



On the Intermediate Line Region in AGNs

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In this paper we explore the intermediate line region (ILR) by using the photoionisation simulations of the gas clouds present at different radial distances from the center, corresponding to the locations from BLR out to NLR in four types of AGNs. We let for the presence of dust whenever conditions allow for dust existence. All spectral shapes are taken from the recent multi-wavelength campaigns. The cloud density decreases with distance as a power law. We found that the slope of the power law density profile does not affect the line emissivity radial profiles of major emission lines: H β , He II, Mg II, C III, and O III. When the density of the cloud at the sublimation radius is as high as $10^{11.5}$ cm⁻³, the ILR should clearly be seen in the observations independently of the shape of the illuminating radiation. Moreover, our result is valid for low ionization nuclear emission regions of active galaxies.

Keywords: active galaxies, emission lines, photoionisation, radiative transfer, accretion disk

1. INTRODUCTION

The emission lines in active galactic nuclei (AGNs) provide a unique opportunity to study the properties of the materials located in the environment of the supermassive black hole (SMBH). From the measurement of full width at half maxima (FWHM) of lines in the observed AGN spectra, it is well understood that there exist two separate regions of the line emission. Lines with FWHM $\geq 3,000$ km s⁻¹ are emitted by materials with densities $\sim 10^{10}$ cm⁻³ in the broad line region (BLR) located closer to the AGN central engine. Whereas, the narrow line region (NLR) with gas densities $\sim 10^5$ cm⁻³, located much farther from the AGN center, emits the narrow lines with FWHM ~ 500 km s⁻¹. A clear spatial separation in between BLR and NLR is present where the emission of lines with intermediate FWHM is not seen in the observations.

Theoretically, the lack of emission from the intermediate line region (ILR) was explained by Netzer and Laor (1993, hereafter NL93) as an effect of dust extinction, both in absorption and scattering of line photons and continuum. The authors considered radially distributed clouds of different density and ionization level. Broad and narrow line regions were separated due to the dust content which cannot be present in BLR since the gas temperatures are so high that the dust grains cannot survive there. Nevertheless, further out from the nucleus there is a boundary radius named sublimation radius, R_d , where temperature drops substantially, and dust can sustain up to the distances where NLR is located. NL93 presented that, for assumed gas parameters, the strong drop of line emission appears at distances where potential ILR is expected. Therefore natural separation between BLR and NLR occurs when the dust is taken into account in photoionisation calculations of cloud's emission. This natural separation disclaims the existence of the ILR. In

the ILR, the ionization parameter is higher than in further located NLR clouds and hence the relative effect of the dust absorption is stronger. The dust suppressed emission in ILR reappears on transition to the NLR when dust absorption becomes negligible due to the low gas temperature. However, in the recent observations of some AGNs, additional intermediate line component of FWHM $\sim 700\text{--}1,200\text{ km s}^{-1}$ is clearly required to fit the lines in their emission spectra (Brotherton et al., 1994; Mason et al., 1996; Crenshaw and Kraemer, 2007; Hu et al., 2008a,b; Zhu et al., 2009; Li et al., 2015). The open questions are: does the ILR exist physically separated from BLR and NLR? What are the mechanisms that give rise to ILR in some sources but not in others?

Recently, Adhikari et al., (2016 hereafter AD16) have shown, using the framework put forward by NL93, that when the density of illuminated clouds is high enough, the dust does not suppress the gap between BLR and NLR and intermediate line emission is clearly visible. The authors performed photoionisation simulations of radially distributed clouds subject to the radiation of four different spectral energy distributions (SEDs), most common types of AGN. The dust content was introduced at the sublimation radius of assumed value: $R_d = 0.1\text{ pc}$. In NL93, the authors assumed a power law with slope -1.5 and normalization $10^{9.4}\text{ cm}^{-3}$ at R_d to describe the variation of the density of gas clouds with distance from the nucleus. These clouds were then illuminated by the mean AGN spectrum. Resulting emissivity profiles contained the suppression of line emission between BLR and NLR as it is commonly observed. AD16 made one step forward, showing that if the density at the sublimation radius is high, of the order of $10^{11.5}\text{ cm}^{-3}$, the ILR is clearly visible. Such result appeared to be independent on SED of illuminated radiation taken into account. Additionally, the authors argued that the low ionization nuclear emission regions (LINERs) should also exhibit the ILR.

