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SPECIALTY SECTION

This article was submitted to Colloidal Materials and Interfaces, a section of the journal Frontiers in Materials

RECEIVED 18 August 2022 ACCEPTED 02 November 2022 PUBLISHED 24 November 2022

CITATION

Kashyap V, Pawar H, Kumar C, Chaudhary N and Saxena K (2022), Analysis of synthesized doped vertical silicon nanowire arrays for effective sensing of nitrogen dioxide: As gas sensors. *Front. Mater.* 9:1022317. doi: 10.3389/fmats.2022.1022317

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Analysis of synthesized doped vertical silicon nanowire arrays for effective sensing of nitrogen dioxide: As gas sensors

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In the present study, the controllable fabrication of silicon nanowires (Si NWs) with vertical alignment was accomplished using metal assisted chemical etching (MACE). The different characteristics, such as structural, morphological, chemical, optical, and dielectric properties were analyzed using X-ray diffraction (XRD), field emission scanning electron microscope (FESEM), Raman spectroscopy, ultraviolet-visible diffuse reflectance spectroscopy (UV-DRS), and LCR [inductance (L), capacitance (C), and resistance (R)] meter (volume of the gas-sensing chamber is 650 mm³). It was revealed from the morphological study i.e., from the FESEM that p-type Si NWs are smaller in size than n-type Si NWs which is attributable to the energy band gap. The optical band gap (E_a) is observed to increase from 1.64 to 1.89 eV with the decreasing of the crystallite size and the optical reflection spectra of the Si NWs show a shift toward a lower wavelength (blue shift). Moreover, Raman spectra verified the red-shifted, asymmetrically broadened Raman lineshapes, which provides information about the size confinement effect in Si NWs. The MACE approach is excellent for synthesizing nanowire structures for use in gas-sensing applications due to its flexibility. The sensitivity of synthesized Si NWs was tested for NO₂ gas. The sensor method is unique based on the testing of the device in the presence of a test gas because the use of the gas-sensing setup has the potential to measure the change in resistance by varying frequency, temperature, and time.

KEYWORDS

FESEM, Raman spectroscopy, dielectric properties, gas sensing, UV-DRS

1 Introduction

The field of electronics and photonics has seen semiconductor nanowires (NWs) emerge as a promising candidate in recent years due to their wide range of possible applications (Xi et al., 2007; Gonchar et al., 2016; Li and Kwok, 2018; Gonchar et al., 2019). Silicon (Si) is gaining importance as a promising approach for an industrial revolution in semiconductor technology due to its technological-based properties.

The enormous growth of microelectronics has benefited Si process technology. Si is potentially the only material that satisfies all of the requirements, including its non-toxicity, strong absorption optics, and high performance (greater than 10%). The manufacturing of effective devices to fuel our economies has shifted toward the more stable Si nano-crystals and NWs. Various methods are available for the synthesis of Si NWs with vast numbers of applications. Si nanostructures (NSs) have been manufactured using a wide range of processes, including chemical etching, epitaxial growth, and electron beam lithography (Hasan et al., 2013; Puglisi et al., 2019). The simplest and lowest cost method for creating NSs is metalassisted chemical etching (MACE) (Li et al., 2013). Si NWs have enormous potential for usage in thermoelectric power, lithium-ion batteries, sensors, and solar energy conversion (Peng et al., 2008; Kim et al., 2010; Liang et al., 2014; Toor et al., 2016; Ghosh et al., 2018; Singh et al., 2018; Kashyap et al., 2021a).

The goal for the scientific community has been to improve human health and the environment by monitoring unsafe and poisonous substances, gases, micro-organisms, and radiation. One-dimensional (1D) material-based sensors are extremely efficient and compatible with recent manufacturing techniques resulting in data that are useful for health and environmental monitoring (Zhu and Zeng, 2017). Nanostructures such as nanowires can be the path to effective and advantageous sensor materials. The scientific community is working to develop sensors that can provide alerts of potential risks in response to a number of pressing issues, including protecting human health, the environment, and safety. Additionally, Si NWs may simply be scaled up for use in gas-sensing applications employing the room temperature MACE approach. The usage of gas-sensing technology is widespread in both household and industrial settings. The most crucial characteristics of gas-sensor devices are their sensitivity, working temperature, selectivity, energy consumption and cost of production (Zhu and Zeng, 2017; Mirzaei et al., 2018; Bhati et al., 2020).

