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Proton aurora and relativistic electron microbursts scattered by electromagnetic ion cyclotron waves

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Charged particle precipitation from Earth's magnetosphere results in stunning displays of the aurora and energy transfer into the atmosphere. Some of this precipitation is caused by wave-particle interactions. In this study, we present an example of a wave-particle interaction between Electromagnetic Ion Cyclotron waves, and magnetospheric protons and electrons. This interaction resulted in a co-located isolated proton aurora and relativistic electron microbursts. While isolated proton aurora is widely believed to be caused by Electromagnetic Ion Cyclotron waves, this unique observation suggests that these waves can also scatter relativistic electron microbursts. Theoretically, nonlinear interactions between Electromagnetic Ion Cyclotron waves and electrons are necessary to produce the intense sub-second microburst precipitation. Lastly, detailed analysis of the auroral emissions suggests that no chorus waves were present during the event. This is in contrast to the most commonly associated driver of microbursts, whistler mode chorus waves, and supports other less commonly considered driving mechanisms.

KEYWORDS

EMIC, SAMPEX, THEMIS, microburst, proton aurora, all-sky-imager

1 Introduction

Since the Van Allen radiation belts were discovered by Van Allen (1959) and Vernov and Chudakov (1960), the loss of radiation belt electrons due to wave-particle interactions has been modeled with quasilinear diffusion theory (e.g., Kennel and Petschek, 1966; Summers et al., 1998; Summers, 2005; Thorne et al., 2005). On long time scales, this model leads to accurate quantification of particle energization and loss (e.g., Lyons and Thorne, 1973; Claudepierre et al., 2020). On shorter time scales, however, quasilinear diffusion does not accurately predict precipitating fluxes on second or sub-second timescales (e.g., Bortnik et al., 2008; Albert and Bortnik, 2009; Saito et al., 2012; Miyoshi et al., 2015; Mozer et al., 2018; Shumko et al., 2018). The importance of nonlinear scattering in regard to particle loss is still unknown. Indeed, even the variety of mechanisms that nonlinearly scatter radiation belt electrons is still unknown (e.g., Grach and Demekhov, 2020; Bortnik et al., 2022). Scattering mechanisms that lead to electron microburst precipitation are of particular interest here; while there is a plethora of observational evidence linking whistler mode chorus waves to electron microbursts, only Douma et al. (2018) shows observational evidence of microbursts scattered by Electromagnetic Ion Cyclotron (EMIC) waves.

Microbursts are intense bursts of electron precipitation that typically last ≈ 100 ms (Shumko et al., 2021). Microburst energies span from tens of keV all the way up to > 1 MeV (Anderson and Milton, 1964; Parks, 1967; O'Brien et al., 2003; Blum et al., 2015; Douma et al., 2017; Shumko et al., 2020b; Kawamura et al., 2021; Shumko et al., 2021; Zhang et al., 2022). The microburst L-MLT distribution peaks in the outer radiation belt L-shells and in the 0–12 MLT region (Lorentzen et al., 2001; O'Brien et al., 2003).

Whistler mode chorus waves share many similarities with microbursts, including similar distribution in MLT and duration (of chorus rising tone elements) (Teng et al., 2017; Meredith et al., 2020; Shumko et al., 2021). These similarities led to a widely-accepted conclusion that microbursts are most often scattered by whistler mode chorus waves (e.g., Lorentzen et al., 2001). This is further supported by theory. Both Miyoshi et al. (2020) and Chen et al. (2020) show that rising tone whistler mode chorus waves can rapidly—and nonlinearly—scatter microburst electrons over the wide range of energies.

EMIC waves are another type of plasma wave that can pitch angle scatter energetic protons and electrons into Earth's atmosphere (e.g., Cornwall, 1965; Summers et al., 1998; Spasojević et al., 2004; Jordanova, 2007; Halford et al., 2016; Yahnin et al., 2021). EMIC waves are typically generated near the magnetic equator by anisotropic ions and are often bounded in frequency by the ion gyrofrequencies into three primary bands: hydrogen, helium, and oxygen (e.g., Gary et al., 1995; Blum et al., 2012; Saikin et al., 2015). In the magnetosphere, EMIC waves are spatially confined in L-Shell and extended in MLT (Mann et al.,

2014; Blum et al., 2017). When some of the EMIC wave power enters the ionosphere, it can duct and be observed over a large geographical area (Woodroffe and Lysak, 2012; Mann et al., 2014; Kim et al., 2018).

