#### **Research Article**

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# Impacts of the gradient-index crystalline lens structure on its peripheral optical power profile

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Abstract: The crystalline lens makes an important contribution to the peripheral refraction of the human eye, which may affect the development and progression of myopia. However, little has been known about the peripheral optical features of the crystalline lens and its impacts on the peripheral ocular refraction. This study aims to investigate the relationship between the structural parameters of the crystalline lens and its peripheral power profile over a wide visual field. The peripheral power profile is defined with respect to the entrance and exit pupil centers along the chief rays. Analysis is performed by three-dimensional ray tracing through the gradient refractive index (GRIN) lens models built from measurement data. It has been found that the vergence of the wavefronts at the entrance and the exit pupil centers of the lens show an approximate linear correlation to each other for each field angle. The exponent parameters of the axial refractive index profile and the axial curvature profile, and the asphericity of the posterior lens surface are found to be the most influential parameters in the peripheral power profiles. The study also shows that there can be significantly different, sometimes unrealistic, power profiles in the homogeneous lens model compared with its corresponding GRIN model with the same external geometry. The theoretical findings on the peripheral lens properties provide a new perspective for both wide-field eye modelling and the design of intraocular lenses to achieve normal peripheral vision.

**Keywords:** crystalline lens; peripheral defocus; peripheral optical power.

## 1 Introduction

In the past decades, the peripheral refractions of the human eye have attracted much attention due to its potential impacts on the progression of myopia [1–3]. As the most complex ocular component, the crystalline lens is a major contributor to the profiles of the peripheral refractions. The lens grows rapidly before adulthood with changes in both structural and optical properties [4, 5]. Meanwhile, the lens structure changes dynamically during accommodation for near vision, which is also a potential factor for the development of myopia [6, 7]. These changes in the lens structure can affect the peripheral ocular refractions and thus may influence the progression of myopia. Therefore, to understand the optical mechanism of myopia, it is fundamental to understand how the lens structure contributes to the distribution of the peripheral refractions of the human eye.

Very few studies have focused on the peripheral optical features of the crystalline lens. Up until now, only one experimental study was found to have measured the peripheral defocus profile of the *in vitro* human lenses [8], but the interpretations of the measured data lacked a further understanding in terms of their relationship with the peripheral ocular refractions. Meanwhile, most of the previous crystalline lens models were built based on a limited number of parameters [9–13]. This results in the restricted capability of the lens model in predicting the peripheral ocular features while maintaining a realistic anatomic structure. Consequently, little is known about the relationship between the lens structure and its peripheral optical properties.

As a further step from our previous work on the physiology-like crystalline lens (PCL) model [14] that was built to reproduce the peripheral physiologic structure of the natural lens, this study aims to investigate the relationship between the lens structure and its peripheral power profile, identifying the impacts of the external lens geometry and the internal gradient refractive index (GRIN)

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structure. This work provides a new perspective for the wide-field eye modelling and the design of intraocular lenses for a natural peripheral visual quality.

## 2 Theoretical approaches

#### 2.1 Establishment of the lens models

The lens models analysed in this study were developed from the PCL model [14], with the posterior external surface patches replaced by one fourth-order polynomial surface to achieve continuity of innumerable derivatives across the entire surface. The modification is based on the fact that a wide zone of the posterior lens surface covers the pathway of the rays for the peripheral visual field. Continuity of the surface geometry ensures the reproduction of realistic peripheral power profiles. As shown in Figure 1, the external lens surface is represented by

$$w^{2} = x^{2} + y^{2}$$

$$= \begin{cases} 2(z + T_{a})R_{a} - (1 + Q_{a})(z + T_{a})^{2}, & -T_{a} \le z \le z_{0} \\ C_{1a}z^{3} + C_{2a}z^{2} + A^{2}, & z_{0} \le z \le 0 \\ 2(-z + T_{p})R_{p} - (1 + Q_{p})(-z + T_{p})^{2} + B_{3p}(-z + T_{p})^{3} \\ + B_{4p}(-z + T_{p})^{4}, & 0 \le z \le T_{p} \end{cases}$$
(1)

where  $C_{1a}$ ,  $C_{2a}$ ,  $B_{3p}$ , and  $B_{4p}$  are derived by the boundary conditions of smooth (first derivative) connection with the adjacent patch. The equations for the internal iso-indicial surfaces are

