



# Conduit Dynamics in Transitional Rhyolitic Activity Recorded by Tuffisite Vein Textures from the 2008–2009 Chaitén Eruption

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The mechanisms of hazardous silicic eruptions are controlled by complex, poorly-understood conduit processes. Observations of recent Chilean rhyolite eruptions have revealed the importance of hybrid activity, involving simultaneous explosive and effusive emissions from a common vent. Such behavior hinges upon the ability of gas to decouple from magma in the shallow conduit. Tuffisite veins are increasingly suspected to be a key facilitator of outgassing, as they repeatedly provide a transient permeable escape route for volcanic gases. Intersection of foam domains by tuffisite veins appears critical to efficient outgassing. However, knowledge is currently lacking into textural heterogeneities within shallow conduits, their relationship with tuffisite vein propagation, and the implications for fragmentation and degassing processes. Similarly, the magmatic vesiculation response to upper conduit pressure perturbations, such as those related to the slip of dense magma plugs, remains largely undefined. Here we provide a detailed characterization of an exceptionally large tuffisite vein within a rhyolitic obsidian bomb ejected during transitional explosive-effusive activity at Chaitén, Chile in May 2008. Vein textures and chemistry provide a time-integrated record of the invasion of a dense upper conduit plug by deeper fragmented magma. Quantitative textural analysis reveals diverse vesiculation histories of various juvenile clast types. Using vesicle size distributions, bubble number densities, zones of diffusive water depletion, and glass H<sub>2</sub>O concentrations, we propose a multi-step degassing/fragmentation history, spanning deep degassing to explosive bomb ejection. Rapid decompression events of ~3–4 MPa are associated with fragmentation of foam and dense magma at ~200–360 m depth in the conduit, permitting vertical gas and pyroclast mobility over hundreds of meters. Permeable pathway occlusion in the dense conduit plug by pyroclast accumulation and sintering preceded ultimate bomb ejection, which then triggered a final bubble nucleation event. Our results highlight how the vesiculation response of magma to decompression events is highly sensitive to the local melt volatile concentration, which is strongly spatially

heterogeneous. Repeated opening of pervasive tuffisite vein networks promotes this heterogeneity, allowing juxtaposition of variably volatile-rich magma fragments that are derived from a wide range of depths in the conduit. This process enables efficient but explosive removal of gas from rhyolitic magma and creates a complex textural collage within dense rhyolitic lava, in which neighboring fused clasts may have experienced vastly different degassing histories.

**Keywords:** volcanic eruption, rhyolite, obsidian, degassing, magma, diffusion, vesiculation, Chaitén volcano

## INTRODUCTION

Explosive-effusive behavioral transitions during silicic volcanic eruptions remain poorly understood, yet are typical of rhyolitic activity (e.g., Eichelberger et al., 1986; Gonnermann and Manga, 2003). The 2008–2009 rhyolitic eruption of Chaitén, in the Chilean Southern Volcanic Zone, is the first to be closely scientifically observed, and is providing unprecedented insights into processes controlling rhyolitic eruptive transitions (Castro et al., 2014). Chaitén was closely followed by the 2011–2012 Cordón Caulle eruption, also in Chile, which is providing additional insights into vent and plume processes (e.g., Schipper et al., 2013; Bonadonna et al., 2015). Notably, both recent eruptions were dominated by hybrid explosive-effusive behavior, which requires a new paradigm to explain the nature of eruptive controls (Schipper et al., 2013; Castro et al., 2014).

Explosive behavior is broadly characterized by magma fragmentation and the transition from a bubble-bearing liquid to a particle-bearing gas (e.g., Gonnermann and Manga, 2007). By contrast, fragmentation is less prevalent during effusive behavior, which is dominated by ductile magma deformation. Volatiles drive explosive volcanism (e.g., Sparks, 1978), yet initial volatile concentrations in explosive and effusive phases are typically similar (Eichelberger et al., 1986). Effusive behavior therefore requires efficient magma outgassing, with a significantly open degassing system, unlike the predominantly closed degassing system operative during explosive behavior (Newman et al., 1988). Outgassing is favored by slower magma ascent, which provides more time for gas separation via fractures or interconnected bubbles. In addition, a spatially variable ascent rate within an individual conduit may permit simultaneous explosive and effusive activity (Schipper et al., 2013). In recently observed Chilean rhyolitic systems, the high magma flux during the first effusive activity requires rapid, efficient gas separation from rising magma (Pallister et al., 2013; Bertin et al., 2015). However, rhyolitic melts are highly viscous, with viscosities increasing from  $\sim 10^5$  to  $\sim 10^{10}$  Pas upon degassing from 4% to 0.1 wt. % H<sub>2</sub>O (Giordano et al., 2008), hindering buoyant bubble rise and thus gas escape. Instead, connected permeable channels must be generated within the conduit, either within deforming, coalescing magmatic foams or fracture networks (Stasiuk et al., 1996; Okumura et al., 2010). The latter, which are generated via shear stresses or gas overpressure, are sometimes preserved as particle-clogged tuffisite veins hosted by ballistic bombs, conduit-filling lava or country rock (Heiken et al., 1988; Stasiuk et al., 1996; Gonnermann and Manga, 2003; Tuffen et al., 2003; Cabrera

et al., 2011; Holland et al., 2011; Castro et al., 2012, 2014; Berlo et al., 2013).

Our understanding of rhyolitic behavioral transitions remains severely limited by a lack of knowledge of the complex processes operating in shallow conduits. Improved constraints on tuffisite vein formation and mechanisms of gas-magma separation are required to refine conduit flow models and identify the thresholds that control hazardous eruptive behavior. In this paper we present a textural study of a tuffisite vein within a rhyolitic obsidian bomb ejected during the transitional phase of the Chaitén 2008 eruption. Pyroclast textures are deeply related to the dynamic characteristics of eruptions and thus provide useful insights into eruptive processes (e.g., Shea et al., 2010; Giachetti et al., 2011; Rust and Cashman, 2011; Gurioli et al., 2015). Here we apply this approach to pyroclasts trapped within the vein, with the aim of quantifying the poorly-constrained fragmentation, decompression and vesiculation processes that accompany tuffisite vein formation. Key uncertainties include the longevity of veins (permeability lifespan), how successive fracture-healing cycles are superimposed, and how veins interact with magmatic foams (Tuffen et al., 2003; Holland et al., 2011; Castro et al., 2012). In particular, we focus on bubble nucleation and growth within diverse juvenile vein-hosted clasts, to reconstruct a multi-step fragmentation and degassing history that culminated in explosive bomb ejection. Results lead to a new model that helps resolve conduit processes during the transitional phase of the Chaitén 2008–2009 eruption.

## SCIENTIFIC AND GEOLOGICAL CONTEXT

### Tuffisite Veins

Tuffisite veins are clast-filled fractures occurring within volcanic conduits and surrounding country rock. Internal, within-conduit tuffisites are networks of linked angular fractures <1 to >10 cm wide that are filled with juvenile clasts (dense lava and pumice) and lithic fragments. Clast sizes span microns-centimeters, and sedimentary structures in tuffisite vein-filling material indicate their defining feature: particle deposition from a mobile fluid phase (Stasiuk et al., 1996; Tuffen et al., 2003). Glassy clasts may be incipiently sintered to densely welded, with the densest welding destroying all porosity and creating viscously-deformed flow bands in silicic lava (Gonnermann and Manga, 2003; Tuffen et al., 2003). Tuffisites occur in dissected conduits, but also in bombs ejected in Vulcanian explosions and are inferred to feed explosive gas release from silicic vents when they intersect with the surface (Stasiuk et al., 1996; Tuffen and Dingwell, 2005;

Cabrera et al., 2011; Castro et al., 2012, 2014; Berlo et al., 2013; Schipper et al., 2013). However, deposition, sintering and compaction of melt-rich particles within veins eventually causes dramatic permeability reduction (Tuffen et al., 2003; Tuffen and Dingwell, 2005; Vasseur et al., 2013; Heap et al., 2015).

The prevalent model for vein generation is shear failure of non-Newtonian magma during shallow ascent in the conduit. Failure is favored by higher strain rates near conduit walls and magma stiffening via degassing and crystallization (e.g., Goto, 1998; Gonnermann and Manga, 2003; Tuffen et al., 2003; Neuberg et al., 2006). When a dense lava plug exists (the lava dome in the case of Chaitén), deeper magma ascent induces shear stresses in the plug, leading to shear failure and plug slip (Neuberg et al., 2006; e.g., Lensky et al., 2008; Holland et al., 2011). Connected fracture networks may briefly allow gas and ash venting at the surface (Schipper et al., 2013), which decreases vein pressure, permitting local degassing at vein walls via diffusion, whereas veins can also receive input of deeper magmatic gas and particles (Cabrera et al., 2011; Castro et al., 2012, 2014; Berlo et al., 2013). Vein wall abrasion may drive further assimilation of host magma particles (Tuffen and Dingwell, 2005).

Clasts within veins present textural evidence, including the presence of truncated flow bands, for repeated fracture-healing cycles. Integrated over time, these cycles permit outgassing and magma densification (Rust et al., 2004; Rust and Cashman, 2007; Cabrera et al., 2011; Castro et al., 2012), potentially assisted by viscous foam collapse and compaction (e.g., Stasiuk et al., 1996). The efficiency of vein-driven outgassing depends on the spatial extent of veins, the timescale over which their permeability is sustained, and the permeability of intercepted magma (Castro et al., 2012; Berlo et al., 2013; Heap et al., 2015). Contrasting clast volatile concentrations within studied tuffisite veins suggest vertical extents exceeding 100 m in shallow conduits (Berlo et al., 2013), whereas H<sub>2</sub>O diffusion gradients around Chaitén veins indicate diffusive lifetimes of ~15 min to 1 day (103–105 s; Castro et al., 2012). Such inferred timescales are consistent with those of vein closure via welding and viscous compaction (Tuffen et al., 2003). Such brief vein opening only permits local diffusive water depletion around veins (~100–200 μm), rendering outgassing inefficient without either unrealistically closely-spaced veins (<1 mm apart) or intersection with foam domains (Castro et al., 2012).

