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A regional classification of time spectral amplitudes in total electron content: Southeastern United States during solar cycle 24

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To investigate the meso-scale structure (100 s of km) of the ionosphere at mid-latitudes, the spectral properties in calculated total electron content (TEC) at a cluster of GPS receivers in and around Florida are analyzed. The ionosphere does not respond exactly the same to periodic solar driving at different locations around the planet, due to the complex electrodynamic interactions of the coupled magnetosphere-ionosphere system. Therefore, at each GPS receiver in the cluster we compare spatio-temporal variations of the spectral amplitudes for diurnal, solar rotation, and seasonal oscillations. The amplitudes for these dominant oscillations are organized with respect to magnetic latitude of the receiver. A low-latitude and high-latitude station are also included to put the mid-latitude ionospheric response into a global context. The amplitudes of diurnal, seasonal, and solar rotation signals are well ordered by magnetic latitude, superposed with meso-scale deviations between stations separated by ~100s of km. The results suggest that spatiotemporal variations of spectral amplitudes in the mid-latitude ionosphere are not dominated by a single process. This conclusion is based on our finding that at high latitude, the shape of the diurnal signal varies significantly less with solar activity compared to low- and mid-latitudes, and additionally, that the ratio of annual to semi-annual amplitudes fluctuates around 1 with time and from station-to-station only at mid-latitudes.

KEYWORDS

ionosphere, total electron content, mid-latitude, global positioning system, meso-scale structure

1 Introduction

Meso-scale (100 s of km) structures form in the ionosphere due to different plasma instabilities (Kelley, 1989a; Kelley, 1989b), some examples of their manifestation including equatorial plasma bubbles (Woodman and La Hoz, 1976), storm enhanced

density plumes (Foster et al., 2021), and (medium-scale) traveling ionospheric disturbances (Hines, 1959, 1960; Hunsucker, 1982). In addition, lightning- and thunderstorm-induced generation of acoustic-gravity waves (Lay et al., 2015) can directly impact the ionosphere at mid-latitudes (Lay et al., 2018) and can be detected from analysis of very low frequency (VLF) pulses (Cheng and Cummer, 2005). The combination of all these processes impact the mid-latitude ionosphere and create irregularities that can disrupt radio-communication (Aa et al., 2019).

Dense arrays of ground-based Global Positioning System (GPS) receivers, such as the National Oceanic and Atmospheric Administration (NOAA) Continuously Operating Reference Stations (CORS) Network (NCN), can be used to probe mesoscale ionosphere structures. In principle, it is possible to resolve strong gradients and ionospheric irregularities at a regional scale using distributed sensors (Lay et al., 2018). A natural use of such a data set would be to train an artificial intelligence (AI) to predict the ionospheric total electron content at a regional scale based on solar wind conditions. But, it is the case that the midlatitude total electron content is dominated by the diurnal signal, with meso-scale structures superposed as fluctuations. Therefore, it is the goal of this study to characterize the ionospheric response to persistent external driving on a regional scale, i.e. what are the amplitudes of strong sinusoidal variations of total electron content (TEC) for the different observatories in a closely separated network of sensors and how do they change in time?

The harmonic response of the ionosphere has received attention for decades through space and ground-based measurements. Davies et al. (1980) compared diurnal variations of TEC during magnetically quiet and disturbed conditions from a total of 5 radio beacons in the United States and Europe. Perevalova et al. (2010) used a widely distributed network of GPS receivers to construct Global Ionospheric Maps (GIMs) with a resolution of 5° longitude and 2.5° latitude. Based on the GIMs, the authors seasonally characterized the heliomagnetically quiet diurnal TEC response for low (0-20°), mid (40-55°), and high (60-87.5°) latitudes, providing a global scale picture of the daily response of the ionosphere. Huang and Roussel-Dupré (2005) quantified the diurnal, seasonal, and solar cycle variations of TEC using Fast On-Orbit Recording of Transient Events (FORTE)-received Los Alamos Portable Pulser (LAPP) signals in Los Alamos, New Mexico. Olwendo et al. (2016) studied the diurnal and annual variations of TEC from 12 receivers in East Africa separated by 25° geographic latitude containing the magnetic equator. They noted the difficulty of accurately computing vertical TEC (vTEC) from slant TEC (sTEC) due to large spatiotemporal gradients of the equatorial ionization anomaly. Guo et al. (2015) constructed TEC time series on a global $5^{\circ} \times 2.5^{\circ}$ grid using more than 350 GPS stations. Morlet wavelet analysis was used to determine periods of TEC variations and the 1-day, 26.5-day, semi-annual and