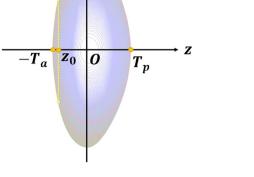
$$\begin{cases} \frac{z^2}{\left(T_a \times \sqrt[p]{\frac{n_{s,e} - n_0}{n_{s,e} - n_0}} \times \gamma\right)^2} + \frac{x^2 + y^2}{(D/2)^2 \gamma^{q+1}} = 1, \ z \le 0 \\ \frac{z^2}{\left((T - T_a) \times \sqrt[p]{\frac{n_{s,e} - n_0}{n_{s,e} - n_0}} \times \gamma\right)^2} + \frac{x^2 + y^2}{(D/2)^2 \gamma^{q+1}} = 1, \ z \ge 0 \end{cases}, \text{ where } \gamma = \sqrt[p]{\frac{n - n_0}{n_{s,e} - n_0}}. \tag{2}$$

The meanings and settings of the lens parameters in Eqs. (1) and (2) can be seen in Table 1. All the values were determined from the measurement data on the human eye. The optical contribution of each parameter is analysed separately within the 'range for investigation' as shown in Table 1, based on the variation among the population.

#### 2.2 Definition of the peripheral lens power

In Gaussian optical theory, the optical power  $\Phi$  is defined for the paraxial region along the optical axis of a rotationally symmetric optical system,

$$\Phi = \frac{n_2}{l_{P2}} - \frac{n_1}{l_{P1}},\tag{3}$$



w

A = D/2

**Figure 1:** Diagram for the modified PCL model with the Z axis pointing towards the retina.

where  $n_1$  and  $n_2$  are the refractive indices for the object and image space, respectively,  $l_{P1}$  is the axial distance from the primary principal point (P1) to the object point, and  $l_{P2}$  is the axial distance from the secondary principal point (P2) to the image point. Based on Eq. (3),  $\Phi$  can be interpreted as the vergence of the emergent wavefront at P2 subtracted by the vergence of the incident wavefront at P1. The difference in the vergence of the wavefronts at these two points is always constant within the paraxial region. However, the paraxial condition can be rarely met in the peripheral visual field of the lens. Thus, new definitions need to be developed for the peripheral optical power.

To understand the contribution of the lens to the peripheral ocular refractions, the definition of the peripheral optical power of the lens should have a direct relationship with the peripheral refraction of the eye, which is defined along the chief rays passing through the pupil center. Therefore, the peripheral optical power of the lens—here named as the 'pupil power',  $\Phi_{pupil}$ —is defined as the difference in the wavefront vergence between the exit and entrance pupil centers. As shown in Figure 2, for a chief ray in the aqueous humor,

$$\Phi_{pupil}(\theta_{aq}, \nu_{in}) = \nu_{out}(\theta_{aq}, \nu_{in}) - \nu_{in}, \qquad (4)$$

where  $\theta_{aq}$  is the angular interval between the chief ray and the lens axis;  $v_{out}$  and  $v_{in}$ , respectively, represent the vergence of the output and input wavefronts at either the exit pupil center or the entrance pupil center, depending on the direction of light propagation. In the definition of the ocular aberrations, the rays are traced from the vitreous

