on Stromboli (Figure 1) were obtained using the same DIAL-lidar described in Aiuppa et al. (2015) and Fiorani et al. (2015, 2016), and realized within the context of the FP7-





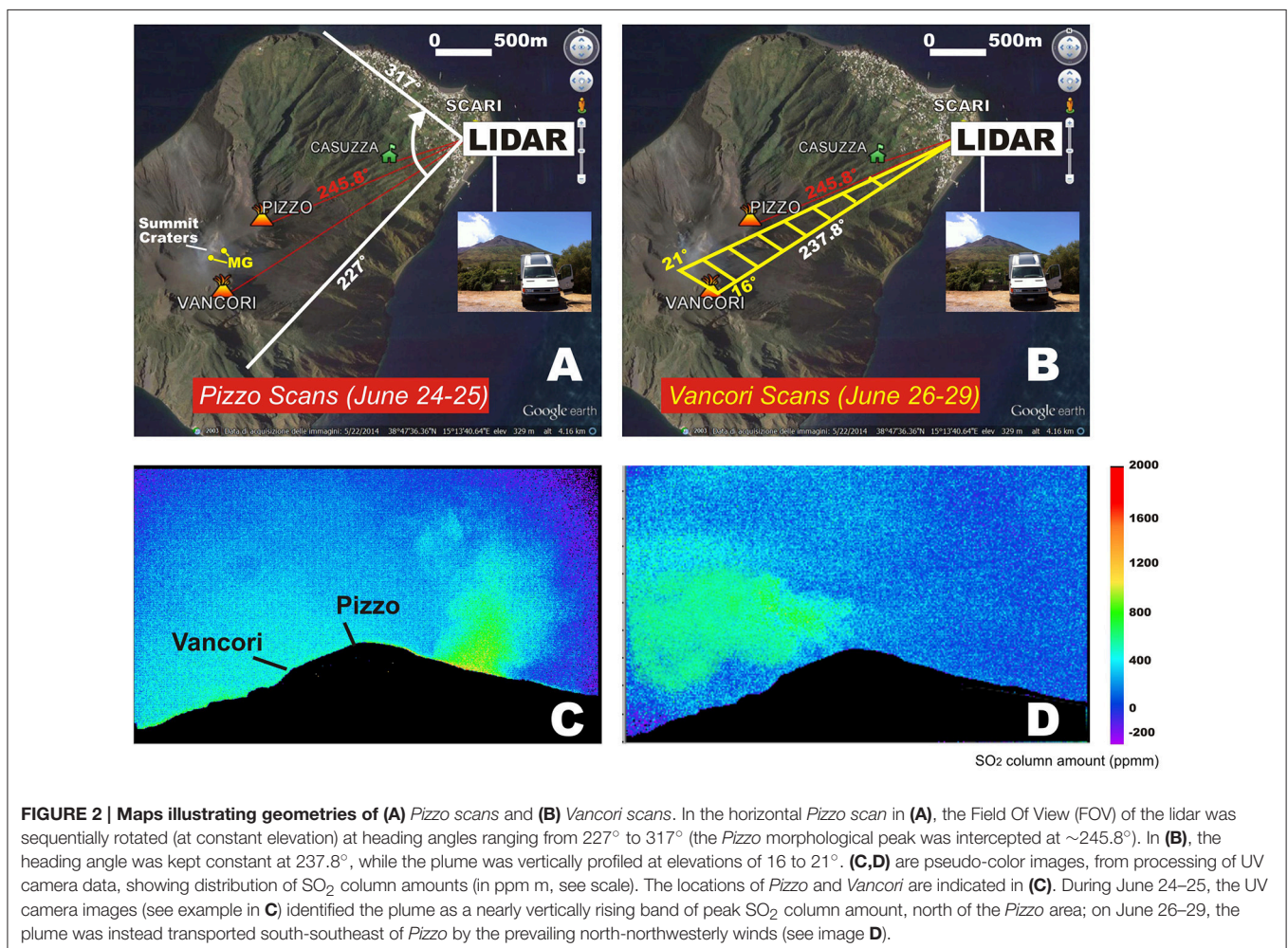
ERC project *Bridge* ([www.bridge.unipa.it](http://www.bridge.unipa.it)). Only key information is reported here, and the reader is referred to previous studies for a detailed description of the instrument. In brief, the Bridge lidar (**Figure 1D**) uses a complex transmitter that integrates (i) an injection seeded Nd:YAG laser with (ii) a double grating dye laser. This transmitter is used to generate laser radiation at  $\sim 2100$  nm, a region of the electromagnetic spectrum absorbed by atmospheric CO<sub>2</sub>, while showing minimal cross-sensitivity to H<sub>2</sub>O (Fiorani et al., 2013). At the ON and OFF wavelengths selected for this experiment, the differential cross section of CO<sub>2</sub> is five orders of magnitude larger than that of H<sub>2</sub>O (Rothman et al., 2013). Considering a CO<sub>2</sub> mixing ratio of 400 ppm, and with the upper and lower ranges of H<sub>2</sub>O mixing ratios used in atmospheric models (Berk et al., 2014), i.e., from 2.59% (tropics, sea level) to 0.141% (high latitude, winter, sea level), the respective CO<sub>2</sub> absorption is 3 and 5 orders of magnitude larger than that of H<sub>2</sub>O. The 2.59% H<sub>2</sub>O mixing ratio is not far from the saturated water vapor pressure at standard atmospheric conditions. We conclude that, even in a condensing volcanic plume, H<sub>2</sub>O absorption is negligible compared to that of CO<sub>2</sub>.

A piezo-electric element is used to sequentially switch the wavelength of the transferred laser beam, from  $\lambda_{ON}$  (2009.537

nm: maximum CO<sub>2</sub> absorption) to  $\lambda_{OFF}$  (2008.484 nm: no CO<sub>2</sub> absorption), at 10 Hz repetition rate. These closely spaced pairs of laser beams are sequentially transmitted into the atmosphere, where they are eventually scattered back by atmospheric back-scatterers (aerosols, water droplets, particles) in either the volcanic plume or the background atmosphere. During their atmospheric propagation, the laser beams are also reflected by any obstacle encountered along the optical path, e.g., in our specific case, the *Pizzo* and *Vancori* walls/rims in front of or behind the volcanic plume (see **Figures 1, 2**). The returned signal is captured by the lidar receiver (a Newtonian telescope, diameter: 310 mm), and then detected and amplified by an InGaAs PIN photodiode module, directly connected with the analog-to-digital converter (ADC).

## Field Operations

During our experiment, the DIAL-lidar operated from a small laboratory truck (**Figure 1C**), positioned in a fixed position at the base of the volcano in the *Scari* area,  $\sim 2$ – $2.5$  km from the degassing vents on the volcano's summit (**Figures 1, 2**). The lidar operated during June 24–29, 2015, including an initial instrumental setup phase. Stable weather conditions



(temperature, 22–26 °C; no rainfall) persisted during the entire measurement window.

