

Analysis of Reconstructed Modified Symmetric Teleparallel f(Q) Gravity

N. Myrzakulov^{1,2}*, S. H. Shekh³, A. Mussatayeva^{1,4} and M. Koussour⁵

¹Department of General and Theoretical Physics, L. N. Gumilyov Eurasian National University, Nur-Sultan, Kazakhstan, ²Ratbay Myrzakulov Eurasian International Centre for Theoretical Physics, Nur-Sultan, Kazakhstan, ³Department of Mathematics, S. P. M. Science and Gilani Arts and Commerce College, Ghatanji, India, ⁴S. Seifullin Kazakh Agro-Technical University, Nur-Sultan, Kazakhstan, ⁵Quantum Physics and Magnetism Team, LPMC, Faculty of Science Ben M'sik, Casablanca Hassan II University, Casablanca, Morocco

The existing analysis reports a reconstruction scheme of the newly proposed gravity say f(Q) gravity through the scale factor of the form $a(t) = a_0 t^n = \frac{1}{1+2}$ by describing the powerlaw cosmology. The reconstructed f(Q) gravity models disclosed how this modified gravity model is capable to replicate dissimilar epochs of the cosmological history. Also, the reconstructed f(Q) gravity models are castoff to develop the expressions for density and pressure and the equation of state parameter. We reconstruct two cases of interacting fluid scenario ghost and pilgrim dark energy with pressureless dark matter. The physical behavior of the models is talked over the evolution of the Universe is accelerated. Moreover, the well-known cosmological planes i.e., $(\omega_D - \omega'_D)$ and (r - s) constructed for our models, also include a comparison of our findings of these dynamical parameters with observational constraints. It is also quite interesting to mention here that the results of the equation of state parameter, $(\omega_D - \omega_D^L)$ and (r - s)-planes coincide with the modern observational data.

Keywords: cosmological evolution, f (Q) gravity, isotropic homogeneous space-time, non-metricity scalar, ghost dark energy, pilgrim dark energy

1 INTRODUCTION

Currently the modern theoretical observational data like type-Ia supernovae (Riess, 1998; Perlmutter, 1999), Cosmic Microwave Background Radiation (CMBR)(Spergel, 2003; Komatsu, 2011), Baryon Acoustic Oscillations (BAO) and Weak Lensing (WL) (Jain and Taylor, 2003; Eisenstein, 2005) Large Scale Structure (LSS) (Tegmark, 2004; Seljak, 2005), have been accounted for an accelerating expansion of the Universe. Fundamentally there are two apparent access for this kind of accelerating expansion of the Universe. First is the modified theories of gravity can also explain such an accelerating expansion of the Universe. In the foregoing epoch, a lot of investigation has been primed in the modified theories of gravitation such as f(R), f(R, T), f(G), f(R, G), and f(T, B) gravity, where R, G, T and T denotes the Ricci scalar, Gauss-Bonnet invariant, trace of energy-momentum tensor and torsion scalar of the Universe respectively. Innumeral works have been established in the framework of this modified theories of gravity and interesting results have been found in (Capozziello et al., 2007; Nojiri and Odintsov 2007; Azadi et al., 2008; Capozziello et al., 2008; Nojiri and Odintsov 2008; Harko et al., 2011; Daouda et al., 2012; Wei et al., 2012; Sharif and Yousaf 2013; Sahoo et al., 2014; Abbas et al., 2015; Chirde and Shekh 2015; Chirde and Shekh, 2016a; Chirde and Shekh, 2016b; Sharif and Fatima 2016; Bhatti et al., 2017; Bhoyar et al., 2017; Chirde and Shekh, 2018a, Chirde and Shekh, 2018b; Chirde and Shekh, 2019; Pawar et al., 2018; Shekh and Chirde, 2019; Shekh and Chirde, 2020; Shekh et al., 2020b; Dagwal and

OPEN ACCESS

Edited by:

Kazuharu Bamba, Fukushima University, Japan

Reviewed by:

Abdul Jawad, COMSATS University Islamabad, Pakistan Vyacheslav Ivanovich Dokuchaev, Institute for Nuclear Research (RAS). Russia

*Correspondence:

N. Mvrzakulov nmyrzakulov@gmail.com

Specialty section:

This article was submitted to Cosmology, a section of the journal Frontiers in Astronomy and Space Sciences

> Received: 23 March 2022 Accepted: 21 April 2022 Published: 04 July 2022

Citation:

Myrzakulov N, Shekh SH, Mussatayeva A and Koussour M (2022) Analysis of Reconstructed Modified Symmetric Teleparallel f(Q) Gravity. Front. Astron. Space Sci. 9:902552. doi: 10.3389/fspas.2022.902552

1

Pawar 2020; Sahoo and Bhattacharjee 2020; Shekh et al., 2020a; Shekh and Chirde 2020; Shekh et al., 2021; Shekh, 2021a; Shekh 2021b). Among the all determinations, one model which is based on the so-called f(Q) gravity theory or symmetric teleparallel gravity, where the nonmetricity scalar Q is works as gravitational interaction. The theory of f(Q) gravity was first introduced by Jimenez et al. (2018) later on Lazkoz, (2019) investigated an interesting restrictions on f(Q) gravity theory by involving the polynomial functions of the redshift in Lagrangian and successfully derived using data from the expansion rate, type-Ia Supernovae, Quasars, Gamma-Ray Bursts, BAO data, and CMBR distance. An examination on f(Q) theory of gravity have been speedily progressed as well as astrophysical data observational constraints to provoke it in contrast to the formulation of standard Einsteinian General Relativity. Mandal et al. (2020a) analyzed the cosmography in f(Q) gravity while the same author in Mandal et al. (2020b) gave a full of energy conditions constraints like weak, strong, dominant and null energy conditions for two models of f(Q) gravity theory. Frusciante, (2021) focused on a specific model of f(Q) gravity theory which at the background level is indistinguishable from ACDM model, while demonstrating measurable and peculiar signatures at linier perturbation level.

Second is dark energy which can be represented by using the equation of state parameter (ω_D) and characterized as $\omega_D = \frac{p_D}{q}$. More about the ω_D parameter is that, if $\omega_D \approx -1$: it behaves like a standard cosmology, if $\omega_D > -1$: the dark energy quintessence model behavior or $\omega_D < -1$: dark energy phantom model behavior and the probability $\omega_D \ll -1$ is governed out by existing cosmological data as well as K-essence, Chaplygin gas and several supplementary limits of ω_D are acquired from observation data results which come from the combination of data SNe-Ia and CMBA as well as Statistics of Galaxy Clustering which are $-1.66 < \omega_D < -0.62$ and $-1.33 < \omega_D < -0.79$ respectively. – 1.44 < ω_D < – 0.92 from luminosity, CMB anisotropy, galaxy clustering statistics and distances of high red-shift SNe-Ia. Notwithstanding the fact that it is preferred by the observational data, the ACDM model tests cosmological constant problems. To overwhelm the theoretical observations mentioned above, various models of dark energy have been proposed in literature. A dynamical dark energy model one, known as ghost dark energy which has a remarkable some non-trivial properties for the expanding Universe having nontrivial topological formation which resolve U(1) problem (Witten, 1979; Kawarabayashi and Ohta, 1980; Nath and Arnowitt, 1981). Sheykhi and Movahed, (2012) observed expansion of the Universe using model parameter constraints in general relativity for an interacting ghost dark energy model. Sadeghi et al. (2013) computed dynamical parameters such as deceleration and equation of state parameters numerically to examine the behavior of the Universe in an interacting ghost dark energy models by varying Λ as well as G. Next is the pilgrim dark energy which devours a phantom-like Universe to prevent the formation of black hole (2012). Sharif and Jawad, (2013) observed the cosmic evolutionary actions of pilgrim dark energy with event horizons and apparent using non-flat Universe while considering the different IR-cutoffs as particle horizon, event horizon and