annual cycles were found to be the dominant signals. Recently, Yasyukevich et al. (2017) employed harmonic spectral analysis to study annual, seasonal, and diurnal variations of the ionosphere. They looked at a full solar cycle of data from 2 GPS receivers in Russia and found oscillations with periods of 1, 1/2 ... up to 1/5 of a day, and 27-day are clearly distinguished in the TEC spectra at high and middle latitudes. During solar maximum, the amplitude of harmonics was twice as high, and the number of clearly manifested harmonics was more than during the solar minimum. They also found using F10.7 solar radio flux that the amplitude of all the ionospheric harmonics follows the solar activity level changes.

Numerous other studies have indicated the presence of periodic signatures in the ionosphere at various locations on Earth (Xingliang et al., 2005; Cai, 2007; Afraimovich et al., 2008; Amiri-Simkooei and Asgari, 2011; Chauhan et al., 2011; Olwendo et al., 2012; Liu et al., 2014; Tariku, 2015; Elemo et al., 2018; Ogwala et al., 2019; Akinyemi et al., 2021). Note that the majority of characterization studies have been performed in low-latitude or high-latitude regions where geomagnetic activity dominates the physical processes. At midlatitudes, tropospheric and geomagnetic phenomena compete in disturbing the ionosphere, and it is not well understood how these multiple sources affect the ionospheric drivers. Even for the relatively lower density night-time mid-latitude ionosphere the Perkins instability produces periodic variations of TEC (Perkins, 1973). The result of these many processes is that a solar cycle spectrum of TEC is composed of signals spanning small-to meso-to larger spatio-temporal scales.

2 Data and TEC calculation

Measurements from the NOAA CORS network can be obtained from the National Geodetic Survey: https:// geodesy.noaa.gov/CORS/. The data rate varies between stations but we use stations taking more than 1 sample per minute and down-sample to 1 sample per minute. The data are daily GNSS Observation files which are provided in the receiver-independent exchange (RINEX) format and must be converted to vTEC. The GPS-TEC analysis software, developed by Gopi Seemala (Ma and Maruyama, 2003; Seemala and Valladares, 2011), uses phase and code values for both L1 and L2 GPS frequencies to eliminate the effect of clock errors and tropospheric water vapor, and thus calculate relative values of sTEC (Sardón et al., 1994; Sardón and Zarraoa, 1997; Arikan et al., 2008). Then, the absolute values of sTEC are obtained by accounting for the differential code biases at the satellites (Valladares et al., 2009). The final vTEC data is then calculated by including the geometric factor. The equations used by the GPS-TEC software can be found in Ma and Maruyama. (2003). It is known that the true value of TEC is difficult to calculate due to the receiver biases, but this is largely



study, and many other CORS stations are given by small squares or circles.

affecting the DC component of the spectrum so it is not relevant to our study.

The data for this study includes 15 mid-latitude stations within 10 degrees of geographic latitude and longitude (corresponding to an area of \sim 700 mi²) of one another (see table

of geomagnetic latitude and longitude in **Figure 1**) covering solar cycle 24 from 2009 to 2019. The station in Jacksonville, Florida (ZJX1), a central location in the cluster, is at magnetic latitude \sim 40°. Stations in Fairbanks, Alaska (magnetic latitude = \sim 65°), and Eusebio, Brazil (magnetic latitude = \sim 4°) are



FIGURE 2

TEC time series (top) and Lomb-Scargle periodogram (bottom) for station ZMA1 during solar cycle 24 (1 January 2009 to 1 January 2020). Annual and semi-annual signatures can be easily identified in the time series, along with higher amplitude diurnal variations. In the spectrum a vertical pink (cyan) dashed line is overlaid at the fundamental frequency of 1/year (365/year) along with some harmonics n = ,2,3, ... The solar rotation peak also appears as a broadband source.



FIGURE 3

Lomb-Scargle periodograms for the TEC time series from ZMA1 during solar minimum (top left), rising (top right), solar maximum (bottom left), and falling (bottom right) phases of solar cycle 24.