During operations, two large motorized elliptical mirrors (major axis: 450 mm) simultaneously aimed the laser beams and the telescope, allowing the laser beam of the lidar to scan the volcanic plume either horizontally (**Figure 2A**) or vertically (**Figure 2B**). In particular, during June 24–25, the volcanic plume was mainly dispersed northwards by gentle southerly winds. From our *Scari* observation point (**Figure 2**), the plume was seen to rise nearly vertically north of the *Pizzo* area (**Figure 2C**). The Line Of Sight (LOS) of the lidar was therefore pointed north of *Pizzo* and the horizontal scan mode was preferred (heading angles: 227–317°; **Figure 2A**). Vertical scans above the *Pizzo* area were also performed. For simplicity, we refer below to these June 24–25 measurements as the *Pizzo scans* (**Figure 2A**).

On June 26–29, the plume was instead transported south-southeast by the prevailing north-northwesterly winds (**Figure 2D**). Vertical scans were therefore preferred that were operated at constant heading angle (237.8°) and at elevation angles from 16 to 21° (**Figure 2B**). In such conditions, the *Pizzo* and *Vancori* peaks were intercepted at elevation angles of 16.98° and 17.78°, respectively, and the volcanic plume was in all cases encountered in the 2300–2700 m range. We hereafter refer to these scans as the *Vancori scans* (**Figure 2B**).

During each profile, 100 lidar returns, 50 at  $\lambda_{\text{ON}}$ , 50 at  $\lambda_{\text{OFF}}$ , and interlaced (OFF after ON, OFF again and so on), were emitted at a 10 Hz rate, then co-added and averaged to increase the signal-to-noise ratio, reducing the signal sampling frequency to 0.1 Hz (temporal resolution of 10 s). The spatial resolution was about 5 m (corresponding to the rise time of the detector module due to its bandwidth). Plume scans, both horizontal and vertical, were retrieved combining about 50 profiles in <10 min. Typically, 10 scans at different elevations were repeated, obtaining a three-dimensional tomography of the volcanic plume.

A cell filled with standard CO<sub>2</sub> gas was periodically used during operations, for check of wavelength accuracy, repeatability and stability. In brief, our calibration procedure involved measuring-by photoacoustic spectroscopy—the absorption of the CO<sub>2</sub> gas cell as a function of wavelength. This calibration, limited to a small interval near the predicted  $\lambda_{\text{ON}}$ , allowed identifying the wavelength at which cell absorption is maximum. The laser system was finally forced to transmit at this radiation. The CO<sub>2</sub> absorption cross-section used in our calculations was based on HITRAN data (Rothman et al., 2013).

## UV Camera

Concurrently with our lidar observations, a dual-UV camera system (Kantzas et al., 2010; Tamburello et al., 2012; Burton et al., 2014) was used to monitor the temporal variations of the SO<sub>2</sub> flux and plume transport speed. A fully autonomous system, similar to that used in other recent work (D'Aleo et al., 2016), was mounted on the roof of the laboratory truck and operated every day from 6 am to 4 pm (Local Time). The UV camera system acquired sequential images of the plume at ~0.5 Hz using two JAI CM 140 GR cameras. Both cameras

had 10-bit digitization and 1392 × 1040 pixels, using an Uka Optics UV lens with a ~37° field of view. Distinct band-pass filters, centered at either 310 nm (where SO<sub>2</sub> absorbs) or 330 nm (no SO<sub>2</sub> absorption), were mounted on the back on the lenses of the two cameras. Each set of co-acquired images from the two UV cameras was processed using the methodology of Kantzas et al. (2010) and integrated into the Vulcamera software (Tamburello et al., 2011, 2012), to calculate an absorbance for each camera pixel. Absorbance was converted into an SO<sub>2</sub> column amount from readings of a co-exposed Ocean optics USB2000+ UV Spectrometer, as outlined in Lübcke et al. (2013). Cameras and spectrometer were both controlled by a mini-pc Jetway. To calculate SO<sub>2</sub> flux time-series, we used Vulcamera to derive temporal records of SO<sub>2</sub> integrated column amounts (ICAs) along a plume cross-section, perpendicular to the plume transport direction. The obtained ICA time-series were then combined with high-temporal resolution (~1 Hz) records of plume transport speed. This latter was derived using an Optical Flow sub-routine using the Lukas/Kanade algorithm (Bruhn et al., 2005; Peters et al., 2015), integrated in Vulcamera. In our specific case, the Lucas-Kanade method was used to track movements of gas fronts (e.g., gas-rich and/or ash-free portions of the plume, having well distinct absorbance features) in consecutive UV camera frames, which allowed us quantifying plume transport speed at 0.5 Hz. We tested performance of this method by using artificial images with known particle velocities, and obtained errors in estimated velocities of <5%. **Table 1** lists daily means ( $\pm 1$  SD: standard deviation) of both SO<sub>2</sub> fluxes and plume transport speed ( $V_p$ ) during our observational period.

## RESULTS

### Characteristics of the DIAL-Lidar Signal

According to lidar theory (Fiorani, 2007), the optical power returned to the lidar receiver at any time  $t$  is produced by back-scattering of the laser beam by an atmospheric layer at distance  $R$  (range) from the source, where  $R = ct/2$  and  $c$  is the speed of light. As such, the lidar offers range-resolved information on atmospheric structure and properties (aerosols, particles and gas molecules) along the laser beam, in the form of an intensity ( $I$ ) vs. range plot (**Figure 3**).

Upon its atmospheric propagation, the beam intensity decreases approximately (a) exponentially, due to atmospheric extinction, according to the Lambert-Beer law; and (b) as  $1/R^2$ , because the solid angle subtended by the receiver is  $A/R^2$ , where  $A$  is the telescope's effective area. The two processes are superimposed. As such, in order to better observe the atmospheric back-scattering, a “range corrected signal,  $S$ ” is commonly used, being given by:  $S = \ln(I R^2)$  (see below). Since the system works in DIAL mode, each intensity profile is in fact acquired at two distinct wavelengths,  $\lambda_{\text{ON}}$ —absorbed by CO<sub>2</sub>—and  $\lambda_{\text{OFF}}$ —not absorbed by CO<sub>2</sub> (**Figure 3**). The two wavelengths are so close that atmospheric behavior, except from CO<sub>2</sub> absorption, is practically identical. The measured intensity contrast between the co-emitted  $\lambda_{\text{ON}}$  and  $\lambda_{\text{OFF}}$  signals allow