conformal age of the Universe. Sharif and Zubair, (2014) noted cosmological evolution of pilgrim dark energy whereas Jawad et al. (2016a) investigated cosmic behavior of pilgrim dark energy in loop quantum cosmology using Hubble horizon in forms IRcutoff for interacting scenario. By assuming the interacting scenario of unified pilgrim ghost dark energy and cold dark matter in the flat FRW Universe framework. Jawad et al. (2017) have constructed the equation of state parameter which exhibits the transition from region of quintessence and then approaches to region of phantom at z = -0.9 whereas Jawad et al. (2016b) illustrated the cosmic acceleration of the Universe under two interacting dark energy models say pilgrim dark energy with Granda-Oliveros cut-off and its generalized ghost version in the DGP braneworld framework and observed that the equation of state parameter behaves like the phantom era of the Universe while the deceleration parameter shows the accelerated expansion of the Universe. Keep in mind an interaction between pilgrim dark energy as a future event and apparent horizons with cold dark matter. Rani et al. (2016) studied the cosmic acceleration in dynamical modified Chern-Simons gravity in the framework of non-flat FRW Universe.

Nevertheless, for the certain geometrical models ascending from the modifications of Einstein's gravitational field equations, the equation of state (ω_D) is no longer playing a vital role and its influence becomes unclear. Consequently, a new diagnosis is required to discriminate all classes of cosmological models. In order to accomplish these classes Sahni et al. (2003) introduced the pair of parameters (r, s) which has no dimension or so-called statefinder parameter of the form:

$$r = \frac{\ddot{a}}{aH^3}, \quad s = \frac{r-1}{3(q-1/2)}$$
 (1)

As, the Universe involving with two component fluid matter Ω_m and exotic form of energy Ω_D . In this situation the statefinder parameters attains the form (Sahni et al., 2003)

$$r = 1 + \frac{9\omega_D}{2}\Omega_D (1 + \omega_D) - \frac{3\dot{\omega}_D}{3H}\Omega_D$$
(2)

$$s = 1 + \omega_D - \frac{\dot{\omega}_D}{3\omega_D H} \tag{3}$$

where ω_D is the equation of state parameter of dark energy.

In the past works so many authors have effectively validated the statefinder diagnostic that it can distinguish a series of cosmological models. In case of Λ CDM and CDM models, the statefinder parameters respectively are fixed as (r, s) = (1, 0) and (r, s) = (1, 1). Also, for quintessence field the trajectories of (r - s)plane lie in the range (s > 0, r < 1) whereas for chaplygin gas which look a lot like to (s < 0, r > 1). In addition to the geometrical diagnostic say statefinder diagnosis, there is another one dynamical diagnostic which was firstly intended by Caldwell and Linder, (2005) and verified the deeds of quintessence scalar field dark energy model through this plane called $(\omega_D - \omega'_D)$ -plane analysis and also used extensively in the literatures. In the $(\omega_D - \omega'_D)$ -plane, ω'_D signifies the advancement of ω_D . Over and done with this plane, the models can be considered as in two different classes as thawing and freezing. The thawing region is described as $\omega'_D > 0$, $\omega_D < 0$ while freezing region as $\omega'_D < 0$, $\omega_D < 0$ on $(\omega_D - \omega'_D)$ -plane.In this paper, we study the correspondence scheme with ghost dark energy and pilgrim dark energy model using reconstruction technique in f(Q) gravity. The paper is prearranged in the following format. Section 2, contains FRW Universe with the source of fluid as an interaction between matter and dark energy from which the equation of state parameter is derived in Section 3 while Section 4, 5 enclose the brief discussion of cosmographic observations through cosmic diagnostic parameters and phase planes by reconstructing f(Q) gravity model with the help of ghost and pilgrim dark energy respectively. Finally, we conclude our results in Section 6.

2 FRW UNIVERSE WITH INTERACTING SOURSE

Consider the spatially homogeneous and isotropic Friedman-Robertson-Walker (FRW) line element of the form:

$$ds^{2} = -dt^{2} + a^{2}(t) \left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta d\phi^{2} \right], \quad (4)$$

where *a* be the scale factor of the Universe. The angle θ and ϕ are the usual azimuthal and polar angles of spherical coordinates. Also, *k* is a constant which represent the curvature of the Universe. If *k* = 1, 0, -1, then this corresponds to closed, flat, open Universe. The energy momentum tensor for dark matter and dark energy are defined as

$$\hat{\mathcal{T}} = \mathcal{T}_{\mu\nu} + \bar{\mathcal{T}}_{\mu\nu} \tag{5}$$

where $\mathcal{T}_{\mu\nu}$ and $\mathcal{T}_{\mu\nu}$ are the energy momentum tensors for pressureless dark matter and dark energy, defined as $\mathcal{T}_{\mu\nu} = \rho_m u_\mu u_\nu$ and $\bar{\mathcal{T}}_{\mu\nu} = (\rho_d + p_d)u_\mu u_\nu + p_d g_{\mu\nu}$ in which ρ_m , ρ_d are the energy densities of dark matter and dark energy respectively and p_d is the pressure of dark energy, u_i is the four velocity of the fluid. Also, $u_i = \delta_4^i$ is a four-velocity vector which satisfies

$$g_{\mu\nu} = u^{\mu}u_{\nu} = -x^{\mu}x_{\nu} = -1$$
 and $u^{\mu}x_{\nu} = 0$ (6)

For a Universe where dark energy and dark matter are interacting to each other the total energy density satisfies the continuity equation as following:

$$\dot{\rho}_D + 3H(\rho_D + p_D) = 0 \tag{7}$$

where ρ_D be the combine energy density of two fluids of the form $\rho_D = \rho_m + \rho_d$. The curvature energy density ρ_k and the critical energy density ρ_{cr} as usual defined by (Yang, 2020):

$$\rho_k = \frac{3k}{8\pi Ga^2} \quad \text{and} \quad \rho_{cr} = 3H^2 \tag{8}$$

Now, the three fractional form of energy densities Ω_D , Ω_m and Ω_k as (Yang, 2020):

$$\Omega_D = \frac{\rho_D}{\rho_{cr}} = \frac{\rho_D}{3H^2},\tag{9}$$

$$\Omega_m = \frac{\rho_m}{\rho_{cr}} = \frac{\rho_m}{3H^2} \tag{10}$$

$$\Omega_k = \frac{\rho_k}{\rho_{cr}} = \frac{k}{H^2 a^2} \tag{11}$$

Then from the above **Eqs 9–11**, the Friedmann equation can then be written as:

$$1 + \Omega_k = \Omega_D + \Omega_m. \tag{12}$$

Consider the interaction between two fluids. So, the energy densities of two fluids do not conserve separately, the continuity of matter of two fluids yields (Yang, 2020)

$$\dot{\rho}_D + 3H(\rho_D + p_D) = -\Gamma \tag{13}$$

where Γ represents the interaction between dark matter and dark energy. In general Γ should be a function with units of inverse of time. For the convenience, choose the following form of interaction term:

$$\Gamma = 3\eta H \left(\rho_m + \rho_D\right) = 3\eta H \rho_D \left(1 + u\right) \tag{14}$$

where η be the coupling parameter. Considering $\eta = 0$, the equation of continuity reduces to the non-interacting case. Here *u* is defined as:

$$u = \frac{\rho_m}{\rho_D} = \frac{\Omega_m}{\Omega_D} = \frac{1 - \Omega_D}{\Omega_D}$$
(15)

