



Inverse Design of Local Solar Flux Distribution for a Solar Methanol Reforming Reactor Based on Shape Optimization

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In this work, an inverse design method that couples the multi-physics model for a solar trough thermochemical reactor (SPTR) and shape optimization model is proposed to find out optimal solar flux distribution for maximizing overall reactor performance. The gradient-based segmentation method is applied to convert the continuous solar flux into step-like flux to guide the concentrator system design. Performance comparisons among uniform flux, linear decreasing flux, and the optimized non-linear flux are also conducted to discuss the reliability of SPTR performance improvement. The results show that the optimized non-linear solar flux can improve the methanol conversion, solar thermochemical conversion, and hydrogen yield of SPTR by 2.5, 3.3, and 2.4%, respectively, compared with the uniform flux. This is attributed to the fact that the optimized non-uniform flux distribution enhances the synergy between temperature and reaction fields, and achieves a better match between spatial solar flux supply and local energy demand by reactions. Also, it is shown that the optimized step-like flux, achieved by regressing the optimized non-linear flux, can perfectly maintain SPTR performance and is effective in boosting SPTR performance under different operating conditions.

Keywords: methanol steam reforming reaction, solar flux distribution, solar thermochemical reactor, shape optimization, multi-physics model

1 INTRODUCTION

The increase in the proportion of solar energy utilization in the energy system will significantly reduce the consumption of fossil fuels, thus protecting the environment and promoting the energy economy (Brockway et al., 2019; Luz and Moura, 2019; He et al., 2020; Mahmoudan et al., 2022). Currently, solar energy utilization technologies include solar thermal power generation (Gilani et al., 2022), photovoltaic (Shakouri et al., 2022), and desalination (Sohani et al., 2022). Among them, solar thermochemical technology to convert solar energy into fuel has the advantages of high energy density, long storage time, and easy transportation (Tang et al., 2022). As an equally clean and high calorific fuel, hydrogen is one of the most desirable energy carriers for future energy systems (Abdalla et al., 2018; Razi and Dincer, 2020). The conversion of solar energy to hydrogen fuel for storing solar energy and hydrogen production can combine the advantages of solar and hydrogen energy, which has already become a focus for scientific research and engineering development (Hosseini and Wahid, 2020).

Currently, hydrogen production by thermochemical conversion driven by solar energy attracts more and more attention (Yadav and Banerjee, 2016; Liu et al., 2019; Zhang et al., 2022). Solar energy can be converted into thermal energy to provide the chemical reaction energy for catalytic hydrogen production such as reforming and pyrolysis of methane and methanol (Ma et al., 2018; Boretti, 2021a, 2021b). Usually, the pyrolysis process and methane reforming require a high temperature of over 800 K, leading to extra structural and cost burden (Wang F. et al., 2014). In contrast, methanol–steam reforming reaction (MSRR) with medium- to low-temperature operation (i.e., 423–573 K) is safer, more stable, and more economical (Kang et al., 2021). Moreover, methanol is a liquid hydrogen carrier that is easy to store and transport, has a high H/C ratio of 4:1, and can produce high-purity hydrogen when needed (Garcia et al., 2021). MSRR shows its significant virtues of high hydrogen selectivity, low carbon monoxide selectivity, and high conversion efficiency (Liu et al., 2016; Garcia et al., 2021), making it a convenient and effective method for hydrogen production in small-scale systems.

Generally, the MSRR system can be applied to parabolic trough collectors (PTC), which are also called solar parabolic trough reactors (SPTR) (Cheng et al., 2019c; Ma et al., 2020; Gharat et al., 2021). Although promising, there are still problems and challenges within the SPTR caused by the non-stationary solar input and non-uniform solar flux distribution, which degrades the reactor performance (He et al., 2019; Ma et al., 2021). There are already some studies on mitigating the effects of non-stationary energy input through active and passive control methods (Saade et al., 2014; Ma et al., 2018; Cheng et al., 2020; Tang et al., 2022). Meanwhile, there are also some studies on enhancing the performance in terms of optic, flow, and heat transfer for boosting SPTR performance (Bellos and Tzivanidis, 2019; Manikandan et al., 2019). For example, Cheng et al. (2019b, 2019a) proposed two kinds of SPTR internal structures to adjust the flow behavior of the reactant mixture inside the SPTR to improve the reactor performance. One is to install a Kenics static mixer inside the reactor tube, which enhances the fluid mixing, reduces the maximum temperature of the reaction bed, and thus increases the methanol conversion (Cheng et al., 2019b). The other one is a catalytic bed with decreasing porosity distribution from top to bottom (Cheng et al., 2019a), which also enhances the reactor's performance. Similarly, Liu et al. (2018) optimized the porosity distribution of the catalytic bed based on the variational method to increase the methanol conversion in a reactor for methanol decomposition.

As for the optical aspect, Gong et al. (2020) proposed a secondary compound parabolic concentrator with enhanced surface flux uniformity, leading to a thermal efficiency improvement of 4.9%. Furthermore, Tang et al. (2021) proposed a broken-line-type secondary concentrator design method based on the flux compensation concept. The newly designed secondary concentrators can achieve over 90% circumferential solar flux uniformity for two kinds of conventional PTCs. Similarly, the secondary uniform reflector proposed by Wang et al. (2014) reduces the circumferential temperature difference. Instead of varying the circumferential

solar flux, Wang et al. (2017) also proposed a concentrator structure with a linear variation of the aperture along the axial direction to vary the axial solar flux distribution to improve the reactor's performance. Kulahli et al. (2019) proposed a new parabolic reflector with varying focal lengths along the length direction of the absorber, which improved the thermal efficiency and net efficiency by 0.21 and 0.63%, respectively.

The aforementioned literature shows that the SPTR performance can be improved to different degrees with different approaches. However, less research has been done to improve the SPTR performance by designing the solar flux spatial distribution of the reactor compared to flow-heat transfer and concentration uniformity improvement (Bellos and Tzivanidis, 2019; He et al., 2020, 2019). There is still a lack of understanding of how surface solar flux distribution affects the performance of SPTR and how to achieve an optimal and reasonable concentrator design. Motivated by these, in this work, we coupled the multi-physics SPTR model with a shape optimization model for solar flux optimization. With the model, the optimal solar flux distribution along the tube with fixed total energy input can be achieved. The gradient-based segmentation method is adopted to convert continuous non-linear solar flux into step-like solar flux distribution, which can be used to guide the design of the concentrator. Finally, the performance enhancement of SPTR by optimized step-like flux is verified under different operating conditions. A summary figure of the entire study to aid understanding can be found in **Supplementary Figure S1**.

2 MATHEMATICAL MODEL AND OPTIMIZATION METHOD

2.1 Model Description

As shown in **Figure 1**, the solar parabolic trough methanol steam reforming reactor (SPTR) in this study mainly comprises three parts: the concentrator system for sunlight collector, the vacuum absorber for solar-to-thermal energy conversion, and the reaction tube for thermal-to-chemical energy conversion, respectively. Our previous work (Tang et al., 2021) showed that a well-designed secondary reflector for PTC can nearly achieve a circumferentially uniform distribution of solar flux on the absorber's surface. As a result, it is reliable and easy to build a 2-D axisymmetric SPTR model, as shown in **Figure 1B** by considering uniform circumferential solar flux distribution. In the SPTR, methanol and steam enter through the inlet, absorb solar heat and react on the reaction particle bed consisting of Cu/ZnO/Al₂O₃ catalyst (Peppley et al., 1999), and finally produce fuel products with solar energy stored. The relevant structural and physical parameters of SPTR are listed in **Table 1**, and a detailed description of the concentrator structure can be found elsewhere (Tang et al., 2021). It should be mentioned that the axial solar flux is not necessarily uniform to boost the reactor performance, which will be optimized in this study.