## Geological Context: The Chaitén 2008–2009 eruption

Volcán Chaitén, located at ~42°50'S in Chilean Patagonia (Figure 1A), produced an explosive rhyolitic eruption on May 1 2008, following only a brief (~24 h) period of detected precursory seismicity (Castro and Dingwell, 2009). Rapid, dyke-controlled magma ascent is inferred, from a magmatic source at 5–10 km depth (Wicks et al., 2011; Pallister et al., 2013). The initial purely explosive activity (phase 1) lasted 10 days, with a ≤19 km plume, followed by the transitional explosive-effusive phase 2 that lasted ~20 days (ending on May 31 2008). This hybrid phase was characterized by the emplacement of rhyolitic lava domes and coulées (Figure 1B) during a succession of

Vulcanian explosions that produced pyroclastic density currents (Alfano et al., 2011, 2012; Castro et al., 2012; Bernstein et al., 2013; Pallister et al., 2013). The ensuing phase 3, with months of exogenous/endogenous lava emplacement, ended in late 2009/early 2010 (Pallister et al., 2013).

The initial lava effusion rate which occurred at the onset of phase 2, 66 m<sup>3</sup>s<sup>-1</sup>, was the highest observed for any silicic eruption (Pallister et al., 2013). It has been since exceeded only by the initial >70 m<sup>3</sup>s<sup>-1</sup> flux at Cordón Caulle in June 2011 (Bertin et al., 2015). This high flux of dense, degassed magma implies that gas separation from magma became highly efficient within, at most, 10 days of the onset of the eruption (Castro et al., 2012). Such high lava fluxes reflect the relatively low viscosity of nearly-aphyric magma (<1 vol % crystals, Castro and Dingwell, 2009), compared with highly-crystalline andesites or dacites, and in total ~0.8 km<sup>3</sup> of lava was emitted (Pallister et al., 2013). The synchronicity of pyroclastic and lava emissions during hybrid phase 2 activity indicate that explosive outgassing processes facilitate emission of a high flux of degassed lava (Castro et al., 2014).

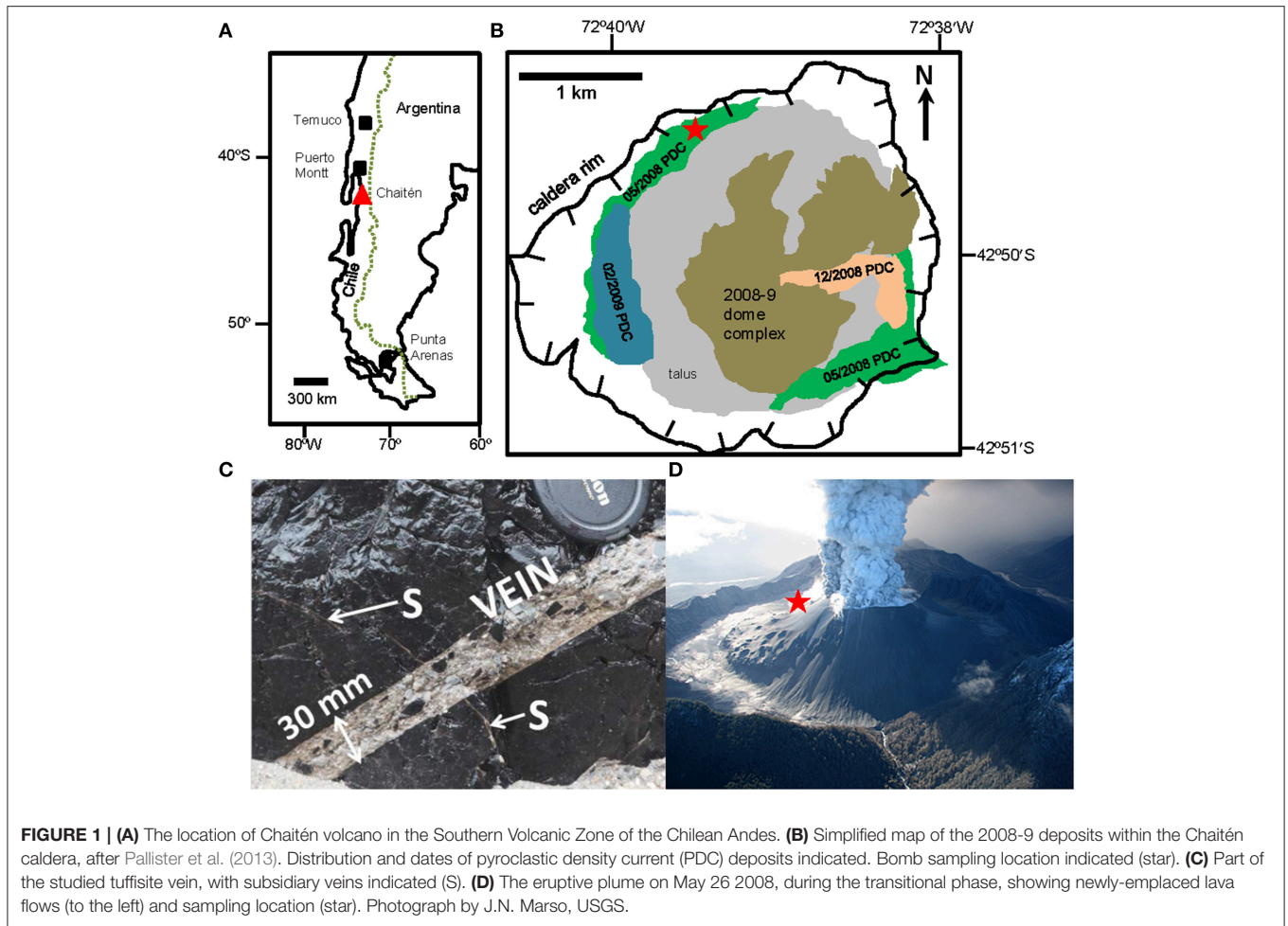
## Studied Tuffisite Vein

A 30 × 40 × 50 cm dense obsidian bomb (Figure 1C) ejected during the hybrid activity (phase 2) in May 2008 (Figure 1D) was sampled in the north-western part of the Chaitén caldera (42°49'33.66''S, 72°39'27.53''W), ~800 m from the 2008 vent. It is cut by a 3 cm-wide, parallel-walled, rectilinear tuffisite vein that is one of the widest and largest documented to date within any bomb. The *main vein* (Figures 1C,2) is filled with centimetric clasts of dense black obsidian similar to the host, alongside a variety of paler clast types, all within a fine-grained matrix. A network of far narrower, strikingly planar, intersecting veins (*subsidiary veins*; tens of microns to ~1 mm wide; Figures 1C,2) occur both sub-parallel to and propagating from the main vein walls into the host obsidian. Similar veins are also found in some of the larger dense obsidian clasts within the vein. Subsidiary veins are filled with very fine-grained, pale clastic material.

## METHODS

### Bulk Rock Analysis and Major Element Compositions

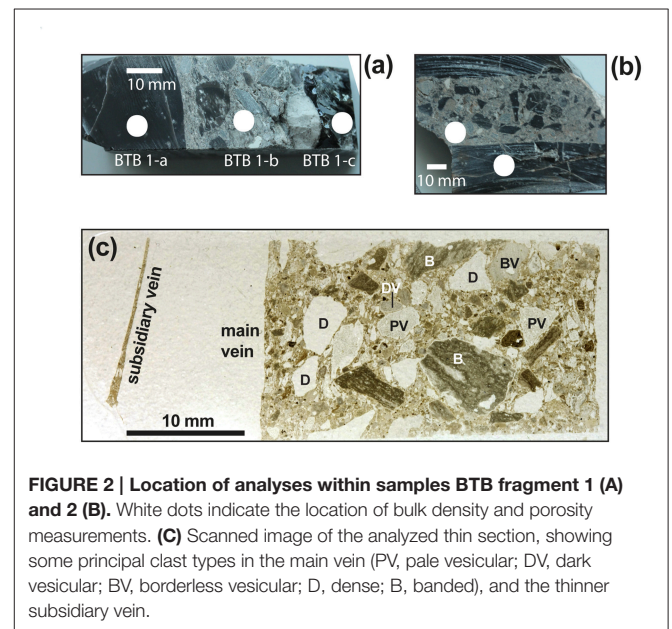
The density and porosity of the main vein and host obsidian were determined following the procedures outlined in Houghton and Wilson (1989), Formenti and Druitt (2003) and Takeuchi et al. (2008) at the Laboratoire Magmas et Volcans (LMV), of Clermont Ferrand, France. The tuffisites and the host obsidian were cut into rectangular prisms (BTB 1-a, BTB 1-b, BTB 1-c; Figure 2A). The porosity was calculated from each fragment using the dense rock equivalent (DRE) measured with the Accupyc 1340 Helium Pycnometer on the same powdered samples (Table 1). Description of the protocol for density and vesicle connectivity measurements is provided in Gurioli et al. (2015), and in the Supplementary Datasheets 1 and 2. The major element composition of the glasses and the mineral phases was characterized using electron probe microanalysis with a



CAMECA SX100 system at LMV, using spot sizes between 5 and 20 microns depending on local clast textures.