**TABLE 1 | Results of volcanic gas plume observations at Stromboli volcano, 24–29 June 2015.**

Date	Plume speed (m/s)		SO <sub>2</sub> flux (kg/s)		CO <sub>2</sub> /SO <sub>2</sub> (Molar ratio)		CO <sub>2</sub> flux <sup>a</sup> (kg/s)		CO <sub>2</sub> flux <sup>b</sup> (kg/s)	
	Mean	(1 SD)	Mean	(1 SD)	Mean	(1 SD)	Mean	(Uncertainty) <sup>c</sup>	Mean	(Uncertainty) <sup>d</sup>
Method	UV camera		UV camera		Multi-GAS		Multi-GAS + UV camera		DIAL-LIDAR	
24/6/15	5.5	0.3	1.5	1.2	8	1.2	8.4	8.0	8.3	2.1
25/6/15	4.8	0.3	n.d.	n.d.	n.d.	n.d.	n.d.	n.d.	18.1	4.5
26/6/15	5.2	0.3	2.5	0.8	7.9	1.2	13.4	6.3	16.5	4.1
27/6/15	6.1	0.3	3.4	2.1	8.6	1.3	20.1	15.3	13.5	3.4
28/6/15	6.1	0.4	2.7	2.3	8.9	1.3	16.7	16.6	8.6	2.2
29/6/15	5.5	0.4	n.d.	n.d.	n.d.	n.d.	n.d.	n.d.	9.2	2.3

Plume speed and SO<sub>2</sub> flux are obtained by processing UV camera images. For both parameters the daily average and its standard deviation (SD) are quoted (the latter is taken as representative of uncertainty). The plume volcanic gas CO<sub>2</sub>/SO<sub>2</sub> ratios are derived from in-situ Multi-GAS observations taken on the volcano’s summit; each quoted ratio is the average (+1 SD) over a 30-min observational period, from 16 to 16:30 Local Time. No successful Multi-GAS plume detection was obtained in other daily observational windows (04-04:30; 10-10:30-22-22:30). Two independent estimates of the CO<sub>2</sub> flux are reported, based on either:

<sup>a</sup>Multiplying the SO<sub>2</sub> flux by the CO<sub>2</sub>/SO<sub>2</sub> ratio, or

<sup>b</sup>Processing of DIAL-LIDAR results. Uncertainties in the derived CO<sub>2</sub> fluxes are from either

<sup>c</sup>Error propagation on SO<sub>2</sub> fluxes and CO<sub>2</sub>/SO<sub>2</sub> ratios (taken as 1 SD), or

<sup>d</sup>Estimated at ±25% (see appendix).

range-resolved CO<sub>2</sub> concentrations in the volcanic plume to be obtained.

An example of a lidar-based atmospheric profile, obtained at Stromboli during a typical *Vancori scan*, is illustrated in **Figure 3**. As described above, the lidar registered one such profile every 10 s, since 100 lidar returns acquired at 10 Hz were co-added and averaged to increase the signal-to-noise ratio. Each of the atmospheric profiles (e.g., **Figure 3**) acquired during the *Vancori scan* contains the following characteristic features:

- 1) at R ~0, a first strong intensity peak is recorded for both λ<sub>ON</sub> and λ<sub>OFF</sub> (**Figure 3A**); this peak, which we refer to as I<sub>0,ON</sub> and I<sub>0,OFF</sub>, is due to scattering inside the laboratory truck of some photons of the transmitted laser pulse. This peak yields the pulse transmission zero-time, and its intensity is proportional to the transmitted energy (used for signal normalization),
- 2) for R between 0 and ~500 m, a weak signal is observed that is returned from atmospheric back-scatterers encountered by the laser beam along the optical path (**Figure 3A**); this signal, as explained before, attenuates with distance and vanishes at R ~500 m,
- 3) a I<sub>P,ON</sub> and I<sub>P,OFF</sub> peak at R ~1900 m (**Figures 3A,B**); this is produced by reflection of the lidar beam by the southeastern margin of the *Pizzo* morphological peak (see **Figure 2**),
- 4) a series of weak but resolvable peaks observed in the range interval 2300–2700 m (**Figure 3B**); in these peaks, the λ<sub>ON</sub> signal appears strongly attenuated relative to the co-acquired λ<sub>OFF</sub> signal, a fact due to laser absorption by CO<sub>2</sub> molecules in the volcanic plume,
- 5) a I<sub>V,ON</sub> and I<sub>V,OFF</sub> peak at R ~2800 m, which is produced by reflection of the laser beam by the *Vancori* peak (**Figure 3B**).

Atmospheric profiles obtained during the *Pizzo scans* of June 24–25 share similar characteristics, except that the *Pizzo* morphological peak is intercepted by the lidar beam at R ~2300 m, and the plume is encountered either before or after the *Pizzo* (**Figure 4**). The *Vancori* peak was obviously not encountered.

## Data Processing and Calculation of CO<sub>2</sub> Concentrations

We processed each acquired atmospheric profile using a Matlab analysis routine, with the aim of calculating the CO<sub>2</sub> concentrations in the atmospheric background and in the volcanic plume. The data processing routine consists of the following steps, all based on the Lambert-Beer law relation:

- a) Initially, the CO<sub>2</sub> concentration in the natural background atmosphere, C<sub>0</sub>, is calculated as:

$$\ln \left[ \frac{I_{P,OFF}/I_{0,OFF}}{I_{P,ON}/I_{0,ON}} \right] = 2 \Delta \sigma C_0 R_P \quad (1)$$

where I<sub>P,ON</sub> (I<sub>P,OFF</sub>) stands for intensity of the ON (OFF) lidar signal [(3) in **Figure 3A**] caused by reflection of the laser beam off the surface of the *Pizzo* wall (R<sub>P</sub> = 2294 m); I<sub>0,ON</sub> (I<sub>0,OFF</sub>) is the intensity of the ON (OFF) lidar peak caused by laboratory scattering of the laser pulse [(1) in **Figure 3A**]; and Δσ is the CO<sub>2</sub> differential absorption cross-section.