Each blue dot is the amplitude of the 365/year signal calculated from a Lomb-Scargle periodogram for the corresponding phase of solar cycle 24 (solar minimum on top and rising phase on bottom). The amplitudes are plotted with respect to the magnetic latitude of the station. Error bars are calculated as the smallest and largest amplitudes from a sliding window periodogram over 90 days sub-intervals of each solar phase. As a general trend the amplitudes align with magnetic latitude and solar activity. Meso-scale deviations from this trend indicate the presence of meso-scale ionosphere structure.

also included to put the mid-latitude cluster of stations into larger-scale context. The magnetic latitude and longitude of each station was calculated from the IGRF-13 dipole model. A map of the mid-latitude stations is shown in **Figure 1**. All of the NOAA CORS receivers in the region are plotted as little squares while the stations in this study have large red pentagons and labels. To calculate power spectra, we use a Lomb-Scargle periodogram (VanderPlas, 2018) which has the advantage over a traditional Fourier decomposition in that the data do not have to be uniformly sampled. From N observations, x_i for $i \in [1, N]$, taken at times t_i , the Lomb-Scargle Periodogram is defined (Lomb, 1976):

$$P(f) = \frac{1}{2\sigma} \left\{ \frac{\left[\sum_{N}^{i=1} (x_i - \bar{x}) \cos(2\pi f(t_i - \tau))\right]^2}{\sum_{N}^{i=1} \cos^2(2\pi f(t_i - \tau))} + \frac{\left[\sum_{N}^{i=1} (x_i - \bar{x}) \sin(2\pi f(t_i - \tau))\right]^2}{\sum_{N}^{i=1} \sin^2(2\pi f(t_i - \tau))} \right\}$$

with $\bar{x} = 1/N\sum_{N}^{i=1}x_i$ the mean of the data and $\sigma^2 = 1/(N-1)\sum_{N}^{i=1}(x_i - \bar{x})^2$ the variance. The offset τ is chosen such that $\sum_{N}^{i=1}\cos(2\pi f(t_i - \tau))\sin(2\pi f(t_i - \tau)) = 0$. The data gaps in the time series used for this study are generally a few days or less but some stations are missing periods of a year or more. The middle panel of **Figure 1** shows the data coverage during solar cycle 24



24 (solar maximum on top and falling phase on bottom). The amplitudes are plotted with respect to the magnetic latitude of the station. Error bars are calculated as the smallest and largest amplitudes from a sliding window periodogram over 90 days sub-intervals of each solar phase. As a general trend the amplitudes align with magnetic latitude and solar activity. Meso-scale deviations from this trend indicate the presence of meso-scale ionosphere structure.

for each GPS station. We have chosen not to sort the data into disturbed and quiet intervals for two reasons. 1) Geomagnetic disturbances can occur at regular intervals when associated with a feature on the Sun that persists for longer than a solar rotation, which is a sinuosoidal signal that we want to characterize. 2) The diurnal variation of day-night photoionization is always the dominant signal for a long time series (20 + days) of TEC, even if it contains storm-time intervals.

In addition to calculating the spectrum for an entire solar cycle of data, we also sequester portions of the time series corresponding to the different phases of solar cycle 24. We have defined these intervals as shown in the bottom panel of Figure 1, which gives the sunspot number time series for 1 January 2009 to 1 January 2020. The phases are defined as Solar Minimum: 1 January 2009–15 May 2010 and 16 August 2018–1 January 2020,

Rising toward Maximum: 16 May 2010–15 February 2013, Solar Maximum: 16 February 2013–15 November 2015, Falling toward Minimum: 16 November 2015–15 August 2018.

3 TEC spectra

The Miami (ZMA1) TEC time series for 1 January 2009 to 1 January 2020 is shown in the top panel of **Figure 2**. Note that the sunspot number (bottom of **Figure 1**) peaks strongly during solar maximum from 2014 to 2015 but also has a peak of similar magnitude during the second half of 2011. On the longest visible time-scale, the TEC time series from Miami reflects the shape of the sunspot number curve. Also easy to pick out are the annual and semi-annual cycles of TEC. The diurnal variation dominates



the signal being too high frequency to easily identify in the 11year time series. However, the second panel of **Figure 2** shows the Lomb-Scargle periodogram for the TEC time series, which indicates the signal with a frequency of 365/year has the largest power at about 20 TECU. The magenta (cyan) vertical dashed lines indicate the location of a fundamental frequency at 1/year (365/year) and the corresponding harmonics for n = 2,3, ... The fundamental and first harmonic of the seasonal cycle have similar power. In the case of the diurnal cycle, many harmonics are well resolved. The fundamental dominates the harmonics in power but it is noteworthy that the first harmonic has similar power to the seasonal signal. The power in the signal at the solar rotation ~ 27 day period (~14/year, see red arrow) is broadband and lacking in discernible harmonics, because it results from multiple sources on the Sun which do not always last a full solar rotation.