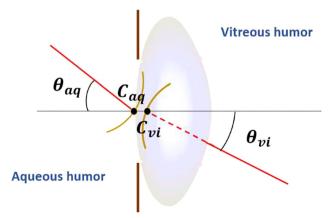
Parameter	Meaning of the parameter	Value	Range for investigation	Source of data <sup>c</sup>
R <sub>a</sub> (mm)	Radius of curvature at the anterior lens vertex	11.580	11-15	Mutti et al. [4]
<i>R<sub>p</sub></i> (mm)	Radius of curvature at the posterior lens vertex	6.303	6-8.5	Mutti et al. [4]
Qa	Asphericity of the anterior external surface	$4.6 \times 10^{-4}$	-22 to 5	Ishii et al. [15], Dubbelman et al. [16]
$Q_p$	Asphericity of the posterior external surface	$2.49\times10^{-4}$	-15 to 3	Ishii et al. [15], Dubbelman and Van der Heijde [16]
<i>T</i> (mm)	Axial lens thickness	3.421	2.9-4.1	Mutti et al. [4]
$T_a/T$	Ratio of axial thicknesses between the anterior lens section and the total lens thickness	0.45	0.4-0.52	Martinez-Enriquez et al. [17], Ishii et al. [15]
D	Lens diameter	8,574	8-8.8	Ishii et al. [15]
$n_0^{b}$	Refractive index at the lens center	1.4012	1.4-1.41	Khan et al. [18]
$n_{s_a}^{b}$	Refractive index at the anterior lens vertex	1.3663	1.36-1.38	Khan et al. [18]
$n_{s_p}^{b}$	Refractive index at the posterior lens vertex	1.3801	1.376-1.382	Khan et al. [18]
$n_{s_e}^{b}$	Refractive index at the equatorial edge	1.3560	1.34-1.362	Khan et al. [18]
р	Exponent parameter for the axial index profile	3.3120	2-4.2	Khan et al. [18]
q	Exponent parameter for the axial curvature profile	1.494	1.2-2.2	Set to fit the paraxial lens power measured by Mutti et al. [4]
θ <sub>a</sub> (°)	Subtended angle to the lens center of the conic zone on the anterior external surface	120	Not available	Set to cover the optical zone

Table 1: Lens parameters<sup>a</sup> for the physiology-like crystalline lens model of age 11 years and the range for investigation.

<sup>a</sup>Elaboration of the parameters can be seen in the article by Li and Fang [14]. <sup>b</sup>All the refractive index parameters are referenced to the wavelength of 589 nm. <sup>c</sup>Source of data for the lens parameters is based on the average and range of data measured in the population, which were applied in this study for investigating their impacts on the peripheral lens power.

humor throughout the lens, thus  $v_{out}$  and  $v_{in}$  correspond to the wavefronts at the iris center ( $C_{aq}$ ) and its conjugate point in the vitreous humor ( $C_{vi}$ ), respectively. It is expected that  $\Phi_{pupil}$  may vary with  $v_{in}$ ; hence  $\Phi_{pupil}$  is formulated as a function of both  $\theta_{aq}$  and  $v_{in}$ .

The vergence of a wavefront is essentially the multiplication of the refractive index and the local curvature of the wavefront surface. In three-dimensional ray tracing, the wavefront surface is often not rotationally



**Figure 2:** Diagram for defining the peripheral pupil power of the lens model along a chief ray (the red line) passing through the entrance and exit pupil centers  $C_{aq}$  and  $C_{vi}$ .

symmetric. As defined in differential geometry, the curvature at a point on the surface can be described by two local principal curvature values—the maximum and minimum curvature ( $\kappa_{max}$  and  $\kappa_{min}$ ) along the two perpendicular directions [19]. Hence, the vergence of a wavefront also has two components, namely the spherical equivalent vergence ( $v_{SE}$ ) and the plus cylinder vergence ( $v_{Cyl}$ ), as shown below:

$$v_{SE} = n_{wave} \cdot \left(\frac{\kappa_{max} + \kappa_{min}}{2}\right), \tag{5.a}$$

$$v_{Cyl} = n_{wave} \cdot (\kappa_{max} - \kappa_{min}), \qquad (5.b)$$

where  $n_{wave}$  is the refractive index of the medium for the wavefront. Similarly,  $\Phi_{pupil}$  can be divided into two components—the spherical equivalent pupil power ( $\Phi_{pupil-SE}$ ) and the plus cylinder pupil power ( $\Phi_{pupil-Cyl}$ ), which are here defined as

$$\Phi_{pupil-SE} = v_{out-SE} - v_{in-SE}, \qquad (6.a)$$

$$\Phi_{pupil-Cyl} = v_{out-Cyl} - v_{in-Cyl}$$
(6.b)

In the definition of the ocular aberrations, the input wavefront at  $C_{vi}$  is emitted from the point on the retina and is thus always spherical. Hence,  $v_{in-Cyl} = 0$  and  $\Phi_{pupil-Cyl}$  is essentially the plus cylinder vergence of the wavefront emerged at  $C_{aq}$ . It should be mentioned that the sign of  $v_{Cyl}$  is

always positive, while the sign of  $v_{SE}$  follows the convention that a positive value corresponds to a converging wavefront.