To build the multi-physics model for SPTR, it is useful and necessary to introduce reasonable assumptions in order to save computational resources while ensuring sufficient realism. The present 2-D multi-physics SPTR model is developed based on the

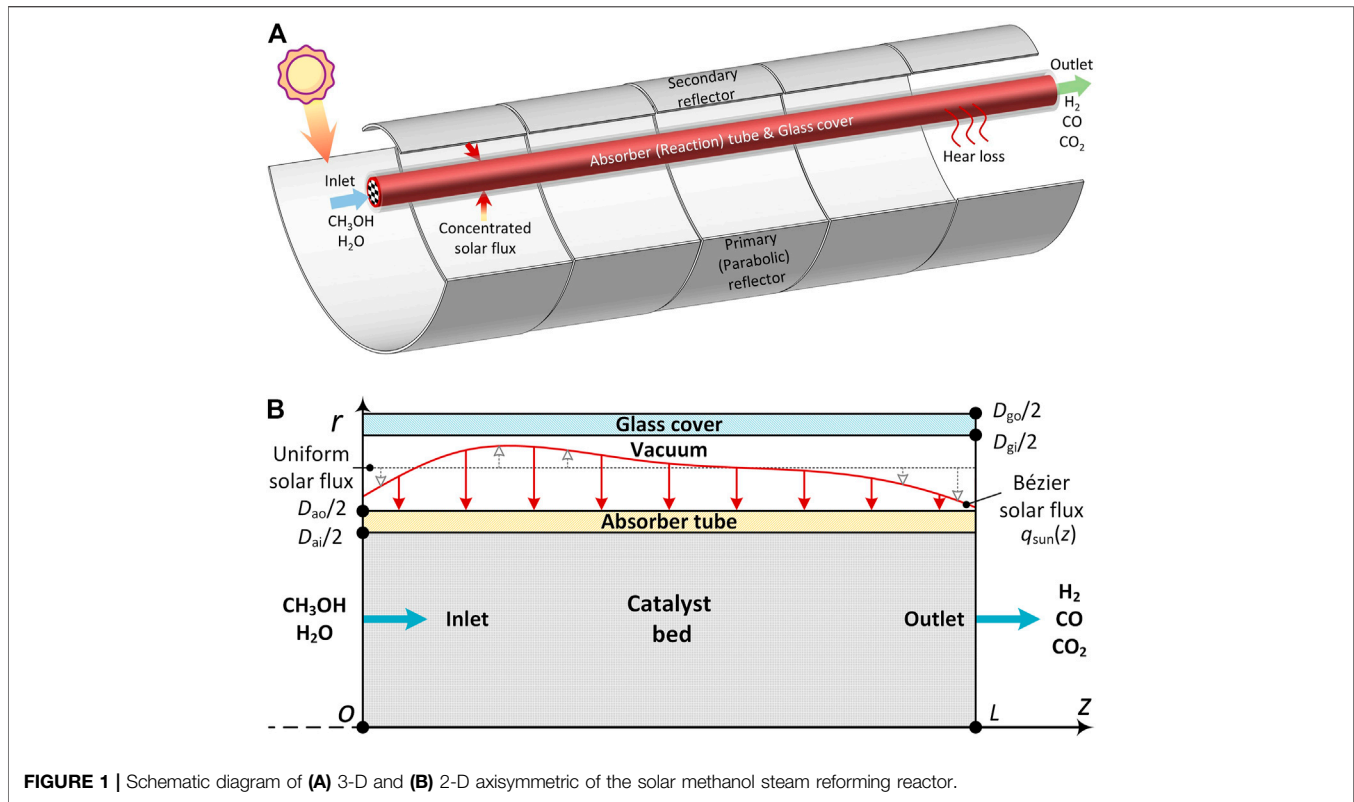


FIGURE 1 | Schematic diagram of **(A)** 3-D and **(B)** 2-D axisymmetric of the solar methanol steam reforming reactor.

TABLE 1 | Relevant structural and physical parameters of SPTR (Ma et al., 2018).

Parameter	Value
ρ_p	1,300 kg m ⁻³
$C_{p,p}$	542 J kg ⁻¹ K ⁻¹
λ_p	20 W m ⁻¹ K ⁻¹
ε_p	0.4
D_{ei}	66 mm
D_{ao}	70 mm
L	5 m
D_{gi}	110 mm
D_{go}	115 mm

following assumptions: 1) the catalytic reaction bed is an isotropic homogeneous porous medium, 2) the gas mixture flows in a steady laminar flow, 3) the heat loss is only for radiation and convection with the atmosphere, and 4) the local thermal equilibrium assumption in the porous reaction bed is valid (Liu et al., 2016). In the following, we briefly present the governing equations for simulating the mass/species transfer, fluid flow, heat transfer, and chemical reactions of the gas mixture.

2.1.1 Mass Conservation

The mass conservation equation for a gas mixture in the porous reaction bed is:

$$\nabla \cdot (\rho_m \vec{U}) = 0 \quad (1)$$

where \vec{U} is the superficial velocity vector and ρ_m is the density of the gas mixture. For simplicity, the gas mixture is usually assumed to be ideal gas expressed as:

$$\rho_m = \frac{p\bar{M}}{R_g T} \quad (2)$$

where p and \bar{M} indicate the pressure and the average molecular weight of the gas mixture, respectively.

2.1.2 Momentum Conservation

The momentum conservation equation for fluid flow in the porous reaction bed can be expressed as:

$$\frac{\rho_m}{\varepsilon_p^2} (\vec{U} \cdot \nabla) \vec{U} = -\nabla p + \nabla \cdot \left\{ \frac{1}{\varepsilon_p} \left[\mu_m (\nabla \vec{U} + (\nabla \vec{U})^T) - \frac{2}{3} \mu_m (\nabla \cdot \vec{U}) \vec{I} \right] \right\} + \vec{F} \quad (3)$$

where \vec{F} is the resistance term and μ_m is the dynamic viscosity of the gas mixture which is described in the Wilke (1950) equation, that is:

$$\mu_m = \frac{\sum_{i=1}^n x_i \mu_i}{\sum_{j=1}^n x_j \phi_{ij}} \quad (4)$$

$$\phi_{ij} = \frac{1}{\sqrt{8}} \left[1 + \left(\frac{\mu_i}{\mu_j} \right)^{1/2} \left(\frac{M_i}{M_j} \right)^{1/4} \right]^2 \left(1 + \frac{M_i}{M_j} \right)^{-1/2} \quad (5)$$

where x_i , μ_i , and M_i refer to the molar fraction, viscosity, and molecular weight of species i , respectively. The species i denotes one of the reaction gases, such as CH_3OH , H_2O , H_2 , CO , and CO_2 .

The resistance term \vec{F} in Eq. 3 can be expressed as (ERGUN, 1952):

$$\vec{F} = -\frac{150\mu_m(1-\varepsilon_p)^2}{\varepsilon_p^3 d_p^2} \vec{U} - \frac{1.75\rho_m(1-\varepsilon_p)}{\varepsilon_p^3 d_p} |\vec{U}| \vec{U} \quad (6)$$

where d_p is the average porous media pore size, which is related to the porosity ε_p :

$$\varepsilon_p = 0.375 + \frac{0.34}{D_{ai}/d_p} \quad (7)$$

2.1.3 Energy Conservation

Considering the assumption of local thermal equilibrium, there is no temperature difference between porous media and gas mixture at the same position. The energy conservation equation for the reaction bed region can be given by:

$$\rho_m C_{p,m} \vec{U} \cdot \nabla T = \nabla \cdot (\lambda_{\text{eff},p} \nabla T) + S_{\text{chem}} \quad (8)$$

where $C_{p,m}$ is the specific heat capacity of the gas mixture, $\lambda_{\text{eff},p}$ is the effective thermal conductivity of the porous reaction bed, and S_{chem} denotes the chemical reaction source term.