Figure 3 shows Lomb-Scargle periodograms for the TEC time series from **Figure 2** split into the different phases of the

solar cycle. The diurnal cycle (365/year), first harmonic, and many more harmonics (not shown), are well resolved during all phases of the solar cycle. The solar rotation is well resolved during all phases except for solar minimum. At solar minimum, the annual signal dominates the semi-annual signal. During the rising phase, the semi-annual signal is slightly more powerful than the annual signal. At solar maximum, the annual signal is slightly more powerful than the semi-annual signal. During the falling phase, the annual signal is significantly more powerful than the semi-annual signal but not so much as during solar minimum. Splitting the data into phases of the solar cycle shows also that the frequency for "annual", "semi-annual", and "solar rotation" is not fixed in time. Only from an entire solar cycle of data are the annual and semi-annual signals present as wellresolved peaks very close to 1/year and 2/year. Later, we will discuss the solar rotation peak for the different solar cycle phases.



Ratio of fundamental diurnal vs. first harmonic amplitudes as a function of magnetic latitude during solar maximum (top) and falling phase (bottom). Error bars are calculated from 90-day sub-intervals. Errors bars approaching a ratio of 1 suggest the presence of a strong semi-diurnal driver.

We construct periodograms for each station similar to Figures 2, 3 and extract the amplitudes for the diurnal, annual, and seasonal oscillations, as well as the relevant harmonics. The power in the fundamental diurnal frequency for each station is plotted in Figures 4, 5 as a function of magnetic latitude of the station. Figure 4 shows solar minimum (top) and rising phase (bottom) and Figure 5 shows solar maximum (top) and falling phase (bottom). The error bars are computed from a 90-day sliding window (with 30-day separations) periodogram for the corresponding phase. The smallest fundamental diurnal amplitude from all sub-intervals is the value of the downwards error bar and like-wise for the upwards error bar. The error bars thus represent the variability of the signal over the corresponding solar cycle phase. The amplitude at all stations corresponds to solar activity level and magnetic latitude. If we assume the stations should fall on a straight line connecting CLGO and

BRFT, deviations from such a line represent the meso-scale ionospheric structure of interest to this study. At solar minimum and falling phase, within the mid-latitude cluster of stations the power deviates from a straight line by 1-2 TECU, while the deviation is 10-20 TECU for rising phase and solar max. It may be alarming to the reader that the error bars are much larger than these variations, but since the error bars represent the timevariation of the amplitudes, it can be expected that all the stations are having their largest or smallest amplitudes at a similar time of the year. Note also that the variation between stations is similar at the top and bottom of the error bars, compared with the variations between stations for the blue points. This is an effect of meso-scale variations (from station-to-station) superposed onto the fundamental signal which does not have a constant frequency in time. A final consideration to take into account when comparing the error bars with the blue points is that the



error bars are computed from the spectrum of a time series that is shorter than the spectrum which produced each blue point.

The diurnal signal mostly results from the day-night cycle of photoionization, and the reason to characterize this signal with high accuracy is because the amplitude is much greater than that of most isolated ionospheric structures. What we aim to predict is the location and extent of ionosphere structures associated with strong TEC gradients generated by geomagnetic disturbances originating from the Sun (i.e., storm enhanced density plume). Of course, these disturbances would not be resolved as peaks in the spectra (unless they are associated with a feature on the Sun that lasts longer than a solar rotation), but by knowing what the diurnal amplitude is at any station and any time, it becomes possible to subtract out this component of the signal and have remaining only the variations due to meso-scale ionospheric structures.