#### 2.3 Optical analysis procedure

In this study, the peripheral power profiles of the lens models as defined in Table 1 were obtained by threedimensional ray tracing through a set of self-developed MATLAB programs. The detailed procedure is listed as follows.

(1) Finding  $C_{vi}$  of the lens model

In this study,  $C_{aq}$  is assumed to coincide with the anterior lens vertex. By tracing rays in the aqueous humor entering the lens that all pass through  $C_{aq}$ , the wavefront emerged out of the posterior lens surface can be derived as a series of points on the rays with the same optical path length from  $C_{aq}$ . The exit pupil center  $C_{vi}$  is then computed as the center of the sphere-fit to this wavefront.

(2) Locating the chief rays

For a given set of field angles  $\theta_{aq}$ , the chief rays can be determined by ray tracing throughout the lens from the aqueous humor to the vitreous humor.

(3) Obtaining the relationship between  $\Phi_{pupil}$  and  $v_{in-SE}$  for each field angle

Along the sections of the chief rays in the vitreous humor, a set of object points were located by a predefined set of  $v_{in-SE}$ , which determine the distances from  $C_{vi}$  with the given refractive index of the vitreous humor. Then, threedimensional ray tracing was performed for each object point separately to obtain the wavefront emerged out of the anterior lens surface at  $C_{aq}$ . The wavefront was fitted to the Zernike polynomials up to the 6th order. The vergence parameters— $v_{out-SE}$  and  $v_{out-Cyl}$ —of the emerging wavefront in the aqueous humor with the refractive index of  $n_{aq}$  were then calculated based on differential geometry and the power vector notation, as [20, 21]

$$v_{out-SE} = \frac{1}{R^2} \cdot \left(4\sqrt{3} c_2^0 - 12\sqrt{5} c_4^0 + 24\sqrt{7} c_6^0\right) \cdot n_{aq}, \quad (7.a)$$

$$v_{out-Cyl} = 2\sqrt{J_0^2 + J_{45}^2},$$
 (7.b)

$$J_0 = -\frac{1}{R^2} \left( 2\sqrt{6} c_2^2 - 6\sqrt{10} c_4^2 + 12\sqrt{14} c_6^2 \right) \cdot n_{aq}, \qquad (7.c)$$

$$J_{45} = -\frac{1}{R^2} \left( 2\sqrt{6} c_2^{-2} - 6\sqrt{10} c_4^{-2} + 12\sqrt{14} c_6^{-2} \right) \cdot n_{aq}$$
(7.d)

Here *R* is the semi-diameter of the wavefront;  $c_n^m$  are the Zernike coefficients;  $J_0$  and  $J_{45}$  are the horizontal/vertical and the oblique power vectors of astigmatism in the power notation, respectively [20]. Note that Eq. (7.b) for  $v_{out-Cyl}$  was derived in a similar way to the derivation of the plus cylinder power of the wavefront refraction of the human eye [20, 22]. In this study, the discussion of the crystalline lens is focused on a rotationally symmetric structure. For the ease of formulation, the peripheral lens power is analysed along the vertical meridian, in which case

$$v_{out-Cyl} = 2|J_0| \tag{8}$$

As shown in the section below, a mathematical pattern can be easily observed between  $J_0$  and  $v_{in}$ , because there is no restriction on the sign of  $J_0$ . Thus, the trendlines of  $J_0$ with respect to  $v_{in-SE}$  and  $\theta_{aq}$  were calculated first, while  $\Phi_{pupil-Cyl}$  can be derived afterwards by Eq. (8).