The effective parameters of the porous media domain can be expressed by the volume averaging method as (Elbahjaoui and El Qarnia, 2019):

$$\lambda_{\text{eff},p} = (1-\varepsilon_p)\lambda_p + \varepsilon_p\lambda_m \quad (9)$$

$$C_{p,\text{eff},p} = (1-\varepsilon_p)C_{p,p} + \varepsilon_p\lambda_m \quad (10)$$

where the effective parameters of the gas mixture are expressed by Poling et al. (2001) as:

$$\lambda_m = \frac{\sum_{i=1}^n x_i \lambda_i}{\sum_{j=1}^n x_j \phi_{ij}} \quad (11)$$

$$C_{p,m} = \frac{\sum_{i=1}^n c_i C_{p,i}}{\sum_{j=1}^n c_j M_j} \quad (12)$$

where the correlation properties of pure species i are defined as temperature-related Eqs 13–15, and the correlation coefficients are listed in **Supplementary Tables S1–2** (Gordon and McBride, 1972).

$$\lambda_i = A_{1,i} + A_{2,i}T + A_{3,i}T^2 + A_{4,i}T^3 \quad (13)$$

$$C_{p,i} = R_g(B_{1,i} + B_{2,i}T + B_{3,i}T^2 + B_{4,i}T^3 + B_{5,i}T^4) \quad (14)$$

$$h_i = R_g\left(B_{1,i}T + \frac{B_{2,i}}{2}T^2 + \frac{B_{3,i}}{3}T^3 + \frac{B_{4,i}}{4}T^4 + \frac{B_{5,i}}{5}T^5 + B_{6,i}\right). \quad (15)$$

2.1.4 Species Conservation

The concentration of each species in the reaction system is represented by the mass transfer equation:

TABLE 2 | SPTR model boundary conditions.

Domain	Porous Reaction Bed	Solid Tube
Inlet ($z = 0$)	$m = m_{\text{in}}, T = T_{\text{in}} = 423.15 \text{ K}, \text{C}_{\text{H}_2\text{O}}: \text{C}_{\text{CH}_3\text{OH}} = 1.1$	$\frac{\partial T}{\partial z} = 0$
Outlet ($z = L$)	$\rho = \rho_{\text{out}} = 101325 \text{ Pa}$, inhibits reflux	
Center ($r = 0$)	Axisymmetric boundary: $\frac{\partial T}{\partial r} = \frac{\partial U}{\partial r} = \frac{\partial c_i}{\partial r} = 0$	
Inner ($r = D_{\text{ai}}$)	No-slip and no-flux boundary: $U = \frac{\partial c_i}{\partial r} = 0$	
Wall ($r = D_{\text{ao}}$)	Radiation q_{rad} and heat flux $q_{\text{sun}}(z)$ boundary	
Glass ($r = D_{\text{go}}$)	Radiation q_{rad} and convection q_{con} boundary	

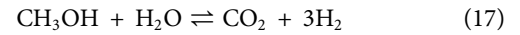
$$\nabla \cdot (-D_i \nabla c_i) + \vec{U} \cdot \nabla c_i = R_i \quad (16)$$

where c_i is the concentration of species i , D_i is the mass diffusion coefficient of species i , and R_i denotes the chemical reaction source term of species i .

2.1.5 Reaction Kinetics

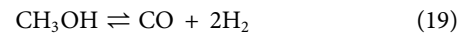
The MSRR reaction system with three reactions is present in SPTR, and the reaction kinetics (Peppley et al., 1999) are shown below:

1) Methanol steam reforming (MSR)



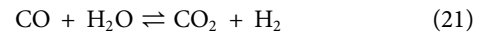
$$R_{\text{MSR}} = \frac{k_R K_{\text{CH}_3\text{O}^{(1)}} \frac{p_{\text{CH}_3\text{OH}}}{p_{\text{H}_2}^{1/2}} \left(1 - \frac{p_{\text{H}_2}^3 p_{\text{CO}_2}}{K_R p_{\text{CH}_3\text{OH}} p_{\text{H}_2\text{O}}}\right) C_{\text{S}_1} C_{\text{S}_{1a}} S_p \rho_p}{\left(1 + K_{\text{CH}_3\text{O}^{(1)}} \frac{p_{\text{CH}_3\text{OH}}}{p_{\text{H}_2}^{1/2}} + K_{\text{HCOO}^{(1)}} p_{\text{CO}_2} p_{\text{H}_2}^{1/2} + K_{\text{OH}^{(1)}} \frac{p_{\text{H}_2\text{O}}}{p_{\text{H}_2}}\right) \left(1 + K_{\text{H}^{(1)}}^{1/2} p_{\text{H}_2}^{1/2}\right)} \quad (18)$$

2) Methanol decomposition (MD)



$$R_{\text{MD}} = \frac{k_D K_{\text{CH}_3\text{O}^{(2)}} \frac{p_{\text{CH}_3\text{OH}}}{p_{\text{H}_2}^{1/2}} \left(1 - \frac{p_{\text{H}_2}^2 p_{\text{CO}}}{K_D p_{\text{CH}_3\text{OH}}}\right) C_{\text{S}_2} C_{\text{S}_{2a}} S_p \rho_p}{\left(1 + K_{\text{CH}_3\text{O}^{(2)}} \frac{p_{\text{CH}_3\text{OH}}}{p_{\text{H}_2}^{1/2}} + K_{\text{OH}^{(2)}} \frac{p_{\text{H}_2\text{O}}}{p_{\text{H}_2}^{1/2}}\right) \left(1 + K_{\text{H}^{(2)}}^{1/2} p_{\text{H}_2}^{1/2}\right)} \quad (20)$$

3) Water–gas shift (WGS)



$$R_{\text{WGS}} = \frac{k_W K_{\text{OH}^{(1)}} \frac{p_{\text{CO}} p_{\text{H}_2\text{O}}}{p_{\text{H}_2}^{1/2}} \left(1 - \frac{p_{\text{H}_2} p_{\text{CO}_2}}{K_W p_{\text{CO}} p_{\text{H}_2\text{O}}}\right) C_{\text{S}_1}^2 S_p \rho_p}{\left(1 + K_{\text{CH}_3\text{O}^{(1)}} \frac{p_{\text{CH}_3\text{OH}}}{p_{\text{H}_2}^{1/2}} + K_{\text{HCOO}^{(1)}} p_{\text{CO}_2} p_{\text{H}_2}^{1/2} + K_{\text{OH}^{(1)}} \frac{p_{\text{H}_2\text{O}}}{p_{\text{H}_2}}\right)^2} \quad (22)$$

where the relevant kinetic parameters are listed in **Supplementary Table S3**, and the detailed definitions are described in Peppley et al. (1999).

Combined with the reaction kinetics, the reaction source term in the energy equation and the species equations can be expressed as:

$$S_{\text{chem}} = \sum_r R_r \Delta H_r \quad (23)$$

$$R_i = \sum_r n_{i,r} R_r \quad (24)$$

where $n_{i,r}$ is the stoichiometric number and R_r is the reaction rate expressions shown in Eqs 18, 20, 22.

2.1.6 Boundary Conditions

In the computational domain of Figure 1B, the corresponding boundary conditions are listed in Table 2.

The surface-to-surface radiation exists between the absorber tube and the glass cover, and the diffuse surface flux is expressed as:

$$q_{\text{rad}} = \varepsilon_e (G - \sigma T^4) \quad (25)$$

where G is the irradiation and ε_e is the emissivity of the diffuse surface. The surface emissivity is set to 0.9 for the glass cover and to $\varepsilon_e = 0.00042T_{\text{ao}} - 0.0995$ for the selective coating on the absorber tube wall (Cheng et al., 2019c).

The radiation from the glass cover to the ambient is calculated by Stefan–Boltzmann’s law. Besides, there are natural convection losses with the ambient, and the heat transfer coefficient is calculated by the horizontal cylinder formula in Bergman et al. (2011), expressed as:

$$q_{\text{con}} = h(T_{\text{amb}} - T) \quad (26)$$

$$h = \frac{\lambda}{D_{\text{go}}} \left(0.6 + \frac{0.387\text{Ra}_{D_{\text{go}}}^{1/6}}{\left[1 + \left(\frac{0.559}{\text{Pr}} \right)^{9/16} \right]^{8/27}} \right)^2 \quad (27)$$

where T_{amb} is the ambient temperature and the air parameters in Eq. 26 come from the COMSOL material library.

2.2 Solar Flux Optimization Method

The solar flux distribution is optimized to get high-quality energy conversion and improve the overall SPTR performance by achieving a better match between the energy field and the reaction field. The primary performance evaluation parameters of SPTR are the methanol conversion efficiency η_m and the solar thermochemical conversion efficiency η_s , given below:

$$\eta_m = \frac{m_{\text{CH}_3\text{OH}, \text{in}} - m_{\text{CH}_3\text{OH}, \text{out}}}{m_{\text{CH}_3\text{OH}, \text{in}}} \times 100\% \quad (28)$$

where $m_{\text{CH}_3\text{OH}, \text{in}}$ and $m_{\text{CH}_3\text{OH}, \text{out}}$ indicate the inlet and outlet methanol flow rate, respectively.

$$\eta_s = \frac{Q_{\text{chem}}}{Q_{\text{sun}}} = \frac{\iiint S_{\text{chem}} dV}{q_{\text{sun}}(z)\pi D_{\text{ao}}L} \times 100\% \quad (29)$$

where Q_{chem} and Q_{sun} denote the thermochemical reaction energy and the total solar input energy, respectively.