In this paper, we expand the work of AD16 and investigate the variation of density profile of radially distributed clouds. All photoionisation simulations are done with the most recent version of the CLOUDY code (Ferland et al., 2017). To accommodate broad types of ionizing SEDs in our calculations, we considered four distinct AGN types: Sy1.5 galaxy Mrk 509 (Kaastra et al., 2011), Sy1 galaxy NGC 5548 (Mehdipour et al., 2015), NLSy1 galaxy PMN J0948+0022 (D'Ammando et al., 2015), and LINER NGC 1097 (Nemmen et al., 2014), each of them obtained from currently available simultaneous multi-wavelength observations. As a result of our simulations, we derived the line emissivity radial profiles for major emission lines: $[\text{H}\beta\ \lambda 4861.36\ \text{\AA}, \text{He II}\ \lambda 1640.00\ \text{\AA}, \text{Mg II}\ \lambda 2798.0\ \text{\AA}, \text{C III}]\ \lambda 1909.00\ \text{\AA}$ and $[\text{O III}]\ \lambda 5006.84\ \text{\AA}$.

Adopting the density normalization to be $10^{11.5}\text{ cm}^{-3}$ at $R_d = 0.1\text{ pc}$, all the power law density distribution yield continuous line emissivity profiles with prominent intermediate line emission component in permitted lines $\text{H}\beta$, He II , and Mg II , independent of the density slopes and the spectral radiation shapes adopted. Below we briefly outline the photoionisation model itself, and discuss the resulting line emissivity profiles focusing mainly on the visibility of ILR in different AGN.

2. PHOTOIONISATION MODEL

The simulation of the photoionisation process is done with the publicly available numerical code CLOUDY version c17 (Ferland et al., 2017), which takes into account all the relevant radiative processes when a gas cloud is subjected to an incident radiation field. A simplistic geometrical set up of gas distributed from BLR further out to NLR is arranged by assuming spherical clouds with varying gas density, n_{H} , and the total column density N_{H} , at each radial distances, r , from the SMBH:

$$n_{\text{H}}(r) = 10^{11.5} (r/R_d)^{-\beta}, \quad N_{\text{H}}(r) = 10^{23.4} (r/R_d)^{-1} \quad (1)$$

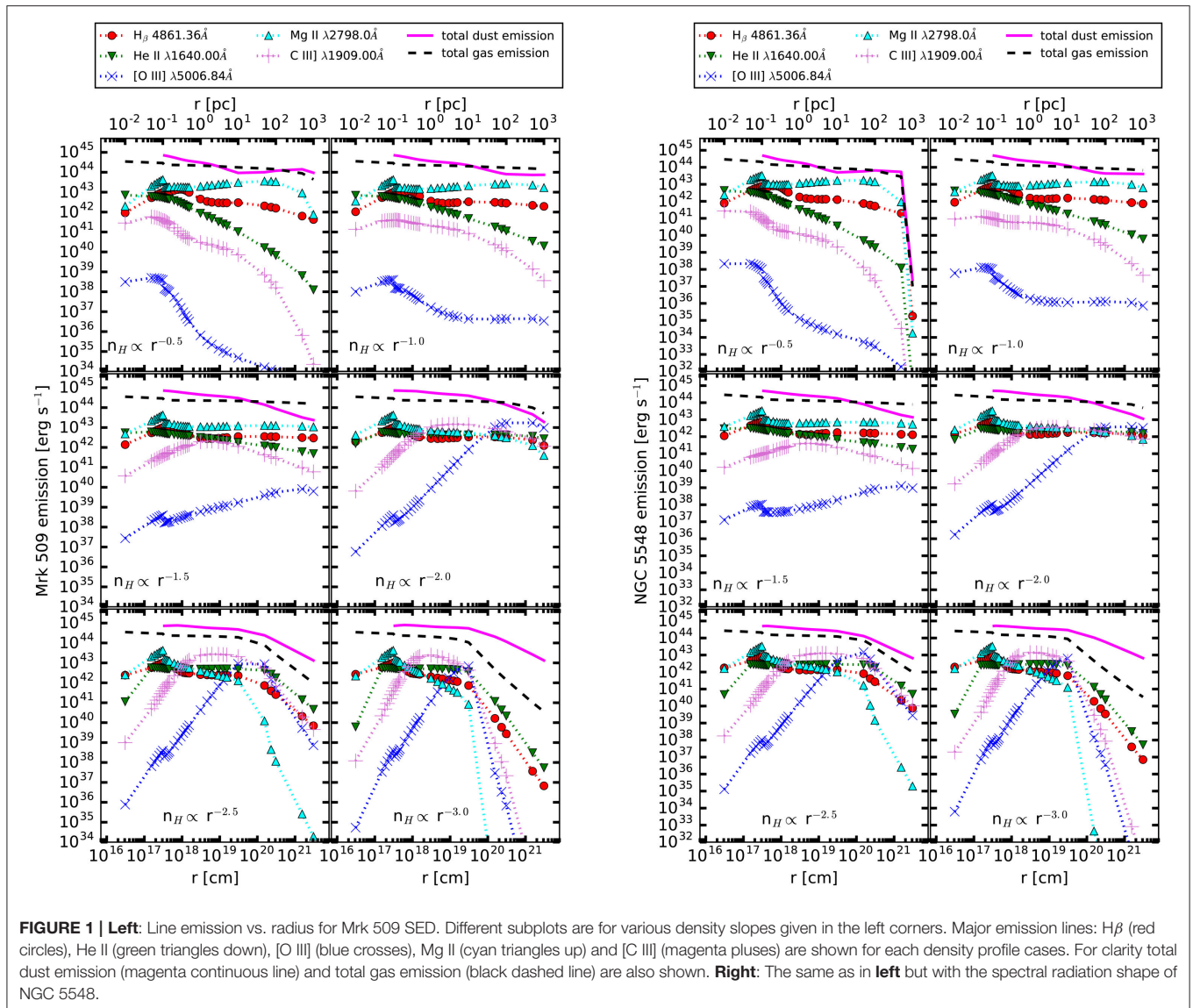
where β is the power law density slope, and R_d is arbitrarily chosen to be equal 0.1 pc (following NL93 and AD16). The total column density of a cloud located at the sublimation radius is assumed after NL93: $N_{\text{H}}(\text{at } R_d) = 10^{23.4}\text{ cm}^{-2}$, and the gas hydrogen density after AD16: $n_{\text{H}}(\text{at } R_d) = 10^{11.5}\text{ cm}^{-3}$. Here, we stress that the density normalizations lower than the value adopted in this paper, do not reproduce the intermediate line emission as shown by AD16.