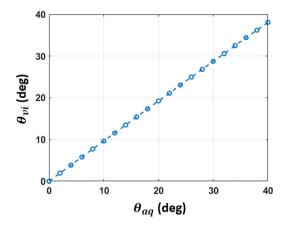
For all the calculations involved in this study, the wavefront diameter is set as 3 mm and the GRIN structure of the lens model is approximated by 200 iso-indicial layers. The refractive indices of the aqueous humor and the vitreous humor are set as 1.333. Around 5000 rays are traced throughout the lens. The boundary of the ray bundle is set wide enough so that the 1.5 mm-radius entrance pupil can be fully covered by the rays even at the most peripheral field angle.

## **3 Results**

#### 3.1 Evaluation of the lens pupil power

Based on the method described above, the peripheral lens power profile was computed on the 11-year-old lens model defined in Table 1.  $C_{vi}$  was calculated at 0.043 mm behind  $C_{aq}$  along the lens axis, suggesting that the iris center and its conjugate point almost coincide with each other. Meanwhile, it has been found that  $\theta_{vi}$  is very close to  $\theta_{aq}$  for each chief ray, as can be seen in Figure 3. Interestingly,  $\theta_{vi}$ is approximately linear to  $\theta_{aq}$  with the slope of 0.95. The *R*-squared value for the linear fit is 0.9998.

The changes in the lens pupil power with respect to the absolute value of  $v_{in-SE}$  are shown in Figure 4. All the trendlines appear approximately linear, which reveals a unique feature of the lens that can be applied to formulate the lens power. The maximum deviation of the linear fit to the relationship between  $\Phi_{pupil-SE}$  and  $v_{in-SE}$  is below 0.025 D, while the maximum deviation for the linear fit between  $J_0$  and  $v_{in-SE}$  is below 0.015 D. Moreover, both



**Figure 3:** The change in the vitreous field angle  $\theta_{vi}$  with respect to the aqueous field angle  $\theta_{aq}$  for the chief rays.

 $\Phi_{pupil-SE}$  and  $J_0$  tend to decrease with larger magnitude of  $v_{in-SE}$  for the peripheral visual field. Meanwhile, for the same value of  $v_{in-SE}$ , both  $\Phi_{pupil-SE}$  and  $\Phi_{pupil-Cyl}$  tend to be larger for the more peripheral field, and the steepness of the trendlines increases with larger eccentricity. This indicates that the peripheral pupil power is more sensitive to  $v_{in-SE}$  of larger eccentricity.

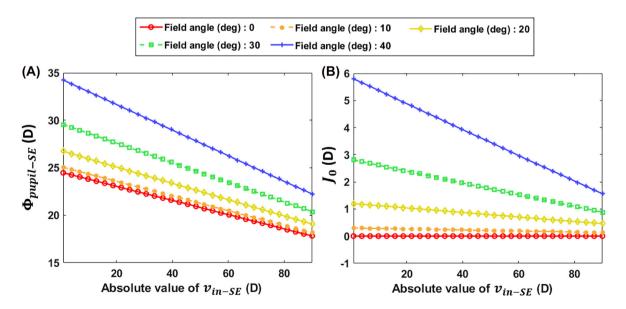
#### 3.2 Impacts of the lens structure

To understand how the structural features of the lens affect its peripheral power profile, 13 major lens parameters were investigated within the intervals measured on the population as listed in Table 1. In addition, the role of the gradient index distribution in the power profiles was also analysed, by comparing each lens model with its corresponding homogeneous model having the same external geometry. The refractive index of the homogeneous model is the equivalent refractive index ( $n_{eq}$ ) of the GRIN model, which is computed by its paraxial optical power:

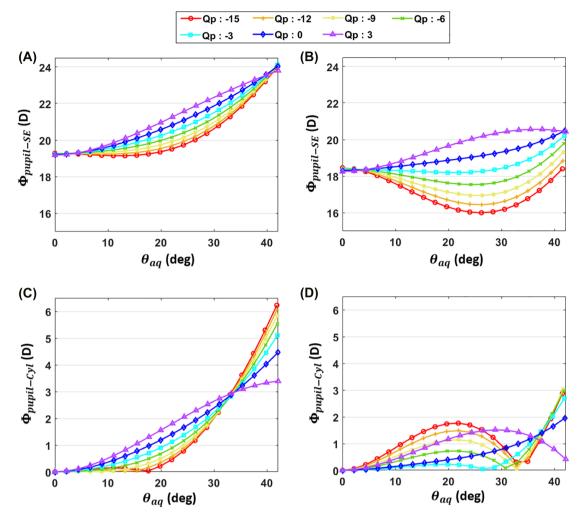
$$P_{paraxial-GRIN} = \frac{n_{eq} - n_{aq}}{R_a} + \frac{n_{vi} - n_{eq}}{R_p} - \frac{T}{n_{eq}} \cdot \frac{n_{eq} - n_{aq}}{R_a}$$
$$\cdot \frac{n_{vi} - n_{eq}}{R_p}$$
(9)

To understand the situation in the eye, the peripheral lens power profile was analysed within a test eye model defined by a spherical retinal contour with a radius of curvature of 12 mm and an axial distance of 19 mm from  $C_{vi}$ . Since  $C_{vi}$  is quite close to the anterior lens vertex for all the lens models, the vitreous chamber depth is around 15–16 mm, which is within the range measured on children [23]. This leads to the value of  $v_{in-SE}$  decreasing from around –70 D at the central field to –81 D at the most peripheral visual field ( $\theta_{vi} = 40^\circ$ ).

The results for Qp are displayed in Figure 5. For most peripheral field locations of the GRIN models, both  $\Phi_{pupil-SE}$ and  $\Phi_{pupil-Cyl}$  increase with eccentricity and tend to be larger with higher Qp. The trendlines of  $\Phi_{pupil-Cyl}$ , in contrast, show a steeper increase with lower values of Qp at  $\theta_{aq}$  >33°. In contrast, a reduction in both  $\Phi_{pupil-SE}$  and



**Figure 4:** Change of  $\Phi_{pupil-SE}$  (A) and  $J_0$  (B) with respect to the absolute value of  $v_{in-SE}$ .



**Figure 5:** Peripheral power profiles of the 11-year lens model calculated in the test eye model with different values of Qp, including: (1)  $\Phi_{pupil-SE}$  for the gradient refractive index (GRIN) (A) and homogeneous (B) models; (2)  $\Phi_{pupil-Cyl}$  for the GRIN (C) and homogeneous (D) models.

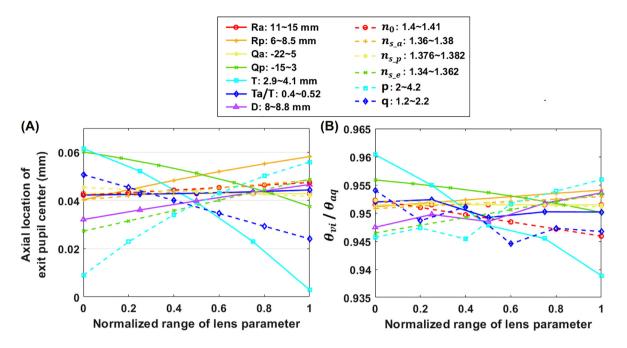
 $\Phi_{pupil-Cyl}$  with larger value of  $\theta_{aq}$  can be observed in certain field ranges of the homogeneous models, and this pattern is more evident for the lower values of Qp. Comparison between the GRIN and the homogeneous models indicates that GRIN structure (1) contributes positively to the peripheral profile of  $\Phi_{pupil-SE}$  and (2) induces opposite values of  $J_0$  compared with the external geometry. These results suggest that having a realistic paraxial optical power alone does not ensure the lens model to have a realistic peripheral power profile.