## Fourier Transform Infra-Red (FTIR) Spectroscopy

Glass water contents were determined using synchrotron-source infra-red spectroscopy (SFTIR) at the MIRIAM beamline, Diamond Light Source, UK, and glolar source infra-red spectroscopy at Lancaster University. The MIRIAM beamline consisted of a Bruker Vertex 80 V FTIR Interferometer, with KBr beamsplitter, coupled to a Hyperion 3000 microscope with a broadband MCT detector and a  $10\ \mu\text{m}$  square aperture. Analyses at Lancaster University employed a Thermo Nicolet IR interferometer at Lancaster University, with KBr beamsplitter, coupled to a Continuum Analytical microscope, with an MCT-A detector and  $75\ \mu\text{m}$  square aperture. Using the SFTIR system, 128 spectra were collected at  $8\ \text{cm}^{-1}$  spectral resolution within the mid-IR range of  $4000\text{--}1000\ \text{cm}^{-1}$ . Using the glolar source, 128 spectra were collected at  $4\ \text{cm}^{-1}$  resolution between  $6000$  and  $1000\ \text{cm}^{-1}$ . Eighteen-point linear baseline corrections were



**TABLE 1 | Bulk textural parameters for host obsidian and main tuffisite vein (see Figure 2A for measurement locations).**

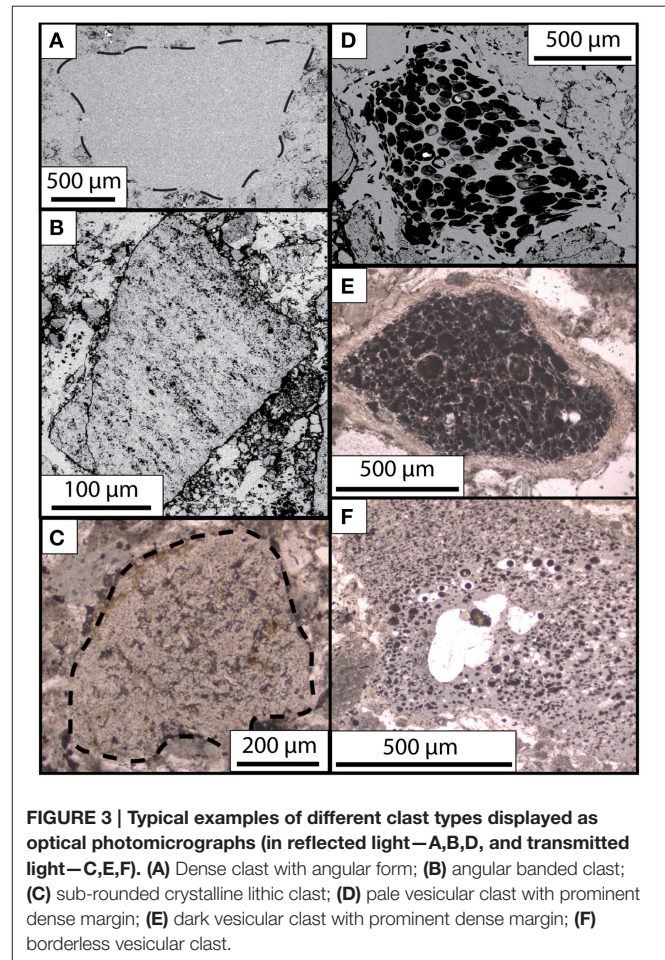
Sample (uncertainty)	Density kgm <sup>-3</sup> (± 30)	Porosity % (± 0.06)	Pore-free density kgm <sup>-3</sup> (±30)	Connectivity % (± 0.03 %)
BTB 1-a obsidian	2281	3.78	2371	90.85
BTB 1-c obsidian	2250	5.05	2370	92.37
BTB 1-b tuffisite vein	2070	15.13	2440	65.86

used to quantify absorption peaks at 3550 cm<sup>-1</sup> (H<sub>2</sub>O<sub>t</sub>) and 1630 cm<sup>-1</sup> (H<sub>2</sub>O<sub>m</sub>). These absorption values were next converted to species concentrations using the Beer-Lambert Law, using the measured glass density of 2281 kg m<sup>-3</sup> (Table 1). Wafer thicknesses were determined with a Mitutoyo digital micrometer (accuracy of ±3 μm) or reflection fringe widths, depending on sample thickness (von Aulock et al., 2014). Respective absorption coefficients are 80 and 55 l mol<sup>-1</sup> cm<sup>-1</sup> for the 3550 and 1630 cm<sup>-1</sup> peaks (Newman et al., 1986; Leschik et al., 2004). We calculated the uncertainty for our IR data using errors in thickness measurement, uncertainties in absorption coefficient, peak height measurement and sample density, and found the combined uncertainty to be less than 10% for all measurements. Thus, we use 10% as a conservative estimate for all FTIR error, as suggested by previous authors (e.g., von Aulock et al., 2014). Analytical results and details of error quantification are provided in Supplementary Datasheet 3.

## 2D and 3D Textural Imaging

Quantitative textural analysis was applied to the vein to provide the componentry, morphology, and size distribution of the different clasts within the tuffisites (Figures 2–4) as well as to quantify their internal vesicularity and vesicle size distribution. Overviews of sawn sample surfaces and a thin section (12 images) were gathered using a high-resolution digital scanner (5 and 10 μm pixel size, respectively). Higher resolution 1600 × 1200 pixel images were captured using a scanning electron microscope in backscatter mode on polished sections, with 0.13 μm pixel size at the highest resolution, and using petrological microscopy. Images were collected over a range of magnifications (25–150×) to obtain a representative particle size range, and processed used Adobe Photoshop, with grayscale thresholds selected to distinguish between clast types. Clasts were separately reconstructed, with edges hand-drawn. Vesicles were hand-drawn where they overlap in vesicular clasts, but automatically isolated within vesicle-poor clasts, using grayscale thresholds, and then manually reconstructed.

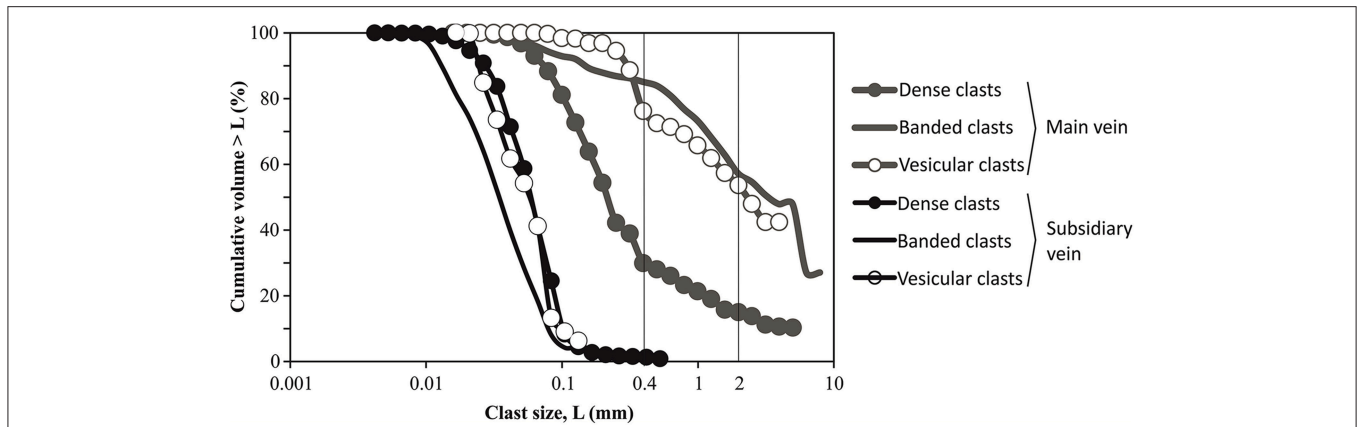
Thresholded grayscale images were analyzed using FOAMS (Shea et al., 2010, 2011). Exploiting the multi-scale functionality of FOAMS, we quantified the number, size and shape of particles (elongation, regularity, eccentricity, and aspect ratio) within the main and subsidiary tuffisite veins (Table 2), as well as the internal 2D vesicle size distributions and number densities. The smallest objects detected were 5 pixels in diameter. Binning for equivalent diameter of clasts and bubbles were calculated from the minimum diameter detected divided by the bin factor (10<sup>0.1</sup>). Prior to using FOAMS, clastic componentry within veins



**FIGURE 3 | Typical examples of different clast types displayed as optical photomicrographs (in reflected light—A,B,D, and transmitted light—C,E,F). (A) Dense clast with angular form; (B) angular banded clast; (C) sub-rounded crystalline lithic clast; (D) pale vesicular clast with prominent dense margin; (E) dark vesicular clast with prominent dense margin; (F) borderless vesicular clast.**

was quantified from overview images using ImageJ. Finally, the widths of vesicular clast dense borders were quantified in Photoshop. To quantify shape difference between clast types, we used two shape parameters calculated with FOAMS: clast elongation and regularity. The elongation comes from the equation  $\epsilon = (a - b)/(a + b)$ , where  $a$  and  $b$  are best-fit ellipse semi-long and semi-short axes respectively (Shea et al., 2010). The regularity is calculated from  $r = A/\pi ab$ , where  $A$  is the clast area, with a perfect circle having a regularity of 1 and an elongation of 0 (Shea et al., 2010).

Overall vein morphology was qualitatively examined in 3D using synchrotron source X-ray microtomography (μ-CT) performed in hutch 3B of the Australian Synchrotron (Clayton, Victoria), using their Ruby detector with X-ray energy of 45 kV. A



**FIGURE 4 | Clast size distributions of the different clast types within the main tuffsite vein and the subsidiary vein.** All components are considerably coarser in the main vein than the subsidiary vein.

**TABLE 2 | Detailed clast classification within the main tuffsite vein.**

Clast type	Name	Proportion	Color	Vesicles	Size	Borders	Fractures
1	Dense clasts	35 ± 3% (52 ± 4%)	Pale	Absent or < 1 %	< 5 mm (< 0.6 mm)	Sub-rounded to angular	Occasional hairline fractures
2	Banded clasts	42 ± 3% (46 ± 5%)	Dark	Very small, organized in bands	< 8 mm (< 0.2 mm)	Angular and vesicle-poor	Fractured margins
3	Crystalline clasts		Beige-orange or very dark	Absent	< 2 mm (< 0.2 mm)	Sub-rounded	-
4	Pale vesicular clasts	22 ± 4% (3 ± 1%)	Pale	Rounded	0.4–4 mm (0.01–0.2 mm)	Sub-rounded to angular, vesicle-poor	-
5	Dark vesicular clasts		Dark	Rounded to irregular	0.4–4 mm (0.01–0.2 mm)	Sub-rounded to angular, vesicle-poor	-
6	Borderless vesicular clasts		Dark	Rounded to irregular	< 4 mm (0.01–0.2 mm)	Sub-rounded to angular, vesicle-rich	-

The size and proportion of the main components are also provided for the subsidiary vein (in brackets). Colors indicated are those observed in petrographic thin section using plane-polarized light.