Under the above defined parameters, the equation of state parameter for dark energy can be derived as (Yang, 2020):

$$\omega_D = -\frac{1}{2 - \Omega_D} \left(1 - \frac{\Omega_k}{3} + \frac{2\eta}{\Omega_D} \left(1 + \Omega_k \right) \right)$$
(16)

For flat Universe after taking k = 0, the equation of state parameter for dark energy from Eq. (16) can be rewritten as (Yang, 2020):

$$\omega_D = -\frac{1}{2 - \Omega_D} \left(1 + \frac{2\eta}{\Omega_D} \right) \tag{17}$$

There exist several dynamical dark energy models, in literature, presented by various authors both in general relativity and in modified theories of gravitation. The most of the authors who have analyzed some cosmological models with dark energy both in general relativity and modified theories of gravitation, some of them are mentioned in Ref. (Chirde and Shekh, 2015, Chirde and Shekh, 2018a; Bhoyar et al., 2017; Shekh and Chirde 2020); Naidu et al., 2012; Sarkar and Mahanta 2013; Kiran et al., 2015; Santhi et al., 2017; Aditya and Reddy, 2018). The reconstruction phenomenon of a well-known PDE model with f(G) gravity in the presence of power law scale factor Jawad and Rani, (2015) have reconstructed f(G) models with respect to two values of PDE parameter; that is, u = 2, -2 and checked the significant cosmological aspects of these reconstructed models while Jawad et al. (2016a); Jawad et al. (2017) studied the cosmological consequences of pilgrim dark energy model in the framework of generalized teleparallel gravity by considering the reconstruction

scheme for f(T) models with power law scale factor taking Hubble horizon and Nojiri-Odintsov length as infrared cut-offs also observed that the Hubble parameter lies within observational suggested ranges while deceleration parameter represents the accelerated expansion behaviour of the Universe and the model corresponds to the quintessence region and phantom region for different cases of pilgrim dark energy parameter *u* under pilgrim dark energy models in fractal Universe while very recently Shekh, (2021a) analyzed the dynamical investigation of different models of holographic dark energy using Friedman-Lemâitre-Robertson-Walker cosmological model in the context of same f(Q) gravity by governing the features of the model in view of the relation between cosmic time and redshift which yields a purely accelerating evolving Universe.In order to obtain the analytic solution, consider the following form of dimensionless scale factor as of the form:

$$a(t) = a_0 t^n \tag{18}$$

where the subscript 0 denote the value of a quantity at present, a_0 is a constant represents the present day value of the scale factor and moreover set a_0 to 1. The deceleration parameter (*q*) of the Universe, reads ast

$$q = -\frac{a\ddot{a}}{\dot{a}^2} = -1 + \frac{1}{n}$$
(19)

The cosmic scale factor in terms of the deceleration parameter may be written as

$$a(t) = t^{1/(1+q)}$$
(20)

From the above Eq. (20), it is observed that q > -1 is the condition for expanding Universe in the power-law expansion cosmological model.

The expansion history of the Universe and the present expansion rate of the Universe are respectively described by the Hubble parameter as:

$$H = \left(\frac{1}{1+q}\right)t^{-1} \quad \text{and} \quad H_0 = \left(\frac{1}{1+q}\right)t_0^{-1} \tag{21}$$

The above Eq. (21), shows that the expansion history of the Universe in power-law cosmology depends on the two parameters say H_0 and q.

Considering the connection between a and z, the time derivatives of H using Eq. (21) are obtained as:

$$H = H_0 (1+z)^{1+q}, \dot{H} = -H_0 (1+z)^{2(1+q)}, \ddot{H} = 2H_0 (1+z)^{3(1+q)} \text{ and } \ddot{H} = -6H_0 (1+z)^{4(1+q)}.$$
 (22)

3 SOME BASICS AND FIELD EQUATIONS OF *F*(*Q*) GRAVITY

Let us consider the action for f(Q) gravity of the form (Shekh, 2021b)

$$S = \left(\frac{1}{2\kappa^2} \int f(Q) + \int \mathfrak{L}_m\right) \sqrt{-g} d^4 x, \qquad (23)$$

where f(Q) is a general function of Q, \mathcal{L}_m is the matter Lagrangian density and g is the determinant of metric g_{uv} .

The non-metricity tensor and its traces are such that

$$Q_{\gamma\mu\nu} = \nabla_{\gamma} g_{\mu\nu}, \qquad (24)$$

$$Q_{\gamma} = Q_{\gamma \mu}^{\mu}, \qquad \tilde{Q}_{\gamma} = Q_{\gamma \mu}^{\mu}. \qquad (25)$$

Moreover, the superpotential as a function of non-metricity tensor is given by

$$4P^{y}_{\mu\nu} = -Q^{y}_{\mu\nu} + 2Q_{(\mu^{y}\nu)} - Q^{y}g_{\mu\nu} - \tilde{Q}^{y}g_{\mu\nu} - \delta^{y}_{(\gamma^{Q}\nu)}, \qquad (26)$$

where the trace of non-metricity tensor Eq. (20) has the form

$$Q = -Q_{\gamma\mu\nu}P^{\gamma\mu\nu}.$$
 (27)

Expression for energy-momentum tensor for the matter, whose definition is

$$T_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta\left(\sqrt{-g} \mathcal{L}_m\right)}{\delta g^{\mu\nu}}.$$
 (28)

Variation of action **Eq. (18)** with respect to metric tensor, one can obtain gravitational equation

$$\begin{aligned} &\frac{2}{\sqrt{-g}} \nabla_{\gamma} \left(\sqrt{-g} f_{Q} P^{\gamma}{}_{\mu\nu} \right) + \frac{1}{2} g_{\mu\nu} f + f_{Q} \left(P_{\mu\gamma i} Q_{\nu}{}^{\gamma i} - 2 Q_{\gamma i\mu} P^{\gamma i}{}_{\nu} \right) \\ &= -\kappa^{2} \hat{T}_{\mu\nu}, \end{aligned}$$

$$(29)$$

where $f_Q = \frac{df}{dQ}$. The variation of Eq. (24) with respect to the connection term, obtain

$$\nabla_{\mu}\nabla_{\gamma}\left(\sqrt{-g}\,f_{Q}P^{\gamma}{}_{\mu\nu}\right) = 0. \tag{30}$$

For isotropic, homogeneous and spatially FRW space-time provided in **Eq.** (4), one can find the modified Friedmann equations for f(Q) gravity as

$$\dot{H} + 3H^2 + \frac{\dot{f}_Q}{f_Q}H = \frac{1}{2f_Q}\left(\kappa^2 p_D + \frac{f}{2}\right)$$
(31)

$$3H^2 = \frac{1}{2f_Q} \left(-\kappa^2 \rho_D + \frac{f}{2} \right) \tag{32}$$

The overhead dot represents the differentiation with respect to cosmic time *t*.