A sequence of harmonics in the spectra can be indicative of a signal that is more fundamentally a square wave than a sinusoid. This is an expected feature of the ionospheric TEC because the steady-state ionospheric density is determined by the balance of ionization and recombination. At night-time photoionization ceases but there is still a residual ionization from cosmic rays Webber. (1962), so the recombination rate will decrease throughout the night until it is in balance with the cosmic-ray ionization rate. To identify a square wave, the spectrum must have a series of harmonics with amplitudes that decrease with *n*. Figures 6, 7 illustrate the presence of mesoscale ionosphere structure based on the ratio of the amplitudes in the fundamental vs n = 1 harmonic, in the same way as Figures 4, 5. In Brazil, the amplitude ratio maximizes at solar maximum and minimizes at solar minimum, and, surprisingly (because it is a log scale plot), has a larger upward error bar



for all phases except the falling phase. In Alaska, the amplitude ratio is mostly constant throughout the solar phases, having the smallest variation from minimum to maximum of all stations in this study. Quantitatively, the standard deviation of amplitude ratio across the 4 solar phases is 1.8 in Alaska, 7.8 in Brazil, and an average of 13.7 from all mid-latitude stations in this study. Similar to low-latitude, at middle-latitude the amplitude ratios correspond well with solar activity (higher ratio for more solar activity). Averaged across all mid-latitude stations the amplitude ratio is 19.1 at solar minimum, 43.9 in rising phase, 53.3 at solar maximum, and 31.9 in falling phase. Note also that from stationto-station at middle latitudes the ratio can vary by 10 or greater during all solar phases, similar to the variations demonstrated in Figures 4, 5. Additionally, during the solar minimum and falling phases, a few of the downward error bars approach 1. This is not the result of a square wave but suggests at these stations there

is a driver at a frequency of twice per day rather than being a harmonic of the diurnal signal, potentially the semi-diurnal atmospheric tide (Oberheide et al., 2015).

For the 1/year signal derived from a periodogram of the entire solar cycle time series only a single harmonic is well resolved and the amplitude ratio between fundamental and first harmonic fluctuates about 1. This indicates a competition between drivers with frequencies of 1/year and 2/year, which are otherwise known as the annual and semi-annual cycles. Figure 8 quantifies the mid-latitude meso-scale structure of these different drivers. The top panel shows the amplitude of the annual signal as a function of magnetic latitude and the bottom panel shows the ratio of annual to semi-annual amplitudes. The error bars are constructed the same as Figure 4 but using a 3-year sliding window. From the top panel, there is a similar general dependence on magnetic latitude as for the fundamental



diurnal signal and also a station-to-station deviation that is maximally 1-2 TECU from a power law fit. The bottom panel contains information about the competition of drivers with 1/year and 2/year frequencies. Within the cluster of mid-latitude stations the ratio of periodogram amplitudes from the entire cycle solar varies from less than 1 in some locations to greater than 1 in others. The error bars also indicate many of the stations go from a ratio that is less than 1 to greater than 1 or *vice versa* throughout the course of the solar cycle. In Brazil and Alaska, the error bars show the ratio is never greater than 1. The source of this meso-scale mid-latitude structure in the spectra is not the result of transient solar driving, as those effects occur on much shorter time scales, but must result from the interaction of the ionosphere with the neutral atmosphere and ground.

Finally, **Figures 9**, **10** quantify the solar rotation amplitude during the different solar cycle phases. Because of the broadband nature of the solar rotation signal, it is intrinsically more difficult

to quantify. We have taken the largest amplitude from a frequency window centered at 1/26.5 = 0.038 days⁻¹ (labeled as solar rotation in Figure 3) with a width 0.01 days⁻¹. For reference, the reader may note in the top left panel of Figure 3 there is a strong peak to the left of the solar rotation peak at a frequency of 1/45 days⁻¹, which is outside of the tolerance window. The error bars in Figures 9, 10 use a 360-day sliding window. The fact that many of the lower error bars are very small is due to the different sources on the Sun, since each sinuosoidal driver will in this case exist for a much shorter time than an entire solar cycle phase. Similar to the diurnal frequency, the amplitudes are generally aligned with solar phase and magnetic latitude. At midlatitude the typical station-to-station deviation is about 0.1-0.2 TECU during solar maximum and rising phases and just slightly smaller during solar minimum and falling phases. The driver for this signal, being mostly coronal hole high speed streams and co-rotating interaction regions, results in a broadband signal because these features are appearing and disappearing on the Sun. Since the solar rotation signal is lower frequency than the diurnal signal but higher frequency than the annual signal, the ionospheric response on these time scales is in between that produced by transient solar wind drivers and the interaction of ionosphere with neutral atmosphere and ground.