3D reconstruction video of the vein is provided (Supplementary Video 1).

### Uncertainties and Errors in Textural and Bulk Analyses

Uncertainties related to quantitative textural analysis include user subjectivity, together with sample variability and possible non-representativeness (see discussion in Gurioli et al., 2015). Grayscale thresholding around vesicles and clasts includes intrinsic errors related to threshold position and manual drawing, which are greatest (~10%) for small vesicles and when lower-resolution images are used (e.g., 5 μm rather than 0.13 μm pixel size), and exacerbated by small-scale spatial heterogeneities (Baker et al., 2011). When converting 2-D data to a 3-D projection, however, the error depends on the stereological model used and is thus harder to estimate (i.e., particle shapes have to be assumed, Cashman, 1988). Absolute sampling errors related to non-representativeness were calculated for FOAMS and ImageJ analyses (Table 2) using population-sample size relations (Van der Plas and Tobi, 1965). Error on equivalent diameter of clasts and vesicles is 5%, associated with

“misrepresentation of 1 pixel in area” within a particle (Sahagian and Proussevitch, 1998; Shea et al., 2010). Analytical precisions and accuracies are 0.01–0.03 and 0.03 g cm<sup>-3</sup> respectively for density determination (Barker et al., 2012), and 0.02 and 0.03 % respectively for connectivity measurements (Formenti and Druitt, 2003).

## RESULTS

### Clast Classification

Componentry analyses permitted classification of six tuffsite clast categories (Table 2, Figure 3): (A) dense, (B) banded, (C) crystalline, (D) pale vesicular, (E) dark vesicular, and (F) borderless vesicular. For simplicity, these six categories were combined into three main types: dense clasts (A), vesicular clasts (D, E, and F), and banded clasts (B and C). Crystalline clasts are grouped with banded clasts because these could not be differentiated in low-resolution overview images; however, the crystalline clasts comprise <0.5 vol. % of the total vein, and thus make a negligible contribution to population statistics.

## Bulk Rock Textural Parameters

The density of the host obsidian slightly exceeds that of the bulk tuffisite vein (2281 vs. 2070 kgm<sup>-3</sup>), because of its lower porosity (0.04 vs. 0.15), although the pore-free tuffisite vein density is marginally higher (Table 1). In the analyzed samples, host obsidian void space is far more connected (>90%) than in the tuffisite (66%), being dominated by a high degree of microfractures with negligible isolated pores. In contrast, a significant part of the porosity in tuffisite veins is isolated, as vesicular clast interiors are enclosed by continuous dense glass borders. These results reflect the final, quenched state of the bomb, whereas porosity values will have been distinctly different earlier in the tuffisite vein evolution (e.g., prior to particle deposition, and vesiculation and sintering of clasts). The measured characteristics are therefore endmember data.

## Clast Shape and Componentry

Clast type proportions were calculated from the weighted average of ImageJ results (Table 2). Componentry in the vein matrix (particles <15 μm) and larger particles appear largely similar, but the former cannot be reliably quantified with the available image resolution. However, the proportions of the three types of clasts were determined from higher-resolution images of the subsidiary vein, and these results, which may provide a reasonable proxy of the main vein matrix (Table 2), reveal a strong decrease in vesicular clast abundance with decreasing clast size. Indeed, vesicular clasts in the main vein are usually >60 μm in diameter, with a very small proportion approaching the lower size limit (20 μm). Within the *subsidiary vein* the grain size distribution of the different types of components is unimodal, with the dense clasts being the smallest in size. In contrast, within the *main vein* the distributions are polymodal, because of the additional presence of medium and large clasts, of which the banded clasts are the largest (Figure 4; see Table 2 for clast size nomenclature). Banded clasts are the most abundant (42.2 ± 3%), followed by dense clasts (35.4 ± 2.5%) and vesicular clasts (22.4 ± 4%) within the *main vein*, whereas the dense clasts are dominant within the *subsidiary vein* and probably in the matrix as well.

Vesicular clasts (0.02–5 mm, Figure 4) are the least elongate (≤0.56 elongation; 0.75 for dense and 0.74 for banded clasts), and most regular (≥0.72 regularity, vs. 0.37 and 0.45 for dense and banded clasts, respectively). They are therefore the most circular (≤0.95 eccentricity; 0.99 for other clasts) and morphologically homogeneous clast population. Banded clasts have the greatest size range (0.005–8 mm) and most diverse shapes, whereas dense clasts (0.004–5 mm) have the most irregular borders.

## Vesicle and Crystal Textures

Vesicular clasts were divided into pale, dark and borderless categories (Figure 3, Table 2), based both on their color in plane-polarized light and whether they display a bubble-free border. Of the 45 vesicular clasts analyzed, 16 were rejected from quantitative analysis because of their insufficiently large populations of resolvable vesicles (<50 vesicles), leaving 29 clasts (21 pale, 5 dark, and 3 borderless). The host obsidian is texturally similar to dense clasts, containing <2% euhedral millimetric plagioclase crystals often surrounded by hairline

fractures. Crystals within the dense clasts are predominantly euhedral, ≤5 mm plagioclase, with minor clinopyroxenes and biotites, consistent with previous mineralogical vein descriptions for juvenile 2008–2009 magma (Castro and Dingwell, 2009).

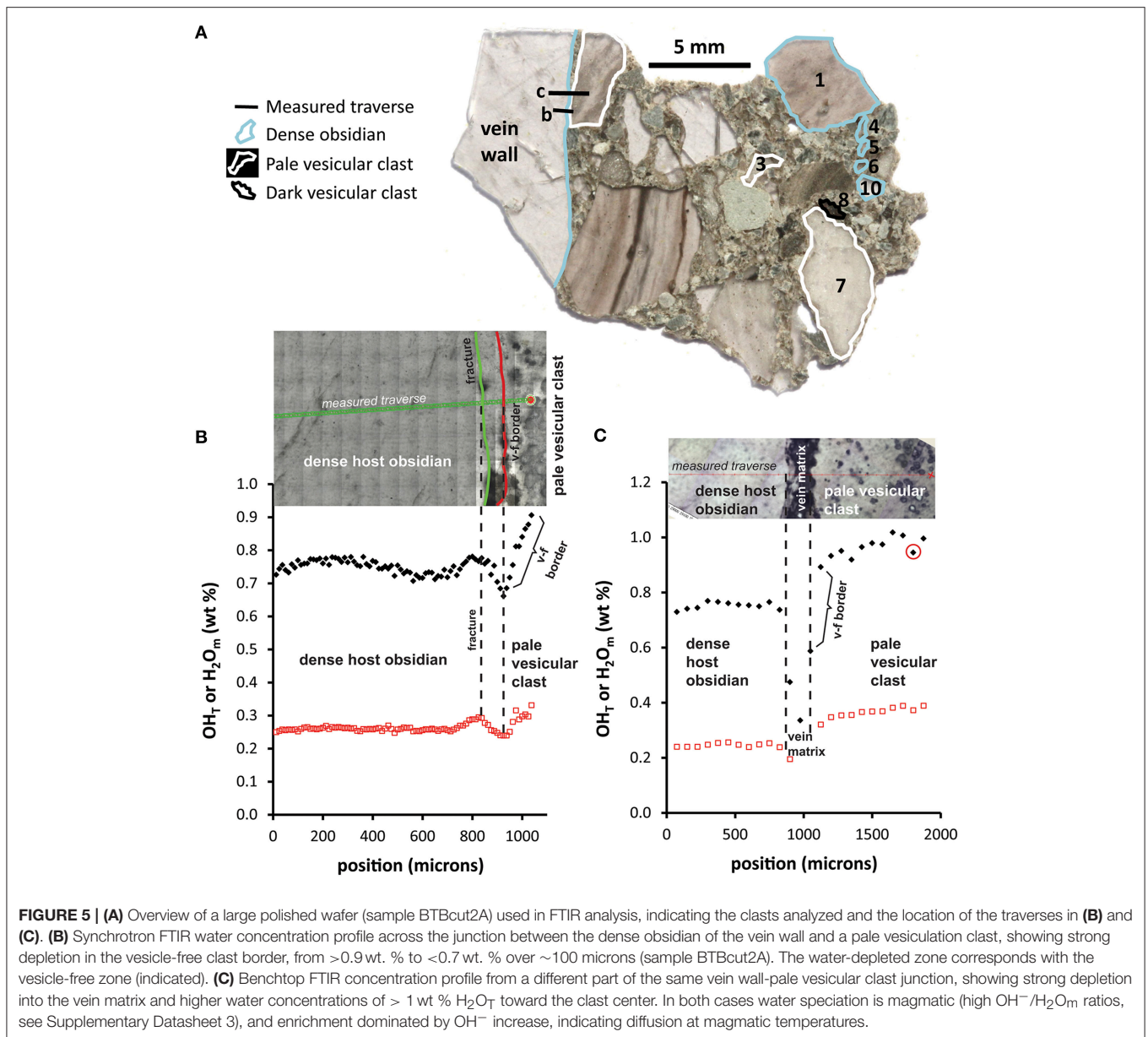
Pale vesicular clasts have a similar gray color in plane polarized light to the dense obsidian clasts. The margins of vesicular clasts (60 μm–4 mm diameter) are generally vesicle-free, except for borderless clasts and two pale vesicular clasts whose margins bear smaller, more elongated vesicles than their centers. Mean vesicle-free margin widths are 50–75 μm (calculated from 10 to 40 thickness measurements within each of 11 clasts), but reach 110 μm for the largest clasts (~4 mm), whose margin thickness is more spatially variable. In some vesicular clasts, dense borders are only locally well defined (e.g., Figures 5B,C).