Also, the modified Friedmann equations enable us to write the pressure and the density for the Universe as

$$\kappa^2 \rho_D = \frac{f}{2} - 6H^2 f_Q \tag{33}$$

$$\kappa^{2} p_{D} = \left(\dot{H} + 3H^{2} + \frac{\dot{f}_{Q}}{f_{Q}} H \right) (2f_{Q}) - \frac{f}{2}$$
(34)

which are the pressure and the density for the f(Q) gravity.For the spatially homogeneous and isotropic FRW Universe, the nonmetricity Q term is defined as $Q = 6H^2$. With the use of **Eq. (22)**, the term Q is observed as





$$Q = 6H_0^2 (1+z)^{(2+2q)}$$
(35)



f(Q) gravity model versus redshift for the appropriate choice of constants.



$$\frac{1}{2}f - 6H^2f_Q = \alpha H \tag{37}$$

The re-arrangement of the above Eq. (37) provide

$$f_{\rm Q} - \left(\frac{1}{12H^2}\right)f = -\frac{\alpha}{6H} \tag{38}$$

which is the first order linear differential equation in Q whose solution is of the form as

$$f(Q) = (c \ln Q + c_1)Q^{1/2}$$
(39)

where *c* and c_1 both are the positive constants of integration.**Eq.** (39), represents the reconstructed ghost dark energy f(Q) gravity model. The plot which shows the behavior of reconstructed ghost dark energy f(Q) gravity model versus redshift and the non-metricity parameter are respectively presented in **Figures 1**, 2 which described that the reconstructed ghost dark energy f(Q)

4 RECONSTRUCTION OF GHOST DARK ENERGY *F*(*Q*) GRAVITY

As, we know the ghost dark energy model is one of the dynamical dark energy model whose energy density is defined as (Pawar et al., 2018)

$$\rho_{GDE} = \alpha H \tag{36}$$

where α is an arbitrary model constant parameter having square dimension. We establish the correspondence between ghost dark energy and f(Q) gravity model by equating the corresponding densities. Using **Eqs 33, 36**, it follows that



gravity model is always positive and increases exponentially with respect to both z and Q. Also, it is noted that the values of positive constants of integration not affect on the behavior of f.

4.1 Cosmographic Observations in Reconstructed Ghost Dark Energy *f*(*Q*) Gravity Model

Using Eq. (39) in Eqs 33 and 34, the expression for energy density and pressure are obtained as:

$$\kappa^2 \rho = (c + c_1 + c \ln Q) Q^{1/2}$$
(40)

$$\begin{aligned} \kappa^2 p &= 2 \Big(\dot{H} + 6H^2 \Big) \Big(c + \frac{c_1}{2} + \frac{c}{2} \ln Q \Big) Q^{-1/2} \\ &+ 2H \Big(\frac{-c_1}{4} - \frac{c}{4} \ln Q \Big) \dot{Q} Q^{-3/2} - \Big(\frac{c_1}{2} + \frac{c}{2} \ln Q \Big) Q^{1/2} \end{aligned} \tag{41}$$

The behavior of both pressure and energy density versus redshift of reconstructed ghost dark energy f(Q) gravity model is clearly shown in **Figures 3**, **4** respectively. **Figure 4** depicted that the energy density of reconstructed ghost dark energy f(Q) gravity model is always positive and exponentially increases (see **Figure 4**) while pressure is always negative and shows negatively decreasing behaviour (see **Figure 3**) for all z = -1 to z > 0. Hence, such a behavior of pressure is the evidence of existance of dark energy.

4.1.1 Equation of State Parameter

$$\omega_D = \frac{-1}{2 - (c + c_1 + c \ln Q)Q^{-1/2}} \left(\frac{(c + c_1 + c \ln Q)Q^{-1/2} - 2\eta}{(c + c_1 + c \ln Q)Q^{-1/2}}\right)$$
(42)

Eq. (42) represents the expression for equation of state parameter of ghost dark energy f(Q) gravity model. **Figure 5** represents the dynamical evolution of the equation of state parameter of ghost dark energy f(Q) gravity model for three consecutive values of η . One can see in **Figure 5** that at late Universe (z < -1) towords $\eta = 0.10, 0.15$



and 0.20 the value of equation of state parameter of ghost dark energy f(Q) gravity model is less than -1 i.e. $(\omega_D)_{GDE} < -1$ which represents the model involve phantom field dark energy whereas for present Universe (z = 0), it is a little bit upper than -1 i.e. $(\omega_D)_{GDE} > -1$ and early Universe (z > 0) it has the value $(\omega_D)_{GDE} > 0$. Hence the present Universe consist of a quintessence field dark energy and early Universe involve barotropic fluid. Also, notice that by increasing the value of interaction parameter η , equation of state parameter takes more negative values, below the -1. It is observed that the equation of state in our framework can cross the phantom divide line as supported by recent astrophysical observations as well as the analysis of holographic dark energy inflation with Hubble's cut-off analyzed by Shekh, (2021a) and also with Jawad and Rani, (2015), Jawad, (2015).

4.1.2 $(\omega_D - \omega_{D'})$ -plane

The derivative of Eq. (42) with respect to $\ell n(a)$, gives

$$\omega_D' = \frac{(c/2 + c_1/2 + (c/2) \ln Q) Q Q^{-3/2}}{2 - (c + c_1 + c \ln Q) Q^{-1/2}} \left\{ 2\eta - \frac{2\eta - (c + c_1 + c \ln Q) Q^{-1/2}}{(c + c_1 + c \ln Q) Q^{-1/2} (2 - (c + c_1 + c \ln Q) Q^{-1/2})} \right\}$$
(43)

The plot of ω'_D with respect to ω_D for the ghost dark energy f(Q) gravity model is shown in **Figure 6**. Figures 5, 6, indicates that when $\omega_D < 0$ then also $\omega'_D < 0$ which represents freezing region.

4.1.3 (r-s)-plane

$$s = 1 + \frac{-1}{2 - (c + c_1 + c \ln Q)Q^{-1/2}} \left(\frac{(c + c_1 + c \ln Q)Q^{-1/2} - 2\eta}{(c + c_1 + c \ln Q)Q^{-1/2}} \right) \left(\frac{(c + c_1 + c \ln Q)Q^{-1/2} - 2\eta}{(c + c_1 + c \ln Q)Q^{-1/2}} \right) \times (c + c_1 + c \ln Q)Q^{-1/2}} \right) \\ + \frac{(c + c_1 + c \ln Q)Q^{-1/2} - \frac{2}{3} \left(\frac{2\eta\dot{Q}}{(2 - (c + c_1 + c \ln Q)Q^{-1/2}} + \frac{(c + c_1 + c \ln Q)Q^{-1/2}}{2(2 - (c + c_1 + c \ln Q)Q^{-1/2}} \right) \times (c + c_1 + c \ln Q)Q^{-1/2}} \right) }{(44)} \\ s = 1 + \frac{-1}{2 - (c + c_1 + c \ln Q)Q^{-1/2}} \left(\frac{(c + c_1 + c \ln Q)Q^{-1/2} - 2\eta}{(c + c_1 + c \ln Q)Q^{-1/2}} \right) + \frac{(\frac{2\eta\dot{Q}}{(2 - (c + c_1 + c \ln Q)Q)Q^{-1/2}} + \frac{(c + c_1 + c \ln Q)Q^{-1/2}}{2(2 - (c + c_1 + c \ln Q)Q^{-1/2}} - 2\eta}}{3\left(\frac{(c + c_1 + c \ln Q)Q^{-1/2} - 2\eta}{(c + c_1 + c \ln Q)Q^{-1/2}}\right)} \right)}$$

$$(45)$$

Figure 7, shows the evolution trajectory for ghost dark energy f(Q) gravity model in (r - s)- plane towards different value of η . From **Figure 7**, the evolution trajectories of (r - s)- plane favors the chaplygin gas model with s < 0 and r > 1. Hence, our results are





consistent with the analysis of Wu and Yu, (2006) and the authors who have achieved the state finder diagnostic for the phantom and quintom dark energy model (Sharif and Zubair, 2014; Sharif, 2018).