Typically, one of the most important SPTR performances is the methanol conversion efficiency η_m , and other evaluation parameters are directly or indirectly related to the conversion rate, such as the hydrogen yield m_{H_2} . Thus, the conversion efficiency η_m is chosen as the maximization target in this study.

For solar flux distribution, it is described by the Bezier curve. This curve is a way that can describe various types of shapes with mathematical expressions, and it has been widely used in scenarios such as airfoil optimization (Wei et al., 2020) and robot motion control (Song et al., 2021). Furthermore, the

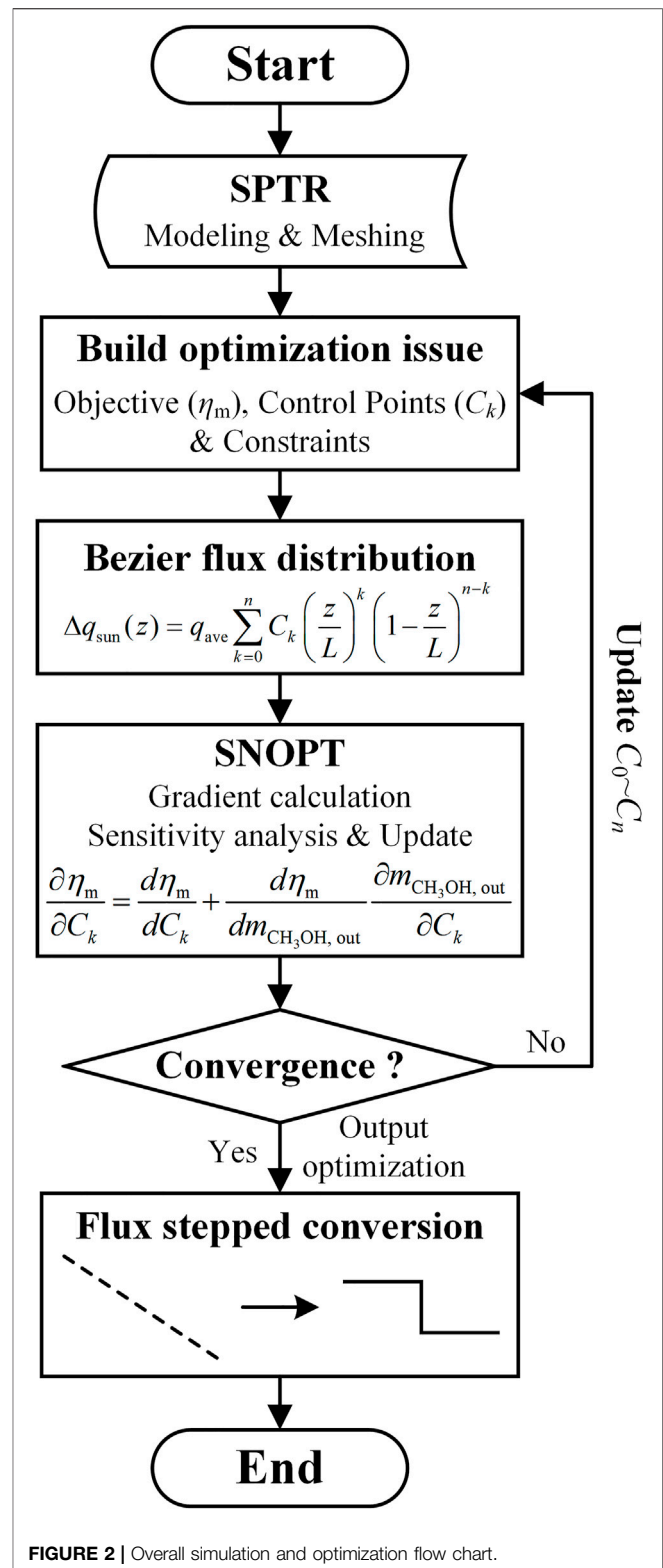


FIGURE 2 | Overall simulation and optimization flow chart.

curve has the property of high-order derivability, which is well suited for the gradient-based optimization algorithm in this study. The corresponding n th-order Bernstein basis function

representing the deviation $\Delta q_{\text{sun}}(z)$ of solar flux relative to the uniform flux q_{ave} is expressed as:

$$\Delta q_{\text{sun}}(z) = q_{\text{ave}} \sum_{k=0}^n C_k \left(\frac{z}{L}\right)^k \left(1 - \frac{z}{L}\right)^{n-k} \quad (30)$$

where C_k is the control point of the curve and is also the control variable.

Since the sintering temperature of the Cu/ZnO/Al₂O₃ catalyst is at 573.15 K (Yong et al., 2013), a maximum temperature constraint is required. Furthermore, a point-by-point constraint of constant surface solar flux greater than 0 is set to prevent unreasonable negative flux values. Also, in order to accelerate the optimization, an advance constraint of greater than 0 at the control points at both ends, that is, $C_0, C_n \geq -1$ is imposed. To sum up, for a certain total solar input, the optimization problem can be concluded as:

$$\left\{ \begin{array}{l} \max_{C_0-C_n} : \eta_m = \frac{m_{\text{CH}_3\text{OH}, \text{in}} - m_{\text{CH}_3\text{OH}, \text{out}}}{m_{\text{CH}_3\text{OH}, \text{in}}} \\ \text{s.t.} : T_{\text{max}} \leq 573.15 \text{ K} \\ C_0, C_n \geq -1 \\ q_{\text{sun}}(z) \geq 0 \\ \oint q_{\text{sun}}(z) dA = \frac{\pi D_{\text{ao}}^2 L q_{\text{ave}}}{4} \end{array} \right. \quad (31)$$

The optimization process of the solar flux distribution is solved by the Sparse Nonlinear Optimizer (SNOPT) based on Sequential Quadratic Programming, an algorithm suitable for solving large-scale constrained optimization problems with smooth non-linear functions in the objective and constraints (Gill et al., 2005). The overall process of SNOPT-based optimization of SPTR solar flux distribution is shown in **Figure 2B**.

COMSOL is a powerful numerical analysis software for multi-physics fields based on the finite element method, covering various fields such as mechanics, fluids, electromagnetism, heat transfer, chemicals, electrochemistry, and acoustics. It is especially suitable for developing models involving multiple physical fields (Fu et al., 2022; Tang et al., 2022). Calculations were performed on a computer with an 8-core 3.6 GHz CPU and 32 GB RAM, and the entire SPTR multi-physics model was solved iteratively by the full-coupling method in COMSOL (about 10 min), and the solar flux distribution is optimally solved by SNOPT (about 133 min). The convergence criterion for the above two solution processes was determined to be 10^{-5} by the tolerance test.

2.3 Model Validation

For grid types, structured grids have the advantages of fast generation, good quality, low memory consumption, and are suitable for regular and simple computational domains (e.g., **Figure 1B**). Thus, custom meshing in COMSOL is used to generate structured grids of different sizes for comparison, and the comparison results are shown in **Figure 3**. The grid group of $56 \times 1,000$ has only a relative error of 10^{-5} compared to the group

of $76 \times 1,500$, which can be considered since the grid system of $56 \times 1,000$ is accurate enough for this study.

To verify the accuracy of the current multi-physics field model, the flow-heat exchange-radiation component is validated with the experimental results of Dudley et al. (1994). As shown in **Table 3**, the results obtained using the present model under different conditions show great agreement with the experimental results.

For chemical validation, the present model was compared with the experimental study by Peppley et al. (1999) and the simulation study by Cheng et al. (2019c). As shown in **Figure 4**, the methanol conversion of the present model matches each other with the results in Cheng et al. (2019c) and Peppley et al. (1999), which especially agrees well with the results of Cheng et al. (2019c). The aforementioned validation shows the reliability of the model.