4 Discussion and conclusion

The strong diurnal signal of ionospheric TEC is mainly due to Earth's rotation. At mid-latitudes, the amplitude of this signal, in addition to the seasonal and solar rotation frequencies, were found to vary on a scale of a hundred to a few hundreds of kilometers. In order to build a predictive capacity for regionalscale TEC at middle-latitudes, which is important at locations such as the Space Coast of Florida or near large airports, it is therefore necessary that whatever model employed be able to reproduce these spectral responses. Without a good characterization of the strong diurnal signal there is no way to predict the strong gradients of TEC which can develop over the Southeastern United States Foster and Rideout. (2007).

Compared to the higher latitude station in Alaska and the lower latitude station in Brazil, the deviation from expected diurnal amplitude (i.e., a straight line fit with respect to magnetic latitude) at mid-latitudes is of the order of a few TECU for each phase of the solar cycle. The number and amplitudes of the diurnal harmonic signals indicate the daily variation is not a perfect sine wave but partially a square wave. This is due to the fact that the night-time recombination rate falls until it reaches the residual ionoziation rate from cosmic rays, defining the trough of a square wave. Our results show that the shape of the diurnal signal varies from station to station and throughout the different phases of the solar cycle. In addition, some stations suggest a significant influence from a twice-daily driver. A key result from the analysis of the diurnal harmonics is that in Alaska the ratio of fundamental and first harmonic amplitudes is constant throughout the solar cycle, while at middle and low latitudes the amplitude ratios have a correspondence with solar activity. Although this could have been expected because the ionospheric drivers are different at those magnetic latitudes, this is the first study to systematically characterize mid-latitude variations of TEC spectral amplitudes for a cluster of observatories with appropriate separation to resolve meso-scale structures.

Based on a ratio of amplitudes that fluctuate around 1 with time, we find that the annual and semi-annual drivers of ionospheric TEC exhibit a competing dominance that varies from station-to-station at mid-latitudes. Olwendo et al. (2016) explain that during equinoxes the sub-solar point is near the equator, where the eastward electric field is often larger, and this intensifies the fountain effect which controls ionization levels at the equator. Near solstices, the subsolar point moves to higher

latitudes, and the fountain is expected to wane at the equator. The peak of ionization occurring when solar radiation hits the ionosphere most perpendicularly throughout the year would be at the summer solstice in high latitudes and at the equinoxes at the equator. Therefore, at mid-latitudes we expect the source of strong annual and semi-annual drivers to be a competition of these two effects. In Alaska and Brazil the ratio can approach but always remains greater than 1, indicating the competition of processes plays out less dramatically than at mid-latitudes, i.e. ionosphere TEC is more strongly controlled by a single process.

The mid-latitude meso-scale station-to-station variations of spectral amplitudes (which represent meso-scale spatial structures) and their ratios results from the complex interaction of solar wind-magnetosphere-ionosphere as well as the neutral atmosphere. This interaction can be a source of ionospheric structure for both short (1-10 s of hours) and longer (months to years) time scales. This is further complicated by the fact that the magnetospheric state is determined by instantaneous as well as the time-history of solar wind conditions Lavraud et al. (2006); Borovsky and Valdivia. (2018). It is also the case that space climate is coupled to the neutral atmosphere climate. For instance, transequatorial winds can change the recombination rate by lifting the ionosphere on either side of the equator (Olwendo et al., 2016). The neutral composition has also been shown to have effects on ionospheric storms (Liou et al., 2005). The meso-scale structure of ground conductivity may also play a role in the ionospheric response. The ground conductivities between the different mid-latitude stations in this study vary between 1 and 8 millimhos/meter. Although the more important factor may be proximity to the ocean, where conductivity is close to 5,000 millimhos/meter. However we have not attempted to classify any of these specific effects in this study, only to characterize what are the sum total sinuosoidal responses of the ionosphere due to the combination of driving forces. In this way, it may not be necessary for a predictive model to include the physics of the many competing affects that drive regular TEC oscillations in the ionosphere, just the observed spectral amplitudes. It is also worth noting that for the solar rotation signal a predictive model would benefit from additional information beyond the spectral amplitude, such as the number of sunspots on the solar disc, for example.

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Data availability statement

Publicly available datasets were analyzed in this study. This data can be found here: National Geodetic Survey https://geodesy.noaa.gov/CORS/.

Author contributions

BB wrote the manuscript and advised student RC on the data analysis. RC analyzed the data and produced Figures. KN provided funding and help revising the manuscript. XM provided the concept idea for this paper and helped revising the manuscript. SD provided expertise on ionospheric physics.

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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