5 RECONSTRUCTION OF PILGRIM DARK ENERGY *F*(Q) GRAVITY

Proceeding the same as it is in section 3. In this section, we reconstructed the pilgrim dark energy f(Q) gravity model. For this consider an interesting model for the description of dark energy is pilgrim dark energy in which the total energy in a box of size L could exceed the mass of a black hole of the same size i.e., $\rho_p L^3 \ge m_{pl}^2 L$, where m_{pl} be the Plank reduced mass. Therefore, the first property of pilgrim dark energy is (Jawad et al., 2016b)

$$\rho_p \ge m_{pl}^2 L^{-2} \tag{46}$$

the simplest way to set the above equation is:

$$\rho_p = 3n^2 m_{pl}^{4-s} L^{-s} \tag{47}$$

where *n* and *s* are the conventional constant and pilgrim dark energy parameter respectively. Thus, from **Eqs 46**, 47, $L^{2-s} \ge m_{pl}^{s-2} = \ell_{pl}^{2-s}$, where $m_{pl} = 1/\sqrt{8\pi G}$ is the reduced Planck length, which is extremely short length. Obviously, since $L > \ell_{pl}$ in general, it is required that

$$s \le 2$$
 (48)

Keeping the box of size L = 1/H towards Hubble's cutoff, the dynamical pilgrim dark energy model whose energy density from **Eq. (47)** is obtained as

$$\rho_{PDE} = 3n^2 m_{pl}^{4-s} H^s \tag{49}$$

We establish the correspondence between pilgrim dark energy and f(Q) gravity model by equating the conforming densities. Using **Eqs 33, 49**, it follows that

$$\frac{1}{2}f - 6H^2 f_Q = 3n^2 H^u \tag{50}$$

The re-arrangement of above Eq. (50) provide

$$f_{\rm Q} - \left(\frac{1}{12H^2}\right)f = -\frac{3n^2H^u}{6H^2}$$
(51)

which is the first order linear differential equation in Q whose solution is of the form as

$$f(Q) = -c_2 Q^{\frac{\mu}{2}} + c_3 Q^{1/2}$$
(52)

where $c_2 = \frac{n^2}{(u-1)6^{u/2-1}}$ and c_3 be the constant of integration.**Eq. (52)** represents the reconstructed pilgrim dark energy f(Q) gravity model. The plot which shows the behavior of reconstructed pilgrim dark energy f(Q) gravity model versus redshift and the non-metricity parameter are respectively presented in **Figures 8**, **9** which described that the reconstructed pilgrim dark energy f(Q) gravity model is always negative and decreases negatively with respect to both *z* and *Q*.





5.1 Cosmographic Observations in Reconstructed Pilgrim Dark Energy f(Q)Gravity Model

Using Eq. (52) in Eqs 33 and 34, the expression for energy density and pressure are obtained as:

$$\rho = \frac{c_2 \left(1 - u\right)}{2} Q^{u/2} \tag{53}$$

$$p = (H + 3H^{2})(c_{2}uQ^{\frac{1}{2}-1} + c_{3}Q^{-1/2}) + \frac{H}{2}(c_{2}u(u-2)Q^{\frac{u}{2}-1} - c_{3}Q^{-1/2})\dot{Q} - \frac{c_{2}}{2}Q^{u/2} - \frac{c_{3}}{2}Q^{1/2}$$
(54)

The behavior of both pressure and energy density versus redshift of reconstructed pilgrim dark energy f(Q) gravity model is clearly seen in **Figures 10**, **11** respectively. The



FIGURE 11 Plot of the energy density of the reconstructed pilgrim dark energy *f*(*Q*) gravity model versus redshift for the appropriate choice of constants.









energy density of reconstructed ghost dark energy f(Q) gravity model is always positive and exponentially increases (see **Figure 11**) whereas pressure is always negative and shows negatively decreasing behevior (see **Figure 10**) for all z = -1 to z > 0 which is the evidance of existance of dark energy.

5.1.1 Equation of State Parameter

$$\omega_D = -\frac{1}{2 - c_2 (1 - u) Q^{\frac{u}{2} - 1}} \left(1 + \frac{2\eta}{c_2 (1 - u) Q^{\frac{u}{2} - 1}} \right)$$
(55)

Eq. (55) represents the expression for equation of state parameter of pilgrim dark energy f(Q) gravity model. Figure 12 represents the dynamical evolution of the equation of state parameter of pilgrim dark energy f(Q) gravity model for three consecutive values of n. One can see in Figure 12 that at late Universe (z < -1) towords $\eta = 0.10, 0.15$ and 0.20 the value of equation of state parameter of pilgrim dark energy f(Q) gravity model is less than -1 i.e., $(\omega_D)_{PDE} < -1$ which represents the model involving phantom field dark energy which is the same as that of ghost dark energy f(Q)gravity model whereas for present (z = 0) and early Universe (z > 0), the equation of state parameter of pilgrim dark energy f(Q) gravity model is a little bit upper than -1 i.e. $(\omega_D)_{PDF} > -1$. Hence the present and early Universe consist quintessence field dark energy. Also, in the pilgrim dark energy f(Q) gravity model it is noticed that by increasing the value of interaction parameter η , equation of state takes more negative values, below the -1. It is observed that the equation of state in our framework can cross the phantom divide line as supported by recent astrophysical observations along with the work holographic dark energy inflation with Granda-Oliveros cut-off and Renyi holographic dark energy in both Hubble's as well as Granda-Oliveros cut-off investigated by Shekh, (2021b) also with the work of Jawad and Rani, (2015), Jawad et al. (2015).

5.1.2 $(\omega_D - \omega_{D'})$ plane

$$\omega_D' = \frac{2\eta c_2 (1-u) (u/2-1) Q^{\frac{u}{2}-2} \dot{Q}}{(2-c_2 (1-u) Q^{\frac{u}{2}-1})} - \frac{c_2 (1-u) (u/2-1) Q^{\frac{u}{2}-2} \dot{Q}}{(2-c_2 (1-u) Q^{\frac{u}{2}-1})^2} + \frac{2\eta (u/2-1)}{(2-c_2 (1-u) Q^{\frac{u}{2}-1})^2} \frac{\dot{Q}}{Q}$$
(56)

The plot of ω'_T with respect to ω_T for the pilgrim dark energy f(Q) gravity modelis shown in **Figure 13**. **Figures 12**, **13**, indicates that when $\omega_T < 0$, $\omega'_T > 0$ it represents thaving region.