3 RESULTS AND DISCUSSION

3.1 Optimization Results of Solar Flux Distribution

To obtain the optimal concentrator combination for improving the SPTR performance, this section first optimized the axial distribution of solar flux on the SPTR at an inlet reaction gas mass flow rate of 0.012 kg s^{-1} , an inlet steam methanol ratio of 1.1, and an average surface solar flux input of $13,000 \text{ W m}^{-2}$. In the optimization process, two types of solar flux distributions are considered, which are linear and non-linear. The linear distribution corresponds to the 1st-order Bezier function, that is, n equals 1 in **Eq. 30**, while the non-linear distribution is set to the 5th-order Bezier function, corresponding to n equals 5.

Figure 5 shows the hydrogen generation rate field in SPTR for uniform solar flux distribution and two optimized distributions, where the relevant control parameters of Bernstein basis functions corresponding to the optimized distributions are listed in **Table 4**. For the distribution characteristics, the final optimization result of the linear flux distribution is decreasing from inlet to outlet, and this result validates each other with the conclusion of Wang et al. (2017), showing the feasibility of the optimization model. However, the optimal non-linear distribution shows an arch-shaped distribution with the solar flux converging in the middle and less on the sides, as in **Figure 5C**. For the hydrogen generation rate, both optimized solar flux distributions produce a higher hydrogen reaction zone ($r_{\text{H}_2} > 60 \text{ mol m}^{-3} \text{ m}^{-1}$) compared to the relatively low uniform generation rate ($r_{\text{H}_2} < 60 \text{ mol m}^{-3} \text{ m}^{-1}$) under the uniform distribution, while it shifts with the location of the energy concentration. Intuitively, the non-linear optimization has a higher and larger high reaction zone, indicating stronger reaction dynamics.

To clearly observe the changes in comprehensive performance before and after flux optimization, **Figure 6** synthesizes the comparative results of SPTR performance under different solar flux distributions. Apparently, both non-uniform optimized solar flux distributions exhibit better SPTR performance than uniform distributions, as reflected by higher methanol conversion, solar thermochemical efficiency, and hydrogen yield. Particularly, the

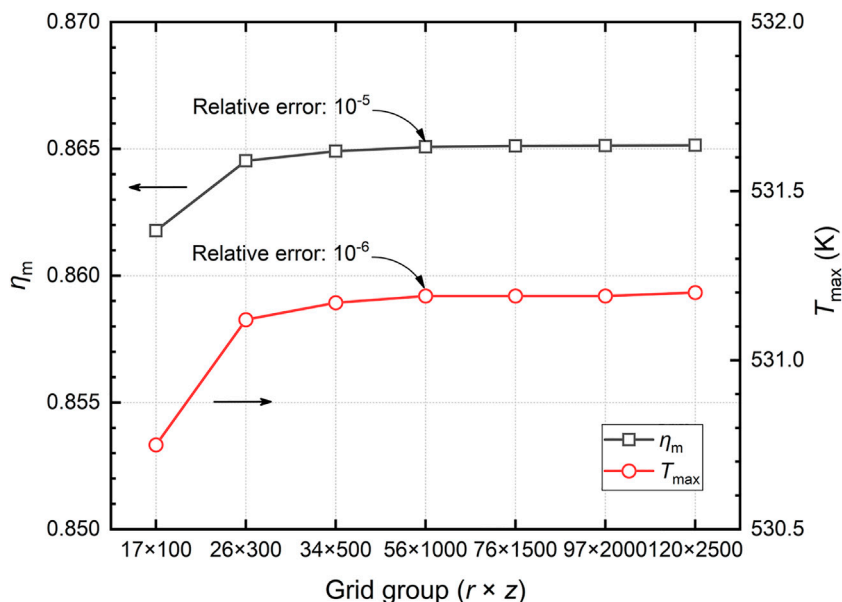


FIGURE 3 | Grid independence verification.

TABLE 3 | Comparative verification of the flow-heat exchange-radiation component with experimental results (Dudley et al., 1994).

Case	u_w (m/s)	T_{in} (K)	T_a (K)	V_{in} (L/min)	DNI (W/m^2)	$T_{o,exp}$ (K)	$T_{o,num}$ (K)	Error (%)
1	2.6	375.35	294.35	47.7	933.7	397.15	397.60	0.11
2	1	570.95	301.95	55.5	937.9	590.05	590.68	0.11
3	2.6	652.65	302.65	56.8	920.9	671.15	671.81	0.10
4	3.7	424.15	295.55	47.8	968.2	446.45	446.99	0.12
5	2.5	470.65	297.45	49.1	982.3	492.65	493.31	0.13
6	3.3	523.85	299.35	54.7	909.5	542.55	542.85	0.06
7	2.9	572.15	300.65	55.6	880.6	590.35	590.40	0.01
8	4.2	629.05	304.25	56.3	903.2	647.15	647.74	0.09

enhancement reaches the maximum under the non-linear flux distribution as shown in **Figure 5C**, where the methanol conversion, solar thermochemical efficiency, and hydrogen yield are enhanced by 2.5, 3.3, and 2.4%, respectively, relative to the uniform flux. Although there is a relative increase of 19.0% in CO yield, the slow methanol decomposition reaction rate and the CO consumption of the water–gas shift reaction result in a CO yield of only about 1/10 of the H₂ yield. Thus, the discussion later focuses on the primary target product H₂ rather than the by-product CO. Moreover, the maximum catalytic bed temperatures are all below the catalyst sintering temperature of 573 K, which implies a long catalyst life and good stability of SPTR performance (Yong et al., 2013).

To clarify the reason for the performance improvement, **Figure 7** shows the average temperature, average hydrogen generation rate, and average methanol concentration distribution along the flow direction for the three solar flux distributions. To facilitate the analysis, the catalytic bed is divided into three zones (Z₁–Z₃) using the intersection of hydrogen generation rates as the boundary, as shown in the

dashed line in **Figure 7**. As seen in **Figure 7B**, both optimized flux distributions have more concentrated fast reaction rate zones due to the concentrated solar energy distribution, which correspondingly appear in the Z₁ and Z₂ zone. With the linearly optimal solar flux distribution, more energy is invested in the front zone (Z₁), which raises the reaction temperature more rapidly and reaches the higher fast reaction zone faster compared to the uniform distribution. However, in the Z₃ zone, as seen in **Figure 7A**, only a small amount of reactants are present, where more energy will be used to heat the fluid (e.g., the uniform case in **Figure 7A**) apart from making the reaction complete. Clearly, neither the linear nor the non-linear optimized distribution allocates much energy input in Z₃. Under the non-linear optimal solar flux distribution, more energy is distributed in the intermediate zone (Z₂). Simultaneous temperature increases with progressively increasing solar input in Z₁ at non-linear flux, at which the temperature level of Z₂ (480–526 K) surpasses the 423–498 K of the optimal reaction zone Z₁ at linear flux. Further, 80% of the reactant molar fraction

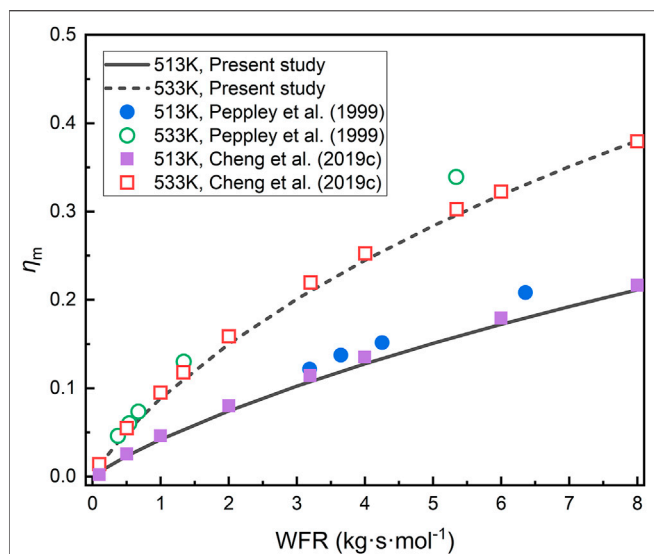


FIGURE 4 | Comparative verification of the chemical part with the experimental results (Peppley et al., 1999) and simulation results (Cheng et al., 2019c).

TABLE 4 | Control parameters of Bernstein basis functions for optimized solar flux distribution.