Besides three previously considered types of AGN incident radiation shapes: Sy1.5 galaxy Mrk 509 (Kaastra et al., 2011), Sy1 galaxy NGC 5548 (Mehdipour et al., 2015), NLSy1 PMN J0948+0022 (D'Ammando et al., 2015), in this paper we also used the shape appropriate for LINER NGC 1097 (Nemmen et al., 2014). This choice of SED covers the general shapes of the radiation emanating from the AGN central engine. Adopted SEDs are the incident spectra used in the photoionisation simulation, where clouds distributed along the range of radii are exposed to the same type of radiation. All the SEDs are normalized to the bolometric luminosity $10^{45}\text{ erg s}^{-1}$ which is an input to the CLOUDY code. This allows to compute the ionizing flux i.e., ionization parameter at each cloud radius.

We adopted the CLOUDY default chemical abundances, which are mostly the Solar values derived by Grevesse and Sauval (1998) for the gas clouds at $r \leq R_d$, whereas the interstellar medium (ISM) composition¹ with dust grains is used for the clouds at $r > R_d$. This assumption is consistent with the observational suggestions that the BLR is devoid of dust whereas the lower temperature in NLR allows its existence. On moving further out from BLR to NLR, the increase in radial distance is accompanied by the decrease in ionization degree and a cloud thickness.

The aim of this paper is to search how the appearance of intermediate line emission is sensitive on the value of density power law slope. Below we present line luminosity radial profiles for six values of $\beta = 0.5, 1.0, 1.5, 2.0, 2.5, 3.0$. We are aware that the power law density distribution of clouds does not reflect realistic situation in AGN, but it is sufficient for the purpose of this paper. In the forthcoming paper (Adhikari et al., in preparation), we plan to use realistic density profile, which is expected where clouds form from outflowing gas above the accretion disk atmosphere. Furthermore, we plan to consider measured values of bolometric luminosities, which give the realistic position of sublimation radius for each type of AGN due to the formula by Nenkova et al. (2008).

¹for details see the Hazy1 CLOUDY documentation.