- b) Secondly, ΔC, the average excess CO<sub>2</sub> concentration in the volcanic plume cross-section between *Pizzo* and *Vancori* [(3,5) in **Figure 3B**], is derived from:

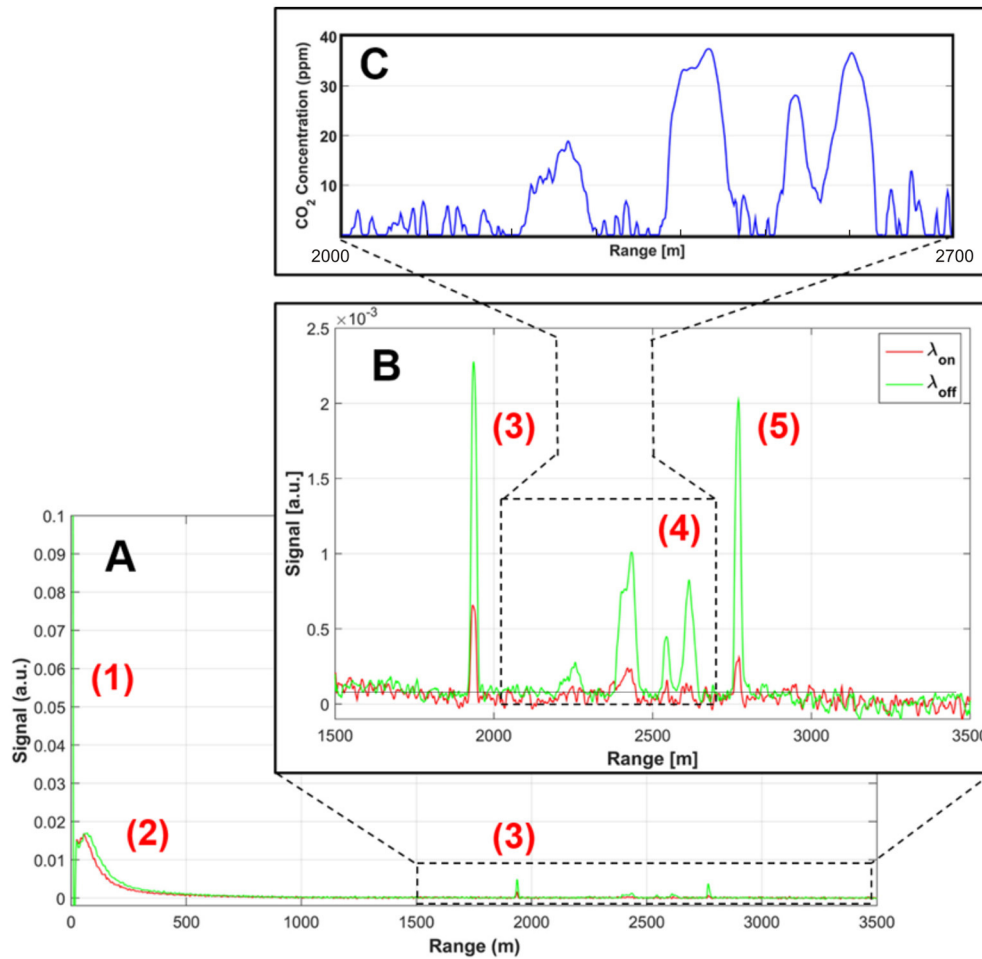
$$\ln \left[ \frac{I_{V,OFF}/I_{P,OFF}}{I_{V,ON}/I_{P,ON}} \right] = 2 \Delta \sigma (C_0 + \Delta C) (R_V - R_P) \quad (2)$$

Where I<sub>V,ON</sub> (I<sub>V,OFF</sub>) is the peak intensity of the ON (OFF) lidar signal caused by reflection of the laser beam off the surface of the *Vancori* rock wall (at R<sub>V</sub> = 2837 m).

- c) Thirdly, C<sub>CO<sub>2</sub>,i</sub>, the excess CO<sub>2</sub> concentration corresponding to each i-th ADC channel of the lidar profile (**Figure 3C**) is calculated from:

$$C_{CO_2,i} = k S_i \quad (3)$$





**FIGURE 3 | (A)** Example of a lidar-based atmospheric profile, obtained at Stromboli during a typical *Vancori scan*, in the form of a range (distance) vs. signal intensity (arbitrary units, a.u.) plot. Peak (1) yields the pulse transmission zero-time (scattering, inside the laboratory truck, of some photons of the transmitted laser pulse); peak (2) is the returned signal from atmospheric back-scatterers along the laser optical path; peak (3) is the returned signal produced by reflection of the lidar beam by southeastern margin of the *Pizzo* morphological peak. **(B)** is a detail of **(A)**, for ranges between 1500 and 3500 m. In this panel, peak (3) is as in **(A)**; the series of peaks observed in the range interval 2300–2700 m (4) are due to back-scattering of the laser beam from the volcanic plume; peak (5) is produced by reflection of the laser beam by the *Vancori* peak; **(C)** a profile of in-plume excess CO<sub>2</sub> concentrations, in the 2000–2700 m range interval, calculated from processing of the lidar signal in **(B)**. See text for the procedure used.

$$k = \frac{\Delta C (R_V - R_P)}{\Delta R \sum_i S_i} \quad (4)$$

$$S_i = \ln (I_{i,OFF} R_i^2) \quad (5)$$

where  $\Delta R$  is the range interval corresponding to each ADC channel, and  $I_{i,OFF}$  and  $R_i$  are the OFF lidar signal and the range of the  $i$ -th ADC channel (the OFF signal has been chosen because its signal-to-noise ratio is higher). **Figure 3C** shows an example of in-plume excess CO<sub>2</sub> concentration profile, obtained by applying the procedure above to the lidar profile of **Figure 3B** (in the 2100–2700 m range interval, where the volcanic plume was detected).

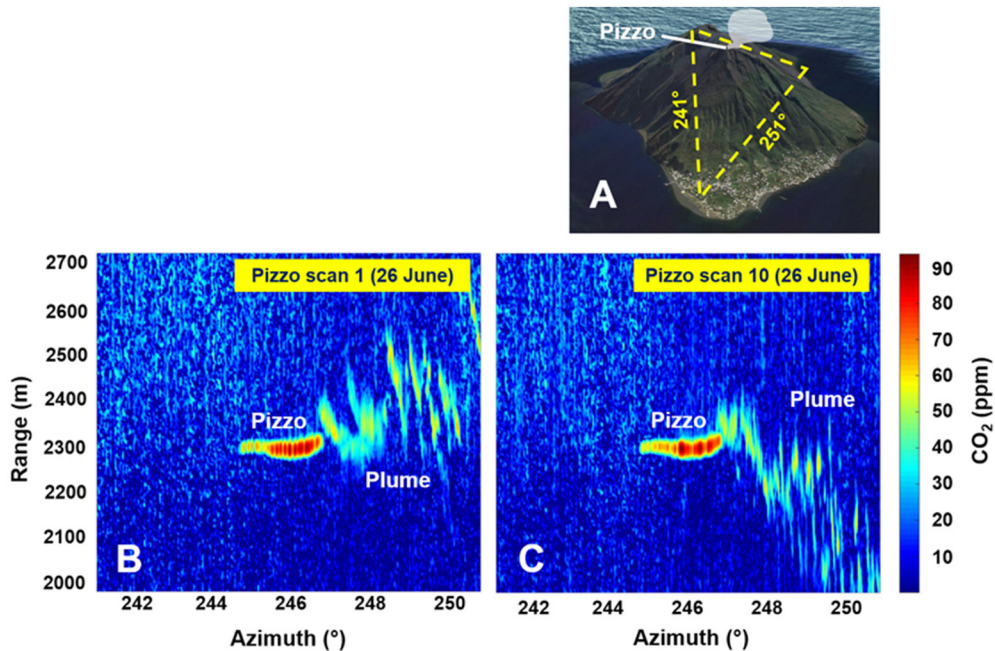
### In-Plume CO<sub>2</sub> Concentration Maps