5.1.3 (r-s)-plane

$$s = \begin{cases} 1 - \frac{9}{2} \frac{c_{2}(1-u)Q^{u^{2}-1}}{2-c_{2}(1-u)Q^{t-1}} \left(1 + \frac{2\eta}{c_{2}(1-u)Q^{t-1}}\right) \left(1 - \frac{1}{2-c_{2}(1-u)Q^{t-1}} \left[1 + \frac{2\eta}{c_{2}(1-u)Q^{t-1}}\right]\right) \\ -\frac{3c_{2}(1-u)}{2H} \left(\frac{2\eta c_{2}(1-u)(u^{2}-1)Q^{t+2}\dot{Q}}{(2-c_{2}(1-u)Q^{t-1})} - \frac{c_{2}(1-u)(u^{2}-1)Q^{t+2}\dot{Q}}{(2-c_{2}(1-u)Q^{t-1})^{2}\dot{Q}} + \frac{2\eta(u^{2}-1)}{(2-c_{2}(1-u)Q^{t-1})^{2}} \frac{\dot{Q}}{\dot{Q}}\right) Q^{u^{2}-1} \end{cases}$$

$$s = \begin{cases} 1 - \frac{1}{2-c_{2}(1-u)Q^{u^{2}-1}} \left(1 + \frac{2\eta}{c_{2}(1-u)Q^{u^{2}-1}}\right) - \frac{1}{3H} \left(\frac{2\eta c_{2}^{2}(1-u)Q^{u^{2}-1}}{2\eta + c_{2}(1-u)Q^{u^{2}-1}}\right) + \frac{1}{3H} \left(\frac{c_{2}^{2}(1-u)Q^{u^{2}-1}}{2-c_{2}(1-u)Q^{u^{2}-1}} - \frac{2\eta c_{2}(1-u)(u^{2}-1)Q^{u^{2}-1}}{2-c_{2}(1-u)Q^{u^{2}-1}}\right) - \frac{1}{3H} \left(\frac{2\eta c_{2}^{2}(1-u)Q^{u^{2}-1}}{2-c_{2}(1-u)Q^{u^{2}-1}}\right) + \frac{1}{3H} \left(\frac{c_{2}^{2}(1-u)^{2}(u^{2}-1)Q^{u^{2}-1}}{2-c_{2}(1-u)Q^{u^{2}-1}} - \frac{2\eta c_{2}(1-u)(u^{2}-1)Q^{u^{2}-1}}{2-c_{2}(1-u)Q^{u^{2}-1}}\right) - \frac{1}{2} \left(\frac{1}{2}\right) \frac{1}{2}$$

Figure 14, shows the evolution trajectory for pilgrim dark energy f(Q) gravity model in (r - s)- plane towards different value of η . From **Figure 14**, the evolution trajectories of (r - s)- plane favor the quintessence field dark energy model with s > 0 and r < 1. Hence, our results are consistent with the analysis of (Sharif and Zubair, 2014; Sharif, 2018).

6 CONCLUSION

In the present analysis, to talk over the evolution of the Universe We have considered the interacting f(Q) gravity with pressureless matter in an FRW Universe via the

reconstruction scheme with power-law form of the scale factor. For this purpose, we have used a newly proposed ghost and pilgrim dark energy model which has a strong repulsive force without a formation of black holes also these are the remarkable models that solves the recent cosmic acceleration problem. The Hubble horizon is taken as the IR cutoff, which provides reliable results with interaction. To establish the correspondence between ghost, pilgrim dark energy and f(Q) gravity model, the corresponding densities have been considered equal. The physical motivation and the consequence of considering this equality show whether the reconstructed model is a realistic one or not. Such methods i.e. equating two densities have been extensively studied in the literature on ghost and pilgrim dark energy in the framework of different modified gravities. To discuss the ghost and pilgrim dark energy f(Q) gravity models, We have explored the evolution trajectories of the equation of state parameter, the $(\omega_D - \omega'_D)$ -phase plane, and the state finder (r - s)- plane. The final results in both ghost and pilgrim dark energy models are respectively concise as follows.

6.1 Ghost Dark Energy f(Q) Gravity Model

- (1) The reconstructed ghost dark energy f(Q) gravity model represents increasing behavior forever with respect to both z and Q which shows that the reconstructed model is a realistic one.
- (2) The equation of state parameter represents the late Universe it involves phantom field dark energy while the present Universe consist of a quintessence field dark energy and early Universe involve barotropic fluid. Also, we should notice that with the different increasing value of interaction parameter η , equation of state parameter takes more negative values below the -1. Hence, our results are consistent with the current accelerated cosmic behavior and hence I conclude that the ghost dark energy f(Q)gravity model favors the dark energy phenomenon. The equation of state in our framework can crosses the phantom divide line as supported by recent astrophysical observations.
- (3) The evolutionary behavior of the (ω_D ω'_D)-phase plane towards all η represents freezing region which confirmed the cosmological expansion is more accelerating in interacting Ghost dark energy f(Q) gravity model.
- (4) The corresponding trajectories of (r s)-plane indicate Chaplygin gas model for all η . Furthermore, it attains CDM limit but cannot achieve Λ CDM limit.

6.2 Pilgrim Dark Energy *f*(**Q**) Gravity Model

- (1) The reconstructed ghost dark energy f(Q) gravity model is always negative and decreases negatively with respect to both z and Q.
- (2) The equation of state parameter at late Universe is less than -1 which represents the model involved phantom field dark energy which is the same as that of Ghost dark

energy f(Q) gravity model whereas for present and early Universe it is a little bit upper than -1. Hence the present and early Universe consists of a quintessence field dark energy. Also, in the Pilgrim dark energy f(Q) gravity model, it is noticed that by increasing the value of interaction parameter η , equation of state takes more negative values, below the -1. It is observed that the equation of state in our framework can cross the phantom divide line as supported by recent astrophysical observations.

- (3) The evolutionary behavior of the $(\omega_D \omega'_D)$ -phase plane towards all η represents thawing region in an interacting Ghost dark energy f(Q) gravity model.
- (4) The corresponding trajectories of (r s)- plane favor the quintessence field dark energy model for all η . Furthermore, it crosses Λ CDM limit.

Hence, the results obtained in both the reconstruction scheme under ghost and pilgrim dark energy are rusumbles with the recent modern theoretical observational data as well as the work analyzed by (sharf and Zuber, 2014; Jawad and Rani, 2015; Jawad et al., 2016a; Wu and Yu, 2006; Shafiz and Saba, 2019).

DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/supplementary material, further inquiries can be directed to the corresponding author.

AUTHOR CONTRIBUTIONS

All of the authors listed have contributed a significant, direct, and intellectual contribution to the work and have given their permission for it to be published.

FUNDING

This research was funded by the Science Committee of the Ministry of Education and Science of the Republic of Kazakhstan (Grant No. AP09058240).

ACKNOWLEDGMENTS

The content of this manuscript has been presented in part at the "Two day International Virtual Conference on Recent Trends in Mathematical Sciences (IVCRTMS-2021)", https://www.ijirmf.com/ivcrtms-oct-2021/. The authors are very much grateful to the honorary referee and the editor for the illuminating suggestions that have significantly improved our work in terms of research quality and presentation. 10.054

Expansion Law. Astrophysics 60, 259-272. doi:10.1007/s10511-017-9480-y Caldwell, R. R., and Linder, E. V. (2005). Limits of Quintessence. Phys. Rev. Lett. 95, 141301. doi:10.1103/physrevlett.95.141301

Abbas, G., Momeni, D., Aamir Ali, M., Myrzakulov, R., and Qaisar, S. (2015).

Aditya, Y., and Reddy, D. R. K. (2018). FRW Type Kaluza-Klein Modified

Gravitation. Eur. Phys. J. C 78, 619. doi:10.1140/epjc/s10052-018-6074-8

Azadi, A., Momeni, D., and Nouri-zonoz, M. (2008). Cylindrical Solutions in

Bhatti, M. Z.-u. -H., Yousaf, Z., and Hanif, S. (2017). Role off(T)gravity on the

Bhoyar, S. R., Chirde, V. R., and Shekh, S. H. (2017). Stability of Accelerating

Anisotropic Compact Stars in F(G) Gravity. Astrophys. Space Sci. 357, 158.

Holographic Ricci Dark Energy Models in Brans-Dicke Theory of

Metric F (R) Gravity. Phys. Lett. B 670, 210. doi:10.1016/j.physletb.2008.

Evolution of Collapsing Stellar Model. Phys. Dark Universe 16, 34-40. doi:10.