According to the type of material, the depletion layer either forms or deforms during the gas-sensing phenomenon which includes the transfer of electrons between the sensing material and the target gases. Zhang et al. (2016) conducted the first study in the field of Si NW based gas sensors. Si NWs for gas sensing have a sensing function that is fairly comparable to a conventional semiconductor gas sensor, since a single Si NW unit has high sensitivity in a nanosensor applications. Peng et al. (2009) created two independent gold electrodes that have a limit of detection (LOD) of 500 ppb against NO₂ and are placed on the top of a vertical Si NW array. To increase the sensitivity of the Si NW array for gas detection, each one should be connected to an electrical circuit (Zhang et al., 2016; Morganti et al., 2021). The problem has been solved by developing a number of methods for fabricating a continuous electrode on top of a Si NW array.

This research focuses on functionalization of Si NWs, synthesis, and the application of Si NWs as gas sensors. In this study, we have fabricated Si NWs using the MACE method and the role of quantum confinement has been established. Techniques such as ultraviolet-visible diffuse reflectance spectroscopy (UV-DRS) was used for the band-gap variation in the nano regime for the different doped Si NWs, and the noble method of gas sensing was introduced to detect NO₂ gas at lower concentrations.

2 Materials and methods

In this study, the MACE technique was used to fabricate Si NWs samples of n-type Si (100) (doped with ~10³ phosphorous atoms/cm³) and p-type Si (100) (doped with ~10³ Boron atoms/ cm³) having resistivity 1–10 Ω cm, thickness 265 ± 25 μ m, 2 \times 2 cm², keeping all other parameters constant. For the deposition process, hydrofluoric acid (HF) (Sigma-Aldrich, assay 99.9%) and silver nitrate (AgNO₃) (Sigma-Aldrich, assay 99.9%) were taken in the proportion of 4.8 M and 5 mM, respectively, for 60 seconds at room temperature to deposit silver nanoparticles (Ag NPs). The etching solution was made by taking HF and hydrogen peroxide (H2O2) (Merck, assay 99.9%) in the proportion of 4.6 and 0.5 M, respectively. After the deposition process, the samples were etched for a duration of 40 min each. The etched wafers were then transferred to HNO₃ to dissolve the remaining Ag NPs and they were then dipped into HF solution to remove the oxide layer (Supporting Information). The samples are named as n_{40} (n-type and etching duration of 40 min) and p_{40} (p-type and etching duration of 40 min).

For characterizing the samples, X-ray diffraction (XRD) was recorded by using Cu (K_{α}) radiation (0.154 nm). Morphological analysis of Si NWs was carried out using a field emission scanning electron microscope (FESEM) (HITACHI SU8010). The Raman spectra were recorded using a HORIBA Jobin Youn iHR-550 spectrometer equipped with a cooled charge-coupled device (CCD) detector. An excitation line with a wavelength of 785.0 nm was produced using an air-cooled Ar⁺ ion laser and the samples were exposed for 10 seconds. The laser power density was less than 10^7 mW/cm². The dielectric and sensing properties were investigated using a LCR [inductance (L), capacitance (C),



and resistance (R)] meter (HIOKI, IM3536) and a gas-sensing system, respectively.

3 Results and discussion

3.1 Structural properties

The XRD patterns of the Si NWs for the n-type and p-type samples are shown in Figure 1. The XRD peak for the n_{40} sample is observed at $2\theta = 69.16^{\circ}$, assigned (004) plane, and for the p_{40} sample is observed at $2\theta = 68.8^{\circ}$, lattice plane (004). The sharp and narrow peak observed in the n_{40} -type indicates improvements in crystallinity (Westra et al., 2010; Ghahramanifard et al., 2018).

The average crystalline size of Si NWs was calculated by the Scherrer equation (Minin and Minin, 1994):

$$D = \frac{0.9\lambda}{\beta\cos\theta}$$
(1)

where D is the crystalline size (average), θ is the diffraction angle, λ is the wavelength (0.154 nm) of radiation used for the diffraction, and β is the full-width at half maximum (FWHM) of the observed peaks (~0.002 for both samples). The estimated crystalline size of the Si NWs is 89 nm and 73 nm for the n_{40} and p_{40} samples, respectively.