During their generation and subsequent propagation away from the magnetic equator, EMIC waves can precipitate 10 s keV protons (e.g., Miyoshi et al., 2008; Shoji and Omura, 2011). When the protons impact the atmosphere, they decelerate via charge-exchange: the bare proton strips an electron from an atmospheric molecule and becomes a neutral hydrogen atom in an excited state (Kivelson et al., 1995). This hydrogen atom, newly decoupled from the magnetic field, then emits hydrogen-specific auroral light, primarily in the Lyman- α (121.57 nm), Balmer- α (653.3 nm), and Balmer- β (486.1 nm) lines (e.g., Gallardo-Lacourt et al., 2021). Davidson (1965) and Fang et al. (2004) show that this charge exchange process will spatially smear a fine beam of precipitating protons into a ≈ 100 km radius proton aurora patch.

Similarly, EMIC waves can also scatter relativistic electrons (e.g., Summers et al., 1998; Khazanov et al., 2014; Kubota et al., 2015; Remya et al., 2015; Zhang X.-J. et al., 2016; Kubota and Omura, 2017; Zhu et al., 2020; Grach et al., 2021). The theoretical work by Omura and Zhao (2013) is most relevant here: they demonstrated that EMIC waves can scatter relativistic microbursts into the atmosphere. There, these precipitating electrons experience bremsstrahlung deceleration and emit X-ray photons which are absorbed before they reach the ground (Winckler et al., 1958; Woodger et al., 2015). Therefore, relativistic electron precipitation can only be observed directly in space and indirectly via the secondary X-rays in the upper atmosphere.

Some satellites in low Earth orbit (LEO) are well-equipped to observe this dual electron-proton precipitation that is often attributed to EMIC wave scattering (e.g., Miyoshi et al., 2008; Nishimura et al., 2014; Zhang J. et al., 2016; Qin et al., 2018; Capannolo et al., 2021). These studies often show EMIC precipitation structure that lasts a few to 10s of seconds. From the vantage point of a high-inclination satellite in LEO, these durations can correspond to either the spatial size or temporal duration (Shumko et al., 2020a,b). Nevertheless, EMIC-driven electron precipitation on faster time scales is seldom observed. In fact, Douma et al. (2018) is the only study to our knowledge which shows sub-second electron precipitation associated with an EMIC wave. These authors presented the results of a magnetic conjunction between the Solar Anomalous and Magnetospheric Particle Explorer (SAMPEX) satellite in the northern hemisphere and a ground-based magnetometer in Halley, Antarctica. During the conjunction, SAMPEX observed three > 1 MeV electron microbursts.

Due to the aforementioned EMIC wave ducting and a lack of auroral observations, it is difficult to co-locate the EMIC wave and microbursts in the Douma et al. (2018) conjunction. However, observing an EMIC wave on the ground, together

with co-located proton and electron precipitation, will provide more convincing evidence of EMIC-driven electron microbursts. We present such a study here.

Early on 20 January 2007, the Canadian Array for Real-time Investigations of Magnetic Activity (CARISMA) magnetometers spread throughout Canada observed an EMIC wave lasting almost 2 hours. This wave generated an isolated proton aurora (IPA) patch that was observed by a Time History of Events and Macroscale Interactions during Substorms (THEMIS) all-sky imager (ASI) and a meridional scanning photometer. Meanwhile, the SAMPEX satellite passed directly through the proton aurora and observed very rapid and intense > 1 MeV electron microbursts.