A series of CO<sub>2</sub> concentration profiles (one every 10 s), all similar to those shown in **Figure 3C**, were obtained as the volcanic plume

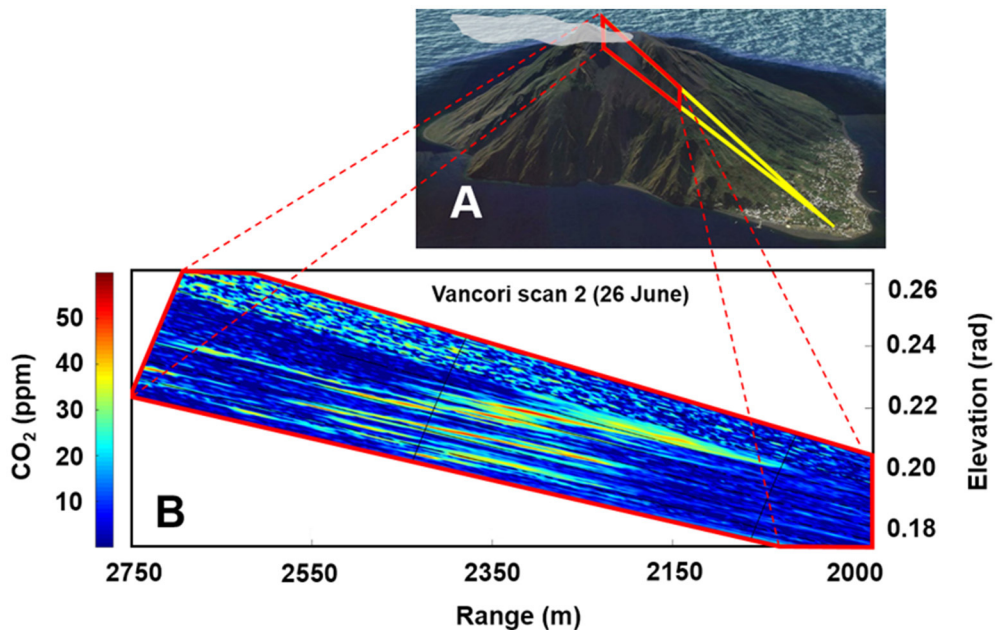
was sequentially scanned by our DIAL-lidar, either horizontally or vertically, during the *Pizzo/Vancori scans*. By interpolating all CO<sub>2</sub> concentration profiles obtained during a single scan, we obtained sequences of CO<sub>2</sub> concentration maps, examples of which are shown in **Figures 4, 5**. Since a full scan of the plume was completed in ~1000–2000 s, each map is in fact obtained from the combination of ~50 to ~100 atmospheric profiles.

The maps illustrate the 2D distribution of CO<sub>2</sub> concentrations as a function of azimuth angle [°] (X axis) and range [m] (Y axis) for horizontal scans (**Figure 4**); or as a function of range [m] (X axis), and elevation angle [°] (Y axis) for the vertical scans (**Figure 5**). In both plots, the color scales (from blue to red) illustrate the level of CO<sub>2</sub> concentrations (in [ppm]) in the investigated space.

**Figures 4, 5** demonstrate the ability of our DIAL-lidar to resolve in-plume volcanic CO<sub>2</sub> from the atmospheric



**FIGURE 4 | Example of two CO<sub>2</sub> concentration maps (B,C)** obtained during a *Pizzo* horizontal scan on June 26. Geometry of the scans and location of the plume are schematically shown in (A). The maps show the distribution of CO<sub>2</sub> concentrations in the lidar's Field Of View (FOV), as a function of heading angle and range. Each map was obtained by interpolation of all CO<sub>2</sub> concentration profiles (e.g., same as 3A), obtained during a given *Pizzo* scan. In the maps, the red colored horizontal bands identify the margin of the *Pizzo* peak (heading angle: 244–245°), while the volcanic plume is the band of peak CO<sub>2</sub> concentration (up to 60 ppm) areas at heading angles of 245–250°.



**FIGURE 5 | Example of a CO<sub>2</sub> concentration map (B)** obtained during a *Vancori* vertical scan on June 26. Geometry of the scan and location of the plume is schematically shown in (A). The map shows the distribution of CO<sub>2</sub> concentrations in the lidar's FVO, as a function of range and elevation, and was obtained by interpolation of all CO<sub>2</sub> concentration profiles (e.g., same as 3A) obtained during a given *Vancori* scan. In the map, the volcanic plume corresponds to the cluster of high CO<sub>2</sub> concentrations (up to 60 ppmv) in the range interval 2200–2500 m and 0.2–0.22 rad (17–19.5°) elevations. The blue colored areas in the 2000–2000 m range correspond to near ambient (<10 ppm above background) CO<sub>2</sub> concentrations.

background CO<sub>2</sub> (blue colors). In the CO<sub>2</sub> distribution maps, clusters of peak CO<sub>2</sub> concentration areas (marked by red, orange and yellow colors) identify the geometry of the plume. The lidar-based plume locations are consistent with visual and UV observations of volcanic plume dispersion (Figure 2). In the *Pizzo* horizontal scans, the plume was intercepted north of the *Pizzo* peak (heading angle: 244–245°), and is identified in the maps of Figure 4 as a cluster of peak CO<sub>2</sub> concentrations (up to 60 ppm above ambient air) at heading angles of 245–250°. The plume was detected over a relative wide range interval ( $R = 2000\text{--}2400$  m), relative to the *Pizzo* peak ( $R \sim 2300$  m). This is consistent with the slightly variable plume transport directions during our June 24–25 observation period that dispersed the plume either toward (Figure 4B) or away from (Figure 4A) the lidar observation point ( $R = 0$ ). A few *Pizzo* vertical scans (not shown) confined the vertical extension of the plume to a diagonal band, extending from  $R = 2300$  m and elevation  $\sim 19^\circ$  (the *Pizzo* area) to  $R \sim 2700$  m and  $\sim 20^\circ$  elevation. Figure 5 is an example of a CO<sub>2</sub> distribution map obtained during a vertical *Vancori* scan. The map exhibits a clear volcanic plume signal, as marked by a cluster of high CO<sub>2</sub> concentrations (up to 60 ppm) in the range interval 2200–2500 m and 17–19.5° elevation. CO<sub>2</sub> remained at background air levels for range distances  $<2000$  and  $>2800$  m.