3.1.1 Micro strain

The XRD peak broadening occurs as a result of a deviation from the ideal crystal. This broadening can be primarily caused by the solution, crystallite size, and micro strain; the micro strain is a major factor among these and estimated as (Samy et al., 2019)

$$\varepsilon = \frac{\beta_{hkl} \cos \theta}{4}, \qquad (2)$$

where ε = micro-strain, β_{hkl} = FWHM, and θ = diffraction angle or Bragg's angle. The micro strain and dislocation density are summarized in Table 1.

3.1.2 Dislocation density

Dislocation is a topological defect; it is the length of dislocation lines in a unit volume (dislocation density) of a crystalline material, i.e., defects present in the samples, mathematically expressed as (Ravichandran and Philominathan, 2009)

$$\delta = \left(\frac{\beta \cos \theta}{k\lambda}\right)^2 = \frac{1}{D^2}$$
(3)

where δ = dislocation density, k = dimensionless shape factor (its value is 0.89), and θ and D are the same as above. The small values of δ obtained in the current investigation support the good crystallinity of the fabricated sample i.e., Si NWs, by using the MACE method. It was found that the dislocation density and micro strain decrease with an increase in crystallite size showing lower imperfections as shown in Table 1. This might be because Si NWs have larger crystallites due to a shift in the dopant (n-type or p-type) which results in a decrease in the occurrence of grain boundaries (Joshi et al., 2003; Begum et al., 2012).

3.2 Morphological properties

Scanning electron micrographs of Si NWs are depicted in Figure 2. Figures 2A–D depict magnified 45°-tilt and cross-sectional FESEM images of Si NWs on n (Figures A,B) and p

S. No.	Sample name	Average crystalline diameter (D) (nm)	Micro strain $\varepsilon(\times 10^{-4})$	Dislocation density $\delta(\times 10^{15} (\frac{\text{lines}}{\text{m}^2}))$	Optical band gap (eV)
1	p ₄₀	73	7.8	3	1.89
2	n ₄₀	89	4.3	2.5	1.64

TABLE 1 Summarized microstructural and elastic parameters of the Si NWs.



(Figures C,D) Si (100) wafers (1–10 Ω cm), keeping other parameters constant, such as chemical concentration and etching time (i.e., 40 min). The length of the Si NWs is calculated using ImageJ software; it is approximately 10 μm and 11 μm for n_{40} and p_{40} Si NWs, respectively, as shown in Figure 3.

Figures 2A–D reveal that, as we deal with different doping (n_{40} and p_{40}), the length of deposited Si NWs increases due to sufficient oxidizing species present in the solution to oxidize the Si (Peng et al., 2011). The p_{40} sample has a denser structure and has a greater length compared to the n_{40} sample, as shown in Figure 3.

The average size (diameter) estimated using ImageJ software, of the Si NWs in n_{40} and p_{40} are 97 nm and 80 nm, as shown in Figures 4A,B, respectively. According

to previous studies the variation in the shape of Si NWs is due to the occurrence of inhomogeneous nucleation of Ag on the Si (100) wafer (Hutagalung et al., 2017; Naffeti et al., 2020; Rouis et al., 2021). We conclude that the lower size of the p_{40} type Si NW compared to the n_{40} -type Si NW is attributable to the energy band gap because as the size reduces the band gap increases. As discussed earlier, the p-type has a band gap of 1.89 eV, while the n-type has a band gap of 1.64 eV (Rouis et al., 2021). Figure 4 depicts the average diameter of n-type and p-type Si NWs, which we calculated using ImageJ software by selecting the Si NWs from the FESEM images. By selecting the diameter of all images of the Si NWs, we take the average length, which comes out as 80 nm and 97 nm for n-type and p-type Si NWs, respectively.