2 Materials and methods

2.1 Instruments

2.1.1 Canadian array for real-time investigations of magnetic activity

We used the CARISMA (Mann et al., 2008) magnetometers to identify EMIC waves. The instruments consist of fluxgate and induction coil magnetometers deployed throughout Canada. Here we use the fluxgate magnetometer from Gillam (CARISMA-GILL) that is close to the ASI that observed the proton aurora. The magnetometer data is collected in two modes: 1- and 8-Hz sample rates. We use the 1-Hz data here.

2.1.2 Solar anomalous and magnetospheric particle explorer

We used the SAMPEX satellite to identify relativistic microbursts. It was launched in July 1992 into a 520, –, 760 km altitude, 82° inclination low Earth orbit (Baker et al., 1993; 2012). The Heavy Ion Large Telescope (HILT; Klecker et al., 1993) observed > 1 MeV electrons. HILT was a large rectangular chamber with the aperture on one end, and 16 solid state detectors on the other. During this event HILT pointed to zenith and the electron counts were accumulated from all of the solid state detectors at a 20 ms cadence.

2.1.3 Time history of events and macroscale interactions during substorms all-sky imager

We use the THEMIS ASIs to study the auroral light. The THEMIS ASIs are an array of charged-coupled device (CCD) auroral cameras spread across Canada and Alaska (Harris et al., 2009; Mende et al., 2009). Each white light imager uses a fisheye lens to expand its field of view to 170°, corresponding to 9° latitudinal and 1 h magnetic local time (MLT) coverage (Donovan et al., 2006; Mende et al., 2009).

The 256×256 pixel images are taken at a 3-s cadence: a 1-s exposure is followed by 2-s processing. To analyze the white light images, we used the THEMIS ASI skymap calibration files that

are provided by the University of Calgary. The calibration data contain, among other things, arrays that map each pixel to (latitude, longitude) coordinates at an assumed auroral emission altitude.

In this study, we use the THEMIS ASI camera stationed in The Pas (THEMIS-TPAS) to study the proton aurora. Furthermore, we briefly use the THEMIS ASI at Gillam (THEMIS-GILL) to identify the proton aurora, but the aurora was overwhelmed by light pollution in that part of THEMIS-GILL's sky.

2.1.4 Meridian scanning photometer

And lastly, we used the Northern Solar Terrestrial Array (NORSTAR) meridian scanning photometers (MSPs) to understand the multispectral properties of the aurora. MSPs are designed to measure the latitudinal location and brightness of aurora at the meridian. We use four channels from the MSP at Gillam (MSP-GILL): 470.9 nm blue-line, 486.0 nm H_{β} (i.e., the Balmer- β emission line), 557.7 nm green-line, and 630 nm red-line (Jackel, 2005). See Unick et al. (2017) for a comprehensive description of the operation and calibration of MSPs.