The peripheral power profiles for the other lens parameters were also analysed by the same procedure. Overall, the distance between  $C_{vi}$  and  $C_{aq}$  is less than 0.06 mm, as can be seen in Figure 6(A). Meanwhile,  $\theta_{vi}$  is approximately linear to  $\theta_{aq}$ , with the value of  $\theta_{vi}/\theta_{aq}$  between 0.94 and 0.96 (Figure 6(B)). These results show that, very interestingly, the sections of the chief rays in the

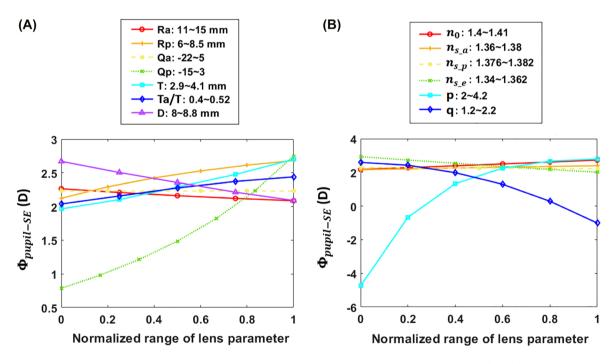
aqueous and vitreous humor are almost collinear to each other.

At around 25° of  $\theta_{aq}$  in the test eye model, the change in the relative  $\Phi_{pupil-SE}$  (peripheral value of  $\Phi_{pupil-SE}$  minus the central value of  $\Phi_{pupil-SE}$ ) with respect to the investigated range of the GRIN lens parameters (as stated in Table 1) are shown in Figure 7. As expected, the value of Qa does not influence the peripheral lens power. For the investigated range of each parameter, the exponent parameters p and qhave the largest effects on the peripheral lens power, followed by Qp and  $n_{s\_e}$ . Furthermore, it can be seen that the increases in q,  $n_{s\_e}$ , Ra and D are associated with a decrease in the relative  $\Phi_{pupil-SE}$ , in contrast with the other lens parameters that are positively correlated with the relative  $\Phi_{pupil-SE}$ .

The results for  $\Phi_{pupil-Cyl}$  can be seen in Figure 8. Similarly, the variation of  $\Phi_{pupil-Cyl}$  is the largest for the



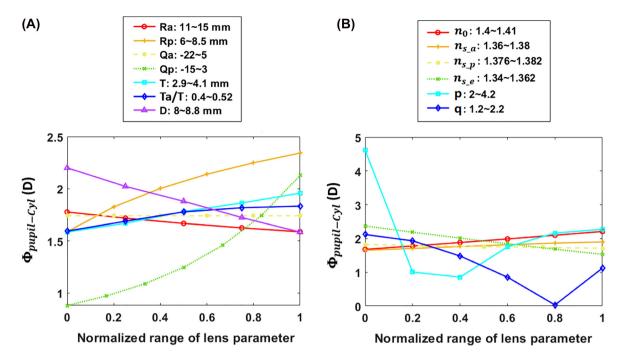
**Figure 6:** Trendlines of the distance between  $C_{aq}$  and  $C_{vi}$  (A) and the field ratio  $\theta_{vi}/\theta_{aq}$  (B) with respect to the normalised range of the lens parameter for all the investigated gradient refractive index lens models.



**Figure 7:** Change of the relative  $\Phi_{pupil-SE}$  at around 25° of  $\theta_{aq}$  with respect to the gradient refractive index (GRIN) lens parameters related to the external lens geometry (A) and the GRIN distribution (B).

investigated intervals of *p* and *q*, while the trendlines of the other lens parameters present a similar pattern to the trendlines of the relative  $\Phi_{pupil-SE}$ . Namely, a larger value of

the relative  $\Phi_{pupil-SE}$  is often associated with a larger value of  $\Phi_{pupil-Cyl}$  regardless of the variations in the lens structure.



**Figure 8:** Change of  $\Phi_{pupil-Cyl}$  at around 25° of  $\theta_{aq}$  with respect to the lens parameters related to the external lens geometry (A) and the GRIN distribution (B).

## 4 Discussion

This study proposed a new way of describing the peripheral optical power of the crystalline lens-as the lens-induced change in the wavefront vergence between  $C_{aq}$  and  $C_{vi}$ along the chief rays. In this way, the contribution of the ocular components (lens, cornea and retina) to the peripheral refraction of the eye can be described separately and quantitatively with mathematical rigor. The ocular refraction is often equivalent to the vergence of the wavefront at the entrance pupil center of the eye, which is also the conjugate point of  $C_{aq}$  in the air. The wavefront has to be emitted from the point on the retina and propagates along the chief ray throughout the eye. Accordingly, the contribution from the retina can be represented by the vergence of the wavefront at  $C_{vi}$ , while the impact of the cornea and its location relative to the lens can be defined as the difference in the wavefront vergence between  $C_{aq}$  and the entrance pupil center of the eye.