3.3 Optical properties

The total reflectance of Si NWs onto n_{40} Si (100) and p_{40} Si (100) in the wavelength range of 250 nm–850 nm is shown in Figures 5A,B. It is observed that the p-type Si NW has better light-trapping than the n-type Si NW. In the n_{40} sample, the observed reflectance was ~9% in the UV region (<300 nm), ~6% in the visible region (300–700 nm), and about 3.5% in the IR region (>700 nm), which is less than the percentage of reflectance given by crystalline Si (c-Si) (reflectance = 66%). In the p_{40} sample, the observed reflectance was ~3% in the UV region (<300 nm), ~4% in the visible region

(300 nm-700 nm), and about 3% in the IR region (>700 nm). Moreover, the total reflectance depends upon on the length of the Si NWs; this corresponds to the reflection, which becomes lower with longer lengths of Si NWs. Total reflectance decreased for the p_{40} sample because the charge carrier gradually increases on the textured surface (Kashyap et al., 2022).

The Kubelka–Munk function was used to transform reflection into absorption spectra (Chiu and Li, 2009; Yogi et al., 2017; Kashyap et al., 2021b):

$$F(R) = \frac{(1-R)^2}{2R} = \frac{K}{S}$$
 (4)





where s denotes the scattering coefficient. If s \rightarrow 0 (no scattering occurs), all the light will be either transmitted or absorbed, while if k \rightarrow 0 (no absorption or transmission occur), all the light will be reflected.

The optical band gap can be estimated by (Shukla and Dixit, 2016; Venkatesan et al., 2019)

$$F(R)h\upsilon = A(h\upsilon - E_g)^n$$
(5)

where F(R) is the Kubelka–Munk function, which is equivalent to the optical absorption coefficient, E_g is the optical band gap of the Si NW, A is a constant, and n is

the power factor which is impacted by several electronic transitions. The plot between $(F(R) hv)^{1/2}$ vs. photon energy (E_g) for the n_{40} type- and p_{40} -type Si NWs is depicted in the inset of Figures 5A,B, given by (i) and (ii), respectively. The estimated values are 1.64 and 1.89 eV for the n_{40} - and p_{40} -typetype samples, respectively. It is observed that the n_{40} sample has a lesser band gap in comparison to the p_{40} sample because the conductivity of the n_{40} sample is higher.

This might be explained by a growth in crystallite size and a reduction in imperfections. The plot between $(F(R) hv)^{1/2}$ vs. photon energy (E_g) for n₄₀-type and p₄₀-types in the inset of



Figure 5 [i.e., (i) and (ii)] illustrates that the Si NWs are linear throughout a large range of photon energy. This suggests that the synthesized Si NWs have a direct optical band gap. The optical band gap of these materials can be calculated by extrapolating the linear portion of the curve to the energy axis. It is noticed that the resulting band gap falls from 1.89 eV (for the p_{40} -type) to 1.64 eV (for the n_{40} -type). According to the usual relationship between band gap and crystallite size, the band gap rises as the crystallite size decreases (summarized in Table 1). Due to the size confinement in the Si NWs, the observed values of E_g are greater than the value of the bulk Si (1.12 eV) (Gao et al., 2007). The variation of different parameters with crystalline size is also confirmed in Figure 6.

3.4 Raman spectroscopy

The Raman spectra for the n_{40} and p_{40} samples are given in Figure 7, along with a comparison with bare Si. A Raman peak at 520 cm⁻¹ with a FWHM of 3.5 cm⁻¹ of c-Si was found due to the Raman active zone center symmetry point $\Gamma = 0$ of the Brillouin zone. The Raman spectra for the n-type and p-type bare Si were found to be identical, with no asymmetry, as

TABLE 2 Calculations of various parameters from the Raman analysis.

shown in Figure 7. Also, the various parameters calculated from the Raman analysis has been given in Table 2

Figure 7 shows that in comparison to their bulk equivalent, the Raman spectra for Si NWs are red-shifted and asymmetrically widened. An asymmetric broadening of the Raman line-shape is caused by the combination of the Fano effect and phonon confinement (Kumar, 2013; Yogi et al., 2016; Saxena et al., 2017; Bhujel et al., 2018). Analysis of the Raman line-shape characteristics such as peak position, peak asymmetry, and FWHM which also reveals the presence of phonon confinement in both samples. The wavenumber was observed at 517.4 \pm 0.26, and 518.4 \pm 0.18 cm⁻¹ for the n₄₀ and p₄₀ samples respectively, with varying asymmetry ratios due to the quantum confinement (QC) effect. This range of wavenumbers of the line is generally defined with a long-range Raman scattering line.