2.2 Methods

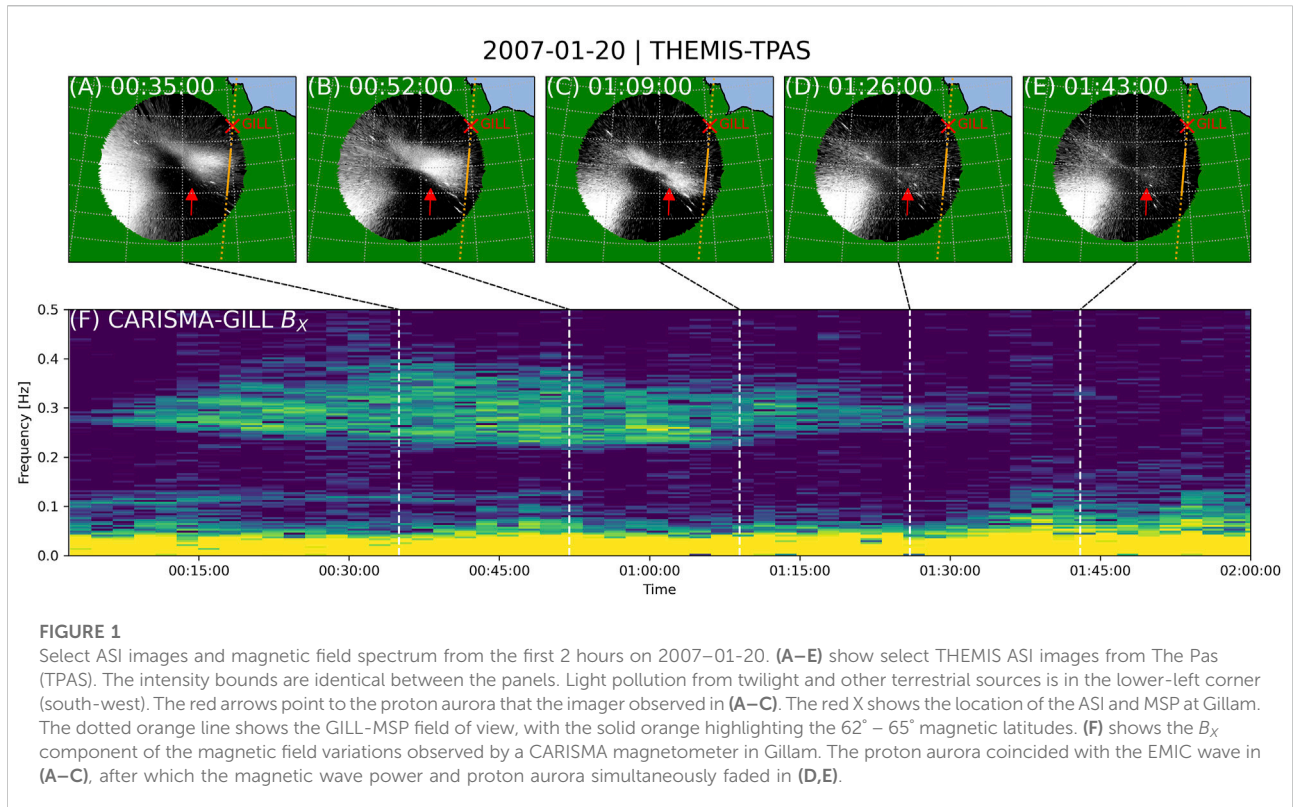
Our analysis consists of two main steps: calculate the CARISMA wave spectrum, and map the THEMIS-TPAS images and SAMPEX footprints to 110 km altitude.

To calculate the CARISMA frequency-time spectrum, we used the Windowed Fast Fourier Transform algorithm implemented by Scipy's spectrogram ()function (Virtanen et al., 2020). We used the Tukey window (also known as the tapered cosine) that had a 256-s length. The windows overlapped by 128 s (50% of the window length).

Then, to accurately compare the two datasets, we mapped the THEMIS-TPAS ASI images and the SAMPEX location to the same altitude. For the images, we mapped the THEMIS-TPAS pixels along their line of sight to their (latitude, longitude) at a 110 km altitude. We used the aforementioned skymap files, together with Python's pymap3d (Hirsch, 2016), and aurora-asi-lib (Shumko, 2022) libraries to do this. For SAMPEX, we mapped its location to its magnetic field footprint at 110 km altitude. For this we used the IRBEM-Lib magnetic field library (Boscher et al., 2012) with the IGRF magnetic field model (Thébault et al., 2015).

3 Results

From approximately 00:00–01:40 UT on 20 January 2007, nine CARISMA magnetometers deployed around central Canada observed EMIC wave power between 0.25–0.4 Hz, classified as Pc



1 (Jacobs et al., 1964). Figure 1F and Supplementary Video S1 show the EMIC wave observed at CARISMA-GILL. The primary band of the EMIC wave spanned 0.25–0.4 Hz. A second weaker band, centered around 0.1 Hz, ended earlier around 01:00 UT. Supplementary Figure S1 shows the spatial extent of the EMIC wave as observed on the ground: it shows a map of Canada with the CARISMA magnetometer locations that observed EMIC waves highlighted. The EMIC wave amplitude was largest at the Island Lake magnetometer (CARISMA-ISLL), 275 km south of CARISMA-GILL. The wave appeared to be almost linearly polarized with the semi-major axis pointing north, suggesting that the EMIC wave entered the ionosphere near ISLL (Woodroffe and Lysak, 2012).