## Major Element Compositions

Electron probe microanalyses show that dense and vesicular glasses contain 74.6 ± 0.5 wt. % SiO<sub>2</sub>, with negligible difference between border and core compositions. Banded clasts are heterogeneous, with some glassy bands similar in composition to other clasts (72.3 wt. % ± 5.9 wt. %), whereas other bands are strongly SiO<sub>2</sub>-enriched (92.8 wt. % ± 3.9 wt. %). These SiO<sub>2</sub>-enriched bands correspond with fractured areas associated with greater crystallinity and porosity, as encountered in the 2011–2012 Cordón Caulle rhyolite lava (Schipper et al., 2015). Other major elements are also heterogeneous in banded clasts, e.g., 6.2–23.2 wt. % Al<sub>2</sub>O<sub>3</sub> and 1.2–8.1 wt. % K<sub>2</sub>O, even when SiO<sub>2</sub>-rich bands are excluded. By contrast, dense clasts are far more compositionally homogenous (e.g., 13.7 ± 0.2 wt. % Al<sub>2</sub>O<sub>3</sub> and 3.2 ± 0.3 wt. % K<sub>2</sub>O). Measured major element concentrations in vesicular clasts are similar to dense clasts, with rather more dispersed values.

## Glass H<sub>2</sub>O Concentrations

Glass H<sub>2</sub>O concentrations reveal significant heterogeneities within and between clasts (Figure 5, Table 3). The interiors of pale vesicular clasts are significantly more water-rich (0.95 ± 0.05 wt % H<sub>2</sub>O) than dense obsidian clasts or the vein walls (0.73 ± 0.02 wt %). Low H<sub>2</sub>O<sub>m</sub> concentrations for given H<sub>2</sub>O<sub>t</sub> values, consistent with high OH<sup>-</sup>/H<sub>2</sub>O ratios, indicate insignificant low-temperature resorption or hydration and total H<sub>2</sub>O variation appears controlled by variation in OH<sup>-</sup> concentration rather than H<sub>2</sub>O<sub>m</sub> (c.f. Denton et al., 2012; Giachetti and Gonnermann, 2013; McIntosh et al., 2014; Giachetti et al., 2015; see Supplementary Datasheet 3 for speciation data). The high-resolution synchrotron FTIR traverse (Figure 5A) shows a ~100 μm-wide zone of significant H<sub>2</sub>O<sub>t</sub> depletion corresponding to the vesicle-free margin of the clast, consistent with previous data for clasts within Chaitén tuffisites and similarly strong depletion at the dense clast margin (Castro et al., 2012). Additional data for a further five dense clasts and the dense borders of three vesicular clasts (two pale, one dark) are provided in Table 3 and Supplementary Datasheet 3. Notably, the highest H<sub>2</sub>O values are only intercepted either by the smaller spot size of the synchrotron-source FTIR, or in the unusually wide vesicle-poor border of one vesicular clast (Figures 5B,C).



## Vesicle Size Distributions

Vesicle size distributions (VSDs) differ between pale, dark and borderless clasts (**Figures 6A–C** respectively). In pale clasts, three peaks occur at 25, 50, and 160  $\mu\text{m}$ , whilst two peaks centered at 42 and 67  $\mu\text{m}$  occur in dark clasts and one peak at  $\sim 22 \mu\text{m}$  in borderless clasts. Such contrasting vesicle size distributions reflect differences in the nucleation history of discrete clast types, as described in Sections Vesicular Clast Textural Characteristics—VSD and Three-Step Vesiculation Model. It is noticeable that the peak associated with the largest vesicles only occurs in the pale clasts.

When full VSDs are plotted as  $\ln(n)$  against equivalent diameter (**Figure 7**), steeper slopes reflect nucleation dominance, whereas shallower slopes indicate protracted growth (e.g.,

Shea et al., 2010). Twenty-nine VSD curves, corresponding to individual vesicular clasts, were created (**Figure 7A**); most can be separated into three segments of distinct slopes (**Figure 7B**). The slope of each segment was calculated and the average derived for each of the three segments, and plotted on **Figure 7C**. A simple three-step bubble nucleation-growth model is therefore proposed (**Figure 7C**), which fits the VSDs of 4 out of 5 dark vesicular clasts, 2 out of 3 borderless vesicular clasts and 9 out of 21 pale vesicular clasts. However, the divergence between the remaining 14 VSDs means that more complex vesiculation histories of individual clasts are both plausible and likely.

Bubble number density (BND) is a sensitive decompression rate meter in magmas, although heterogeneous nucleation can influence BND-decompression rate relationships (e.g.,



**TABLE 3 | Summary of clast H<sub>2</sub>O concentrations as measured via FTIR.**

Clast number	Textural type	Number of analyses	Mean H <sub>2</sub> O <sub>t</sub> (wt.%)	St. Dev. H <sub>2</sub> O <sub>t</sub> (wt.%)	Range of H <sub>2</sub> O <sub>t</sub> values (wt.%)	Mean H <sub>2</sub> O <sub>m</sub> (wt.%)	St. Dev. H <sub>2</sub> O <sub>m</sub> (wt.%)	Range of H <sub>2</sub> O <sub>m</sub> values (wt.%)
1	Dense obsidian clast	116	00.79	00.04	0.65–0.84	00.29	00.01	0.23–0.31
4	Dense obsidian clast	5	00.78	00.02	0.76–0.81	00.26	00.01	0.24–0.28
5	Dense obsidian clast	5	00.76	00.04	0.71–0.80	00.26	00.01	0.24–0.28
6	Dense obsidian clast	4	00.66	00.03	0.61–0.68	00.21	00.01	0.20–0.22
10	Dense obsidian clast	6	00.70	00.03	0.67–0.75	00.21	00.01	0.20–0.22
3	Pale vesicular clast margin	10	00.68	00.05	0.61–0.74	00.23	00.02	0.20–0.25
7	Pale vesicular clast margin	8	00.65	00.02	0.64–0.70	00.22	00.01	0.20–0.22
8	Dark vesicular clast margin	5	00.67	00.06	0.59–0.74	00.23	00.03	0.20–0.26
Traverse B	Dense host obsidian	72	00.75	00.02	0.70–0.78	00.26	00.01	0.25–0.29
Synchrotron source	Pale vesicular clast margin	10	00.79	00.08	0.66–0.91	00.28	00.03	0.24–0.33
Traverse C	Dense host obsidian	11	00.75	00.01	0.73–0.77	00.25	00.01	0.24–0.26
Globar source	Vein matrix	3	00.47	00.13	0.34–0.59	00.20	N/A	0.20–0.20
	Pale vesicular clast margin	11	00.96	00.04	0.89–1.02	00.36	00.02	0.32–0.39

Refer to **Figure 5A** for the position of individual analyzed clasts and traverses. Complete data, H<sub>2</sub>O speciation and error analysis is provided in Supplementary Datasheet 3.

Toramaru, 2006; Hamada et al., 2010; Houghton et al., 2010; Shea et al., 2011, 2012; Alfano et al., 2012). BND values are high and broadly similar for all clast types, indicating high decompression rates. Pale vesicular clast BND values of  $10^{12.7}$ – $10^{15}$  m<sup>-3</sup> occur over a range of vesicularities up to 70%. Dark vesicular clast BND values are similar, but over more restricted vesicularities (30–50%), whereas borderless clast BNDs range from  $10^{13.3}$  to  $10^{14.3}$  m<sup>-3</sup>, but vesicularity is strictly <10%.

## DISCUSSION

### Glass H<sub>2</sub>O Contents and the Origin of the Different Clast Types

Similar H<sub>2</sub>O contents at the margins of diverse clast types indicate some local diffusive re-equilibration occurred in these narrow marginal zones (**Table 3**). However, widely-varying, non-equilibrated H<sub>2</sub>O concentrations at clast centers (~0.66 to ~1 wt %) indicate sourcing from a broad depth range. The average H<sub>2</sub>O value for six analyzed dense clasts is 0.74 wt %, very similar to the average H<sub>2</sub>O content of vein wall obsidian of 0.75 wt %. Using H<sub>2</sub>O solubility-pressure relations for CO<sub>2</sub>-free rhyolitic melt at the eruptive temperature of 825 °C (Newman and Lowenstern, 2002; Castro and Dingwell, 2009), vein wall H<sub>2</sub>O equates to a pressure of 4.6 MPa, corresponding to 210 m depth in a conduit filled with dense vein wall material (2281 kgm<sup>-3</sup>). This depth estimate is a strict minimum as significantly greater mean porosity is likely in the overlying melt column. The equivalent values for 1 wt % H<sub>2</sub>O are 8.1 MPa and 360 m depth. Clasts with H<sub>2</sub>O-rich cores are therefore interpreted to be particles of fragmented, volatile-rich deeper magma that have risen and become trapped within the vein. By contrast, the dense obsidian clasts with similar H<sub>2</sub>O contents to the host obsidian are likely locally sourced from the vein walls. The clast SiO<sub>2</sub> contents broadly coincide with juvenile glass compositions from the 2008–2009 eruption (Amigo et al., 2013; Pallister et al., 2013;  $75.3 \pm 0.3$  wt. % SiO<sub>2</sub>). Silica-enriched portions of banded

clasts are interpreted as zones of local cristobalite precipitation within lava, potentially from an earlier eruptive episode. Banding in these lava clasts probably results from shear deformation (Schipper et al., 2015).