The asymmetry arises due to the presence of nano-scale size confinement (Li et al., 2005; Kumar, 2013). Phonon confinement in Si NWs is thought to be the cause of the red-shifted, asymmetrically widened Raman line-shapes, as shown in Figure 7. The Raman line-shapes are wider on the lower energy side of the peak for n_{40} semiconductors (asymmetry ratio 1.3), but wider on the higher energy side of the peak for p_{40} semiconductors (asymmetry ratio 1.5), as there is more electron-phonon confinement in the n_{40} sample.

In vertically aligned Si NW arrays, metal electrodes must be linked to the ends of the Si NWs. Copper wire was used for the connecting wires and silver epoxy was used to join the wires after they had been tested and chosen for ohmic contact formation. The Si nanowire tips must be in complete contact with a metal electrode; otherwise, a short circuit would result. Figure 8 is a representation of the ohmic contact made with the fabricated samples of Si NWs.

3.5 Dielectric properties

The change of electrons in the presence of the electric field is responsible for the change in the capacitance of the Si NWs. The dielectric properties of the Si NWs were investigated in the form of a thin film of dimensions 2 mm \times 2 mm with 0.00265 mm thickness. The electrodes were deposited on thin film using silver

Sample name	FWHM (cm ⁻¹)	Asymmetric ratio	Peak position (cm ⁻¹)
n ₄₀	21	1.3	517
p ₄₀	15	1.5	518
Bare n-type Si	3.5	1	520
Bare p-type Si	3.5	1	520
	Sample name n ₄₀ P ₄₀ Bare n-type Si Bare p-type Si	Sample name FWHM (cm ⁻¹) n ₄₀ 21 p ₄₀ 15 Bare n-type Si 3.5 Bare p-type Si 3.5	Sample name FWHM (cm ⁻¹) Asymmetric ratio n ₄₀ 21 1.3 p ₄₀ 15 1.5 Bare n-type Si 3.5 1 Bare p-type Si 3.5 1



FIGURE 8 Ohmic contacts with Si NWs and electrodes.

metal to make ohmic contact. The dielectric study of Si NWs was performed in air atmosphere. The dielectric permittivity was calculated by the following formula (Pawar et al., 2021):

$$\varepsilon' = \frac{\complement t}{\varepsilon_0 A} \tag{6}$$

where ε' is the dielectric permittivity of material, ε_0 is the dielectric permittivity of the free space, <u>C</u> is the capacitance, t is the thickness of the thin film, and A is the area of the thin film. The dielectric permittivity of Si NWs as a function of frequency with varying temperature is shown in Figures 9A,B. The maximum dielectric permittivity of both types of Si NWs was found at low frequency, decreasing rapidly with increasing frequency then becoming almost constant after the 1–8 MHz frequency range. The Si NWs had a more conductive nature; hence, the dielectric permittivity of the samples was higher at low frequency because of the Debye orientational polarization



FIGURE 9

Dielectric permittivity and dielectric loss with varying temperature of (A,C) n_{40} Si NWs and (B,D) p_{40} Si NWs as a function of frequency, respectively.



(Rasool et al., 2012). As the temperature decreased from 110°C to 30°C, the dielectric permittivity decreased and became lowest at 30°C, nearly 965 and 1,128 for n-type and p-type Si NWs, respectively, because the dipoles tend to rearrange and follow the direction of the applied electric field. The optimum value of dielectric permittivity at room temperature is a strong indication for the sensing performance of the fabricated device.

The dielectric loss showed a very low loss of energy in operation with varying frequency and time for both devices, which indicates that the devices will show enhanced sensing phenomenon for the detection of NO₂ gas. The dielectric loss as a function of frequency with variation of temperature is depicted in Figures 9C,D. The dielectric loss decreased rapidly as the frequency initially increased until 1 MHz and 0.5 MHz for n-type and p-type Si NWs, respectively, and became constant but as the temperature increased from 30°C to higher temperatures the dielectric loss increased with an increase in frequency because of the ion exchange anisotropy phenomenon (Khan et al., 2020), where at high temperature the ions tend to move from the barrier and produce polarization at a large scale.