The 0.25–0.4 Hz EMIC wave band coincided with an auroral patch that was observed concurrently by two THEMIS ASIs located at Gillam (GILL) and The Pas (TPAS). The ASIs observed the aurora from just after the cameras turned on at sunset until 01:40 UT. THEMIS-TPAS is 500 km south-west of THEMIS-GILL and the auroral patch was observed in between the two imagers. While THEMIS-TPAS observed the patch in the region of the sky away from sunset, THEMIS-GILL observed the aurora light superposed with a strong background light consisting of twilight and terrestrial light. Thus, our analysis focused on THEMIS-TPAS. At TPAS the sunset was at 23:03 UT the previous day and twilight lasted until 01:09 UT on January 20th. The TPAS imager turned on at 00:22:06 UT and was

initially saturated by twilight. The image in Figure 1A, taken at 00:35 UT, was the earliest time when the auroral patch was clearly visible. Panels (E) and (F) in Figure 1 show that both the auroral patch and EMIC waves vanished simultaneously around 01:45 UT, indicating that the auroral patch was driven by the EMIC wave.

Next, we investigate what particles created the auroral patch light. Figure 2 shows a multi-spectral keogram from MSP-GILL in panels (A)–(D) and a THEMIS-TPAS keogram along MSP-GILL's field of view in panel (E). We show MSP-GILL's field of view at 110 km altitude in Figures 1A–E with the dotted orange line. Figures 2A–D shows that the green and H_β MSP channels observed an intensity enhancement corresponding to magnetic latitudes $\lambda = 62^\circ - 65^\circ$ that we highlighted in Figures 1A–E with the solid orange line. The H_β MSP channel is sensitive to precipitating 10 s keV protons, while the green MSP channel is sensitive to secondary electrons that were generated by the protons (Sakaguchi et al., 2008). While the aurora intensity was dim in the H_β channel, we confirmed that it was not observed in the MSP's 480 and 495 nm background channels.

Since we're using THEMIS-TPAS images, we need to compare its light to MSP-GILL to identify the proton emission. Figure 2E shows a keogram constructed along the MSP-GILL field of view using data extracted from the THEMIS-TPAS ASI images. The field of view grazed the auroral patch that THEMIS-TPAS observed—enough that the two can be

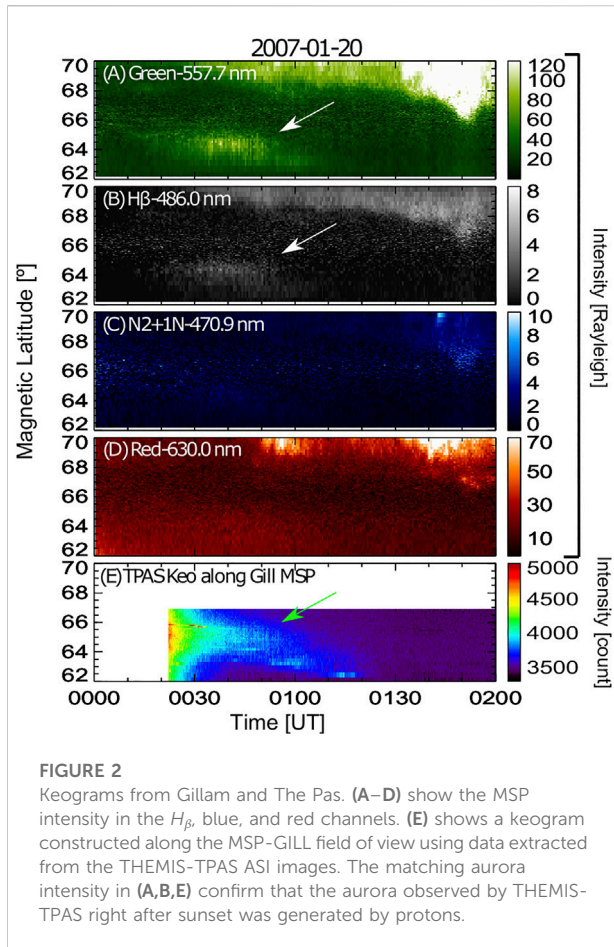


FIGURE 2
Keograms from Gillam and The Pas. (A–D) show the MSP intensity in the H_{β} , blue, and red channels. (E) shows a keogram constructed along the MSP–GILL field of view using data extracted from the THEMIS–TPAS ASI images. The matching aurora intensity in (A, B, E) confirm that the aurora observed by THEMIS–TPAS right after sunset was generated by protons.