## CO<sub>2</sub> Flux

The CO<sub>2</sub> concentration maps served as basis for calculating the CO<sub>2</sub> flux. To this aim, and by analogy with previous work (Aiuppa et al., 2015), we integrated the background-corrected (excess) CO<sub>2</sub> concentrations over the entire plume cross-sectional area covered by each scan, and multiplied this integrated column amount by the plume transport speed. Mathematically, the CO<sub>2</sub> flux ( $\Phi_{\text{CO}_2}$ , in  $\text{kg}\cdot\text{s}^{-1}$ ) was obtained from:

$$\Phi_{\text{CO}_2} = v_p \cdot \frac{PM_{\text{CO}_2}}{10^3 N_A} \cdot N_{\text{molCO}_2\text{-total}} \quad (6)$$

where  $v_p$  is the plume transport speed (in [m/s]) obtained from processing of UV camera images (Table 1);  $N_{\text{molCO}_2\text{-total}}$  is the total-plume CO<sub>2</sub> molecular density (expressed in molecules  $\text{m}^{-3}$ ); and  $PM_{\text{CO}_2}$  and  $N_A$  are, respectively, the CO<sub>2</sub> molecular weight and Avogadro's constant. The term  $N_{\text{molCO}_2\text{-total}}$  was obtained by integrating the effective average excess CO<sub>2</sub> concentrations ( $\overline{C_{\text{exc},i}}$  [ppm]) over the entire plume cross section, according to:

$$N_{\text{molCO}_2\text{-total}} = N_h \cdot 10^{-6} \cdot \sum_i \overline{C_{\text{exc},i}} \cdot A_i \quad (7)$$

where:  $N_h$  is the atmospheric number density (molecules  $\text{m}^{-3}$ ) at the crater's summit height, the term  $10^{-6}$  converts  $\overline{C_{\text{exc},i}}$  into a dimensionless quantity, and  $A_i$  represents the  $i$ -th effective plume area, given by:

$$A_i = l_i \cdot \Delta R \quad (8)$$

where  $\Delta R$  is the spatial resolution of the lidar (1.5 m) and  $l_i$  is the  $i$ -th arc of circumference (Figure 6B):

$$l_i = R_i \cdot \theta \quad (9)$$

In relation (9),  $R_i$  is the  $i$ -th distance vector (in meters) and  $\theta$  is the angular resolution of the system expressed in radians (ranging from  $0.04^\circ \pi/180 = 1.75 \cdot 10^{-4}$  rad to  $0.1^\circ \pi/180 = 0.00175$  rad) (Figure 6B).

Our obtained CO<sub>2</sub> fluxes, shown in Figure 6A, range from 1.8 to 32.1 kg/s. The lidar-based CO<sub>2</sub> flux time-series (Figure 6A) has maximum temporal resolution of 16–33 min (the time required to complete a full scan of the plume for our instrumental configuration). Temporal gaps in the dataset are caused by decreases in the signal-to-noise ratio (SNR) that prevent us from accurately detecting a clear CO<sub>2</sub> excess. These SNR decreases are likely caused by reduction of the backscattering coefficient of the probed air parcel, reflecting temporal variations in condensation extent of the volcanic gas plume. Visual (and UV camera) observations confirmed that the plume was variably condensed during our measurement interval, possibly due to slight changes in atmospheric conditions.

We evaluate the overall uncertainty in our derived CO<sub>2</sub> fluxes at  $\pm 25\%$  at 1s (see appendix).