Universe with Linear Equation of State in F (T) Gravity Using Hybrid

- Capozziello, S., Stabile, A., and Troisi, A. (2007). Spherically Symmetric Solutions in F (R) Gravity via the Noether Symmetry Approach. Cl. Quantum Grav. 24, 2153-2166. doi:10.1088/0264-9381/24/8/013
- Capozziello, S., Martin-Moruno, P., and Rubano, C. (2008). Dark Energy and Dust Matter Phases from an Exact F(R)-cosmology Model. Phys. Lett. Sect. B 664, 12. doi:10.1016/j.physletb.2008.04.061
- Chirde, V. R., and Shekh, S. H. (2015). Dark Energy Cosmological Model in a Modified Theory of Gravity. Astrophysics 58, 106-119. doi:10.1007/s10511-015-9369-6
- Chirde, V. R., and Shekh, S. H. (2016a). Isotropic Background for Interacting Two Fluid Scenario Coupled with Zero Mass Scalar Field in Modified Gravity. Bul. J. Phys. 43, 156.
- Chirde, V. R., and Shekh, S. H. (2016b). Plane Symmetric Dark Energy Models in the Form of Wet Dark Fluid in F (R,T) Gravity. J. Astrophys. Astron. 37, 1-16. doi:10.1007/s12036-016-9391-z
- Chirde, V. R., and Shekh, S. H. (2018a). Dynamic Minimally Interacting Holographic Dark Energy Cosmological Model in F(T) Gravity. Indian J. Phys. 92, 1485-1494. doi:10.1007/s12648-018-1236-y
- Chirde, V. R., and Shekh, S. H. (2018b). Transition between General Relativity and Quantum Gravity Using Quark and Strange Quark Matter with Some Kinematical Test. J. Astrophys. Astron 39, 56. doi:10.1007/s12036-018-9555-0
- Chirde, V. R., and Shekh, S. H. (2019). Dynamics of Magnetized Anisotropic Dark Energy in *f*(*R*, *T*) Gravity with Both Deceleration and Acceleration. *Bulg. J. Phys.* 46, 94,
- Dagwal, V. J., and Pawar, D. D. (2020). Two-fluid Sources in F (T) Theory of Gravity. Mod. Phys. Lett. A 35, 1950357. doi:10.1142/s0217732319503577
- Daouda, M. H., Rodrigues, M. E., and Houndjo, M. J. S. (2012). Static Anisotropic Solutions in F(T) Theory. Eur. Phys. J. C 72, 1890. doi:10.1140/epjc/s10052-012-1890-8
- Eisenstein, D. J. (2005). Detection of the Baryon Acoustic Peak in the Large-Scale Correlation Function of SDSS Luminous Red Galaxies. Astrophys. J. 633, 560. doi:10.1086/466512
- Frusciante, N. (2021). Signatures of f(Q)-Gravity in Cosmology. Phys. Rev. D 103 (4), 044021. doi:10.1103/PhysRevD.103.044021
- Harko, T., Lobo, F. S. N., Nojiri, S., and Odintsov, S. D. (2011). f(R,T) Gravity. Phys. Rev. D. 84, 1. doi:10.1103/physrevd.84.024020
- Jain, B., and Taylor, A. (2003). Cross-Correlation Tomography: Measuring Dark Energy Evolution with Weak Lensing. Phys. Rev. Lett. 91, 141302. doi:10.1103/ physrevlett.91.141302
- Jawad, A., and Rani, S. (2015). Cosmological Evolution of Pilgrim Dark Energy in F(G) Gravity. Adv. High Energy Phys. 2015, 952156. doi:10.1155/2015/952156
- Jawad, A., Debnath, U., and Batool, F. (2015). Generalized Ghost Pilgrim Scalar Field Models of Dark Energy. Commun. Theor. Phys. 64, 590-596. doi:10.1088/ 0253-6102/64/5/590
- Jawad, A., Chattopadhyay, S., and Rani, S. (2016a). Viscous Pilgrim F (T) \$F(T)\$ Gravity Models. Astrophys. Space Sci. 361, 231. doi:10.1007/s10509-016-2814-0

Analysis of Reconstructed Modified Symmetric

- J. Plus 131, 236. doi:10.1140/epjp/i2016-16236-x Jawad, A., Rani, S., Salako, I. G., and Gulshan, F. (2017). Pilgrim Dark Energy Models in Fractal Universe. Int. J. Mod. Phys. D. 26, 1750049. doi:10.1142/ s0218271817500493
- Jawad, A. (2015). Cosmological Analysis of Pilgrim Dark Energy in Loop Quantum Cosmology. Eur. Phys. J. C 75, 206. doi:10.1140/epjc/s10052-015-3430-9
- Jimenez, J. B., Heisenberg, L., and Koivisto, T. (2018). Coincident General Relativity. Phys. Rev. D. 98, 044048. doi:10.1103/physrevd.98.044048
- Kawarabayashi, K., and Ohta, N. (1980). The η Problem in the Large-N Limit: Effective Lagrangian Approach. Nucl. Phys. B 175, 477-492. doi:10.1016/0550-3213(80)90024-3
- Kiran, M., Reddy, D. R. K., and Rao, V. U. M. (2015). Minimally Interacting Holographic Dark Energy Model in Brans-Dicke Theory. Astrophys. Space Sci. 356, 407-411. doi:10.1007/s10509-014-2213-3
- Komatsu, E. (2011). Seven-year Wilkinson Microwave Anisotropy Probe (WMAP) Observations: Cosmological Interpretation. Astrophys. J. Suppl. 192, 18. doi:10. 1088/0067-0049/192/2/18
- Lazkoz, R. (2019). Observational Constraints of F(Q) Gravity. Phys. Rev. D. 100, 104027. doi:10.1103/PhysRevD.100.104027
- Mandal, S., Sahoo, P. K., and Santos, J. R. L. (2020a). Energy Conditions in F (Q) Gravity. Phys. Rev. D. 102, 024057. doi:10.1103/physrevd.102.024057
- Mandal, S., Wang, D., and Sahoo, P. K. (2020b). Cosmography in F(Q) Gravity. Phys. Rev. D. 102, 124029. doi:10.1103/physrevd.102.124029
- Naidu, R. L., Satyanarayana, B., and Reddy, D. R. K. (2012). LRS Bianchi Type-II Dark Energy Model in a Scalar-Tensor Theory of Gravitation. Astrophys. Space Sci. 338, 333-336. doi:10.1007/s10509-011-0935-z
- Nath, P., and Arnowitt, R. (1981). U(1) Problem: Current Algebra and Theovacuum. Phys. Rev. D. 23, 473-476. doi:10.1103/physrevd.23.473
- Nojiri, S. i., and Odintsov, S. D. (2007). Modified Gravity and its Reconstruction from the Universe Expansion History. J. Phys. Conf. Ser. 66, 012005. doi:10. 1088/1742-6596/66/1/012005
- Nojiri, S., and Odintsov, S. (2008). Future Evolution and Finite-Time Singularities in F(R) Gravity Unifying Inflation and Cosmic Acceleration. Phys. Rev. D. 78, 046006. doi:10.1103/PhysRevD.78.046006
- Pawar, D. D., Bhuttampalle, G. G., and Agrawal, P. K. (2018). Kaluza-Klein String Cosmological Model in f(R, T) Theory of Gravity. New Astron. 65, 1-6. doi:10. 1016/j.newast.2018.05.002
- Perlmutter, S. (1999). Measurements of Ω and Λ from 42 High-Redshift Supernovae. Astrophys. J. 517, 565. doi:10.1086/307221
- Rani, S., Jawad, A., Ines, G., and SalakoAzhar, N. (2016). Non-Flat Pilgrim Dark Energy Frw Models in Modified Gravity. Astrophys. Space Sci. 361, 386. doi:10. 1007/s10509-016-2868-z
- Riess, A. G. (1998). Observational Evidence from Supernovae for an Accelerating Universe and a Cosmological Constant. Astron. J. 116, 1009. doi:10.1086/ 300499
- Sadeghi, J., Khurshudyan, M., Movsisyan, A., and Farahani, H. (2013). Interacting Ghost Dark Energy Models with variableGand A. J. Cosmol. Astropart. Phys. 2013, 031. doi:10.1088/1475-7516/2013/12/031
- Sahni, V., Saini, T. D., Starobinsky, A. A., and Alam, U. (2003). Statefinder-A New Geometrical Diagnostic of Dark Energy. Jetp Lett. 77, 201-206. doi:10.1134/1. 1574831
- Sahoo, P. K., and Bhattacharjee, S. (2020). Revisiting the Coincidence Problem in F(R) Gravitation. New Astron. 77, 101351. doi:10.1016/j.newast.2019.101351
- Sahoo, P. K., Mishra, B., and Chakradhar, R. (2014). Axially Symmetric Cosmological Model in f(R, T) Gravity. Eur. Phys. J. Plus. 129, 1. doi:10. 1140/epjp/i2014-14049-7
- Santhi, M. V., Rao, V. U. M., and Aditya, Y. (2017). Bianchi type-VI0 Modified Holographic Ricci Dark Energy Model in a Scalar-Tensor Theory. Can. J. Phys. 95, 179-183. doi:10.1139/cjp-2016-0628
- Sarkar, S., and Mahanta, C. R. (2013). Holographic Dark Energy Model with Quintessence in Bianchi Type-I Space-Time. Int. J. Theor. Phys. 52, 1482-1489. doi:10.1007/s10773-012-1468-0
- Seljak, U. (2005). Cosmological Parameter Analysis Including SDSS Ly a Forest and Galaxy Bias: Constraints on the Primordial Spectrum of Fluctuations, Neutrino Mass, and Dark Energy. Phys. Rev. D. 71, 103515. doi:10.1103/ PhysRevD.71.103515