Parameter	C_0	C_1	C_2	C_3	C_4	C_5
Value	-0.92395	-8.1315	20.5375	25.14	-5.088	-0.99995

remained when the best reaction zone Z_2 reached under non-linear flux, thus being sufficient to show higher reaction performance in terms of temperature and reactant-dominated reaction kinetics.

The relationship between the variable distributions in **Figure 7** shows that there is a degree of similarity under the non-linear case, so **Figure 8** compares the synergy between the fields by the correlation index described by the Pearson correlation coefficient P_r of Eq. 32:

$$P_r = \frac{\sum_{k=1}^n (x_k - \bar{x})(y_k - \bar{y})}{\sqrt{\sum_{k=1}^n (x_k - \bar{x})^2 \sum_{k=1}^n (y_k - \bar{y})^2}} \quad (32)$$

where x and y denote the two discrete variables for comparison, n is the number of discrete variables, and $\bar{}$ represents the average of the corresponding variables.

Usually, the closer P_r is to 1, the higher the positive correlation is, especially when the sample of variables is constant, P_r is 0 (e.g., the solar flux in the uniform case in **Figure 8** is constant). As seen in **Figure 8**, the correlations under the non-uniform optimal distribution are all the highest, showing the strongest synergistic effect to enhance the SPTR performance. Here, the solar flux as an external variable directly affects the temperature change, while the temperature and the reaction rate interact with each other, which ultimately manifests as the effect of the solar flux and the reaction rate, that is, the relationship between the solar energy supply and the reaction energy demand. It should be noted that due to the dominant role of MSR in the SPTR reaction system, the hydrogen reaction rate was found to almost represent the reaction heat consumption in this study.

Overall, the rational distribution of non-linear optimization achieves hysteresis matching of temperature distribution, which in turn achieves energy flow matching of energy input and consumption (i.e., reaction rate) through the synergistic relationship between reaction rate and temperature with P_r up to 0.99.

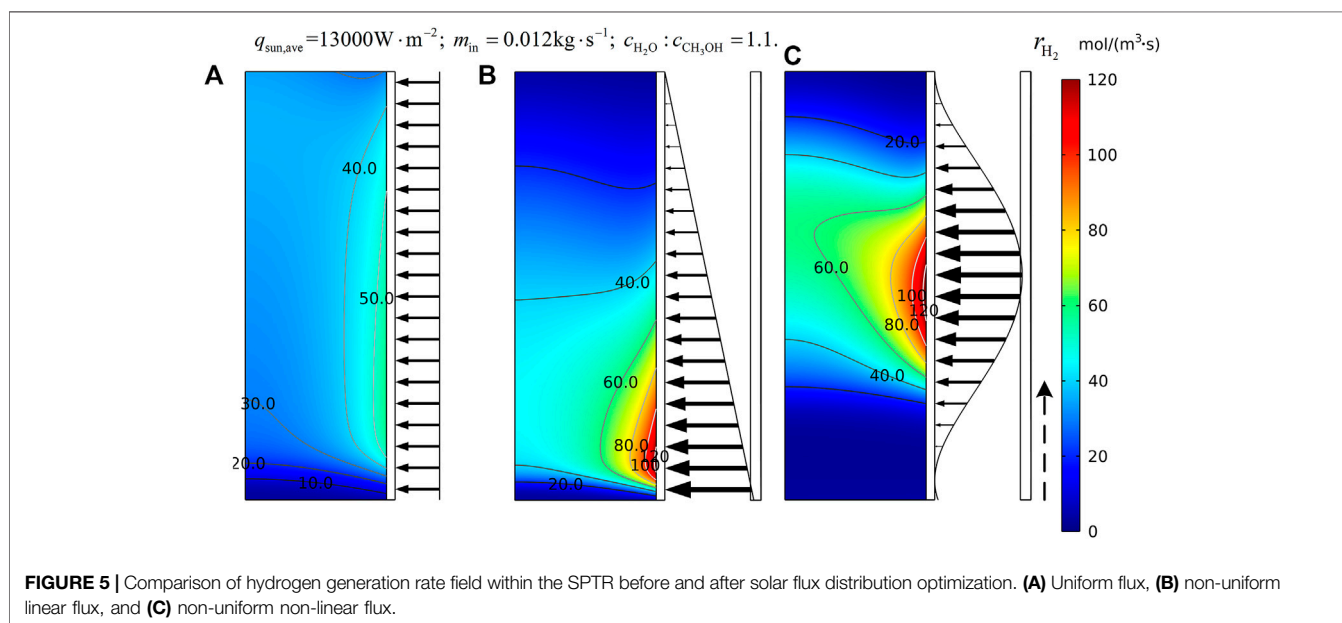


FIGURE 5 | Comparison of hydrogen generation rate field within the SPTR before and after solar flux distribution optimization. (A) Uniform flux, (B) non-uniform linear flux, and (C) non-uniform non-linear flux.

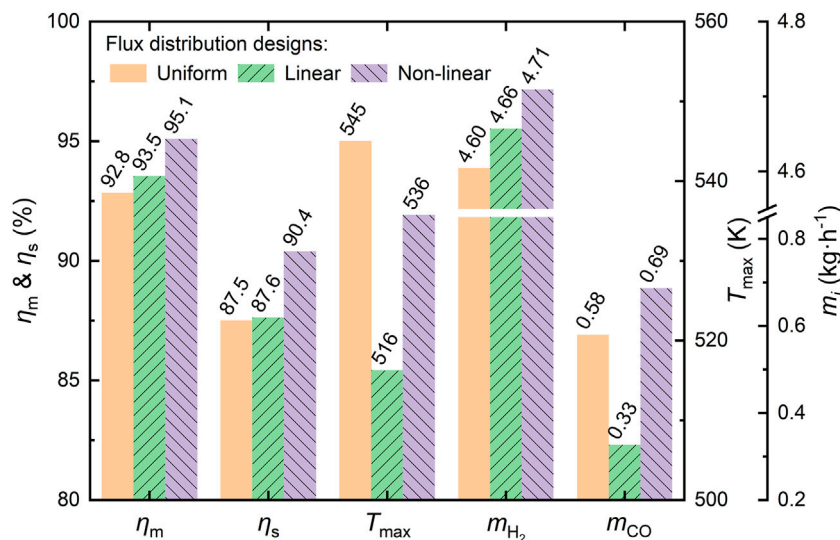


FIGURE 6 | Comparison of SPTR comprehensive performance before and after solar flux distribution optimization.

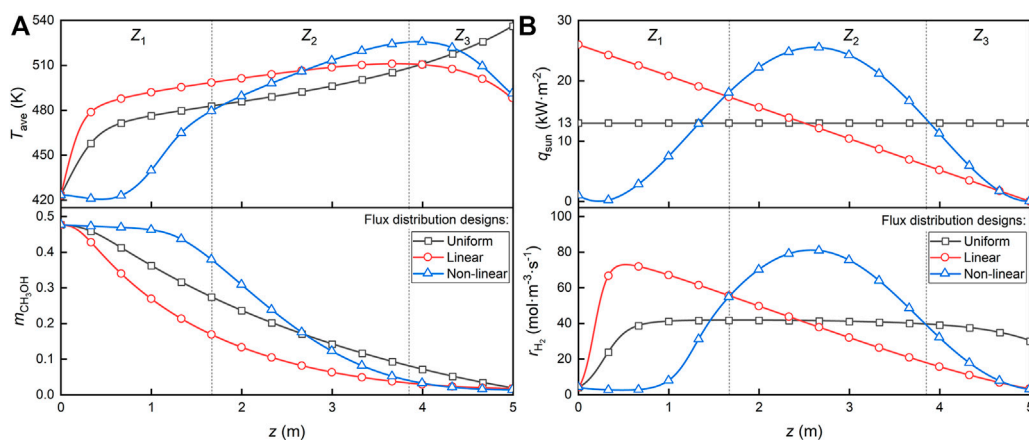


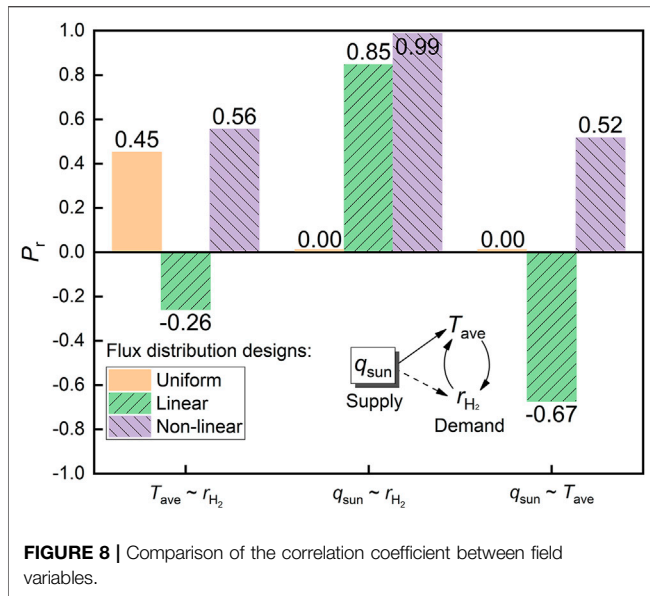
FIGURE 7 | Distribution of the average values of the variables along the flow direction. (A) Temperature and methanol molar fraction and (B) solar flux and hydrogen generation rate.