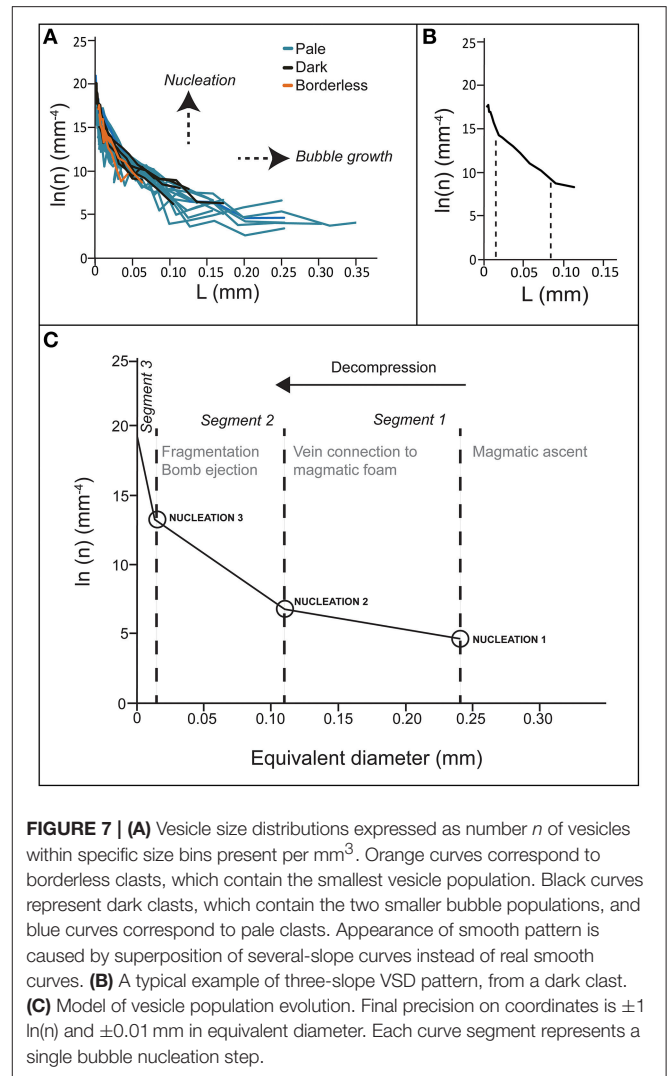
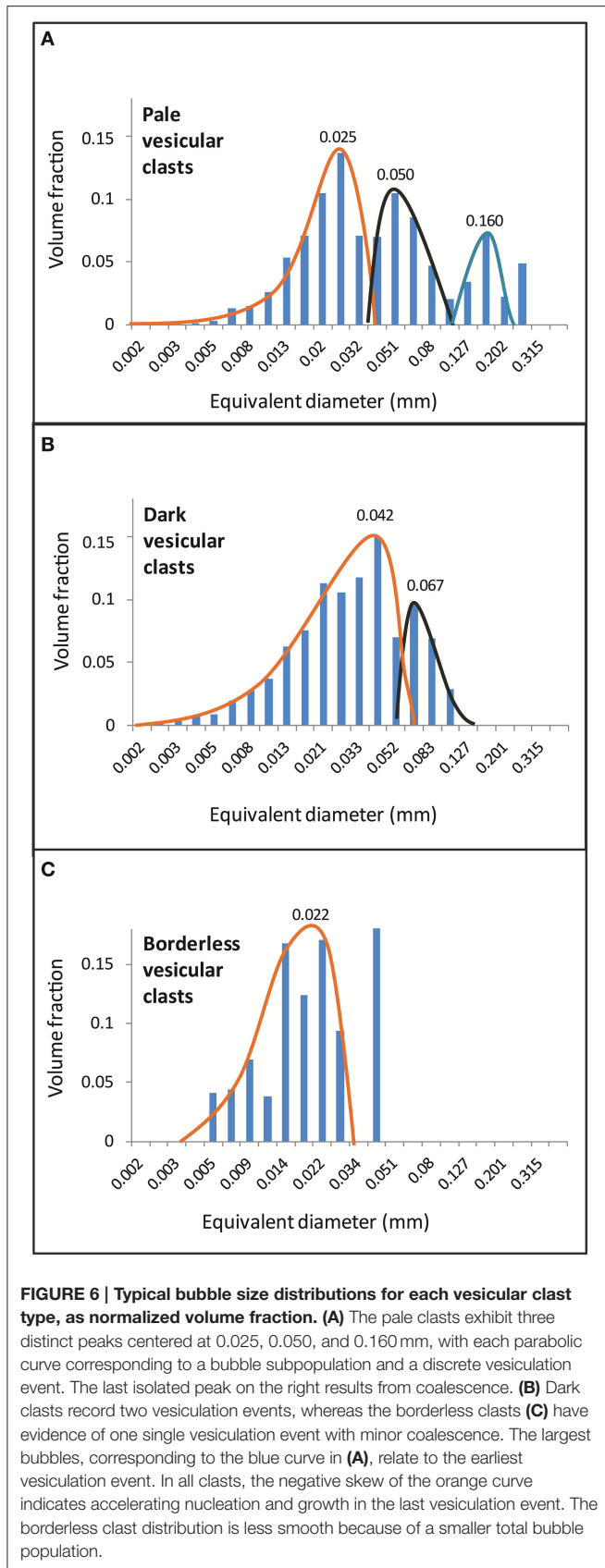
The inferred depth of magma excavation in the vein ( $\geq 360$  m) likely exceeded the dome height during phase 2, as it only attained 350 m on 30/09/2008, well after Vulcanian activity ceased (Global Volcanism Program, 2008; Bernstein et al., 2013). Therefore banded obsidian is likely non-juvenile material from previous rhyolitic activity within the caldera, which has become entrained into the vein (e.g., Rust et al., 2004). The color difference between pale/borderless and dark clasts may relate to differing nanocrystal populations and/or varying redox conditions (Castro et al., 2009).

### Bulk Rock Textural Parameters

The higher porosity-free density of the tuffsite vein (**Table 1**) compared with the host obsidian reflects the presence of more mafic non-juvenile clasts that are denser than the 2008 rhyolite, whereas higher porosity gives the vein lower overall density than the host. In the host obsidian porosity is dominated by discrete micro-fractures, e.g., offshoots from the main vein and subsidiary veins, which are almost completely connected, despite the very low overall porosity. These micro-fractures, interpreted as damage in the main vein walls as the tuffsite vein opened, are significantly interconnected on the sample scale investigated by the He pycnometer.

### Clast and Vesicle Shapes

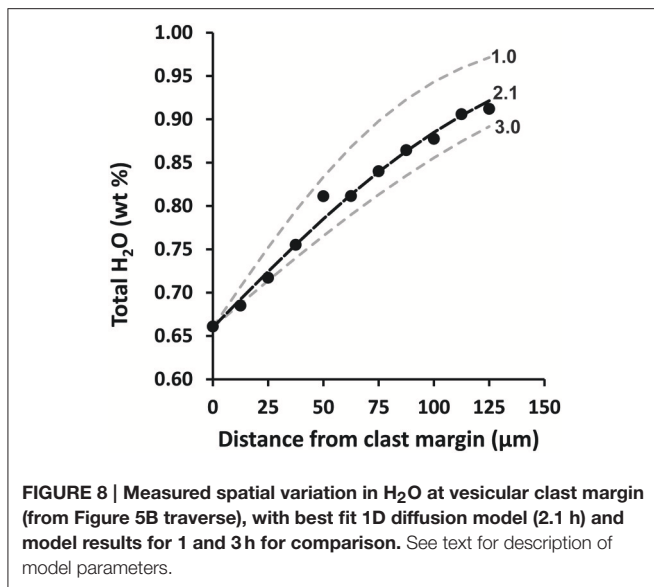
Shape factors such as regularity, elongation and eccentricity provide useful comparison of clast and vesicle forms (Shea et al., 2010; Moitra et al., 2013), and reveal that vesicular clasts are more rounded than dense clasts (for the main vein: average regularity 0.94 vs. 0.90, elongation 0.27 vs. 0.31, eccentricity 0.78 vs. 0.81). This reflects clast inflation by vesiculation, despite local embayments forming through impingement by denser neighboring clasts (**Figure 3D**). By contrast, banded and



dense clasts have more diverse, angular forms, with some clast margins bearing fractures associated with initial clast formation. Vesicle shapes are mostly near-circular, with coalescence creating some less regular forms (Figure 3D; Moitra et al., 2013). Coalescence has largely failed to create pervasive connected channels or open porosity because of enclosure by dense clast margins.

### Vesicular Clast Textural Characteristics—Dense Borders and H<sub>2</sub>O Diffusion

Significant textural differences exist between vesicular clast types, especially in their borders. Dense border thicknesses are mostly 50–75  $\mu\text{m}$  in near-spherical pale/dark clasts, reaching 110  $\mu\text{m}$  in the largest clasts ( $\sim 4 \text{ mm}$ ), and are locally smaller (20–40  $\mu\text{m}$ ) in more inhomogeneous clasts. Some highly texturally heterogeneous vesicular clasts only locally display well-developed dense borders (Figure 5). The minimum diameter of vesicular



vein clasts (60 µm) is approximately twice the minimum commonly seen dense border thickness, and a small proportion (~5% volumetric fraction) of vesicular clasts are <200 µm in diameter, corresponding to a radius of <100 µm, close to the maximum measured thickness of well-developed dense borders. We therefore infer that size-dependent clast vesiculation was influenced by the same diffusive degassing process that formed the dense clast borders, supported by this coincidence between H<sub>2</sub>O-depleted zones and vesicle-free textures (Figure 5B).

In order to estimate the timescale of diffusive water loss from vesicular clasts, we have modeled measured H<sub>2</sub>O concentration gradients using an error function solution to the one-dimensional diffusion equation

$$\frac{c_x - c_b}{c_0 - c_b} = 1 - \operatorname{erf} \left[ \frac{x}{2\sqrt{Dt}} \right],$$

where  $c_x$  is the H<sub>2</sub>O concentration at distance  $x$  from the surface (m),  $c_b$  is the bulk (far-field) H<sub>2</sub>O concentration,  $c_0$  is the (lower) surface H<sub>2</sub>O concentration,  $D$  is the H<sub>2</sub>O diffusivity ( $\text{m}^2\text{s}^{-1}$ ), and  $t$  is time (s). H<sub>2</sub>O<sub>m</sub> is assumed to be the sole diffusing species, and values of concentration- and temperature-dependent H<sub>2</sub>O diffusivity in rhyolitic melt are employed, as defined by Zhang (1999). A magmatic temperature of 825°C is assumed (Castro and Dingwell, 2009), with a fixed melt H<sub>2</sub>O concentration of 1 wt % at the clast center (the highest value measured in this clast; Figure 5C) and 0.66 wt % at the clast margin, as measured. Using these assumed values, the best-fit diffusion time is 2.1 h (Figure 8), but this is a maximum value, as the orientation of the measured traverse with respect to the clast border is unknown. If it were an oblique cut, then the true diffusion zone width may be significantly smaller. For example, if the measured concentration gradient occurred over a spatial scale of 30 µm, the best-time diffusion time would be 12 min.