The AC conductivity as a function of frequency with varying temperature is depicted in Figure 10. It is clearly observed from Figures 9A,B that as the frequency increased from 100 Hz to 8 MHz and the AC conductivity also increased in linear order and became highest at 8 MHz for both types of sensing materials. The AC conductivity of the n_{40} Si NW increased with increasing temperature from 30°C to 110°C and became highest at 110°C, while this was found to be lower for the p_{40} Si NW. The change in conductivity with increasing temperature is observed due to the hopping conduction mechanism (Saidi et al., 2017).



FIGURE 11 Schematic of a dielectric setup including an equivalent circuit, a furnace, a gas-sensing system, and a LCR meter (parallel equivalent circuit is shown in the inset).

3.6 Gas-sensing properties

The instrument used for the gas-sensing analysis is shown in Figure 11, which comprised a LCR meter furnace followed by a gas-sensing arrangement.

The sensing mechanism behind the response toward oxidizing gas (NO_2) is based on the change in resistance or conductance of the semiconductor devices in the presence of



air and an oxidizing gas environment. The change in resistance as a function of frequency at 30°C temperature for n₄₀-type and p₄₀-type Si NWs devices is depicted in Figures 12A,C. It is observed from Figures 12A,C that as the frequency increased the resistance decreased rapidly until 1 MHz and then became almost constant at higher frequencies. Furthermore, the resistance of both the devices decreased as the temperature increased from 30°C to 110°C, which is depicted in Figures 12B,D. The reason behind the decrement of sensor resistance in an air atmosphere at higher temperatures is due to adsorbed oxygen molecules losing electrons at the surface of the sensor, which further move through the sensor and increase the flow of current; hence, the resistance of the sensor decreases simultaneously. The maximum resistance of both sensor devices was observed at 30°C in air; hence, the gas-sensing experiments were performed at 30°C with varying concentrations of oxidizing gas from lower to higher concentrations (10 ppm-250 ppm). The ratio of change in resistance matter for the identification of gas response was calculated using Eqs 7, 8. The objective of the present study is to investigate the change of resistance at 30° C with varying concentrations of NO₂ gas for the fabrication of room-temperature detection devices.

Figures 11B,D show the change in resistance as a function of frequency with varying temperatures of the sensor devices in the presence of oxidizing gas. The change in resistance or conductance also depends on the n-type or p-type behavior of the devices. Usually, the resistance of n_{40} type Si NWs in the presence of oxidizing gas such as NO₂ increases with increasing the concentration of the target gas (10 ppm-250 ppm) while decreasing for p_{40} -type Si NWs owing to the dual modulation of space charge constituency between the adsorbed oxygen with the sensing element and NO₂ gas atmospheres. This occurs because of the physically powerful electron-withdrawing capabilities of NO₂; the accumulation of a hole in p_{40} -type Si NW based sensors



occurs due to the withdrawing of electrons, which increases the conductivity of the sensor (Cao et al., 2013) with increasing concentrations, while decreasing for n_{40} -type Si NW based sensors.

The gas-sensing behavior of the sensors was investigated by observing the change in resistance in air and oxidizing gas atmosphere using the following relationship (Pawar et al., 2020):

$$R (\%) = \left(\frac{\Delta R}{R_g}\right) \times 100 = \left(\frac{R_a - R_g}{R_g}\right) \times 100i R_g > R_a$$
(7)

$$R(\%) = \left(\frac{\Delta R}{R_a}\right) \times 100 = \left(\frac{R_a - R_g}{R_a}\right) \times 100 \text{ if } R_a > R_g \qquad (8)$$

where R (%) is the sensitivity of the target gas, ΔR is the change in resistance, and R_a and R_g are the resistance of the air and the target gas, respectively. n_{40} -type and p_{40} -type Si NW based sensors were used for the detection of oxidizing gas such as NO₂.

The gas-sensing response as the function of frequency with varying gas concentrations from 10 ppm to 250 ppm is shown in Figures 13A,B. The sensing response of the n_{40} -type Si NW based sensor showed that the response toward NO2 gas initially decreased with increasing frequency, reaching a minimum at a certain frequency and then increasing again at higher frequencies. The concentration of the oxidizing gas introduced in the gas-sensing unit, varying from 10 ppm to 250 ppm, exhibited an excellent gas response at lower concentrations, which is at 50 ppm, and a minimum response at higher concentrations at all frequency ranges because the fabricated sensor had a small dimension suitable to detect gas precisely at lower concentrations, but at higher concentrations, the adsorbed layer of the O⁻² and gas-interacting surface was saturated; hence, the free electrons were not available in the sensor element, which can propagate current through the device, and as a result the resistance increased rapidly, which reduced the gas response of the sensor, while decreasing in the case of the p_{40} -type Si NW sensor. In the presence of NO₂ gas, as the frequency increased, the resistance of the sensing device decreased due to saturation polarization which improved the gas sensitivity. A chemiresistive sensor's sensitivity is a resistance-dependent mechanism as the maximum change in resistance was observed in the presence of NO₂ gas, while in the air atmosphere, the gas response increased.