3. LINE EMISSIVITIES

As the results of photoionisation simulations, we compute line luminosities emitted from clouds located at the each radii. Therefore, by presenting line luminosity dependence on the distance from SMBH, i.e., line luminosity radial profile, we can check if the emission from ILR is comparable to the BLR and NLR or substantially lower. This is our basic test for the existence of ILR in all considered types of AGN.

We derived the line luminosity radial profiles for the major emission lines : [H β λ 4861.36 Å, He II λ 1640.00 Å, Mg II λ 2798.0 Å, C III] λ 1909.00 Å, and [O III] λ 5006.84 Å. The resulting line emissivity profiles for the four cases of SED are presented in the **Figures 1, 2**. In all cases of density power law slopes, we recovered a continuous line emission, with a small enhancement of the permitted lines H β and He II at the radial distance around 0.1 pc corresponding to the intermediate region, independent of the

shape of the SEDs in consideration. There is a small reduction of Mg II line at 0.1 pc though not very significant as compared to the suppression presented by NL93. The semi forbidden line C III contribution to the intermediate emission component becomes the most prominent for the density profile with $\beta = 1.5$. These results corroborate with the conclusion of AD16 that when the density of the emitting gas is high enough, the extinction effect of dust grains on line production is negligible.

The most noticeable effects of different density profiles on line emissivities occur in the NLR range, i.e., for $r > 50$ pc. This behavior is quite obvious since for those radii, differences in densities between profiles are the biggest. For $\beta = 0.5$ and 1.0, density falls slowly and remains moderately high across the radii causing the strong suppression of forbidden line [O III]. [O III] is effectively produced in low density environment and becomes prominent when the density around the radius 10 pc becomes low enough, the cases for the profiles with $\beta \geq 2.0$. Narrow