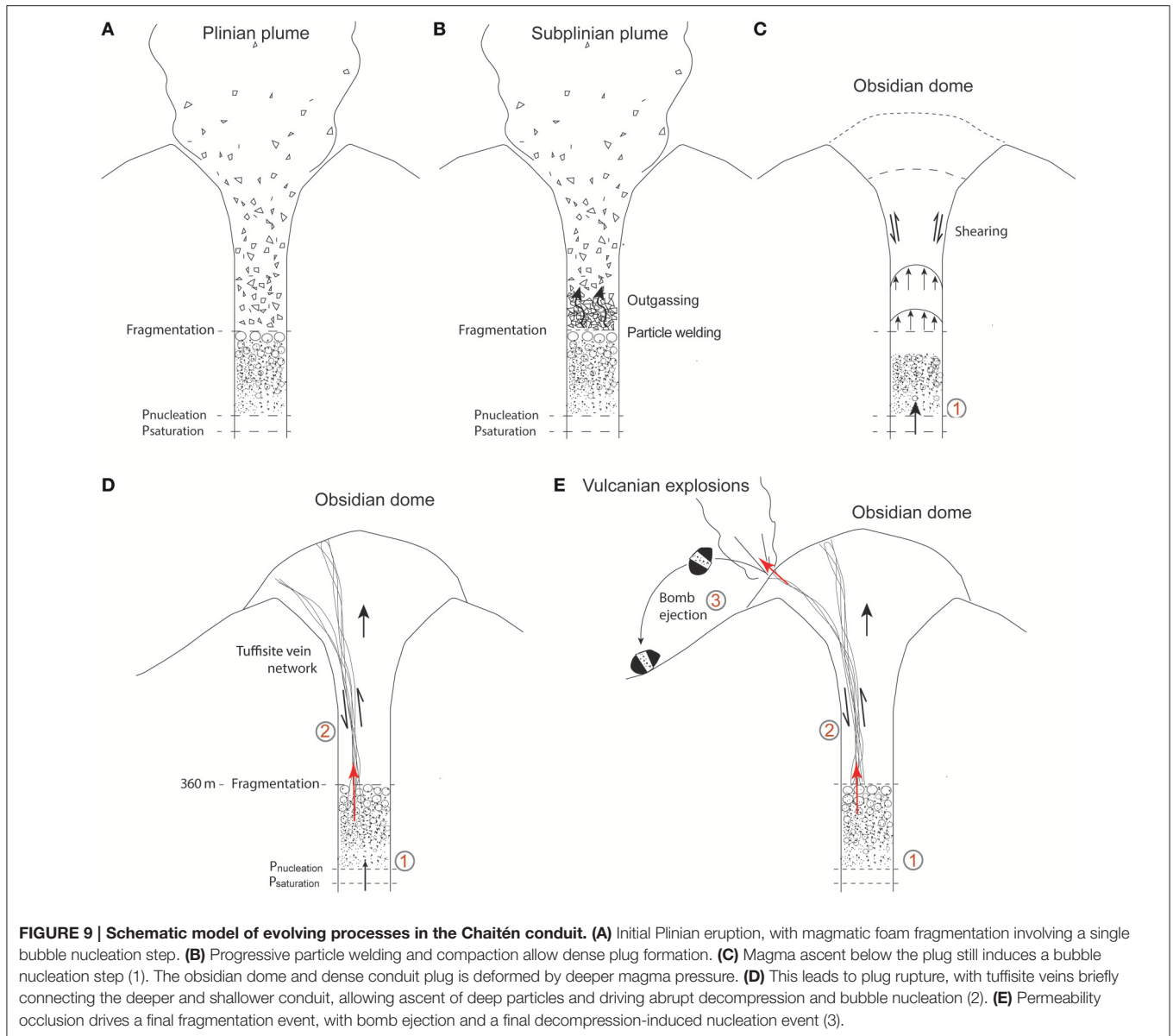
As dense border thickness broadly coincides with the zone of diffusive H<sub>2</sub>O depletion, measured widths of 20–110 µm therefore relate to diffusion timescales of ~10 min to >2 h. However, the zone of suppressed bubble nucleation may not precisely coincide with the width of the diffusively depleted zone, as demonstrated by the unusually heterogeneous vesicular clast analyzed in Figure 5, in which the bubble-poor border width differs greatly (~100 µm to >500 µm) between two measured FTIR traverses at its margin (Figures 5A,B), despite both traverses showing similar depleted zone widths of ~100 µm. This ~10 min to >2 h timescale indicates the total time elapsed from the initial deep fragmentation of the pale vesicular clast, through its residence in the vein, to final bomb quenching.

### Vesicular Clast Textural Characteristics—VSD

At a given magmatic temperature, bubble nucleation requires attainment of a critical supersaturation for stable bubble nuclei to persist (e.g., Navon et al., 1998). In a conduit, this can relate to magma decompression driven by ascent, plug slip or edifice failure, as well as an increase in magma temperature (e.g., Sparks, 1978; Alibidirov and Dingwell, 1996; Lensky et al., 2008; Lavallée et al., 2015). The coincidence between dense borders of vesicular clasts and H<sub>2</sub>O depletion reflects an H<sub>2</sub>O concentration-dependent nucleation response to a decompression event (Figure 5B; Castro et al., 2012). Only H<sub>2</sub>O-rich melt zones exceed the supersaturation threshold for bubble nucleation (e.g., Hoblitt and Harmon, 1993; Wright et al., 2007). Therefore, if all clast types experienced a common decompression history with three main decompression events (Figure 7), only the H<sub>2</sub>O-rich pale vesicular clasts were sufficiently decompression-sensitive to undergo bubble nucleation each time. In the most H<sub>2</sub>O-poor material (dense clasts and host obsidian) bubbles failed to nucleate, indicating that dense clasts have been “passively” plucked and abraded from vein walls.

Decompression events can trigger bubble nucleation or brittle fragmentation, depending on the decompression rate, magma porosity and volatile concentration (e.g., Spieler et al., 2004; Toramaru, 2006). Individual decompression events can therefore simultaneously induce varied magmatic responses—bubble nucleation, brittle failure, or no response—within different parts of the conduit, and the varied vesiculation histories of tuffsite components documented here reflect such divergent magma response over repeated events. Only one of twenty-nine analyzed clasts displays a smooth ln( $n$ ) curve thought to reflect progressively accelerating nucleation and bubble growth (Shea et al., 2010). Other curves comprise a succession of segments separated by breaks in slope. Most of these breaks in slope occur at similar equivalent diameter values, and three segments with similar slopes were identified across a large number of clasts. The average slope value increases from segment 1→2→3 (Figure 7C), reflecting increasingly strong bubble nucleation over successive decompression events, and consequently larger pressure drops that culminated in explosive bomb ejection.

Pale vesicular clasts have undergone significantly greater bubble growth than dark or borderless clasts, and higher H<sub>2</sub>O



in clast centers points toward derivation from greater depth in the conduit. The population of larger bubbles may reflect more prolonged bubble growth caused by more protracted ascent from deeper in the conduit, or faster bubble growth related to higher melt diffusivity and lower melt viscosity at higher  $H_2O$  concentrations (Zhang, 1999; Giordano et al., 2008; Browning et al., in review).

### Three-Step Vesiculation Model

Here we integrate textural and volatile data to propose a model of magmatic processes within the Chaitén conduit (Figure 9). The first bubble nucleation event (1) corresponded to magma decompression and foaming at  $>360$  m depth in the conduit (8.1 MPa vapor pressure). Fractures then opened that created a connection with the shallower conduit ( $\sim 210$  m

depth, vapor pressure  $\sim 4.6$  MPa). This opening may have driven fragmentation of deep magmatic foam as it experienced  $\sim 3.5$  MPa of abrupt decompression, which is sufficient to fragment highly inflated magma (Spieler et al., 2004). Fragments rising into the shallow vein system would have experienced simultaneous decompression-triggered bubble nucleation (2), viscous relaxation, and diffusive degassing from clast and fracture surfaces into the vein. Particles accumulated, blocking the vein, followed by vesiculation and sintering that colluded to compact the fragmented magma and create largely isolated porosity. This created conditions for gas overpressure that led to a final fragmentation event and bomb ejection (3), with accompanying sudden decompression of  $\sim 4.5$  MPa to atmospheric pressure. This last nucleation step was followed by localized bubble growth during bomb cooling. These three steps are

represented in **Figure 9**, with magmatic processes presented on **Figure 7C**.

The smallest vesicle subpopulation was therefore nucleated during bomb ejection, and the borderless vesicular clasts contain this single bubble population. We attribute intermediate bubble sizes to vesiculation driven by deep magma fragments ascending into the newly-opened vein, and the largest vesicle subpopulation to early, deep degassing prior to vein opening. Vesicle growth timescales following bomb ejection can be estimated from 1 atm vesiculation experiments, and Browning et al. (in review) found growth rates of  $0.6 \mu\text{m s}^{-1}$  for Chaitén melt with 0.97 wt %  $\text{H}_2\text{O}$ . This is appropriate for  $\text{H}_2\text{O}$ -rich pale clast centers and will be a slight overestimation for clasts with lower  $\text{H}_2\text{O}$  concentrations, or at elevated pressures where the pressure differential driving bubble growth is lower.

The timescale of post-bomb ejection vesicle growth can be roughly estimated from bubble sizes associated with the final nucleation event (20–40  $\mu\text{m}$ ) and these experimental bubble growth rates, yielding a minimum time of 0.55–1.1 min for bubble growth. This is significantly shorter than the inferred timescale for earlier diffusive degassing of pale vesicular clasts (~12 min). The dwell time of bordered clasts within the vein prior to final ejection was therefore longer than the final phase of vesicle growth within the bomb (for which realistic quenching times are minutes-tens of minutes, depending upon the exact position within the bomb). In contrast, borderless clasts could have experienced a shorter residence time within the vein, preventing significant diffusive degassing prior to the final vesiculation event.