compared. The auroral patch observed by THEMIS–TPAS was constrained to $\lambda = 62^{\circ}$ – 65° along the MSP’s field of view. Thus, the auroral patch seen in the ASI images was caused by the precipitating 10 s keV protons, confirmed with the MSP–GILL, and is an IPA.

At 00:40 UT SAMPEX orbited above the IPA and directly observed precipitating relativistic electrons. [Supplementary Video S1](#) and [Figure 3](#) show the conjunction that was at approximately $L = 5$ and $MLT = 17$ h. [Figures 3A–D](#) shows the mapped THEMIS–TPAS image with the SAMPEX orbit footprint superposed as a dotted red line and the instantaneous footprint as a red circle. [Figure 3E](#) shows that HILT observed a handful of >1 MeV electron microbursts. We used the fitting method described in [Shumko et al. \(2021\)](#) to estimate microburst durations. Some of the microbursts lasted ≈ 100 ms, while three had a ≈ 300 ms duration.

Two more instruments onboard SAMPEX, the Proton/Electron Telescope (PET; [Cook et al., 1993](#)) and the Low-Energy Ion Composition Analyzer (LICA; [Mason et al., 1993](#)), observed precipitation at this time that we show in [Supplementary Figure S2](#). PET is sensitive to >400 keV

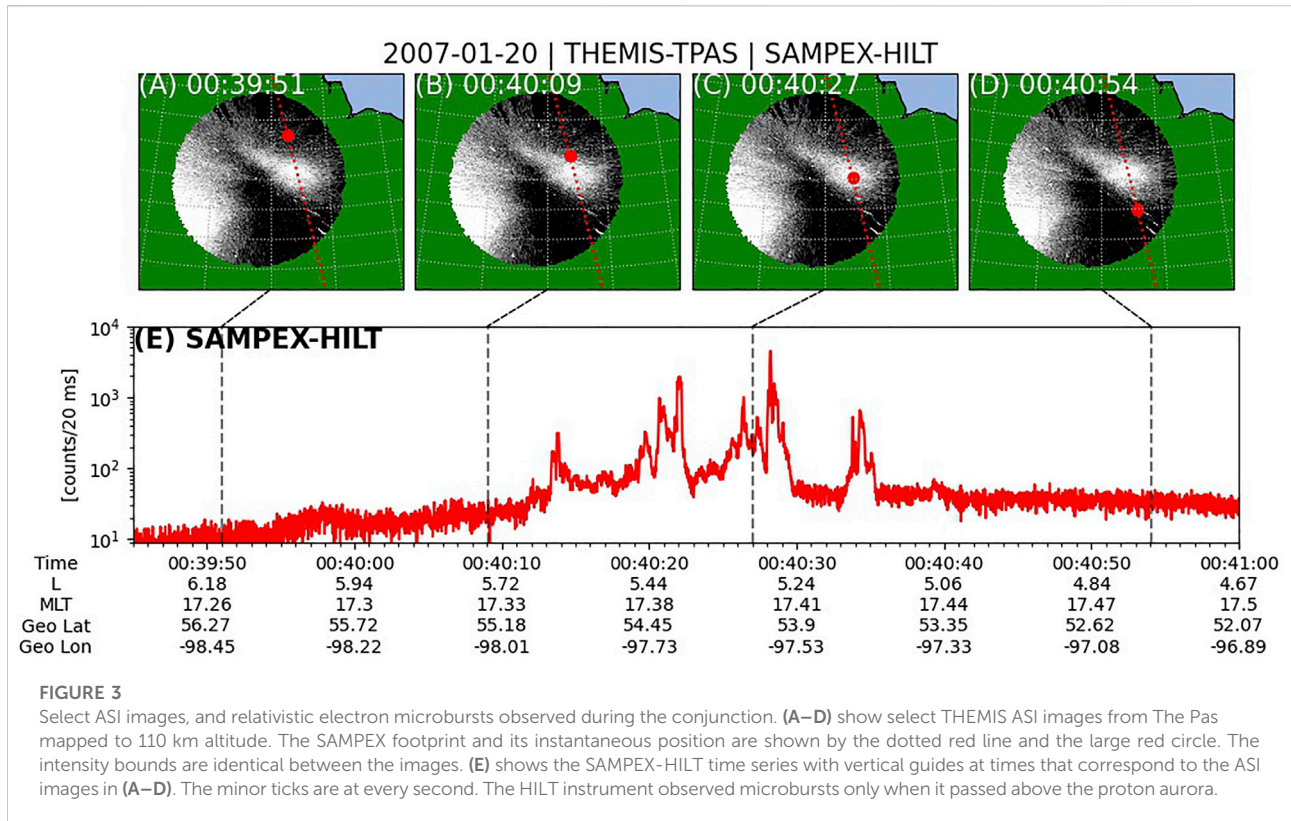
electrons and clearly observed many microbursts at a 100 ms cadence. Electron precipitation was also observed by LICA’s stop microchannel plate that was sensitive to >25 keV electrons. However, LICA’s time resolution was 1 s, so microbursts are difficult to identify. The clearest exception is the microburst observed at 00:40:30 UT shown in [Supplementary Figure S2](#). The three SAMPEX instruments’ integral response make it difficult to infer the energy of the precipitating electrons—we can only claim that some of the electrons had a >1 MeV energy. Furthermore, the SAMPEX instruments show signs of time-energy dispersion. However, we are unable to reliably verify dispersion because the time synchronization between instruments was of order 1 second—longer than the theoretical sub-second energy dispersion.