The gas response R (%) of n_{40} -type and p_{40} -type Si NW sensors as a function of concentrations at specific frequencies of 1, 4, and 8 MHz is depicted in Figures 13C,D.

Figure 13C shows the gas response of the n_{40} -type Si NW based gas sensor, with the maximum sensing performance at 50 ppm found to be 35%, 63%, and 68% at 1, 4, and 8 MHz, respectively, while this was 75%, 68%, and 62% for the p_{40} -type Si NW based sensor, as illustrated in Figure 10D. The aforementioned results revealed that the p_{40} -type Si NW based sensor is the best for the detection of NO₂ gas at room temperature and at lower frequencies, while the n_{40} -type Si NW-based sensor is the best for the detection of NO₂ gas at room temperature and higher frequencies.

Conclusion

The low-cost MACE method was utilized for synthesizing Si NWs. The structural analysis of the synthesized Si NWs confirmed that as the crystallite size increased, the dislocation density decreased simultaneously. The morphological analysis showed that the array of Si NWs had a length of approximately 10 \pm 0.5 and 11 \pm 0.06 μm and an approximate average size (diameter) of 97 \pm 5 nm and 80 \pm 0.001 nm for the n_{40} and p_{40} samples, respectively. The band gap of the synthesized Si NWs was observed to decrease from 1.89 eV (for the p-type) to 1.64 eV (for the n-type) which is a remarkable outcome for gas-sensing applications. Due to greater electron-phonon confinement in the n₄₀ sample, the Raman line-shape was wider on the lower energy side of the peak for the n_{40} -type semiconductor (asymmetry ratio: 1.3) but was wider on the higher energy side of the peak for the p₄₀-type semiconductor (asymmetry ratio: 1.5). The dielectric study confirmed the lower dielectric values of the synthesized Si NWs at room temperature, which is a strong indication for the sensing performance of the fabricated device. Regarding the gassensing response of the n₄₀-type Si NW based gas sensor, the maximum sensing performance at 50 ppm was found to be 35%, 63%, and 68% at 1, 4, and 8 MHz, respectively, while this was 75%, 68%, and 62% for the p_{40} -type Si NW based sensor, which confirms that the p40-type Si NW based sensor is the best for the detection of NO2 gas at room temperature and lower frequencies, while the n₄₀-type Si NW-based sensor is the best for the detection of NO2 gas at room temperature and higher frequencies.

Data availability statement

The raw data supporting the conclusions of this article will be made available by the authors, without undue reservation.

Author contributions

VK: conceptualization, methodology, writing—review and editing, validation, resources, and data curation. HP: writing—review and editing, validation, resources, and data curation. CK: formal analysis and visualization. NC: supervision. KS: supervision and investigation.

Acknowledgments

We would like to thank to Dr. Deepshikha Rathore (ASAS, Amity University, Jaipur, Rajasthan, India) for the use of the dielectric and gas-sensing facilities. we also acknowledge financial support via Senior Research Fellowship (SRF) from the University Grants Commission (UGC), India. The author acknowledges the Sophisticated Analytical Instrumentation Facility (SAIF) and Prof. G. S. S. Saini (Department of Physics, Panjab University, Chandigarh, India) for using and recording the Raman spectroscopy facility.

Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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Supplementary material

The Supplementary Material for this article can be found online at: https://www.frontiersin.org/articles/10.3389/fmats. 2022.1022317/full#supplementary-material Begum, A., Hussain, A., and Rahman, A. (2012). Effect of deposition temperature on the structural and optical properties of chemically prepared nanocrystalline lead selenide thin films. *Beilstein J. Nanotechnol.* 3, 438–443. doi:10.3762/bjnano.3.50

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