## Magnitudes and Rates of Decompression

As BND values allow estimation of decompression rates in explosive eruptions (e.g., Toramaru, 2006), we compare our results with fast decompression experiments on pure rhyolitic melt (Hamada et al., 2010). The measured BND of vesicular clasts span  $10^{12.7}$ – $10^{15} \text{ m}^{-3}$ , which corresponds with the range of BND values in experimental samples decompressed at ~1–10 MPa/s at 800°C. This provides an estimate of the magma decompression rate, although it is an upper estimate, as any heterogeneous nucleation would generate a given BND value at lower decompression rates. Estimated decompression events are 3.5 MPa (upon vein opening from 8.1→4.6 MPa), and 4.5 MPa (on bomb ejection, from 4.6 MPa saturation pressure to 0.1 MPa). Therefore decompression times were <1–4 s, and thus, ignoring gas compressibility, particle velocities (from 360 m→210 m depth in the conduit, and 210 m depth to the surface) of the order of tens to  $>100 \text{ ms}^{-1}$ . Such values are plausible, being in broad accordance with observations of the longevity of Vulcanian explosions (seconds) and bomb ejection velocities (~100  $\text{ms}^{-1}$ ) during similar transitional rhyolitic activity at Cordon Caulle in January 2012 (Schipper et al., 2013). These results also match estimated gas rise velocities of  $50 \text{ ms}^{-1}$  within fracture networks in the shallow conduit at the dacitic Santiaguito lava dome, Guatemala (Johnson et al., 2014) and an estimate of the gas velocity in a different tuffisite vein from Chaitén of at least meters per second (Berlo et al., 2013).

## Comparison with the First Stage of Chaitén Eruption

The vesicle textures described here, from the transitional phase 2 of the Chaitén 2008 eruption, contrast with those from pyroclasts ejected in the initial sub-Plinian phase 1 (Alfano et al., 2012). The latter clasts have a higher mean vesicularity (69%), unimodal VSD profile with peak at 50–80  $\mu\text{m}$ , and BND of  $1.3 \times 10^{14} \text{ m}^{-3}$ . Phase 1 therefore involved a single, rapid homogeneous nucleation event within water-saturated magma during ascent from >5 km depth, whereas up to three vesiculation steps are recorded in phase 2 magma. BND values indicate broadly similar decompression rates (~1–10 MPa/s), which are lower in some phase 2 clasts. Clasts ejected during phase 1 were predominantly non-juvenile fragments (80%), alongside juvenile obsidian fragments (10%) and pumices (10%; Alfano et al., 2011, 2012). Similar particle types occur within phase 2 tuffisite veins, in which the non-juvenile proportion is still large (42%), but with dense obsidian and vesicular glassy fragments more dominant in this phase (35 and 23%). Dissolved  $\text{H}_2\text{O}$  concentrations decreased from 1.3 to 2.3 wt. % in phase 1 to ~0.6–1 wt % in phase 2, reflecting shoaling of the magma fragmentation depth during progressive magma stalling and dense plug formation. As the eruption progressed from phase 1 to phase 2, the rapidly-healing plug of dense magma acted as an increasingly effective trap for particles of deeper fragmented magma, despite briefly permitting escape of batches of exsolved gas. By contrast, in phase 1 the entire conduit effectively acted as a single, open tuffisite vein that provided an uninhibited path for delivery of magmatic fragments and gases to the surface. Trapping of particles allowed individual bodies of magma to experience multiple degassing events in phase 2, with the superimposed effect of facilitating efficient outgassing from the deeper conduit via periodically-blocked tuffisite vein networks. Interactions between veins and variably foamed magma domains (**Figure 7**) boosted the outgassing efficiency, which ultimately drove the complete transition from explosive (phase 1) to effusive (phase 3) activity.

## Broader Implications for Shallow Conduit Processes

The data presented in this paper provide new constraints on fluid and pyroclast transport through tuffisite veins, which controls the episodic, explosive gas-magma separation that characterizes transitional rhyolite eruptions (Schipper et al., 2013; Castro et al., 2014). We infer that pressure changes related to vein opening triggered either fragmentation or vesiculation of neighboring magma, depending upon the rate and amount of decompression, and the porosity-controlled fragmentation threshold (Spieler et al., 2004). Magma porosity within the shallow conduit was heterogeneous over short spatial scales, resulting from the integrated effects of strain localization and trapping of volatile-rich deeper magma fragments within tuffisite veins (Schipper et al., 2013; Castro et al., 2014; Cassidy et al., 2015). As a result, individual decompression events could simultaneously elicit fragmentation and vesiculation responses in juxtaposed, decimeter to meter-scale magma domains, with preferential “quarrying” of vesicular domains promoting release

of batches of exsolved gas, and boosting the efficiency of gas-magma separation (Castro et al., 2012, 2014). Our inference that individual magma batches can experience multiple vesiculation and fragmentation events, which act to progressively extract exsolved magmatic gas in multiple steps, contrasts with existing models of Vulcanian eruptions in which ejected plug material is sequentially replaced by a fresh batch of low-vesicularity magma (Clarke et al., 2002). At Chaitén magma was neither fully replaced nor homogenized, allowing the textural and degassing signature of earlier decompression events to dictate subsequent magmatic responses.

Models of gas mobility within shallow silicic conduits require estimates of permeability and pressure gradients (e.g., Clarke et al., 2002; Collinson and Neuberg, 2012), which may be time-variant. Our results provide constraints on the evolving gas pressure within one tuffsite vein. However, the permeability will have drastically changed between high values in the initially-open vein to lower values upon clogging by clastic particles, to still lower values as particle sintering destroyed pore space (Heap et al., 2015), aided by partial vesiculation of some clasts. The rate of sintering-driven tuffsite compaction is strongly dependent on the particle size (Vasseur et al., 2013), and as our results reveal significant spatial variations in particle size distribution within a tuffsite vein, defining an appropriate rate of sintering for the entire vein system is far from straightforward, and requires modeling of polydisperse particle populations (Wadsworth et al., 2014). Timescales of compaction and permeability destruction are additionally influenced by strain, which can be highly localized within ascending magma. The range and scale of textures observed in the sampled vein is also controlled by the size of the vein itself, and a linear upscaling of these features is unlikely to yield a comprehensive description of the entire shallow conduit. Adequately accounting for evolving shallow conduit permeability and pressure will therefore require observations across a broad population of glassy and pumiceous tuffsite bombs, ideally coupled to direct observation and sampling of these bombs (as erupted) in the field.

As the parameters that most strongly influence the magnitude and duration of gas flux evolve markedly during tuffsite vein formation, quantification of the permeability and volatile concentration heterogeneities within quenched glassy tuffsite veins may not be the most robust method for determining time-averaged fluxes. Direct measurement of exsolved gas fluxes and ejected particle velocities during silicic eruptions, combined with petrological characterization (Cassidy et al., 2015), and investigation of trace metal and short-lived radionuclides within tuffsites (Berlo et al., 2013) may be better approaches. The glassy tuffsite veins, however, provide rich textural and volatile records that belie a multi-scale evolutionary complexity that cannot be ignored in the description of silicic conduits that are in transitional explosive-effusive activity.

## CONCLUSIONS

Our study of a bomb-hosted rhyolitic tuffsite vein ejected from Chaitén in 2008 provides new insights into the enigmatic

mechanisms of gas-magma separation that control silicic eruption dynamics. Vein-filling clasts are magma fragments that ascended transient fracture networks cutting the shallow conduit during transitional explosive-effusive activity. Diverse clast textures indicate that multiple phases of magma fragmentation and vesiculation occurred beneath and within a dense lava plug. Varied clast water concentrations indicate that fracture networks were vertically extensive, with clasts derived from a wide range of depths within the conduit.

We envisage that three vesiculation steps occurred, progressing toward a final Vulcanian bomb-ejecting explosion. First, magma decompression within the deep conduit (~360 m), triggered an early bubble nucleation event. Vein opening then connected the shallow conduit (~210 m) to deeper magma (~360 m), triggering decompression of deep-derived melt fragments as they rose within the vein. These volatile-rich clasts diffusively lost H<sub>2</sub>O into the vein, creating relatively degassed, dense selvages tens of microns in thickness. Blockage of the vein by vesiculating, sintering clasts may have then drastically reduced its permeability, allowing an increase in gas pressure. Pressurization ultimately triggered bomb ejection from a depth of ~210 m in a Vulcanian explosion. This induced decompression of ~4.5 MPa and a final bubble nucleation event in the deep-derived magma clasts, but not the degassed clast selvages or shallower vein host magma. An H<sub>2</sub>O diffusion gradient at a vesicular clast margin indicates that up to ~2 h elapsed between the initial fragmentation and ascent of deeper magma within the vein, and the final quenching of the ejected bomb. Bubble number densities of 10<sup>12.7</sup>–10<sup>15</sup> m<sup>-3</sup> indicate decompression rates of ~1–10 MPa s<sup>-1</sup>, implying decompression timescales of ~seconds and so particle velocities of tens to 100 ms<sup>-1</sup> occur in the tuffsite vein system.

Together, our results provide a uniquely detailed reconstruction of degassing processes within a shallow silicic conduit. Repeated, abrupt pressure changes involved the opening of vertically pervasive tuffsite vein systems, whose permeable lifespan was then severely limited by the accumulation and rapid sintering of clasts. The contrast with the simpler vesiculation history of clasts from the initial purely explosive phase reflects the formation of a dense magma plug during transitional activity. This plug trapped ascending clasts, introducing considerable textural heterogeneity, and allowed multiple decompression events to effectively outgas magma. Accounting for such extreme spatial and temporal variations in upper conduit permeability will constitute a key challenge for the next generation of eruption models.

## AUTHOR CONTRIBUTIONS

The study was devised by ES, HT, and LG. Sample collection was carried out by HT, JC, and CS. ES and LG carried out imaging, textural analysis, and compositional analysis. HT led H<sub>2</sub>O analysis, assisted by JO, JC, KB, EM, and KW. All authors contributed to model development and interpretations. Manuscript preparation was led by ES and HT,

with contributions from all other authors, and figures drafted by ES and HT.

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## SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <http://journal.frontiersin.org/article/10.3389/feart.2016.00059>

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**Conflict of Interest Statement:** The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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