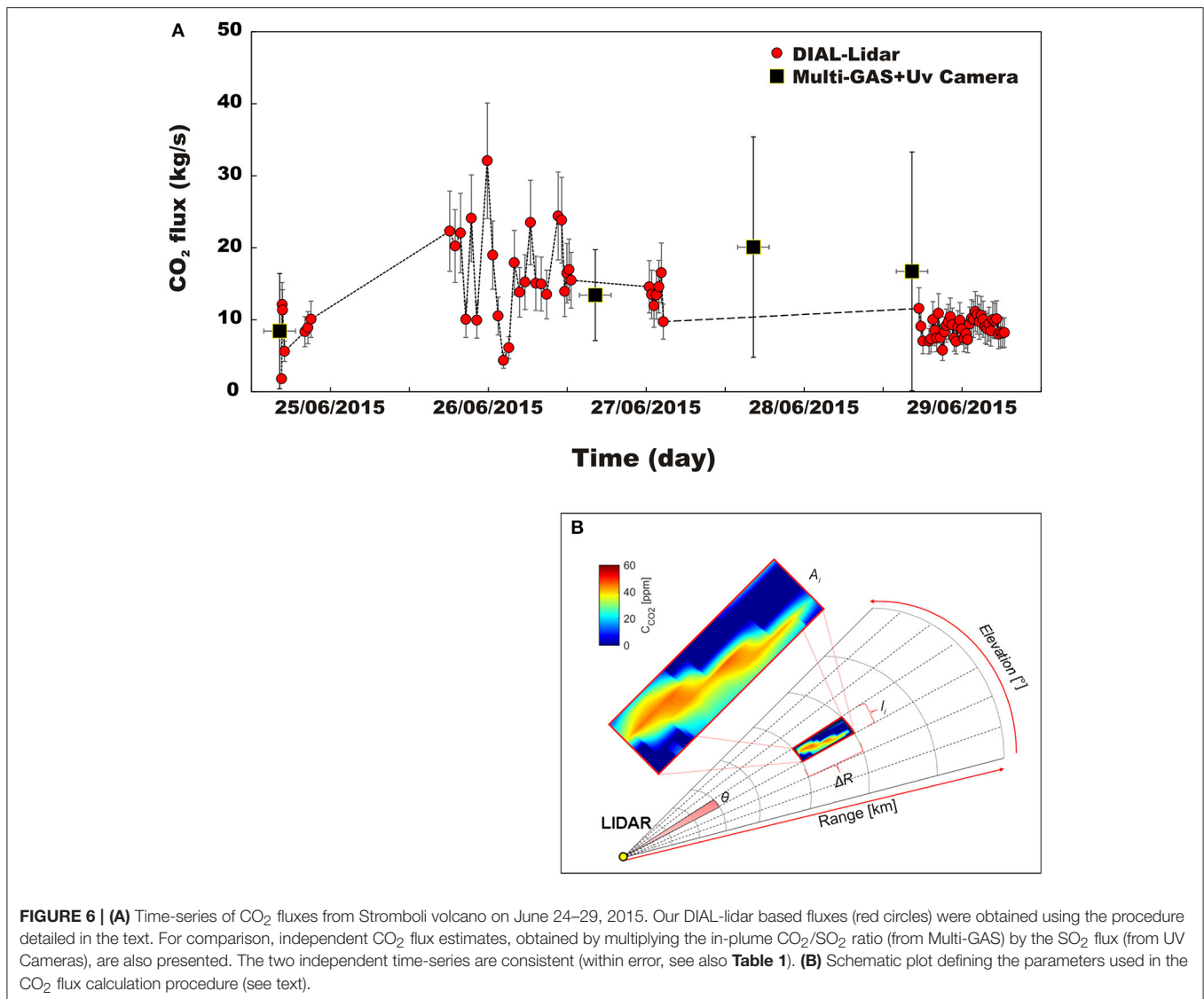
## DISCUSSION

The scarcity of volcanic CO<sub>2</sub> flux data in the geological literature (see Burton et al., 2013 for a recent review) is a direct consequence of the technical challenges in resolving the volcanic CO<sub>2</sub> signal from the large atmospheric background ( $\approx 400$  ppmv). In contrast to SO<sub>2</sub>, which is present at the part per billion level in the background atmosphere, allowing the volcanic flux to be routinely measured from ground and space using UV spectroscopy (Oppenheimer, 2010), remote sensing of volcanic CO<sub>2</sub> has only been achieved during eruptions of mafic volcanoes. In such circumstances, magma/hot rocks can effectively be used as a light source for ground-based Fourier Transform Infra-Red (FTIR) spectrometers (Allard et al., 2005; Burton et al., 2007; Oppenheimer and Kyle, 2008). In contrast, measurement of the far more common "passive" CO<sub>2</sub> emissions from quiescent volcanoes has required access to hazardous volcano's summit craters for direct sampling of fumaroles (Fischer and Chiodini, 2015) or *in-situ* measurement of plumes via either Multi-GAS instruments (Aiuppa, 2015) or active-FTIR (Burton et al., 2000; La Spina et al., 2013; Conde et al., 2014).

A major breakthrough has recently arisen from the possible application of lidars to remote volcanic CO<sub>2</sub> sensing (Fiorani et al., 2013, 2015, 2016; Aiuppa et al., 2015; Queißer et al., 2015, 2016). Aiuppa et al. (2015) were the first to report on a DIAL-lidar-based remote measurement of the volcanic CO<sub>2</sub> flux at Campi Flegrei volcano, but their observations were limited to short ( $<200$  m) measurement distances. Here, we have extended this earlier work to demonstrate that DIAL-lidars can successfully detect volcanic CO<sub>2</sub> at tens of ppmv above the atmospheric background over optical paths up to  $\approx 3$  km (Figures 4, 5). Similar results have recently been obtained at Campi Flegrei volcano by Queißer et al. (2016), suggesting that lidar may soon become an important operational tool in volcanic-gas research.

Our results constrain the CO<sub>2</sub> flux at Stromboli during June 24–29, 2015 (Figure 6A). Averaging all successful results during



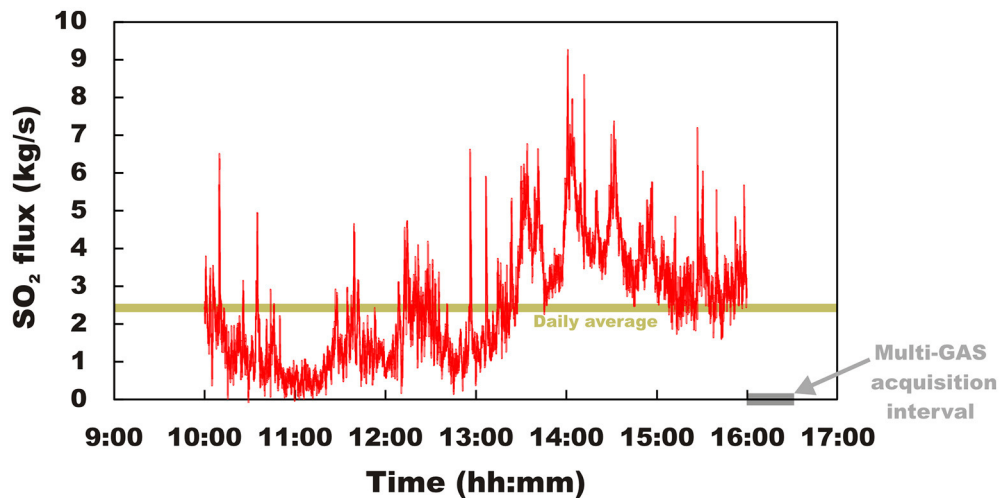


each measurement day, we obtain daily averages of the CO<sub>2</sub> flux between  $8.3 \pm 2.1$  (June 24) and  $18.1 \pm 4.5$  (June 25) kg/s, which correspond to cumulative daily outputs of 718 and 1565 tons, respectively. These results fall well in the range of previous CO<sub>2</sub> measurements on Stromboli. Aiuppa et al. (2010a, 2011) found that the CO<sub>2</sub> flux exhibits large temporal oscillations on Stromboli, from as low as 60 tons/day to as high as 11,000 tons/day, the highest values being observed in the days prior to *paroxysmal* and/or *major* explosions. The time-averaged CO<sub>2</sub> flux from Stromboli has been evaluated at 550 tons/day (Aiuppa et al., 2011) and at 1040–1200 tons/day (Allard, 2010). Our lidar-based CO<sub>2</sub> flux for the entire (June 24–29) measurement period is reasonably close, averaging at  $1050 \pm 250$  tons/day (mean of 80 individual measurements).

**Figure 6A** offers further confirmation to the robustness of our results. In the figure, we compare our lidar-based CO<sub>2</sub> fluxes with independent estimates, in which the CO<sub>2</sub> flux was derived by multiplying the CO<sub>2</sub>/SO<sub>2</sub> ratio of the plume by the SO<sub>2</sub> flux. This

latter approach has been used at volcanoes for years (Aiuppa, 2015), and at Stromboli involves use of two fully automated Multi-GAS instruments, operating on the volcano's summit to measure the in-plume CO<sub>2</sub>/SO<sub>2</sub> ratio (**Figure 2A**; Aiuppa et al., 2009, 2010a,b; Calvari et al., 2014). This is combined with SO<sub>2</sub> fluxes, delivered from either the FLAMES network of scanning UV spectrometers (Burton et al., 2009) or from UV camera observations (Tamburello et al., 2012), to obtain the CO<sub>2</sub> flux.

Problems with this Multi-GAS + SO<sub>2</sub> flux approach include issues of different temporal resolutions and poor temporal alignment of the two time-series. Successful Multi-GAS measurements of plume composition on Stromboli (Aiuppa et al., 2009, 2010b) are restricted to periods when the volcanic plume is dispersed by the local wind field into the *Pizzo* area, where the instruments are deployed (see **Figure 2A**). In addition, the Multi-GAS cannot operate continuously, but only during four equally spaced measurement cycles per day, each being 30 min long (Aiuppa et al., 2009). As such, the temporal resolution of the



**FIGURE 7 |** Temporal record of the volcanic SO<sub>2</sub> flux from Stromboli on June 26<sup>th</sup>, 2015, as derived from our UV camera observations. The figure exemplifies misalignment between Multi-GAS and SO<sub>2</sub> flux time-series; plume CO<sub>2</sub>/SO<sub>2</sub> ratios on June 26<sup>th</sup> were successfully measured only during the 1600–1630 local time Multi-GAS acquisition period, immediately after the end of the SO<sub>2</sub> flux acquisition window (0900–1600 local time). Poor temporal alignment is a flaw in the technique of estimating the CO<sub>2</sub> flux through a combination of Multi-GAS and UV camera records.