3.2 Regression of Real Concentrator

Considering the engineering reality, the current PTC cannot achieve continuous concentrated solar flux, so it is necessary to convert the continuous solar flux distribution into a step-like one in order to get a concentrator group with different apertures and finally get the design guidance of concentrator in SPTR. Since the step-like distribution can be viewed as having an infinite gradient at each junction, the gradient-based segmentation method is used in this study to segment the continuous solar flux distribution to ensure that the features of the distribution are fully preserved.

After comparative consideration, 100 and 75% of the maximum gradient are selected as the splitting points to split the non-linear flux into a step-like flux, as in Figure 9A.

Figure 9B summarizes the corresponding field variable changes and performance diagrams of SPTR under the step-like solar flux distribution. The characteristic of variable change can be observed, where after each flux change, the temperature and reaction rate first have a large rate of change and then gradually level off. In Figure 9B, there is an extremely high correlation of over 0.99 for the results of variable changes between the corresponding step-like and continuous type flux. Also, the SPTR performance is only slightly degraded by about 0.2% due to the loss of some continuous curve properties by the step-like shape, which also shows the effectiveness of the gradient-based segmentation method in this study.



3.3 Comparison of Performance Enhancement Under Different Conditions

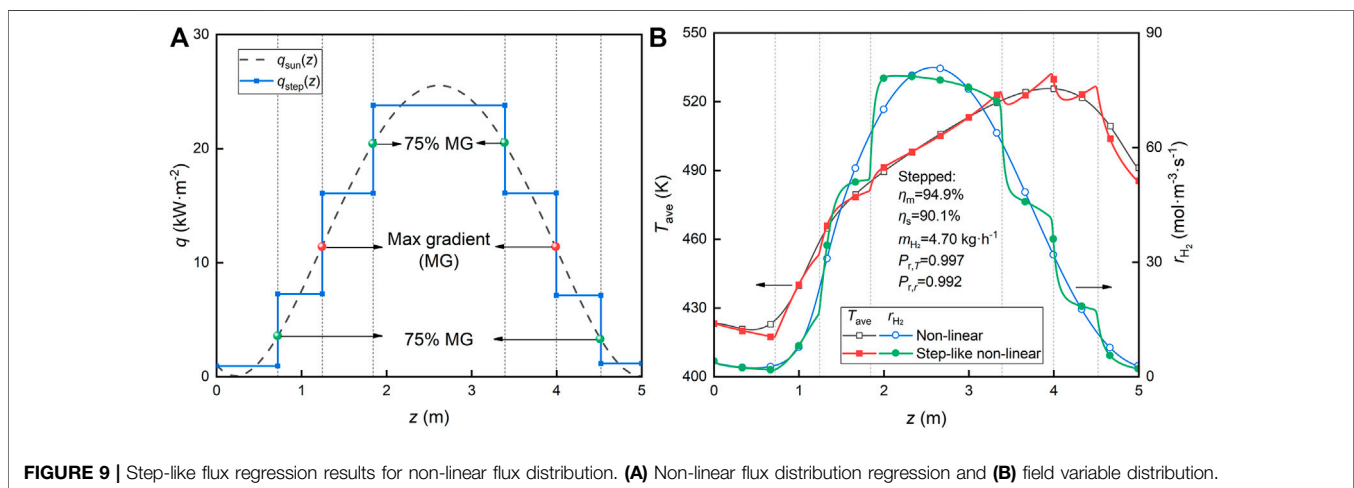
To check the performance effect of the optimized solar flux distribution form (i.e., the step-like solar flux) under different conditions, this section explores the SPTR performance at different reactant mass flow rates and DNIs, which map the demand side and supply side in SPTR, respectively. Different DNIs correspond to different solar average fluxes, where $DNI = 1 \text{ kW m}^{-2}$ corresponds to a solar average flux of 13 kW m^{-2} , implying an average concentration ratio of 13. Here, the solar average flux of $4\text{--}14 \text{ kW m}^{-2}$ and the reactant mass flow rate of 0.012 kg s^{-1} are applied.

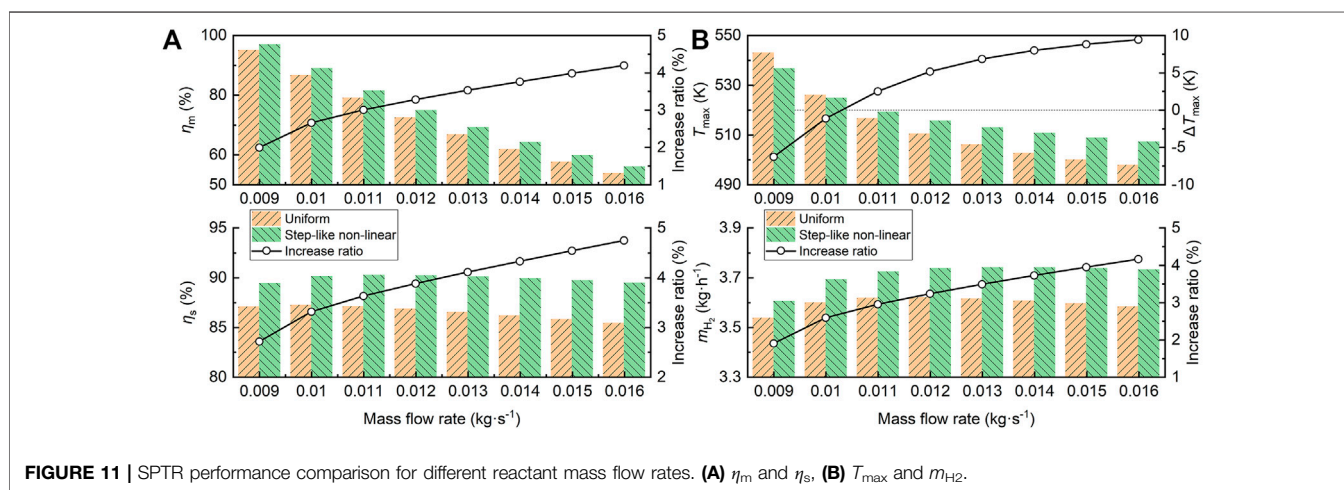
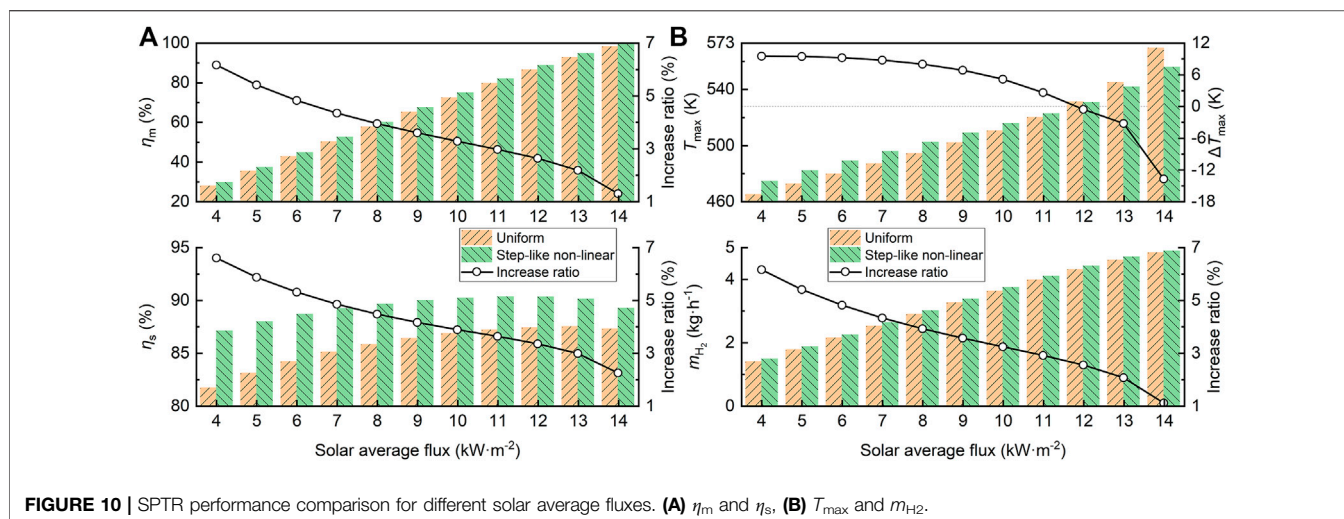
The performance comparison for different solar average fluxes is shown in **Figure 10**, where the methanol conversion η_m , hydrogen generation rate m_{H_2} , and maximum catalytic bed temperature T_{max} all increase with increasing input solar