REFERENCES

doi:10.1007/s10509-015-2392-6

1016/j.dark.2017.04.003

- Sharif, M., and Fatima, H. I. (2016). Built-in Inflation in F(G) Gravity. Int. J. Mod. Phys. D. 25, 1650011. doi:10.1142/s0218271816500115
- Sharif, M., and Jawad, A. (2013). Pilgrim Dark Energy with Apparent and Event Horizons in Non-flat Universe. Eur. Phys. J. C 73, 2600. doi:10.1140/epjc/ s10052-013-2600-x
- Sharif, M., and Yousaf, Z. (2013). Dynamical Instability of the Charged Expansionfree Spherical Collapse in F(R) Gravity. Phys. Rev. D. 88, 024020. doi:10.1103/ physrevd.88.024020
- Sharif, M., and Zubair, M. (2014). Cosmological Evolution of Pilgrim Dark Energy. Astrophys. Space Sci. 352, 263–272. doi:10.1007/s10509-014-1889-8
- Sharif, M., and Saba, S. (2018). Pilgrim Dark Energy in f(G, T) Gravity. Mod. Phys. Lett. A 33, 1850182. doi:10.1142/s0217732318501821
- Sharif, M., and Saba, S. (2019). Ghost Dark Energy Model in f(G) Gravity. Chin. J. Phys. 58, 202–211. doi:10.1016/j.cjph.2018.12.023
- Shekh, S. H., and Chirde, V. R. (2019). Analysis of General Relativistic Hydrodynamic Cosmological Models with Stability Factor in Theories of Gravitation. *Gen. Relativ. Gravit.* 51, 87. doi:10.1007/s10714-019-2565-7
- Shekh, S. H., and Chirde, V. R. (2020). Accelerating Bianchi Type Dark Energy Cosmological Model with Cosmic String in \$F(T)\$ Gravity. Astrophys. Space Sci. 365, 60. doi:10.1007/s10509-020-03772-y
- Shekh, S. H., Arora, S., Chirde, V. R., and Sahoo, P. K. (2020a). Thermodynamical Aspects of Relativistic Hydrodynamics in f(R,G) Gravity. Int. J. Geom. Methods Mod. Phys. 17, 2050048. doi:10.1142/s0219887820500486
- Shekh, S. H., Chirde, V. R., and Sahoo, P. K. (2020b). Energy Conditions of the f(T, B) Gravity Dark Energy Model with the Validity of Thermodynamics. *Commun. Theor. Phys.* 72, 085402. doi:10.1088/1572-9494/ab95fd
- Shekh, S. H., Katore, S. D., Chirde, V. R., and Raut, S. V. (2021). Signature Flipping of Isotropic Homogeneous Space-Time with Holographic Dark Energy in f(G)gravity. *New Astron.* 84, 101535. doi:10.1016/j.newast. 2020.101535
- Shekh, S. H. (2021a). Models of Holographic Dark Energy in F(Q) Gravity. Phys. Dark Universe 33, 100850. doi:10.1016/j.dark.2021.100850
- Shekh, S. H. (2021b). Dynamical Analysis with Thermodynamic Aspects of Anisotropic Dark Energy Bounce Cosmological Model in f(R, G) Gravity. *New Astron.* 83, 101464. doi:10.1016/j.newast.2020.101464

- Sheykhi, A., and Movahed, M. S. (2012). Interacting Ghost Dark Energy in Nonflat Universe. Gen. Relativ. Gravit. 44, 449–465. doi:10.1007/s10714-011-1286-3
- Spergel, D. N. (2003). First-Year Wilkinson Microwave Anisotropy Probe (WMAP)* Observations: Determination of Cosmological Parameters. Astrophys. J. Suppl. 148, 175. doi:10.1086/377226
- Tegmark, M. (2004). Cosmological Parameters from SDSS and WMAP. *Phys. Rev.* D. 69, 103501. doi:10.1103/PhysRevD.69.103501
- Wei, H., Qi, H.-Y., and Ma, X.-P. (2012). Constraining F(T) Theories with the Varying Gravitational Constant. Eur. Phys. J. C 72, 2117. doi:10.1140/epjc/ s10052-012-2117-8
- Witten, E. (1979). Current Algebra Theorems for the U(1) "Goldstone Boson". Nucl. Phys. B 156, 269–283. doi:10.1016/0550-3213(79)90031-2
- Wu, P., and Yu, H. (2006). Statefinder Parameters for Phantom Dark Energy. *Mod. Phys. Lett. A* 21, 1305–1311. doi:10.1142/s0217732306019293
- Yang, L. (2020). Universal Bounds on the Size of a Black Hole. Eur. Phys. J. C 80, 1204. doi:10.1140/epjc/s10052-020-08521-7

Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

Publisher's Note: All claims expressed in this article are solely those of the authors and do not necessarily represent those of their affiliated organizations, or those of the publisher, the editors and the reviewers. Any product that may be evaluated in this article, or claim that may be made by its manufacturer, is not guaranteed or endorsed by the publisher.

Copyright © 2022 Myrzakulov, Shekh, Mussatayeva and Koussour. This is an openaccess article distributed under the terms of the Creative Commons Attribution License (CC BY). The use, distribution or reproduction in other forums is permitted, provided the original author(s) and the copyright owner(s) are credited and that the original publication in this journal is cited, in accordance with accepted academic practice. No use, distribution or reproduction is permitted which does not comply with these terms.