CO<sub>2</sub>/SO<sub>2</sub> ratio time-series is 6 h at best. In contrast, the temporal resolution of UV spectrometers/cameras is higher, from ~10 to 20 min (Burton et al., 2009) to 0.5 s (Tamburello et al., 2012), but observations are intrinsically limited to daylight hours and to good meteorological conditions (no clouds).

Figure 7 exemplifies the issue related to misalignment between Multi-GAS and UV observations. In the June 26<sup>th</sup> example, the only successful Multi-GAS acquisition period (from 1600 to 1630 h local time) clearly did not overlap with the SO<sub>2</sub> flux acquisition window (0900 to 1600 h local time). To overcome this problem, the common practice is to average out available Multi-GAS and UV spectroscopy data to obtain *daily means* of the CO<sub>2</sub> flux (Aiuppa et al., 2010a). Owing to the large inter-daily variability of SO<sub>2</sub> flux (e.g., Figure 7), however, large uncertainties are associated with these derived CO<sub>2</sub> fluxes (see Table 1, and errors bars in Figure 6A).

In spite of the issues above, we find overall consistency between the lidar-based and the traditional (Multi-GAS + UV spectroscopy-based) CO<sub>2</sub> fluxes (Figure 6A). This provides mutual validation for both quantification approaches. Our lidar-based CO<sub>2</sub> flux time-series (Figure 6A) are manifestly more continuous and of better temporal resolution (16–33 min). In addition, the lidar as with other remote sensing techniques is intrinsically safer. We caution, however, that further development is required before the lidar can become an operative tool for volcano monitoring. Improvements will need to occur in portability (the prototype weighs ~1100 kg), and reduced power requirement (6.5 kW) and costs (300 kUS \$). In addition, the current measurement protocol is complex and thus requires great familiarity with the technique. Efforts are now being made to make the lidar more simple, user-friendly and fully automated, including development of an on-line remote control system and of a self-checking routine of the laser's wavelength settings.

Electro-optics and laser/lidar private manufacturers need to be directly involved to transition the prototype into a more widely accessible, commercial instrument.

## CONCLUSIONS

Our proof-of-concept study demonstrates the ability of DIAL-lidars to remotely ( $\approx 3$  km distance) measure the volcanic CO<sub>2</sub> flux. Our reported lidar-based CO<sub>2</sub> fluxes at Stromboli volcano ( $1.8 \pm 0.5$  to  $32.1 \pm 8.0$  kg/s) are in the same range as those obtained using standard techniques that require in-situ observations and are intrinsically more risky for operators. Our results, with those of QueiBer et al. (2016), open new prospects for the use of lidars for instrumental remote monitoring of volcanic CO<sub>2</sub> flux. Further work is warranted in order to standardize and widen potential applications of Lasers in volcanic gas studies.

## AUTHOR CONTRIBUTIONS

AA, LE, and SS conceived the idea. AA, LE, SS, GM, and MN conducted the lidar/UV camera experiment. SP and RD processed the data, with help from AA, LE, and SS. ML provided the Multi-GAS results. AA drafted the manuscript with help from all co-authors.

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**Conflict of Interest Statement:** The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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## APPENDIX—UNCERTAINTY AND ERROR ANALYSIS

Our lidar-based CO<sub>2</sub> fluxes are affected by the following error sources:

- i. systematic error in CO<sub>2</sub> concentration measurement,
  - ii. statistical error in CO<sub>2</sub> concentration measurement,
  - iii. error in plume transport speed,
  - iv. error in identifying the integration area.
- i. *Systematic error of the CO<sub>2</sub> concentration measurement* - It is well known that the DIAL-lidar systematic error is dominated by imprecision in wavelength setting (Fiorani et al., 2015), leading to inaccuracy in differential absorption cross section and thus in gas concentration. To minimize this error, we implemented a photo-acoustic cell filled with pure CO<sub>2</sub> at atmospheric pressure and temperature, close to the laser exit, in order to control the transmitted wavelength before each atmospheric measurement. This procedure allows us to set the ON/OFF wavelengths with better accuracy than the laser linewidth (Fiorani et al., 2016). Assuming that the error in the wavelength setting is  $\pm 0.02 \text{ cm}^{-1}$  (half laser linewidth), in the wavelength region used in this study, the systematic error of the CO<sub>2</sub> concentration measurement is 5.5%.
  - ii. *Statistical error of the CO<sub>2</sub> concentration measurement* - The statistical error has been calculated by standard error propagation techniques from the standard deviation of the lidar signal at each ADC channel. As discussed in Fiorani and Durieux (2001), the statistical error of the lidar signal increases with range. As a consequence, the uncertainty associated with the derived CO<sub>2</sub> concentrations also

increases with range. In the distance range between *Pizzo* and *Vancori*, representing a mean measurement range, and at typical atmospheric and plume conditions encountered during this study, the statistical error of the CO<sub>2</sub> concentration measurement was about 2%. The statistical error exceed 5% at 4 km (well beyond our measurement range).

- iii. *Error in plume transport speed* - The standard deviation and the average value of the wind speed have been calculated for each measurement session, and the corresponding relative error was evaluated (by error propagation technique) at 3%.
- iv. *Error in identifying the integration area* - The integration area in which an excess CO<sub>2</sub> concentration is actually present is probably the most difficult parameter to retrieve accurately, and therefore represents the main error source in our calculated volcanic CO<sub>2</sub> fluxes. The following procedure has been followed. For each CO<sub>2</sub> concentration map (e.g., **Figure 5**), we initially measured: 1) A<sub>15</sub>, the area where the excess CO<sub>2</sub> concentration was larger than 15 ppm; and 2) A<sub>25</sub> the area where the excess CO<sub>2</sub> concentration was larger than 25 ppm. Then, the average between A<sub>15</sub> and A<sub>25</sub> was taken as the best-estimated area, and their semi-difference as the error (~25%). The above thresholds have been chosen because below 15 ppm noise becomes significant, while above 25 ppm the plume area is reduced to its core. Use of a 15 ppm threshold likely underestimates the area (and thus the flux) of the order of magnitude of the measurement error, i.e., 10–20%.

Assuming that each error source is statistically independent, we can quadratically sum all the errors and obtain a cumulative error of ~ 25% (dominated by the area error).