energy, due to the increased total input allowing more chemical and thermal energy conversion. However, the solar thermochemical conversion efficiency η_s decreases as the input energy increases when it reaches a high conversion efficiency η_m , mainly because the chemical energy that can be converted reaches the limit. Furthermore, at different solar flux, the optimized flux all have better performance improvement compared to uniform flux, with the increase ratio (relative percentage improvement of corresponding performance indicator in present figure) increasing as the solar flux decreases. Together with the temperature variation in **Figure 10B**, where the reaction is weak when the energy input is small, the central concentration feature of the optimized solar flux distribution results in a higher temperature level and thus a more significant performance improvement. However, when the energy input is sufficient, the reaction proceeds completely and the energy input at the end flows to the thermal energy; thus, the optimized solar flux with less end distribution allows more energy to be used in the high reaction zone and reduces the temperature level.

For different mass flow rates, which represent the energy consumed on the demand side, $0.009\text{--}0.016 \text{ kg s}^{-1}$ are chosen for the analysis, with a solar average flux of 10 kW m^{-2} . The performance comparison summarized in **Figure 11** shows that as the reactant mass flow rate increases, η_m and T_{max} decrease due to the increase in the flow rate and the decrease in the reaction residence time. However, m_{H_2} and η_s are essentially stable because the energy input determines the hydrogen production when the total amount of reactants is greater than the energy input. Similarly, the optimized solar flux is enhanced at different flow rates, and the increase ratio increases with increasing flow rate.

Overall, the aforementioned results indicate that the form of optimal solar flux distribution with spatial energy supply and demand matching can improve SPTR performance regardless of the relationship between solar input and total reactants, and thus the present study is considered valid and general.





4 CONCLUSION

The primary purpose of this study is to propose an inverse design method for solar flux distribution to achieve optimal SPTR performance enhancement. First, a 2-D axisymmetric model of SPTR was coupled with an optimized solar flux distribution model, which applied the SNOPT algorithm to optimize the axial solar flux distribution described by Bezier curves to achieve better SPTR performance. Further, the energy matching and performance enhancement mechanism of SPTR were analyzed. Then, the continuous solar flux was regressed to a step-like flux with maintained performance to provide design guidance for the concentrator. Finally, the performance enhancement of the step-like flux distribution form with different solar energy input and different reactant flow rates was investigated. The particular conclusions were as follows:

(1) Optimized linear and non-linear solar flux distributions achieve higher and larger fast reaction zone compared to

uniform flux, resulting in significantly improved SPTR performance. Especially at non-uniform flux, the methanol conversion, solar thermochemical conversion, and hydrogen yield are, respectively, enhanced by 2.5, 3.3, and 2.4% relative to the uniform flux. Furthermore, the optimized non-uniform flux enhances the synergy between the temperature field and the reaction field in SPTR, ultimately achieving a spatial correlation of 0.99 between solar energy supply and reaction demand, resulting in energy matching.

(2) To effectively guide the design of the concentrating system, the step-like solar flux distribution based on the gradient-based segmentation method is obtained. The simulation results between the step-like distribution and the continuous distribution have a high correlation of 0.99 and only about 0.2% performance loss, which maximally maintains the performance-enhancing capability of the original distribution shape. The feasibility of the segmentation method and design has also been verified.

(3) Optimized step-like solar flux distribution forms are always effective and suitable for improving SPTR performance under different solar inputs versus total reactants (e.g., different solar flux inputs and reactant mass flow rates). The performance improvement is especially significant when the total solar energy supply is insufficient.

Thus, the method of this study is instructive and informative for the optimal solar flux distribution and concentrator design that enhance the performance of concentrated solar energy utilization systems. Furthermore, in the future, we will consider experiments to further refine and validate the design method of this study.

DATA AVAILABILITY STATEMENT

The raw data supporting the conclusion of this article will be made available by the authors, without undue reservation.

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AUTHOR CONTRIBUTIONS

XT: Conceptualization, methodology, validation, software, writing—original draft, formal analysis, writing—review and editing, and visualization. WY: Conceptualization, formal analysis, funding acquisition, writing—review and editing, and supervision. ZD: Investigation and writing—review and editing. YY: Visualization and writing—original draft.

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SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fenrg.2022.881822/full#supplementary-material>

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NOMENCLATURE

Abbreviations

DNI direct normal irradiance
MSRR methanol steam reforming reaction
SNOPT Sparse Nonlinear Optimizer
SPTR solar trough thermochemical reactor
WFR ratio of catalyst weight to methanol inlet molar flow rate

Nomenclature

c concentration ($\text{mol}\cdot\text{m}^{-3}$)
 C_p specific heat ($\text{J kg}^{-1} \text{K}^{-1}$)
 d_p mean catalyst particle size (m)
 D_{ai} inner diameter of absorber tube (m)
 D_{ao} outer diameter of absorber tube (m)
 D_{gi} inner diameter of glass cover (m)
 D_{go} outer diameter of glass cover (m)
 D_i mass diffusion coefficient of species i ($\text{m}^2 \text{s}^{-1}$)
 h heat transfer coefficient ($\text{W m}^{-2} \text{K}^{-1}$)
 h_i molar enthalpy of species i (J mol^{-1})
 ΔH_r enthalpy change of the r th reaction (J mol^{-1})
 L length of reactor (m)
 m_i mass flow rate of species i (kg h^{-1})
 M molar weight (kg mol^{-1})
 $n_{i,r}$ stoichiometric number of species i in the r th reaction
 p pressure (Pa)
 p_i partial pressure of species i (bar)
 q energy flux (W m^{-2})
 Q power (W)

R chemical reaction rate ($\text{mol m}^{-3} \text{s}^{-1}$)
 R_g universal gas constant ($\text{J mol}^{-1} \text{K}^{-1}$)
 R_r chemical reaction rate of the r th reaction ($\text{mol m}^{-3} \text{s}^{-1}$)
 S_{chem} chemical reaction source (W m^{-3})
 T temperature (K)
 \vec{U} velocity vector (m s^{-1})
 r, z axisymmetric coordinates (m)

Greek symbols

ϵ_e emissivity of the diffuse surface
 ϵ_p porosity of catalytic bed
 η_m methanol conversion efficiency
 η_s solar thermochemical conversion efficiency
 λ thermal conductivity ($\text{W m}^{-1} \text{K}^{-1}$)
 μ dynamic viscosity (Pa s)
 ρ density (kg m^{-3})
 σ Stefan–Boltzmann constant ($\text{W m}^{-2} \text{K}^{-4}$)

Subscripts

amb ambient
ave average value
eff effective value
 i species ($\text{CH}_3\text{OH}, \text{H}_2\text{O}, \text{H}_2, \text{CO}, \text{CO}_2$)
in inlet
m mixed reaction gas
max maximum value
out outlet
p particle catalytic bed