4 Discussion

This event shows strong evidence of EMIC-driven IPA concurrent with >1 MeV electron microbursts. The EMIC wave was observed throughout central Canada for almost 2 hours. The wave precipitated 10 s keV protons in a localized patch which then generated the IPA light—confirmed by the white-light observed by THEMIS–TPAS and the H_{β} light observed by MSP–GILL. The wave also precipitated intense and rapid >1 MeV electron microbursts that SAMPEX observed directly above the IPA.

While [Miyoshi et al. \(2008\)](#) and [Jordanova et al. \(2008\)](#), and others show that quasilinear diffusion can model the gradual EMIC-driven relativistic electron precipitation, it is unable to model the intense sub-second relativistic microbursts studied here. Nonlinear EMIC–electron scattering models are probably needed. [Albert and Bortnik \(2009\)](#) show that nonlinear interactions between electrons and monochromatic EMIC waves lead to rapid electron transport in pitch angle (and negligible transport in energy). The scattering efficiency of relativistic electrons is strongly dependent on the gyrophase difference between the wave and particle and on the inhomogeneity parameter. [Albert and Bortnik \(2009\)](#) describe the inhomogeneity parameter as quantifying the strength of the EMIC wave relative to the inhomogeneity of the magnetospheric plasma and magnetic field.

[Omura and Zhao \(2012\)](#) extended this work from a monochromatic wave to rising tones, generated by the EMIC triggered emission mechanism (e.g., [Grison et al., 2013](#); [Sakaguchi et al., 2013](#)). The authors found that the most effective EMIC–MeV electron scattering occurs with EMIC rising tones and repeated encounters with the wave as the electrons bounce: the EMIC wave is relatively stationary compared to a bouncing MeV electron. [Omura and Zhao \(2013\)](#) applied this theory to model EMIC-driven MeV microbursts. Here, the authors found that EMIC rising tones are very efficient at trapping relativistic electrons and



transporting them into the loss cone when the inhomogeneity parameter is less than one. The authors simulated an EMIC wave that propagated away from the magnetic equator to higher magnetic latitudes for 37 s. As the electrons bounced (with a sub-second bounce period), they repeatedly passed through and resonated with the wave for ≈ 20 ms. This led to efficient pitch angle transport of the resonant electrons into the loss cone. Figure 8b in Omura and Zhao (2013) shows an envelope of precipitating electrons with intense subsecond electron enhancements superposed; thus, the authors classified them as microbursts.