As found by ray tracing through all the investigated lens models, the peripheral lens power (in terms of both  $\Phi_{pupil-SE}$  and  $J_0$ ) at the same visual field presents a highly linear correlation with the wavefront vergence at  $C_{vi}$  for all the field angles. Due to the complexity of the GRIN structure, an analytical explanation of this phenomenon cannot be easily derived. However, calculations of all the lens models show that  $C_{vi}$  is close to  $C_{aq}$  and  $\theta_{aq}$  is close to

 $\theta_{vi}$  for the same chief ray. These findings indicate that the chief rays on both sides of the lens model are almost collinear to each other—having the feature of the optical axis defined in the paraxial optics. In the paraxial optical theory, the wavefront vergence at two conjugated points on the optical axis are always linear to each other [24], which could partially explain the linear relationship between  $\Phi_{pupil-SE}$  and  $v_{in-SE}$ .

The proposed approach for analysing the peripheral lens power was then applied to evaluate the contribution of each lens structural feature within the range measured on the population. In this study, 13 lens parameters were examined based on a 11-year-old lens model constructed from measurement data. In particular, we found that the asphericity of the posterior lens surface plays a significant role in the power profile. This finding is helpful for widefield eye modelling. Since Qp has nearly no impacts on the paraxial lens power, it can be adjusted to reproduce the targeted peripheral refraction profile. Furthermore, this study shows that the peripheral lens power changes the most with the variation of the exponent parameters for the axial index profile (p) and the curvature profile (q) within the examined interval, followed by Qp and  $n_{s_e}$ . These parameters should be given special attention in the construction of the accommodative eye models that aim for reproducing the measured change in the peripheral ocular refraction during accommodation.

The role of the GRIN structure as a whole has been investigated by comparing the GRIN and the homogeneous models of the same external geometry. Although both types of the lenses share the same value of the paraxial optical power, there can be large deviations in the peripheral optical profiles. Our analysis has shown that for lower values of Qp, negative correlations of  $\Phi_{pupil-SE}$  and  $\Phi_{pupil-Cyl}$  with respect to  $\theta_{aq}$  are present in the homogeneous models, which are very rare for the GRIN models. This finding is especially helpful for building the physical eye models, where the crystalline lens is often modelled by a homogeneous lens. To achieve a realistic distribution of the peripheral refraction, the external geometry of the lens model should be manually modified rather than following the external geometry measured on the real lenses.

One limitation of this study is the assumption of a rotationally symmetric lens model, while studies have found some degree of the axial astigmatism in the human lens [25]. In such case, the wavefronts propagating along most of the peripheral chief rays will be distorted due to the torsion in the refracting surfaces. Moreover, the iris center is assumed to locate at the anterior lens vertex, while a slight decentration of the iris center with respect to the lens can exist in many eyes and during pupil constriction. Although these situations are insignificant in most eyes and thus ignored in many eye models, more measurement data of the peripheral optical and structural features of the human lens are needed to develop a further understanding.

## 5 Conclusions

This study proposed a new method for describing the peripheral optical power profiles of the crystalline lens model for the entire visual field. Based on this method, the relationship between the lens structure and its peripheral power profile was systematically investigated on a series of lens models constructed from measurement data. It has been found that the vergence of the wavefronts at the entrance and the exit pupil centers of the lens shows an almost linear correlation to each other for each field location. This can be partly explained by the high degree of collinearity between the sections of the chief rays in the vitreous and the aqueous humor as found in all the lenses. Among the 13 lens parameters, *p*, *q*, *Qp*, and  $n_{s_e}$  were found to have the largest impacts on the peripheral power profiles. The study also shows that there can be significantly different, sometimes unrealistic, power profiles in the homogeneous lens model compared with its corresponding GRIN model with the same external geometry. These findings can be helpful for building eye models and the design of intraocular lenses that aim to reproduce the peripheral ocular refraction profiles with high accuracy and efficiency.

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