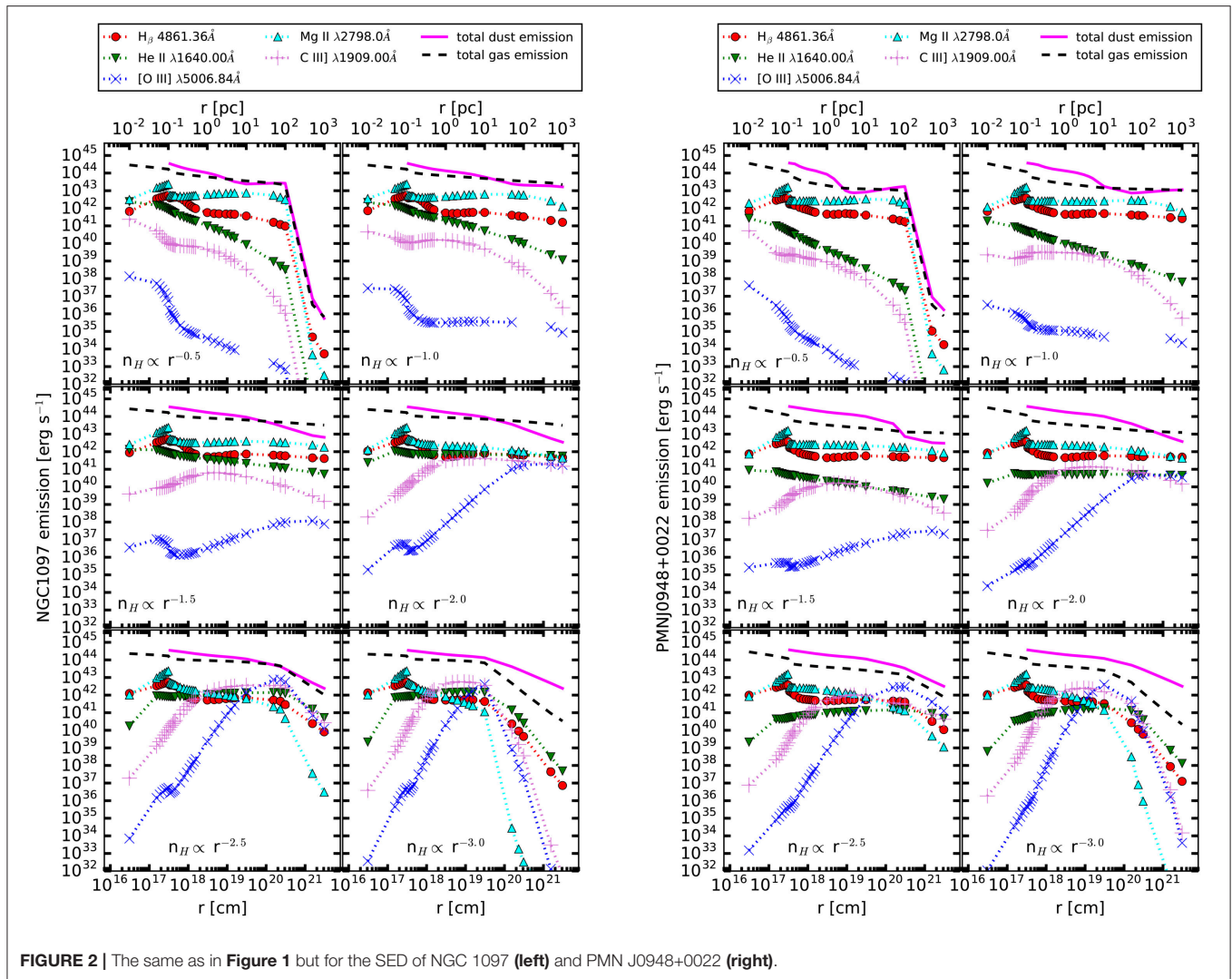


FIGURE 2 | The same as in **Figure 1** but for the SED of NGC 1097 (**left**) and PMN J0948+0022 (**right**).

line emission is dominated by the [C III] components when the density distribution is given by the profiles with $\beta \geq 2.0$. We found that the derived line emissivities for all cases of power law density slopes, particularly in the region of intermediate line emission, do not strongly depend on the shapes of the SED used. There are subtle differences in emissivities corresponding to the BLR and NLR due to the different amount of UV and soft X-ray photons among the SEDs. This result is in agreement with the conclusion of AD16, that the presence of ILR emission is not determined by the shape of the incident radiation.

4. DISCUSSION

The results above confirm the conclusion of AD16, that the dust extinction of the emission lines in AGN introduced by NL93 is important only when the gas density is low. In AD16, the authors adopted only a density power law of slope $\beta = 1.5$ whereas this work demonstrates that the different slopes of the density distribution do not matter significantly as long as the gas density at the sublimation radius is high, in this case being $10^{11.5}$

cm^{-3} . In all cases of considered density profiles, we obtained an intermediate line emission around 0.1–1 pc, mostly manifested in permitted lines H β , He II, and Mg II, and weakly present in the semi forbidden line [C III]. This indicates that the high density and the low ionization environment favors the intermediate line emission rather than the high ionization environment where the forbidden line [O III] is produced. So, in the AGNs where the ILR is seen in observations, the emitting region is composed of the dense and less ionized gas.

The physical reason for the disappearance of the effect of dust is connected with the size of H⁻ ionized front in the gas. At very high value of ionization parameter, i.e., for the low density case, the volume of the H-ionized region is very large, if not the full cloud volume. When the density of cloud increases, the ionization decreases, and a cloud consists of two regions: H-ionized region and H-neutral region. The line emission comes from the H-ionized region, and only the dust in this region competes with the gas for the photons. In other words, the dense clouds have much smaller geometrical thickness of H⁻ ionized layer, smaller dust column density in the region with abundant

photons, and therefore the dust absorption is negligible. Our simulations are not yet aimed to make a quantitative statements about the studied objects. For that, we would need to do more extensive study, representing the bolometric luminosity and the position of the inner radius of the dust distribution appropriate for a given object. However, the grid of results shows a clear trend.

In the recent years, there has been promising claims that, broad line emission clouds in AGN are connected with the wind from the upper part of an accretion disk atmosphere (Gaskell, 2009; Czerny and Hryniewicz, 2011). As shown in the Fig. 6 of AD16, the density profiles computed in the upper part of the standard disk atmosphere can be quite dense with values up to $\sim 10^{15} \text{ cm}^{-3}$ at the assumed position of sublimation radius. Those density values depend on the mass of the black hole and the disk accretion rate. Because of the high gas density, photoionisation simulations outcome with continuous line luminosity radial profile for the reason discussed in the previous paragraph. As the consequence, ILR should be observed together with BLR and NLR. The use of realistic density profiles expected from the accretion disk atmosphere is very important in

the aim to understand the nature and origin of the ILR observed in some AGNs. This work is in progress and will be presented by Adhikari et al. (in preparation).

AUTHOR CONTRIBUTIONS

TPA, AR, and KH were responsible for developing the idea, doing the simulations, analyzing the results and writing the text for the manuscript, BC provided the concept and GF discussed the results.

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