Thus, our observations are theoretically supported by Omura and Zhao (2013)'s results. Nonlinear wave-particle interactions between EMIC rising tones and electrons are necessary to generate these microbursts. Other studies including Nakamura et al. (2019), and Zhu et al. (2020) also presented evidence of nonlinear scattering of electron by EMIC rising tones that resulted in sub-minute flux variations observed by the Van Allen Probes (Mauk et al., 2012). That said, given that EMIC rising tones are necessary for nonlinear microburst scattering, our observations do not show clear EMIC rising tones. This is unsurprising, because ground-based observations of EMIC rising tones are rare—likely explained by the smoothing of the frequency dispersion during the propagation to the ground. In fact, Nomura et al. (2016) claims to be the first to have find EMIC rising tones on the ground.

We consider one final scenario: collocated EMIC and chorus waves. In this case, the EMIC waves were responsible for the IPA while the chorus waves were responsible for microbursts. While there were no direct equatorial wave measurements in this L-MLT region at this time, the MSP-GILL can infer the presence of 10 s keV electrons that are associated with lower-band chorus (Thorne et al., 2010). The green-to-blue MSP ratio is a good proxy for the precipitating electron energy spectrum: higher green-to-blue ratios correspond to lower energy electron precipitation (Rees and Luckey, 1974; Shepherd et al., 1996). The keograms in Figure 2 show an emission in the green but not the blue line. The large green-to-blue ratio indicates an absence of 10 s keV electron precipitation and thus suggest an absence of lower-band chorus waves that are capable of scattering MeV electrons.

The chorus-microburst model is admittedly very simple and widely accepted. But, our results suggest that there are multiple ways to scatter microbursts. Perhaps this is unsurprising, given that microbursts are broadly defined by their intensity and sub-second duration—so we would classify any wave-particle interaction capable of producing rapid precipitation as microbursts.

In conclusion, our observations show EMIC-driven precipitation of both 10 s keV protons that resulted in IPA light, and >1 MeV electron microbursts. Since relativistic electrons were not observed above the full extent of the

proton aurora, this means that the relativistic microburst electrons are episodic: a very specific triggering condition must have been met. Douma et al. (2018) showed an example of EMIC-driven microbursts, but the waves and MeV electrons were observed in opposite hemispheres, and no proton aurora was observed to localize the EMIC wave. Thus, our event is arguably the strongest evidence yet of nonlinear interactions between EMIC waves and electrons, leading to microburst precipitation. Our results are theoretically supported, despite the scant observational evidence in the published literature. Further, this study shows the indispensable utility of ground-based observations to gain a more comprehensive understanding of wave-particle interactions.

Data availability statement

Publicly available datasets were analyzed in this study. This data can be found here: CARISMA: <https://www.carisma.ca/carisma-data-repository>, THEMIS: https://data.phys.ucalgary.ca/sort_by_project/THEMIS/asi/MSP; https://data.phys.ucalgary.ca/sort_by_project/GO-Canada/GO-Storm/msp/IDLsav/SAMPEX: <https://izw1.caltech.edu/sampex/DataCenter/>.

Author contributions

MS, BG-L, KM, and AH first conceived the idea to look for SAMPEX-THEMIS ASI conjunctions. Both KH and YM first highlighted the uniqueness of this conjunction. AH wrote the draft introduction and MS wrote the rest of the manuscript. BG-L made Figure 2 and MS made the other plots. JL, ED, ES, and DG provided the THEMIS ASI and MSP data and expertise. JB and ML provided the SAMPEX data and analysis support. LB helped analyze the SAMPEX microburst observations. IM provided the CARISMA data. All authors contributed to manuscript revision, read, and approved the submitted version.

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Supplementary material

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fspas.2022.975123/full#supplementary-material>

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